



$$\sin^2 \theta_W = 0.238$$

Precision measurement of the Weinberg angle

 θ_W

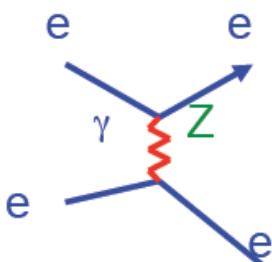
Frank Maas

(Helmholtz Institute Mainz,

Johannes Gutenberg University Mainz)

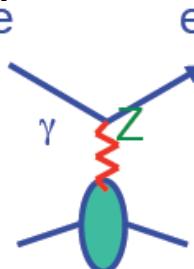
“The Future of Non-Collider-Physics”, Mainz, April 27./28., 2017

Møller Scattering



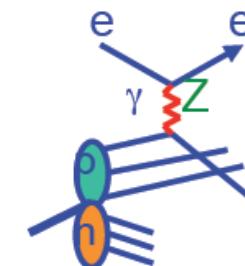
- Purely Leptonic

Q-Weak (JLab) P2 (Mainz/MESA)



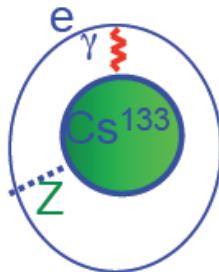
- Coherent quarks in p
- in operation now
- $2(2C_{1u} + C_{1d})$

e-DIS



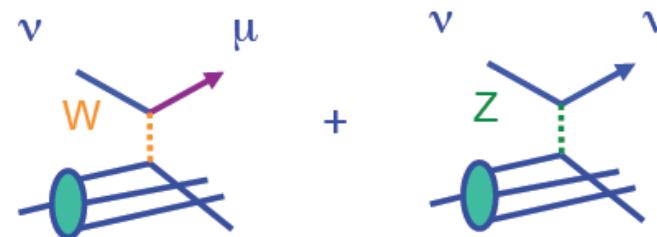
- Isoscaler quark scattering
- $(2C_{1u} - C_{1d}) + Y(2C_{2u} - C_{2d})$

Atomic Parity Violation



- Coherent quarks in entire nucleus
- Nuclear structure uncertainties
- $-376 C_{1u} - 422 C_{1d}$

Neutrino Scattering



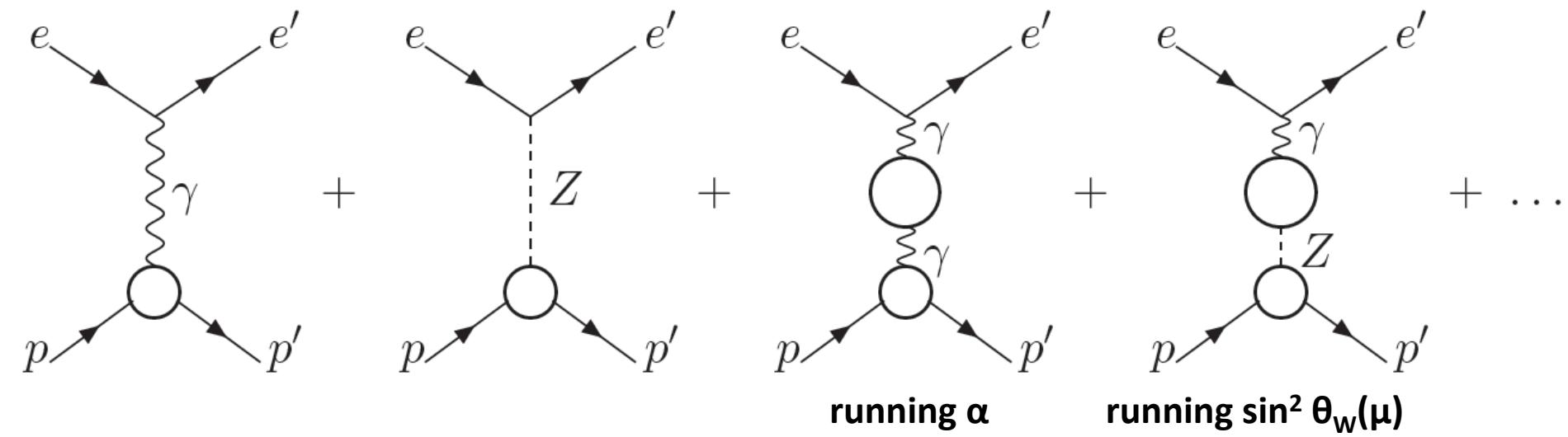
- Quark scattering (from nucleus)
- Weak charged and neutral current difference



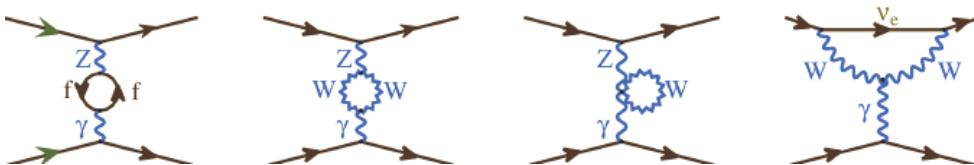
„running“ $\sin^2 \theta_{\text{eff}}$ or $\sin^2 \theta_W(\mu)$

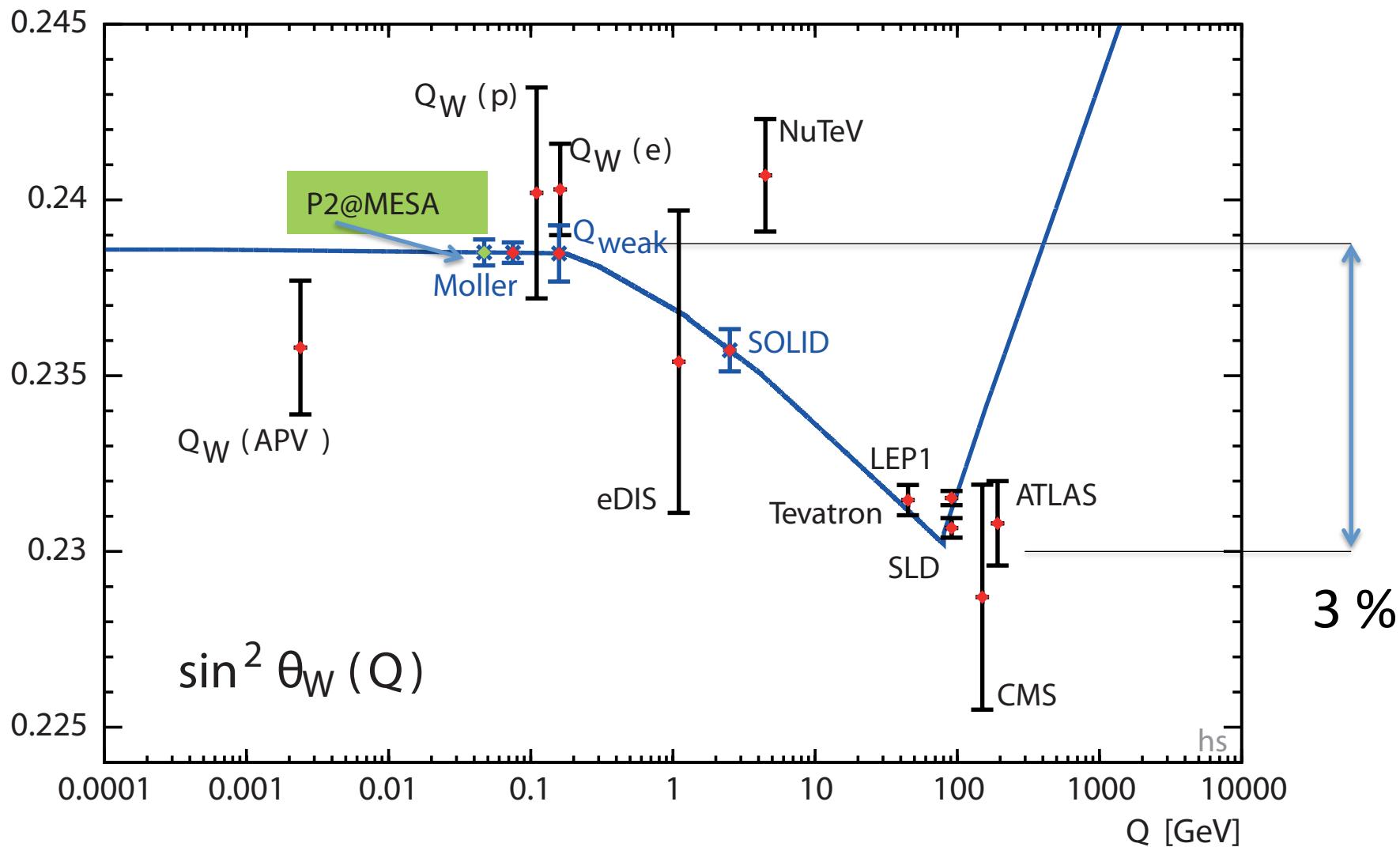


Precision measurements and quantum corrections:



Universal quantum corrections: can be absorbed into a
scale dependent, „running“ $\sin^2 \theta_{\text{eff}}$ or $\sin^2 \theta_w(\mu)$



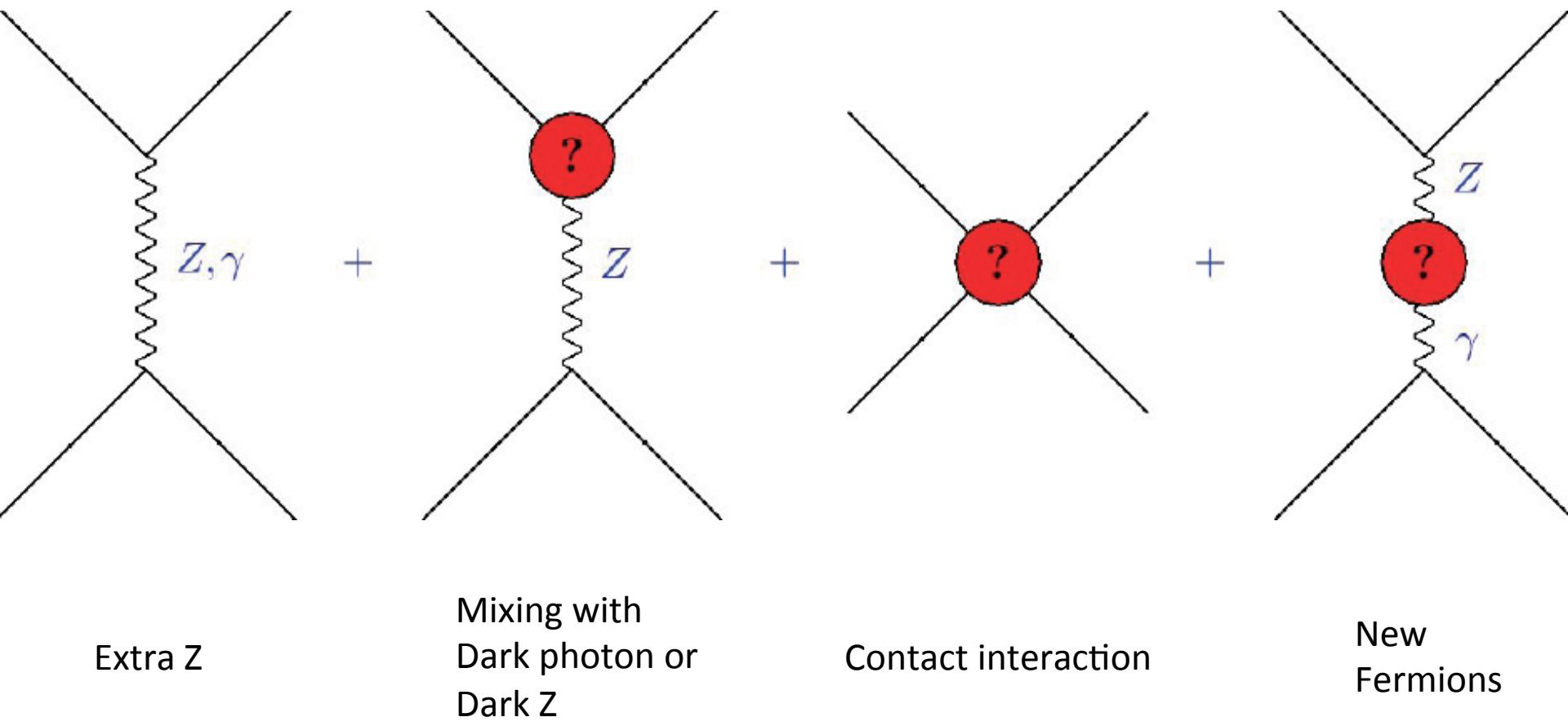




Sensitivity to new physics beyond the Standard Model



Sensitivity to new physics beyond the Standard Model

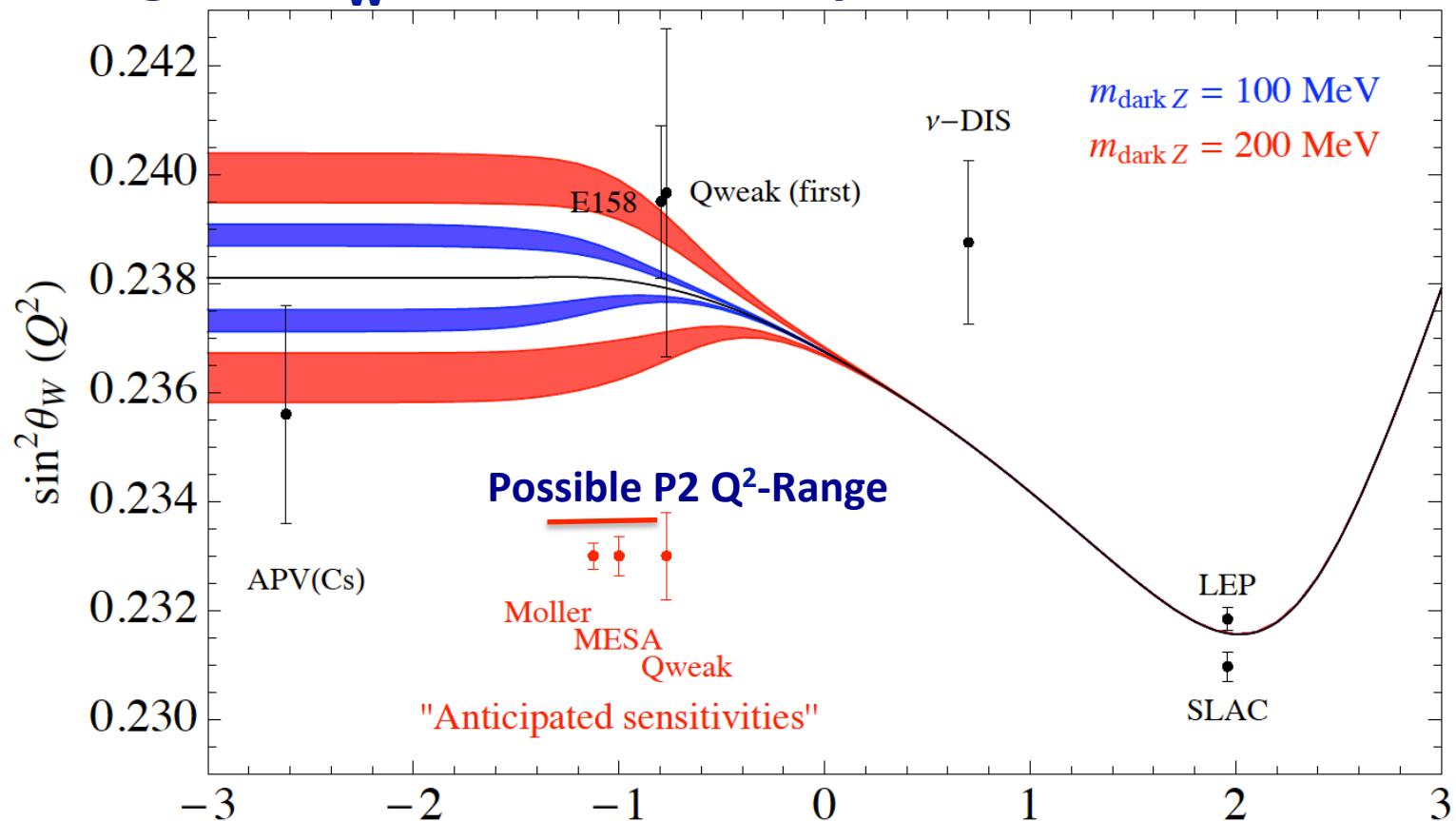




Dark Photon, Z-Boson



Running $\sin^2 \theta_W$ and Dark Parity Violation



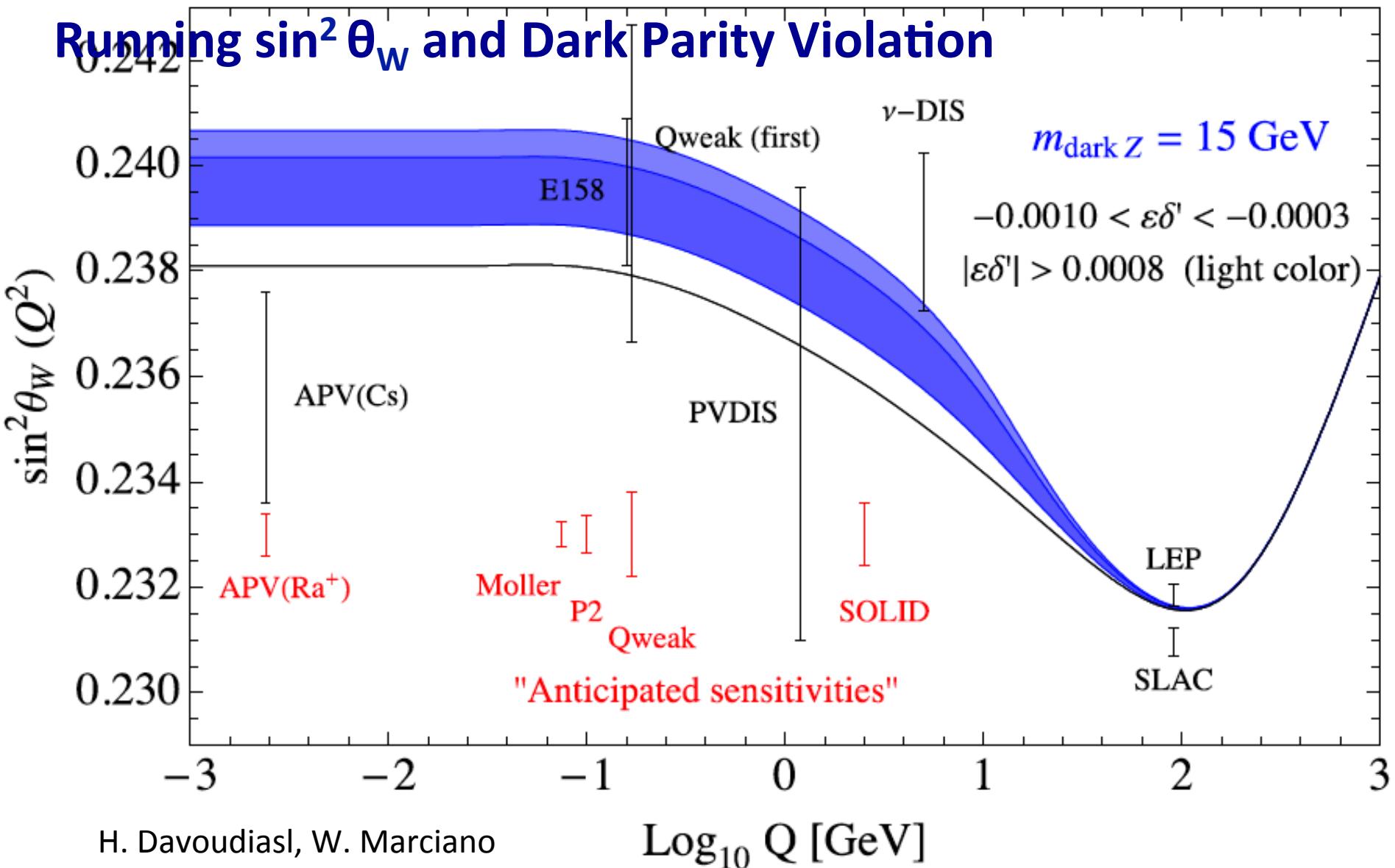
$$Z = \cos \theta_W W_3 - \sin \theta_W B$$
$$A = \sin \theta_W W_3 + \cos \theta_W B$$

 $\log_{10} Q [\text{GeV}]$

Bill Marciano

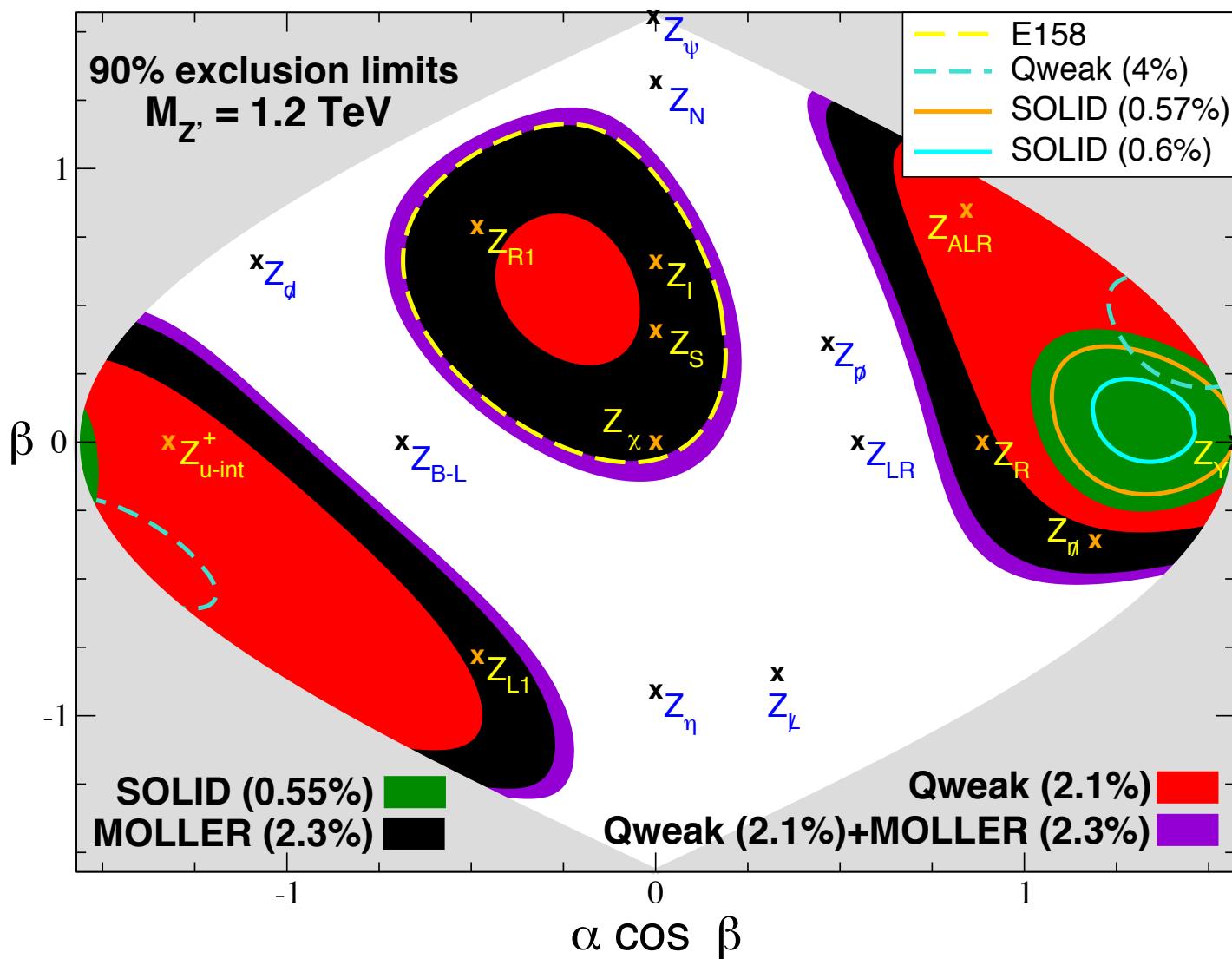


Running $\sin^2 \theta_W$ and Dark Parity Violation





Extra Z-Boson



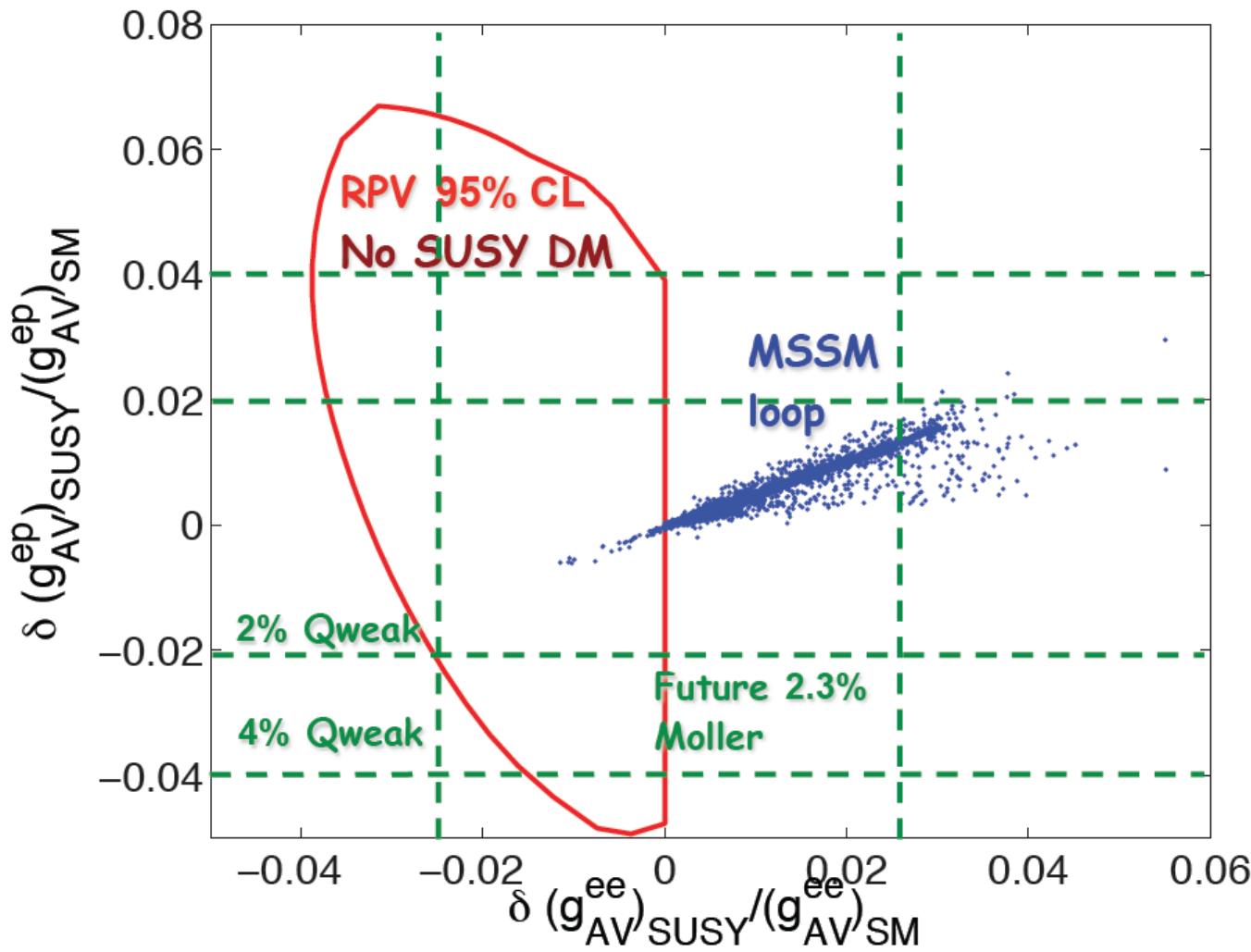


Supersymmetry



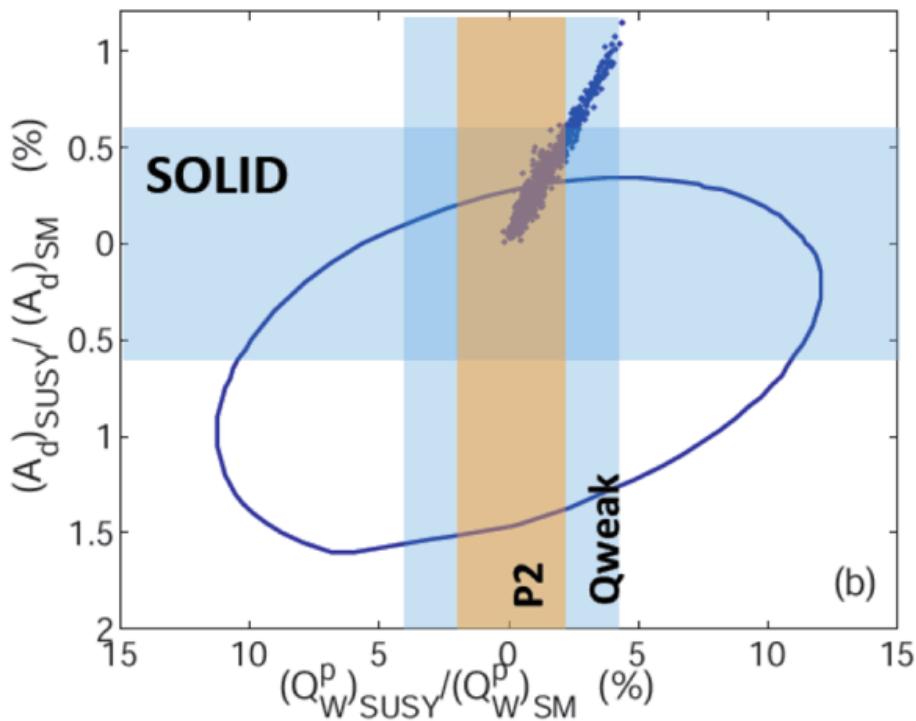
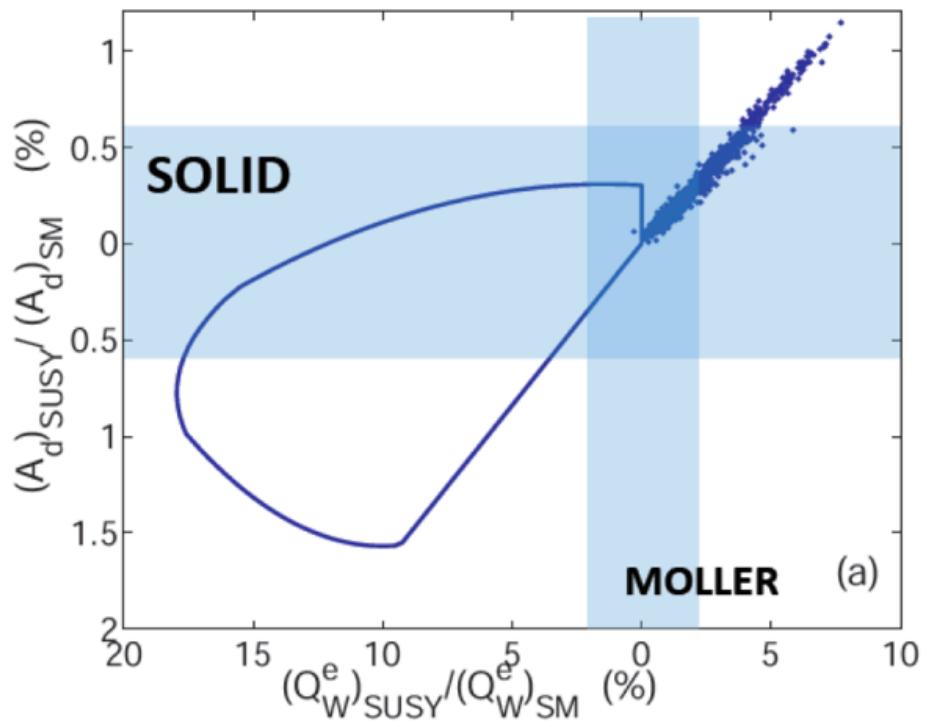
Example: Supersymmetric standard model extensions

Kurylov, Ramsey-Musolf, Su (2003), updated





Ramsey-Musolf and Su, *Phys. Rep.* 456 (2008)





- Complementary access by weak charges of proton and electron

Weak charge of the proton:

$$Q_W^p = 0.0716$$

A horizontal line with a central dot and two vertical error bars extending to the left and right, labeled ± 0.0029 .

Weak charge of the electron:

$$Q_W^e = -0.0449$$

A horizontal line with a central dot and two vertical error bars extending to the left and right, labeled ± 0.0051 .

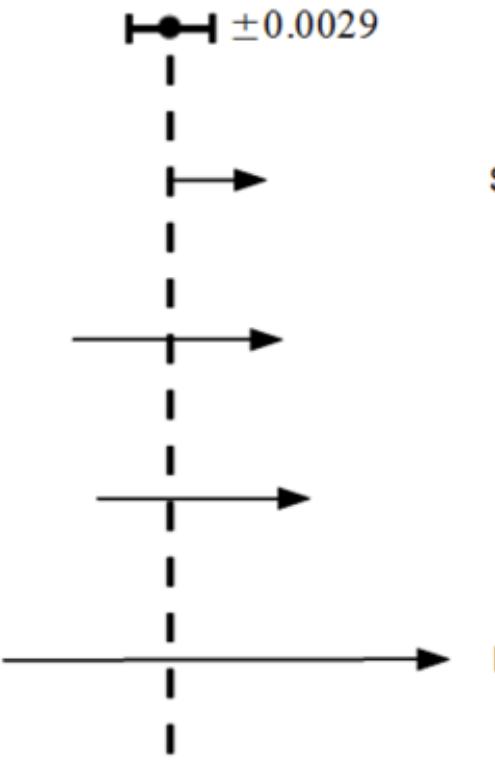
Experiment

SUSY-Loops

$E_6 Z'$

RPV SUSY

Leptoquarks



SM

(Jens Erler, Ramsey-Musolf, 2003)

SM



Weak
Charge
Of
Proton:
**Qweak (Jlab),
P2 (MESA)**

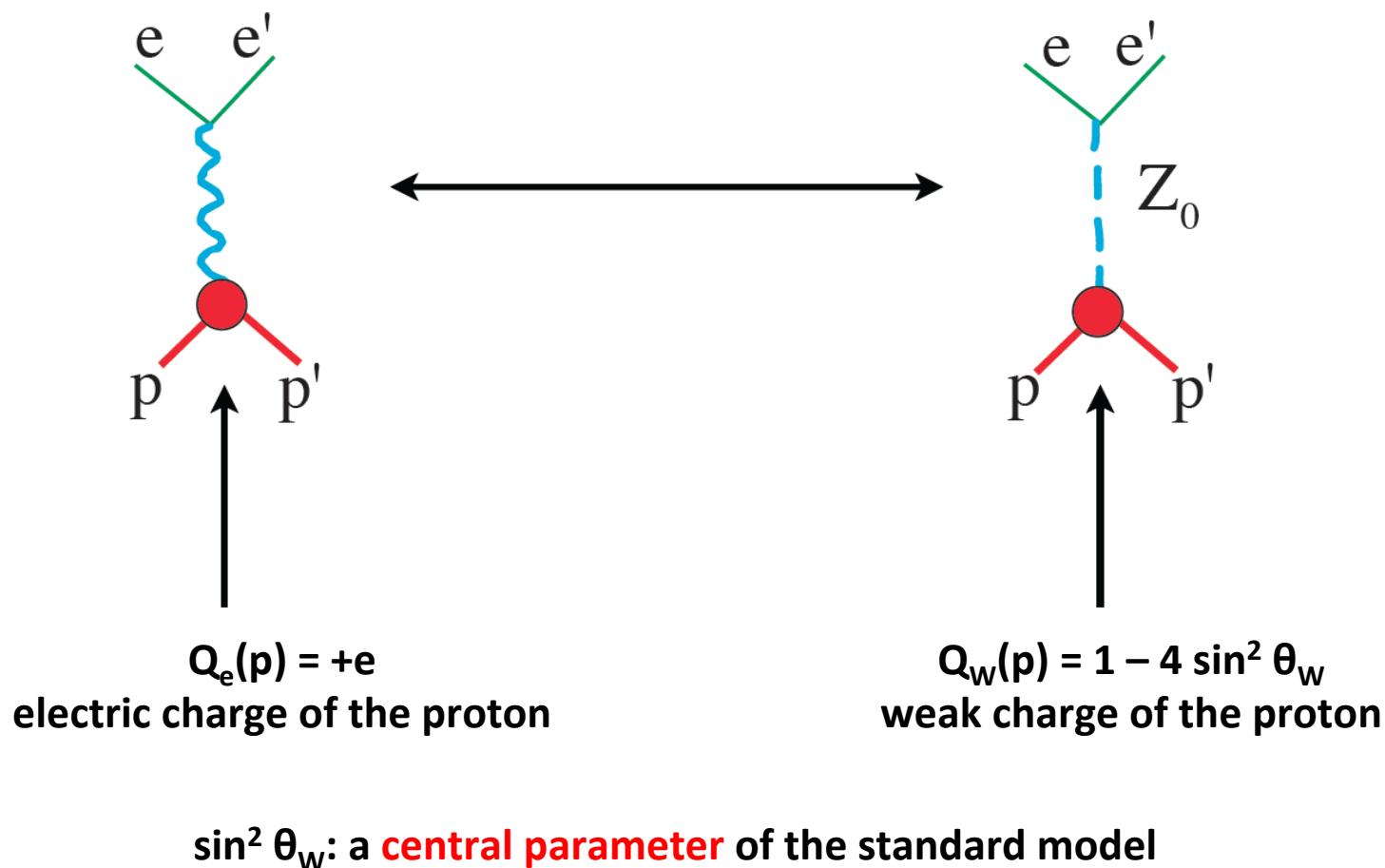
Weak
Charge
Of
Electron:
**MOELLER
(JLAB)**

Weak
Charge
Of
Quarks:
**SOLID (PVDIS)
(JLAB)**



The role of the weak mixing angle

The **relative strength** between the weak and electromagnetic interaction is determined by the **weak mixing angle**: $\sin^2(\theta_w)$





Proton: special case

$$\text{Proton Weak charge: } Q_w(p) = 1 - 4 \sin^2 \theta_w$$

$$\text{Error: } \Delta Q_w(p) = 4 \Delta \sin^2 \theta_w$$

$$\text{Rel. error: } \Delta Q_w(p)/Q_w(p) = 4/((1/\sin^2 \theta_w) - 4) \quad (\Delta \sin^2 \theta_w / \sin^2 \theta_w)$$

$$\text{Rel. error } \Delta \sin^2 \theta_w / \sin^2 \theta_w = ((1/\sin^2 \theta_w) - 4) / 4 \quad \Delta Q_w(p)/Q_w(p)$$

$$\text{Example: } \sin^2 \theta_w (50 \text{ MeV}) = 0.238$$

$$4/((1/\sin^2 \theta_w) - 4) \sim 20$$

$$\Delta Q_w(p)/Q_w(p) = 2\% \quad \text{from Experiment}$$

$$\Delta \sin^2 \theta_w / \sin^2 \theta_w = 0.1 \% \quad \text{same precision as LEP, SLAC}$$

Neutron Weak charge:

$$\Delta Q_w(p)/Q_w(n) = \Delta \sin^2 \theta_w / \sin^2 \theta_w$$



Physics sensitivity from contact interaction (LEP2 convention, $g^2 = 4\pi$)

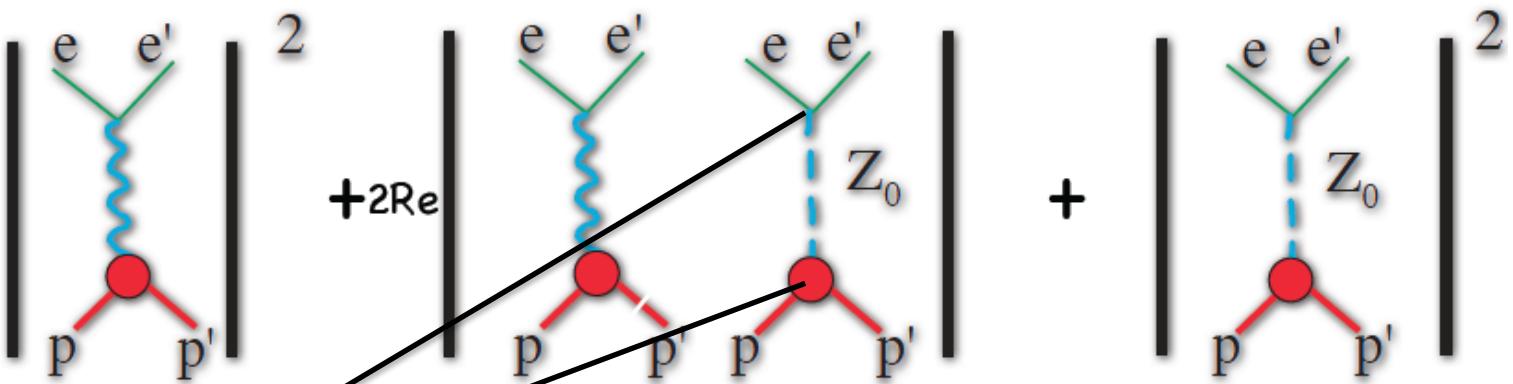
	precision	$\Delta \sin^2 \bar{\theta}_W(0)$	$\Lambda_{\text{new}} \text{ (expected)}$
APV Cs	0.58 %	0.0019	32.3 TeV
E158	14 %	0.0013	17.0 TeV
Qweak I	19 %	0.0030	17.0 TeV
Qweak final	4.5 %	0.0008	33 TeV
PVDIS	4.5 %	0.0050	7.6 TeV
SoLID	0.6 %	0.00057	22 TeV
MOLLER	2.3 %	0.00026	39 TeV
P2	2.0 %	0.00036	49 TeV
PVES ^{12}C	0.3 %	0.0007	49 TeV



Experimental Method: Parity Violating Electron Scattering



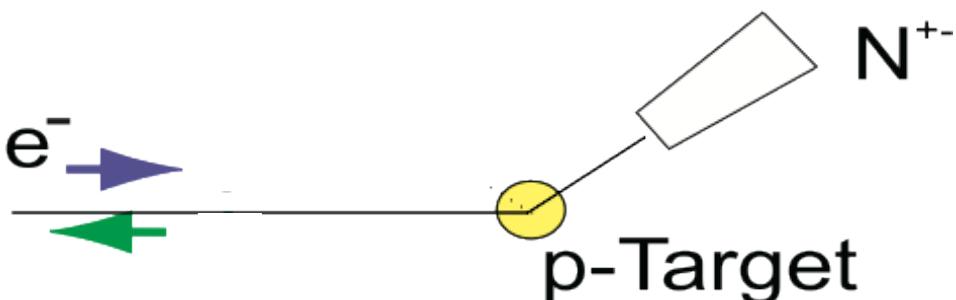
Parity Violating Asymmetry in elastic electron proton scattering

 $\sigma \approx$ 

$$(V-A)_e(V-A)_p$$

$$A_e V_p + V_e A_p$$

V-A coupling:
parity-violating
cross section asymmetry A_{LR}
longitudinally pol. electrons
unpolarised protons





Parity violating cross section asymmetry

$$A_{LR} = \frac{\sigma(e \uparrow) - \sigma(e \downarrow)}{\sigma(e \uparrow) + \sigma(e \downarrow)} = -\frac{G_F Q^2}{4\sqrt{2}\pi\alpha} (Q_W - F(Q^2))$$

$$Q_W = 1 - 4 \sin^2 \theta_W(\mu)$$

polarisation measurement

$$F(Q^2) = F_{EM}(Q^2) + F_{Axial}(Q^2) + F_{Strange}(Q^2)$$

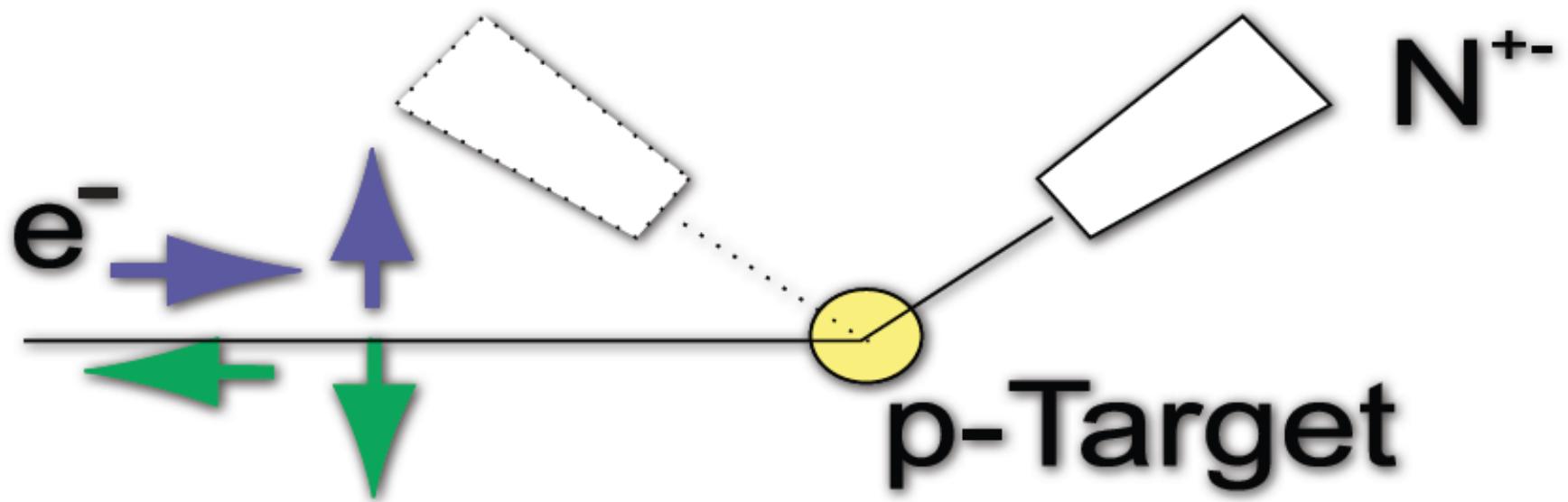
tracking system

weak charge

hadron structure



Conceptually very simple experiments



$$A = (N^+ - N^-) / (N^+ + N^-) \quad \Delta A = (N^+ + N^-)^{-1/2} = N^{-1/2}$$

$$A = 20 \times 10^{-9} \quad 2\% \text{ Measurement} \quad N = 6.25 \times 10^{18} \text{ events}$$

Highest rate, measure Q^2 : Large Solid Angle Spectrometers

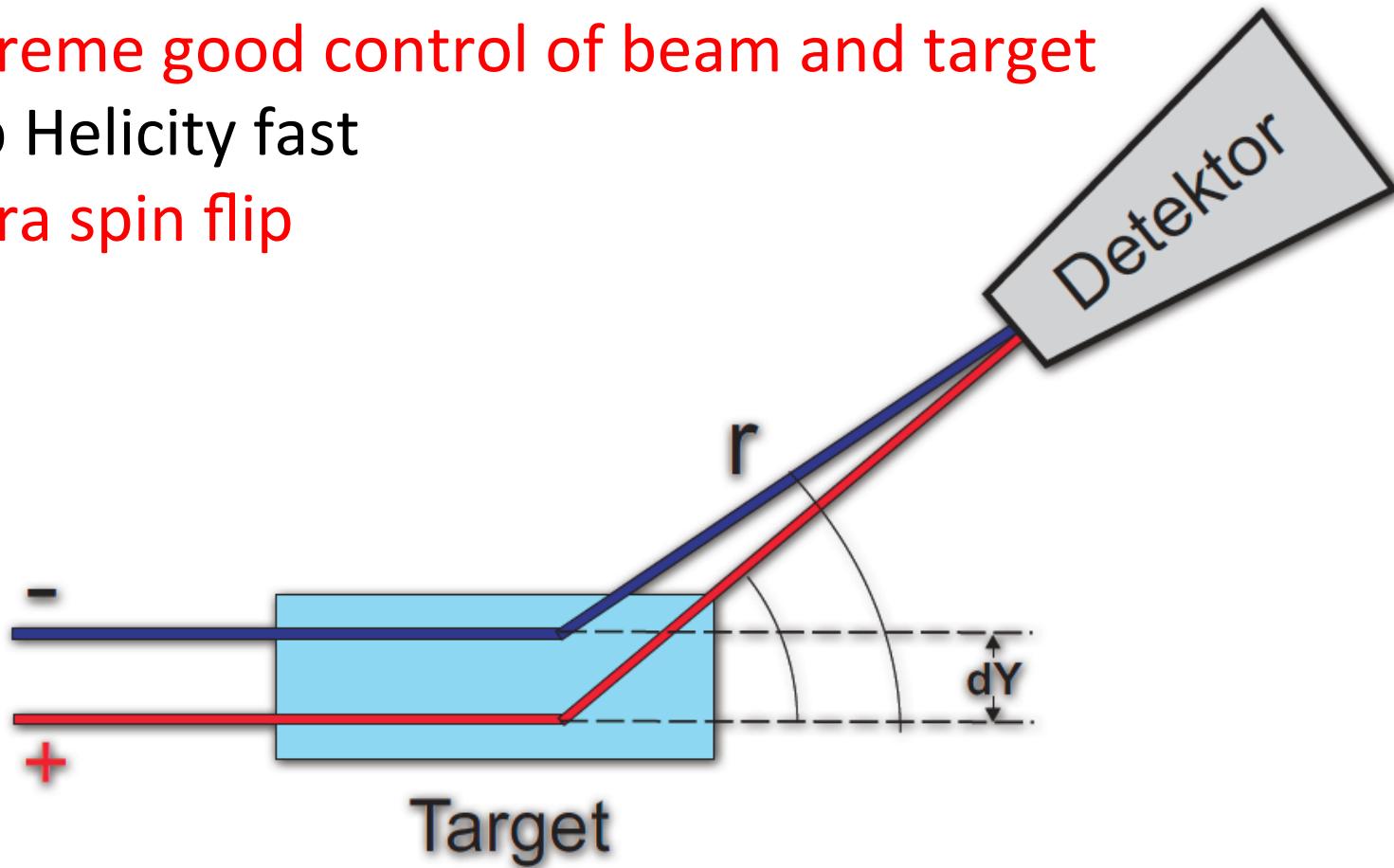


Apparative (false) asymmetries:

Extreme good control of beam and target

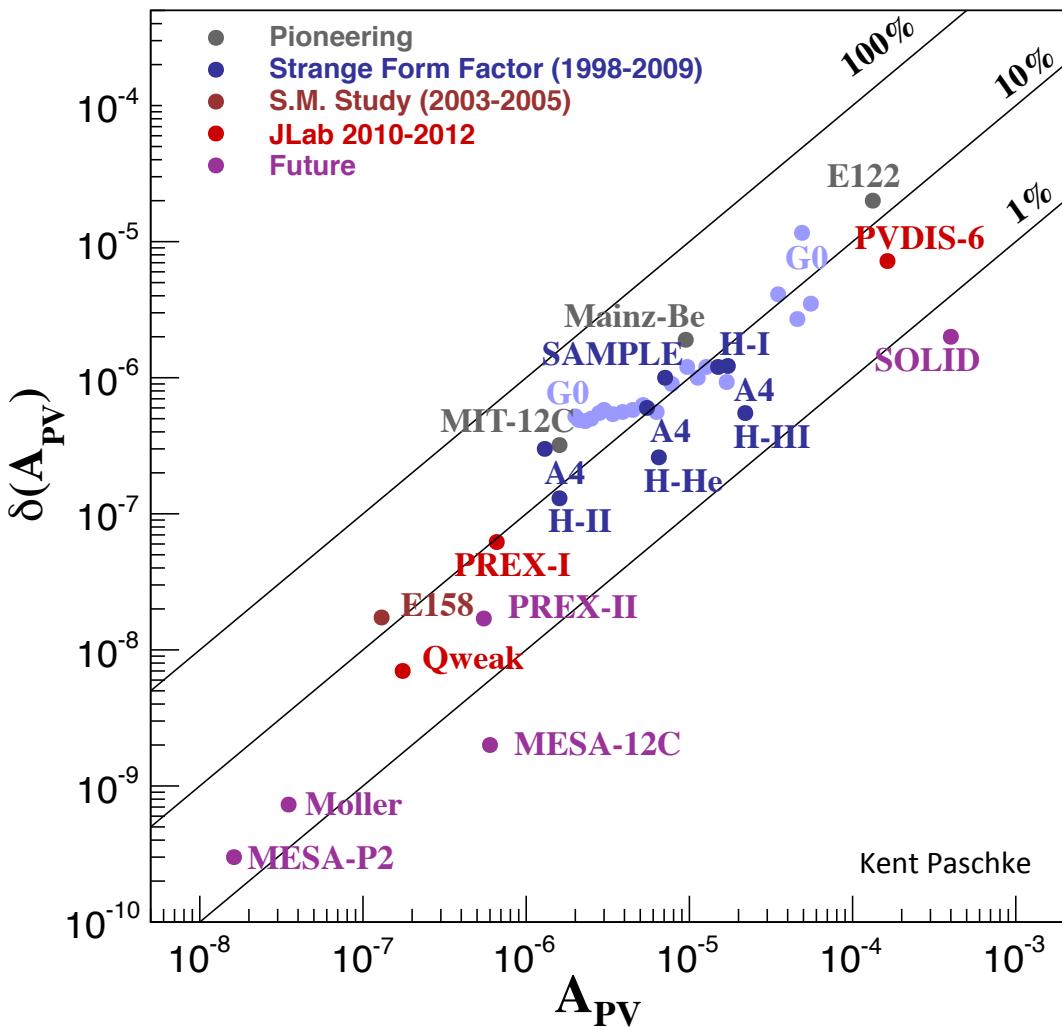
Flip Helicity fast

Extra spin flip



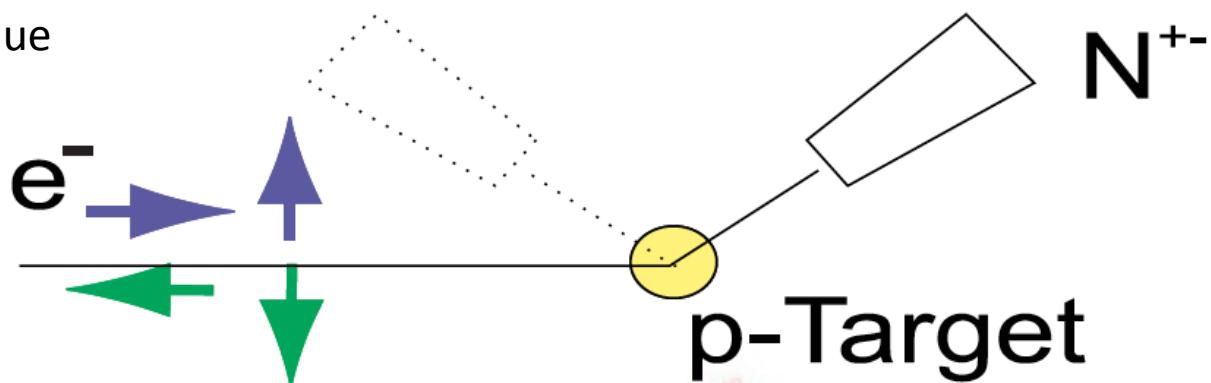


PVeS Experiment Summary





Counting Technique

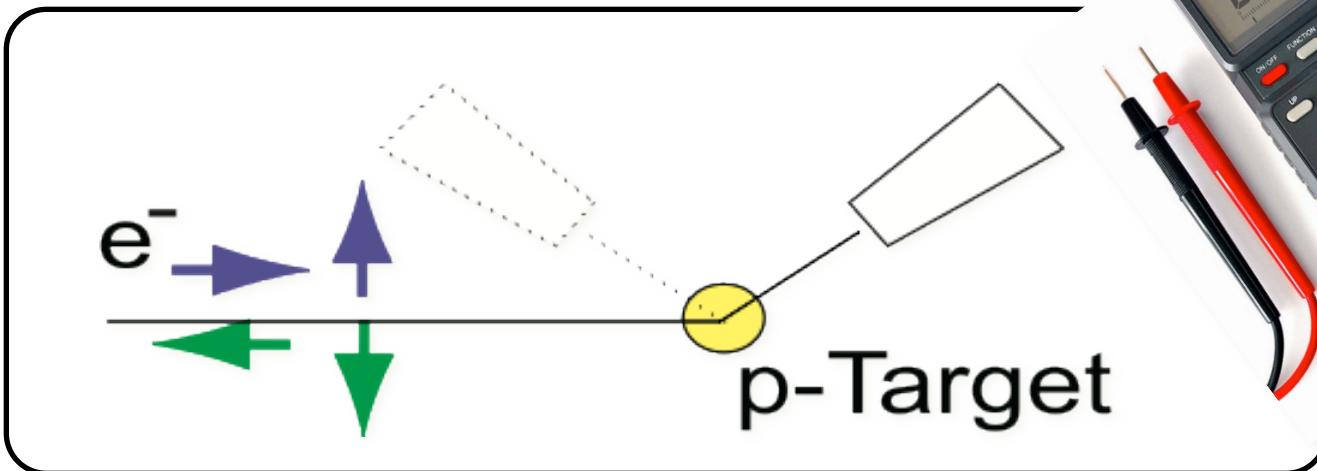


Count scattered electrons:

- pile-up (double count losses)
- Background Asymmetry
- Very Fast Counting (MHz)
- Measure TOF or Energy



Analogue Technique



Measure Flux of Scattered electrons:

- no pile-up (double count losses)
- sensitive to small electr. fields.
- no separation of phys. process



- Collaboration was enlarged during the past years
- Integrating detectors, magnetic spectrometer, theory, electronics, target, polarimetry

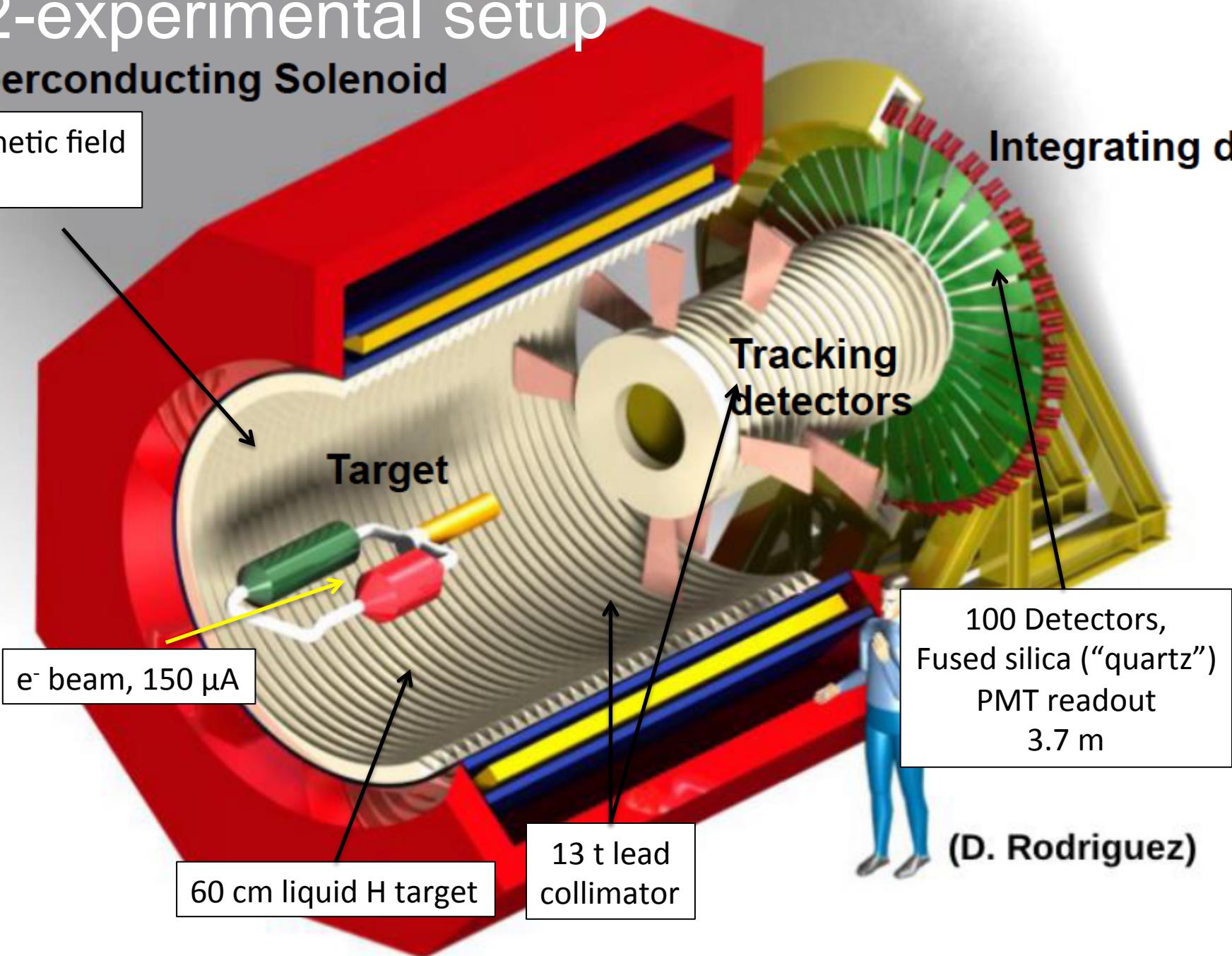
Hubert Spiesberger	Matthias Molitor
Razvan Bucoveanu	Niklaus Berger
Stephan Wezorke	Iurii Sorokin
Frank Maas	Alexey Tyukin
Sebastian Baunack	Marco Zimmermann
Dominik Becker	Silviu Covrig
Kathrin Gerz	Sandesh Gopinath
Thomas Jennewein	Krishna Kumar
Kurt Aulenbacher	Paul Souder
Dr. Valery Tioukine	Michael Gericke

Institut für Kernphysik, Universität Mainz
Institut für Physik, Universität Mainz
Department of Physics and Astronomy,
Stony Brook, NY
Jefferson Lab, Newport News, VA
University of Manitoba, Winnipeg
Syracuse University, Syracuse, NY

P2-experimental setup

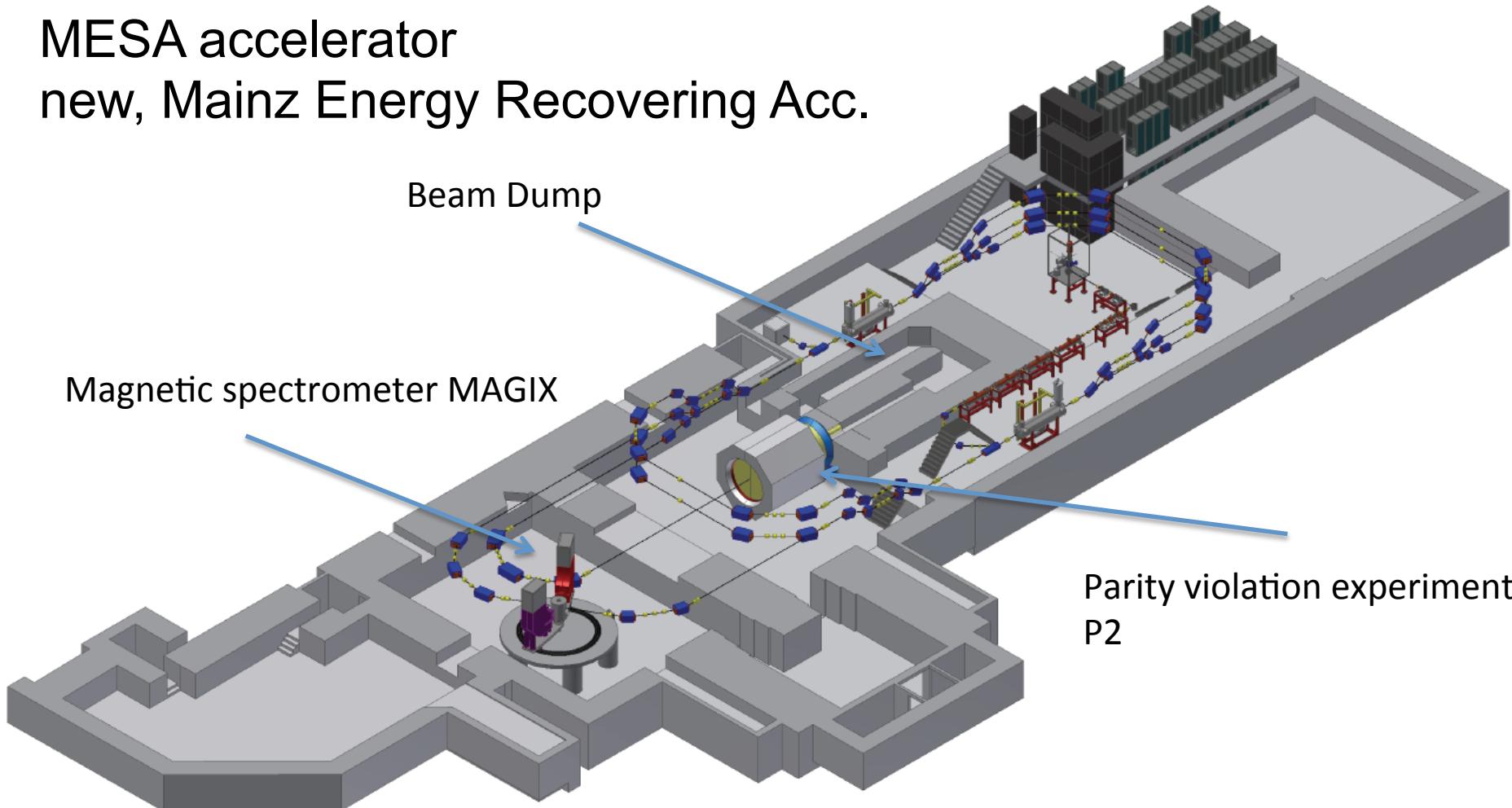
Superconducting Solenoid

Magnetic field
0.6 T





MESA accelerator new, Mainz Energy Recovering Acc.



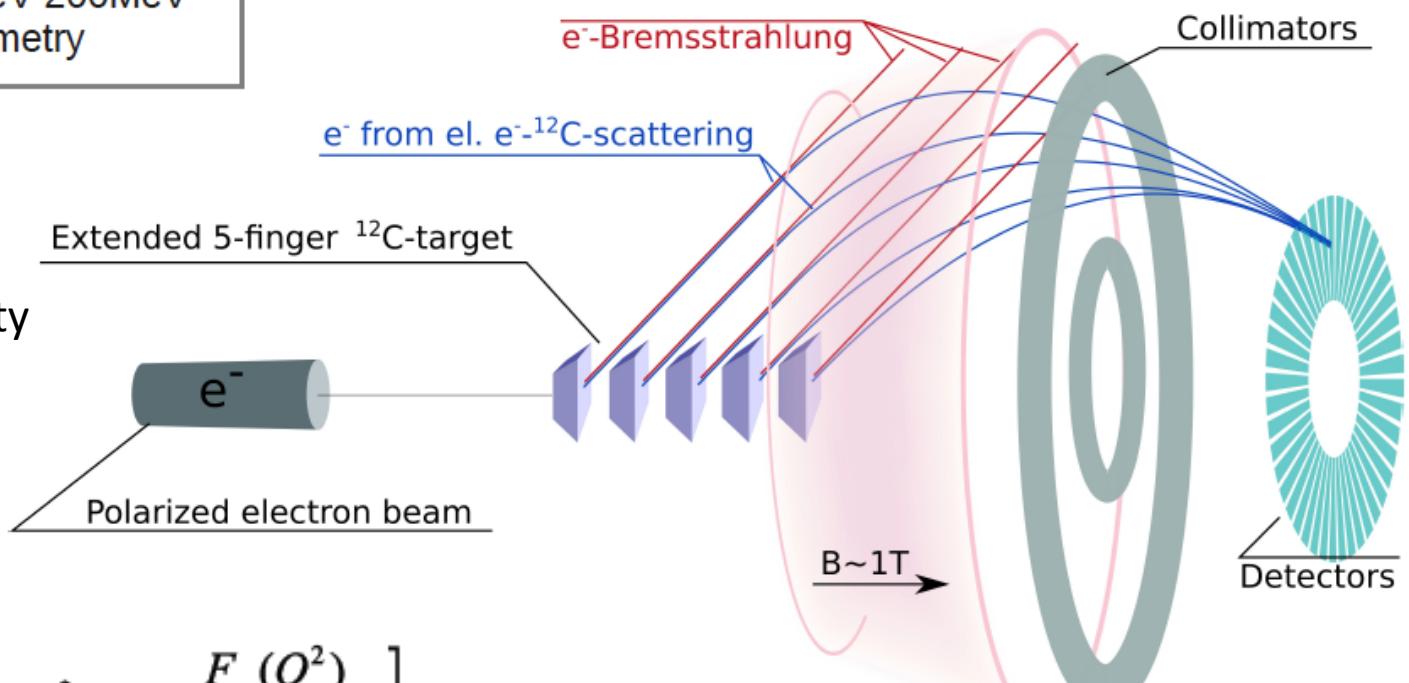


Other Measurements:
Carbon, Lead

EXPERIMENTAL REALIZATION

- MESA:
- 150 μ A
 - 150MeV-200MeV
 - Polarimetry

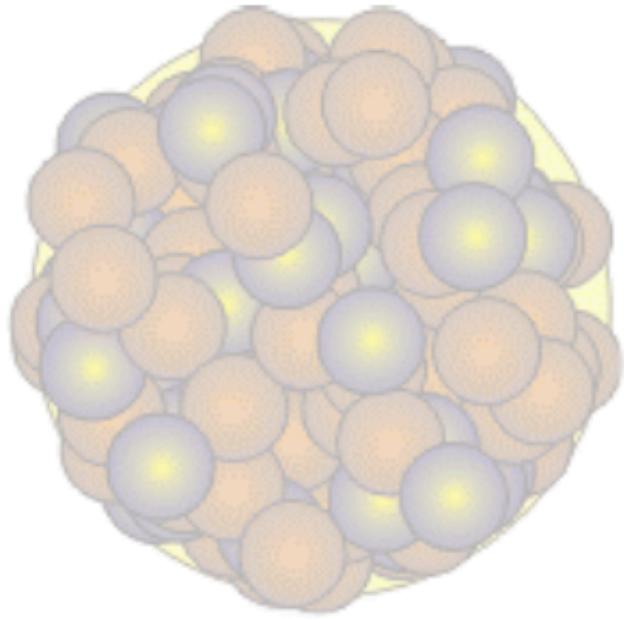
Enhanced sensitivity
To new physics



$$-N = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \left[\underbrace{\sin^2 \theta_W}_{\approx 0} - \frac{F_n(Q^2)}{F_p(Q^2)} \right]$$

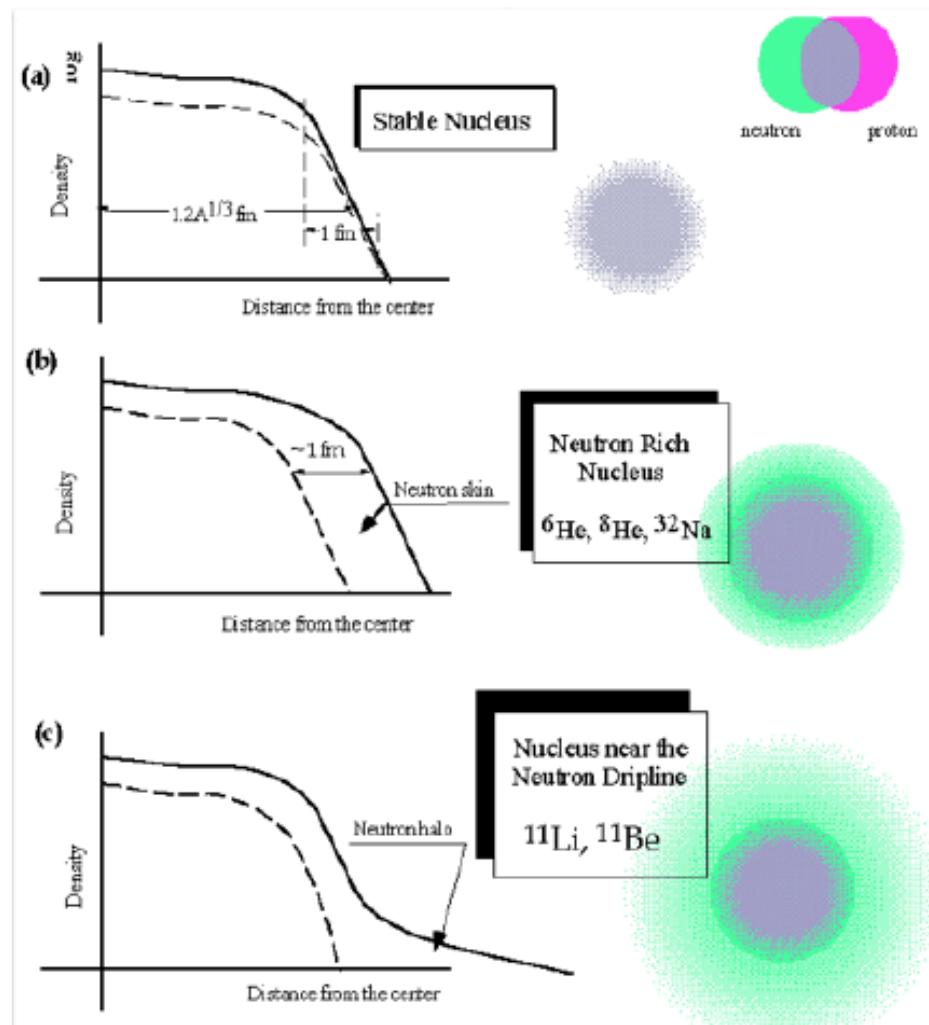
Neutron Skin for beginner

Where do the neutrons go?



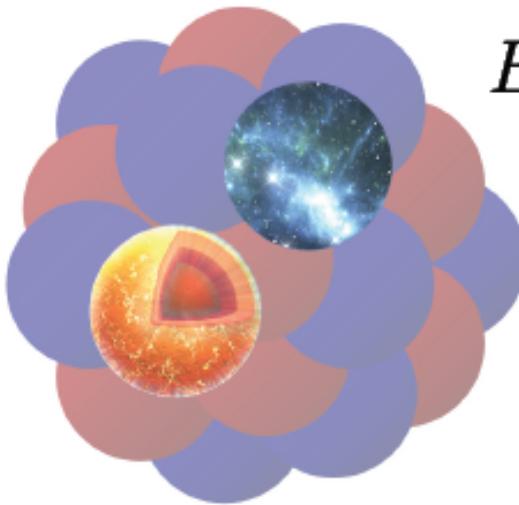
Pressure forces neutrons
out against surface tension

→ EOS



Measurement of neutron distribution in nuclei decisive for Neutron star properties

WHY?



$$E(\rho, \delta) = E(\rho, 0) + E_{sym}(\rho) \delta^2 + \mathcal{O}(\delta)^4$$

↓
symmetry energy

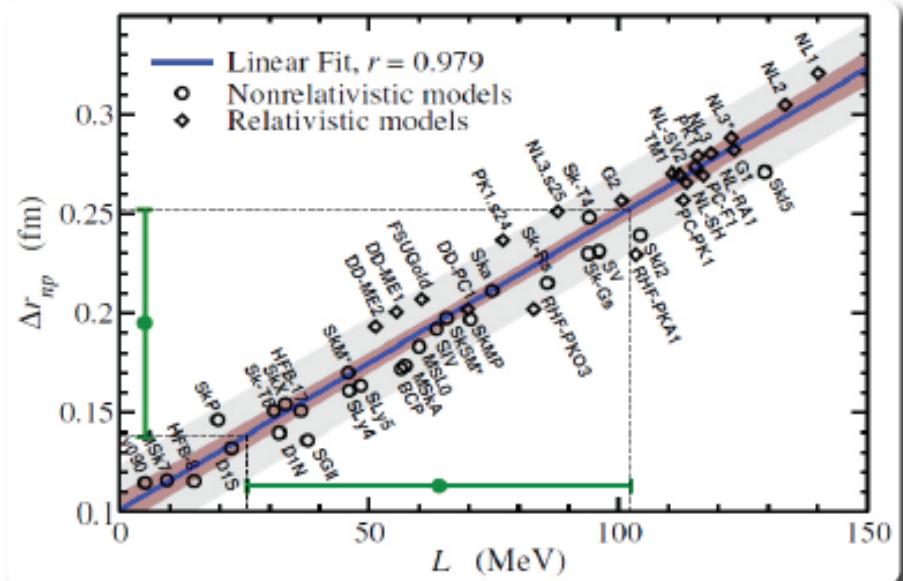
$$E_{sym}(\rho) = \left[S_v + \frac{L}{3} \left(\frac{\rho - \rho_0}{\rho_0} \right) + \frac{K_{sym}}{18} \left(\frac{\rho - \rho_0}{\rho_0} \right)^2 \right] + \dots$$

slope parameter

$$L = 3\rho_0 \frac{\partial E_{sym}(\rho)}{\partial \rho} \Bigg|_{\rho_0}$$

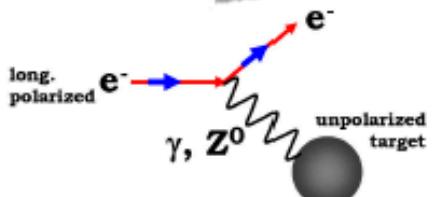
curvature parameter

$$K_{sym} = 9\rho_0^2 \frac{\partial^2 E_{sym}(\rho)}{\partial \rho^2} \Bigg|_{\rho_0}$$



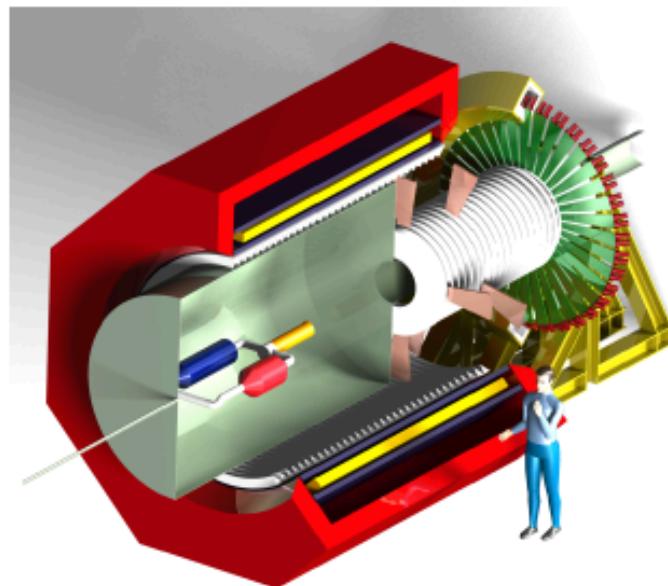
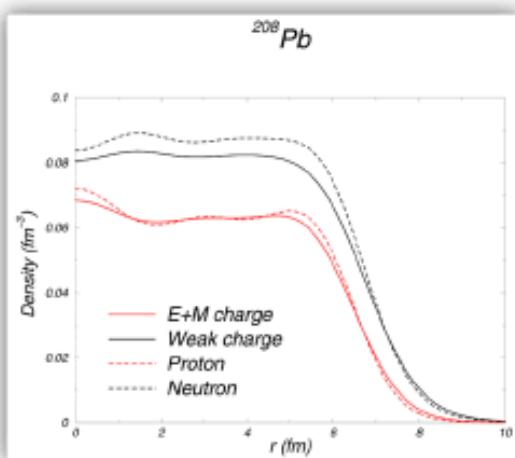
PVES

Enormously Clean ...
Extraordinarily Expensive!

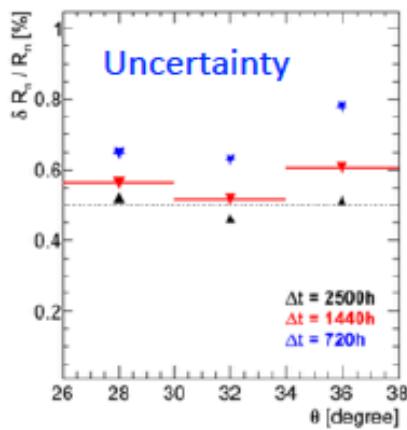
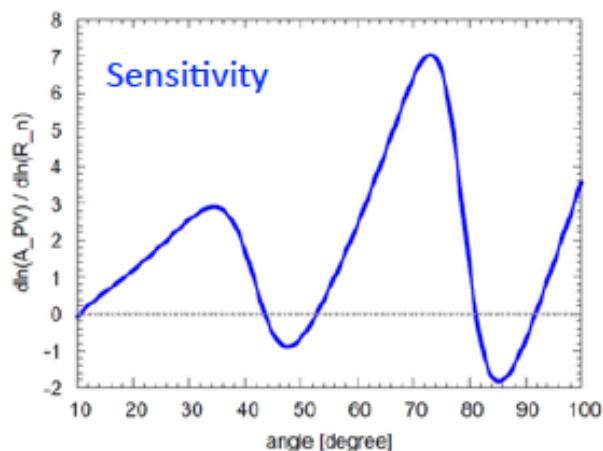


$$A_{PV} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \left[1 - 4\sin^2\theta_W - \frac{F_n(Q^2)}{F_p(Q^2)} \right] \approx 0$$

$$F_{n,p}(Q^2) = \frac{1}{4\pi} \int d^3r \ j_0(qr) \rho_{n,p}(r)$$

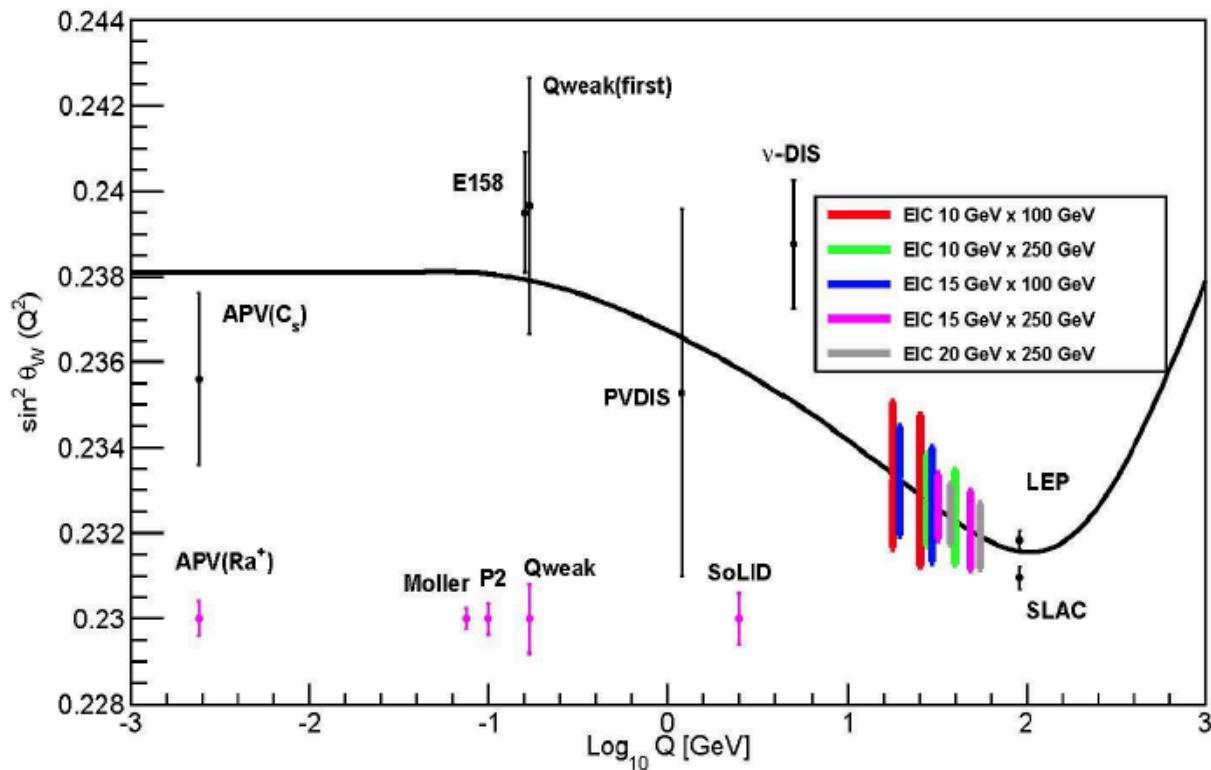


Full azimuthal coverage $\Leftrightarrow 4\times$ stat



Assuming same PREX luminosity:
 $\Delta\theta=4^\circ$: Rate=8.25 GHz, $A_{PV}=0.66$ ppm
1440h $\rightarrow \delta R_n / R_n = 0.5\%$
 (assuming 1% syst. $\delta A_{PV} / A_{PV}$)

World data of $\sin^2 \Theta_w$ including EIC projections



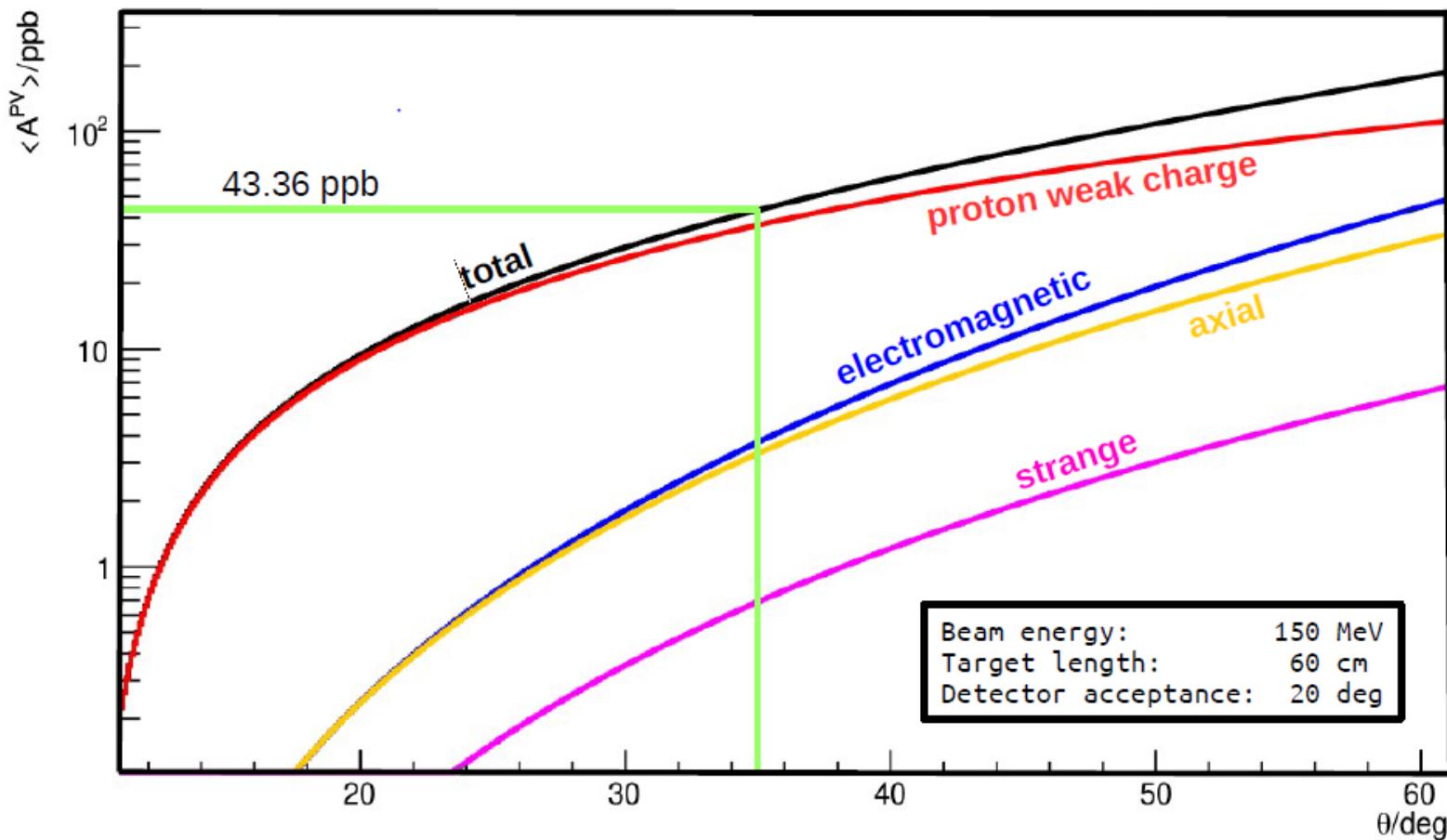
- 200 days of dedicated run
- Can reach similar precision to SoLID measurement
- Interesting Q^2 region never been measured or planned



- Parity violating electron scattering: “Low energy frontier” comprises a sensitive test of the standard model complementary to LHC
- Determination of $\sin^2(\theta_W)$ with high precision (same as Z-pole)
- P2-Experiment (proton weak charge) in Mainz under preparation
New MESA energy recovering accelerator at 155 MeV, target precision is 1.7% in Qweak i.e. 0.13% in $\sin^2(\theta_W)$, Sensitivity to new physics up to a scale of 49 TeV
- Much more physics from PV electron scattering
- Together with Moeller@Jlab (electron weak charge) and SOLID@Jlab (quark weak charge) very sensitive test of standard model and possibility to narrow in on Standard Model Extension



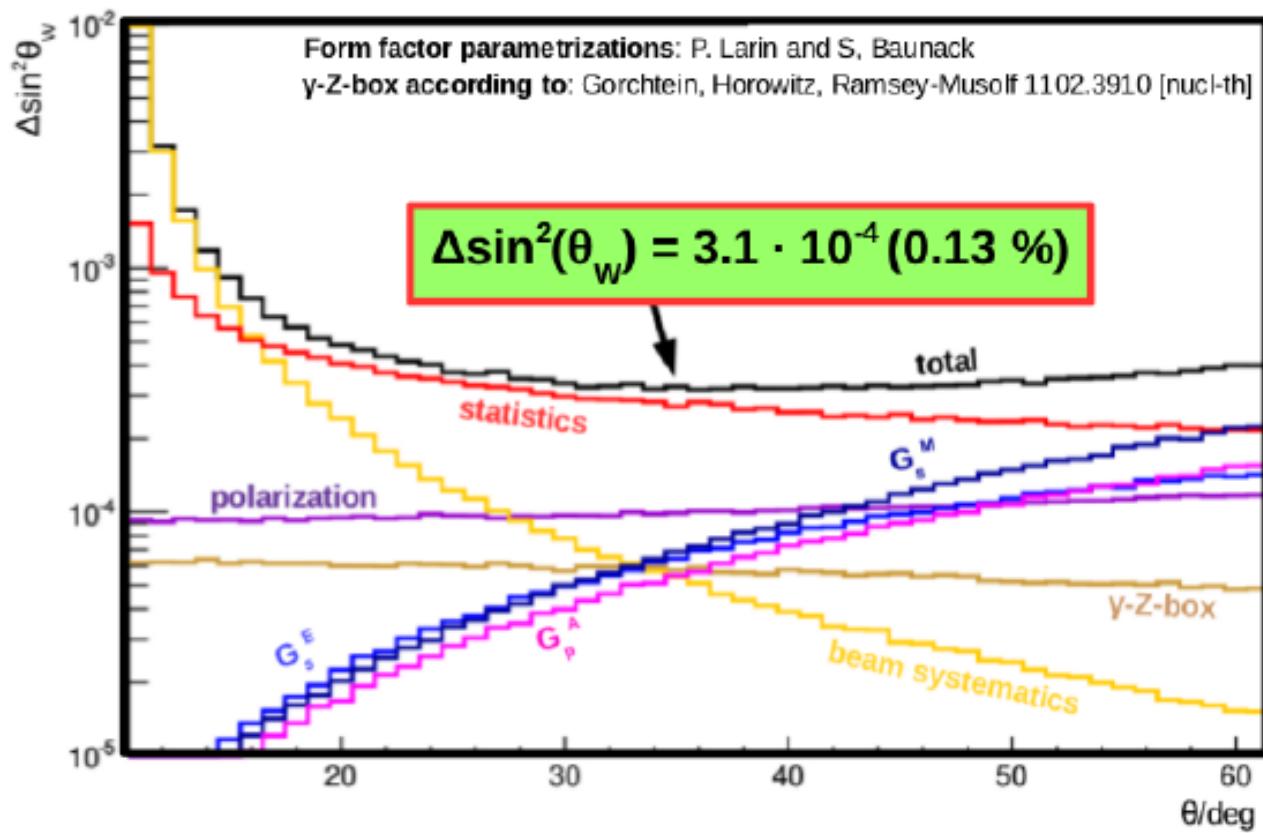
- Contributions to $\Delta \sin^2 \Theta_W$ for 35° central scattering angle, $E=150$ MeV, 10000 h of data taking





JG|U

P2-Precision in $\sin^2 \theta_W$



Beam energy: 155 MeV
 Beam current: 150 μA
 Polarization: $(85 \pm 0.5)\%$
 Target: 60 cm lH₂
 Acceptance: $2\pi \cdot (35^\circ \pm 10^\circ)$
 Rate: 0.5 THz
 Runtime: 10000 h
 ΔA^{app} : 0.1 ppb

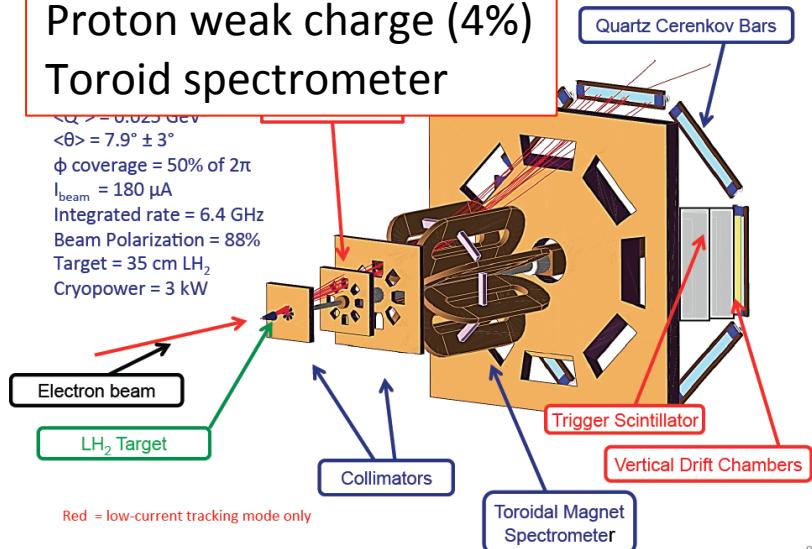
	Total	Statistics	Polarization	Apparative	FF	$\text{Re}(\square_{yzA})$
$\Delta \sin^2(\theta_W)$	3.1e-4 (0.13 %)	2.6e-4 (0.11 %)	9.7e-5 (0.04 %)	7.0e-5 (0.03 %)	1.4e-4 (0.04 %)	6e-5 (0.03 %)
$\Delta A^{\text{exp}}/\text{ppb}$	0.44 (1.5 %)	0.38 (1.34 %)	0.14 (0.49 %)	0.10 (0.35 %)	0.11 (0.38 %)	0.09 (0.32 %)



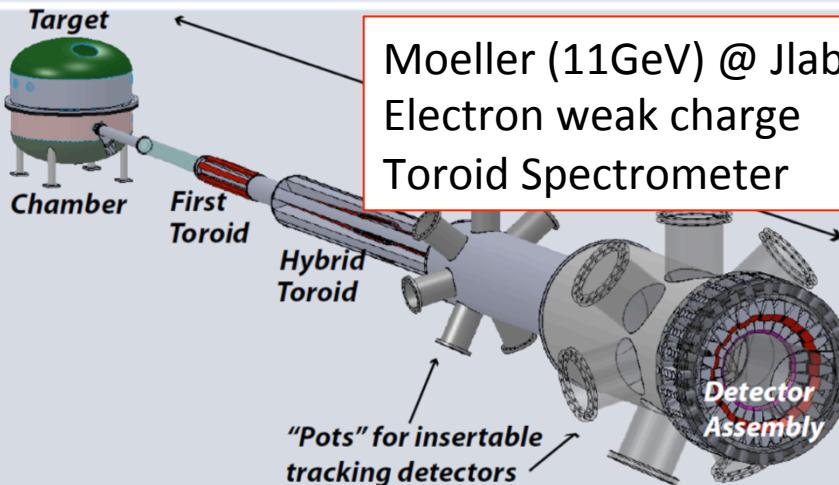
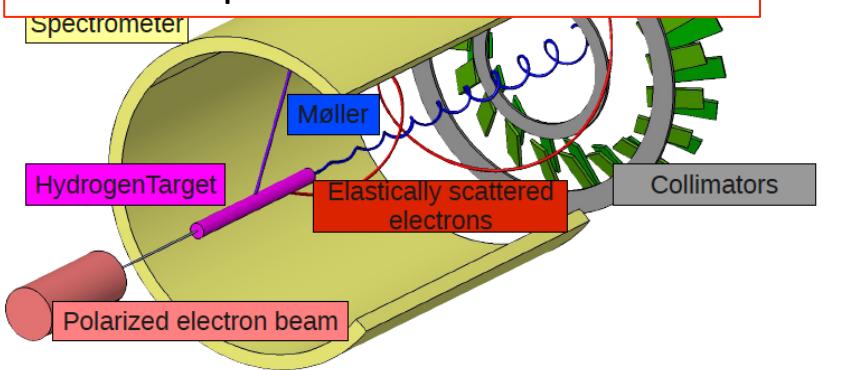
Beam Energy ERL/EB [MeV]	105/155 (105/205)
Operation mode	1300 MHz, c.w.
Elektron-sources	1.) Polarised : NEA GaAsP/GaAs superlattice , 200keV (?) 2.) unpolarised KCsSb, 200keV
Bunch Charge EB/ERL [pC] $7.7\text{pC}=\text{10mA@1300MHz}$	0.15/0.77 (0.15/7.7)
Norm. Emittance EB/ERL [μm]	0.1/<0.5 (0.1/<1)
Spin Polarisation (EB-mode only)	> 0.85
Recirculations	2 (3)
Beampower at Exp. ERL/EB [kW]	100/22.5 (1050/30)
R.f.-Power installed [kW]	140 (180)



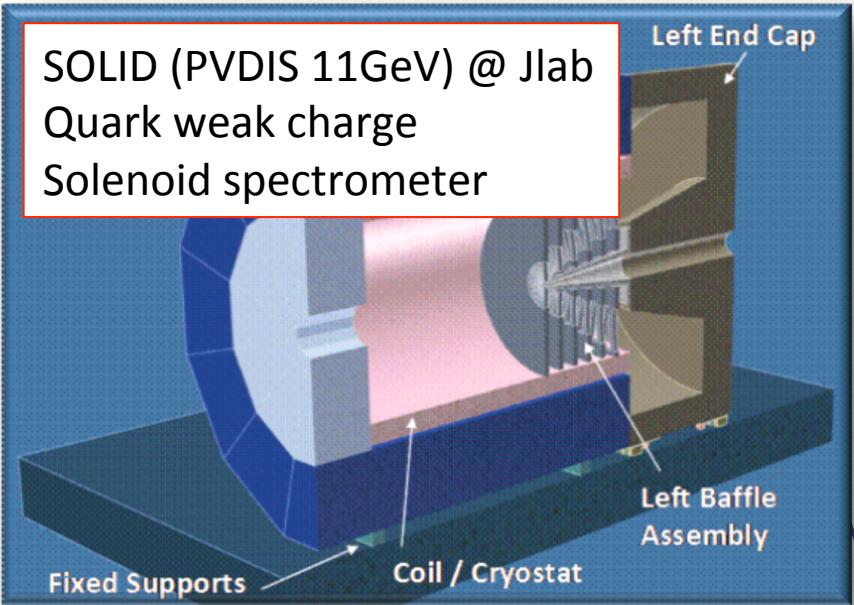
Qweak (1GeV) @ Jlab
Proton weak charge (4%)
Toroid spectrometer



P2@MESA (0.150 GeV) @ Mainz
Proton weak charge (1.7%)
Solenoid spectrometer



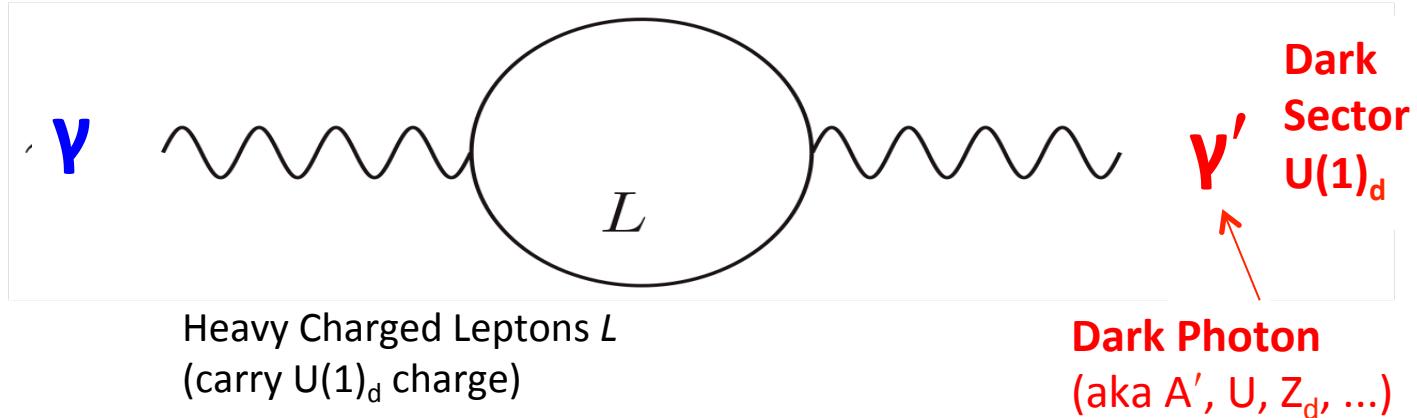
SOLID (PVDIS 11GeV) @ Jlab
Quark weak charge
Solenoid spectrometer





A portal to relate the dark sector to the SM world (coupling $\sim \epsilon^2$)

Standard
Model
Sector
 $U(1)$

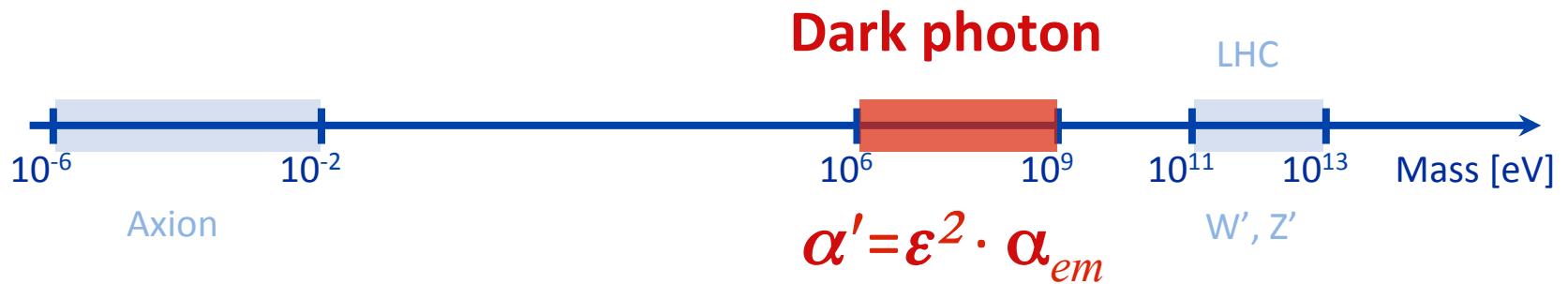


Features à la Arkani-Hamed: A theory of Dark Matter

- More than one Dark Matter particle \rightarrow Dark Sector
- $dm + dm \rightarrow e^+e^-$ explains positron excess
- Astrophysical anomalies (PAMELA, FERMI, DAMA/LIBRA, INTEGRAL, ...) suggest dark photon mass on GeV mass scale (and lighter than $2M_p$)

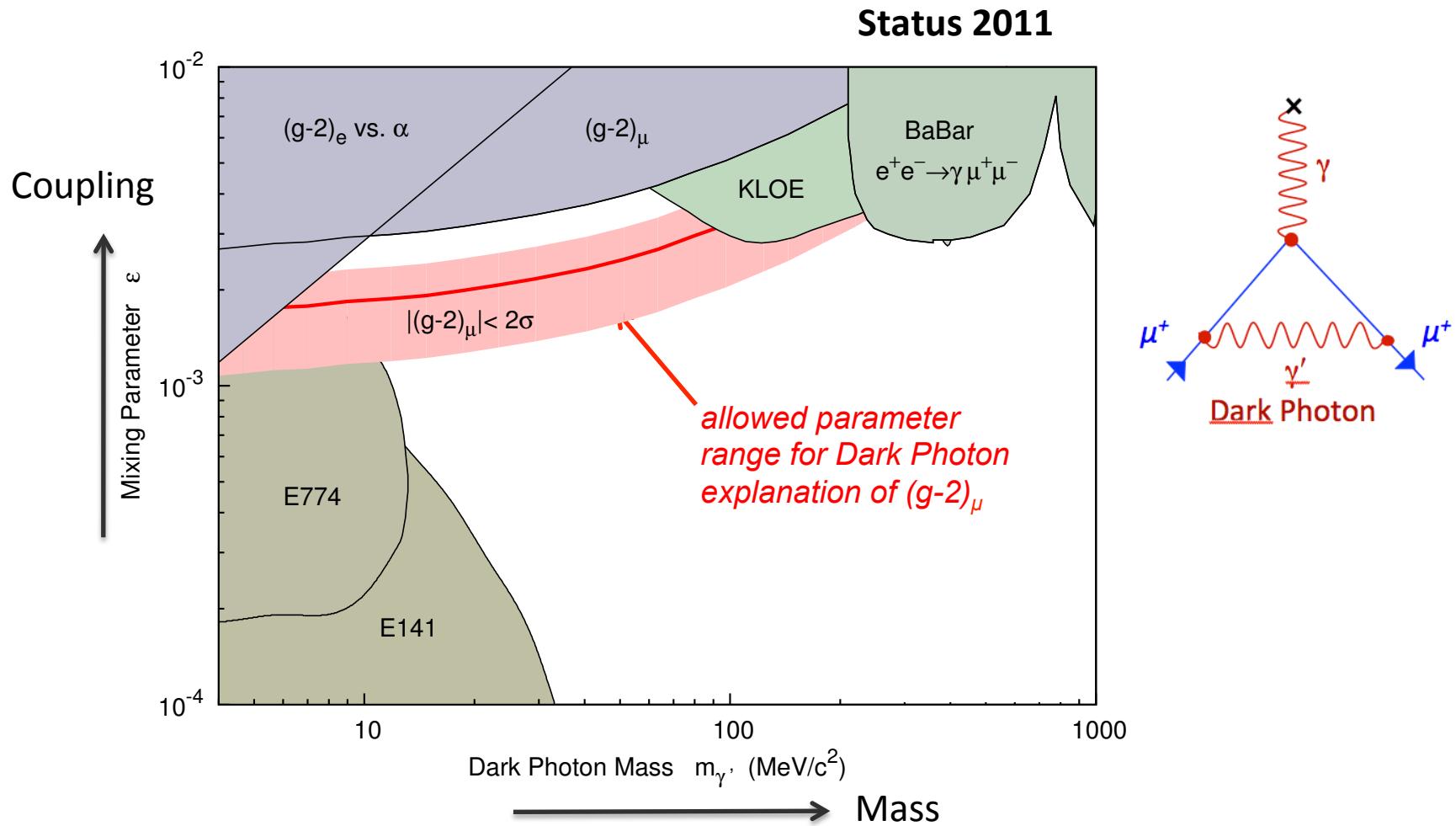


New massive force carrier of extra $U(1)_d$ gauge group; predicted in almost all string compactifications



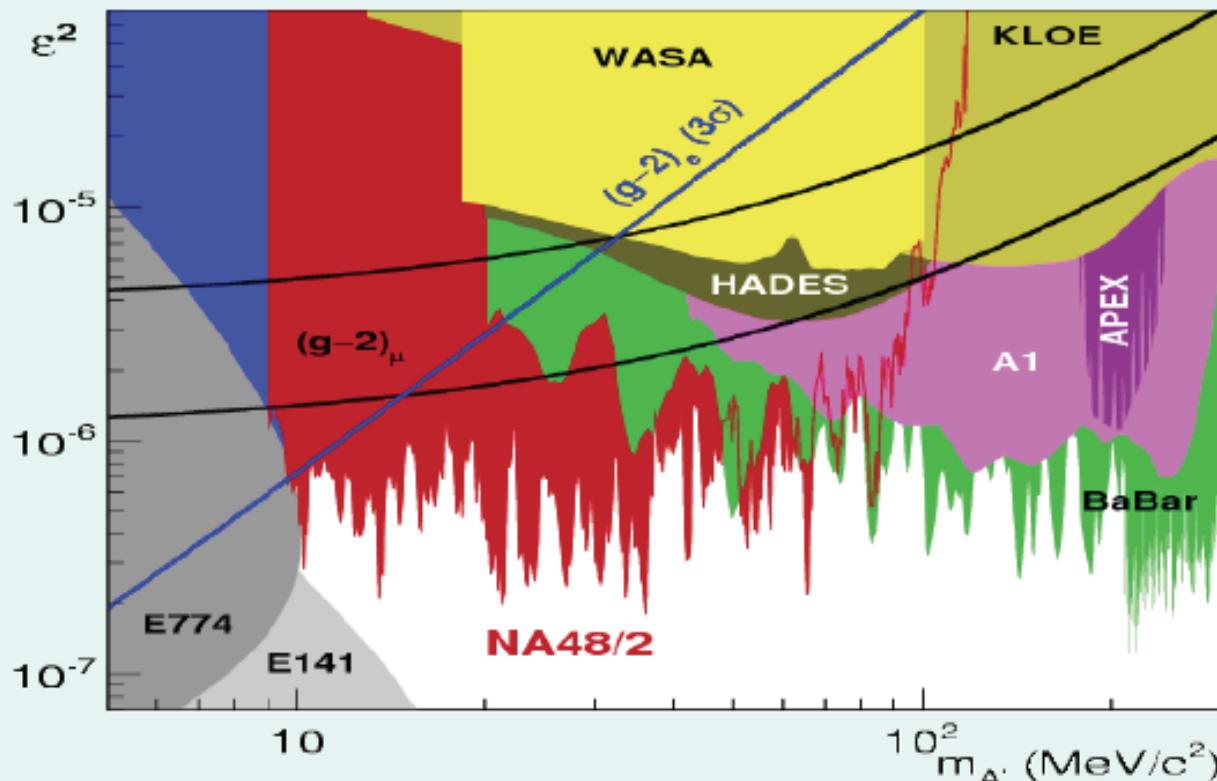
Search for the $O(\text{GeV}/c^2)$ mass scale in a world-wide effort

- Could explain large number of **astrophysical anomalies**
*Arkani-Hamed et al. (2009)
Andreas, Ringwald (2010); Andreas, Niebuhr, Ringwald (2012)*
- Could explain presently seen **deviation of $>3\sigma$** between $(g-2)_\mu$
Standard Model prediction and direct $(g-2)_\mu$ measurement
Pospelov (2008)





NA48/2 Updated Bounds on Dark Photon
 $g_\mu - 2$ discrepancy solution ruled out
Assumes $\text{BR}(Z_d \rightarrow e^+e^-) \sim 1$





High Precision Determination of $\sin^2(\theta_w)$

Running of $\sin^2(\theta_w)$

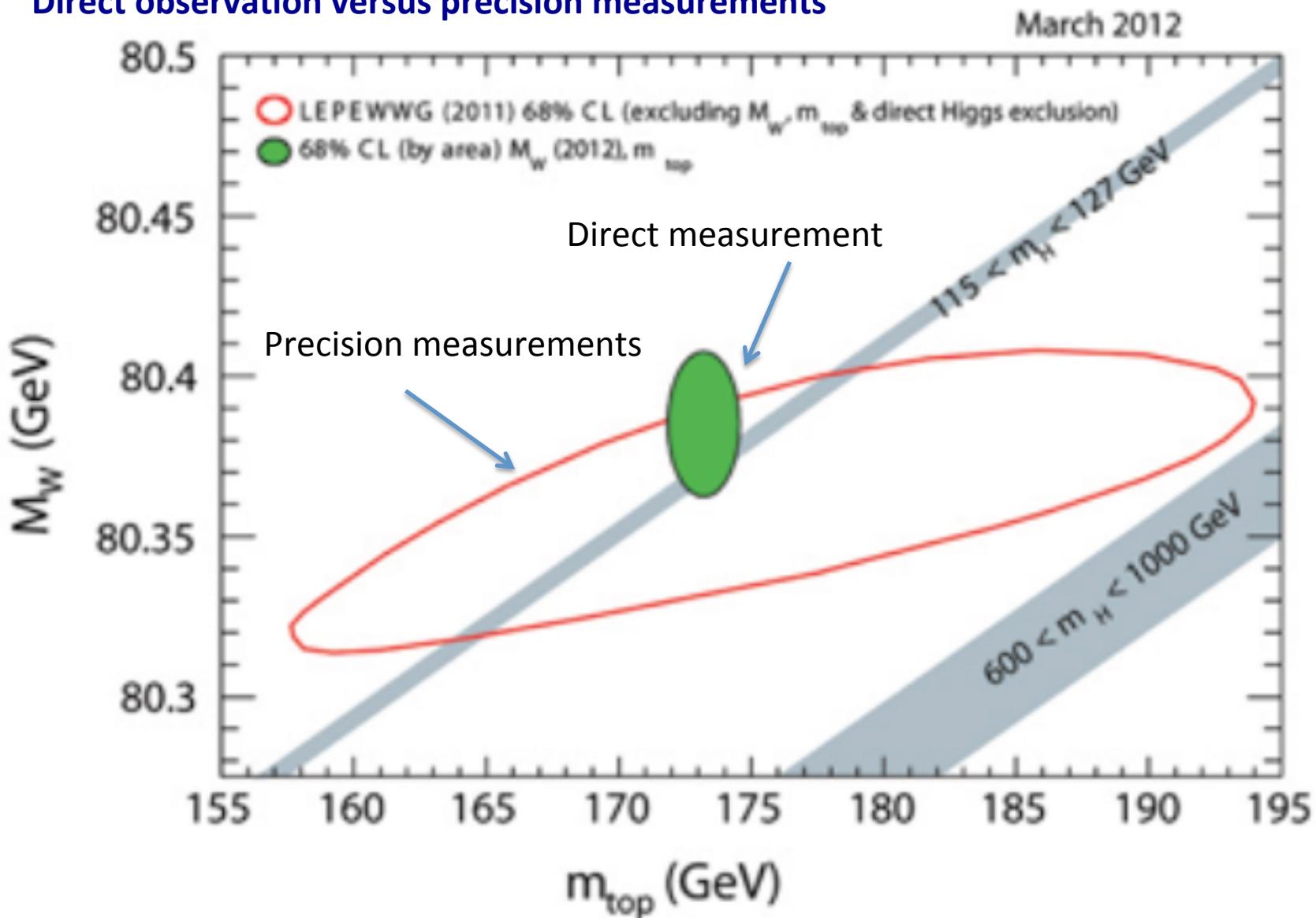
Sensitivity to new physics

Experimental Method

ECT*-Workshop: Physics beyond the standard model and precision nucleon structure Measurements, 01 Aug 2016 to 05 Aug 2016, Villa Tambosi, Trento, Organisers: F. Maas, K. Kumar, P. Souder, M. Vanderhaghen <http://www.ectstar.eu/node/1682>

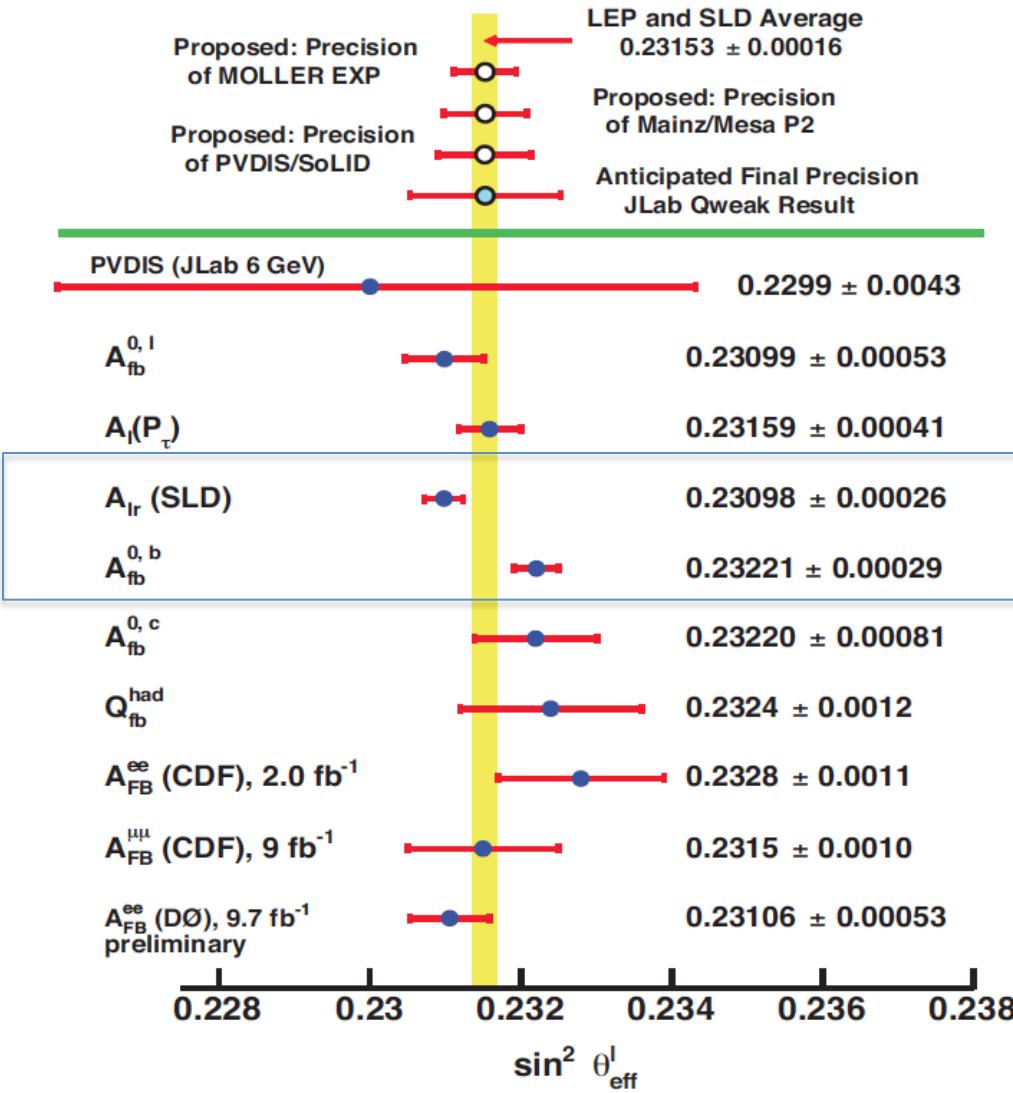


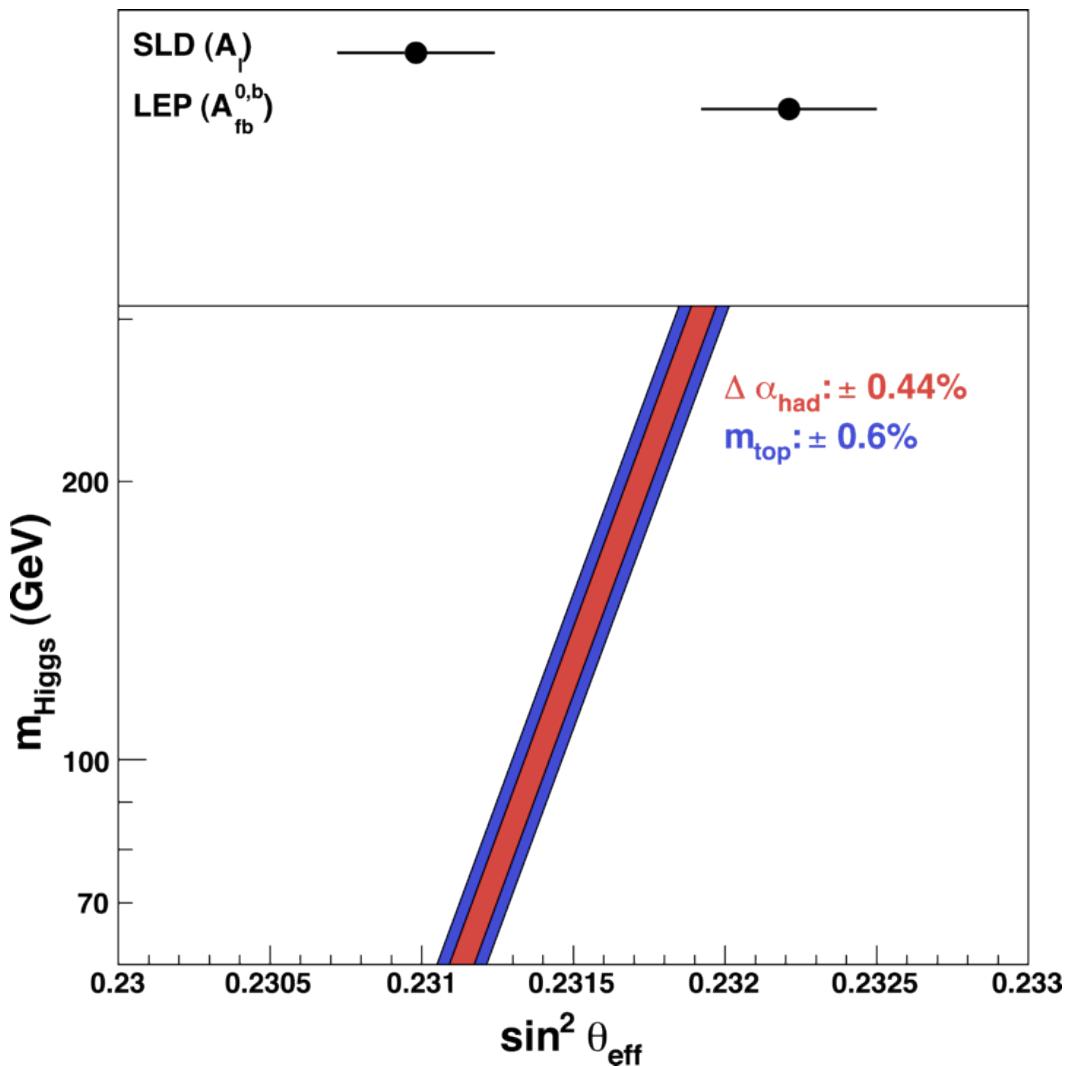
Direct observation versus precision measurements

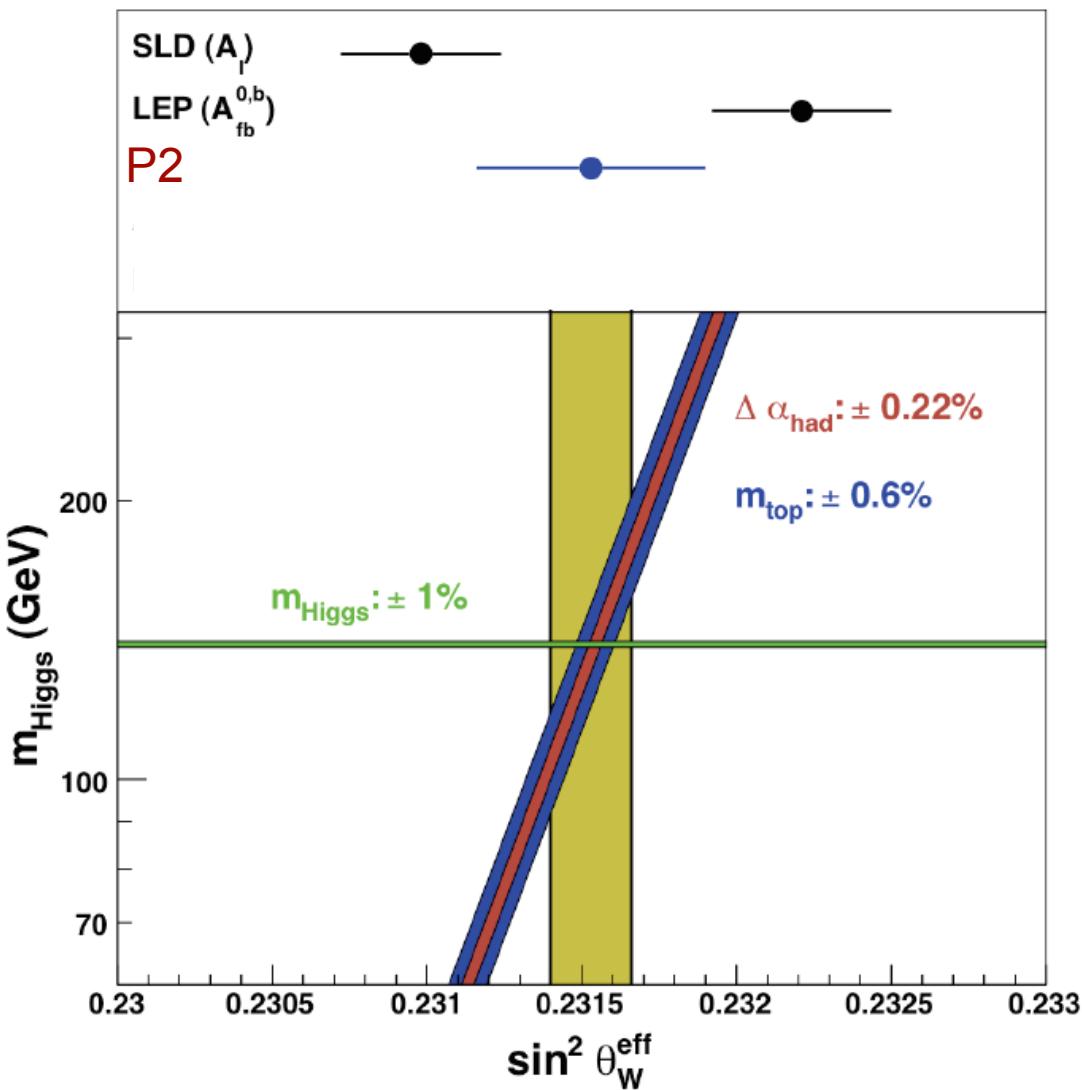


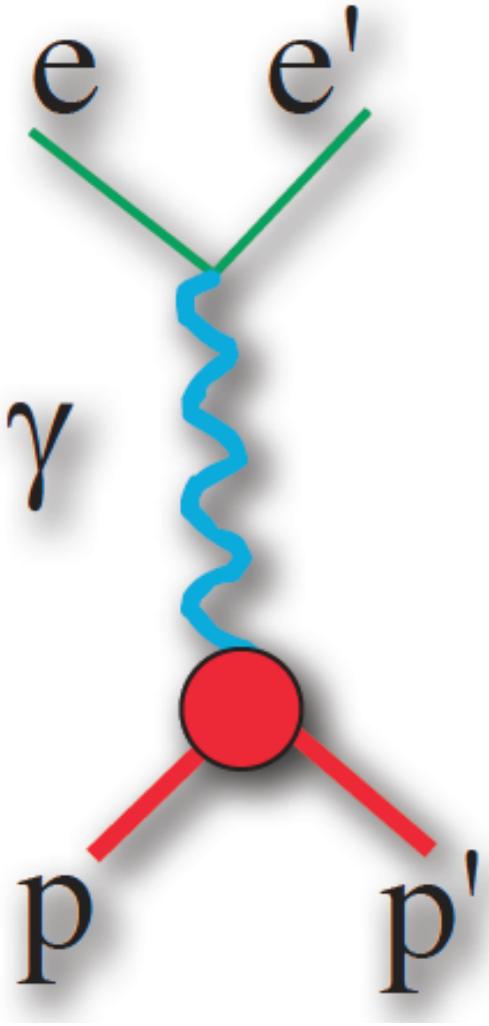


Summary: Measurements of $\sin^2 \theta_{W(\text{effective})}^l$



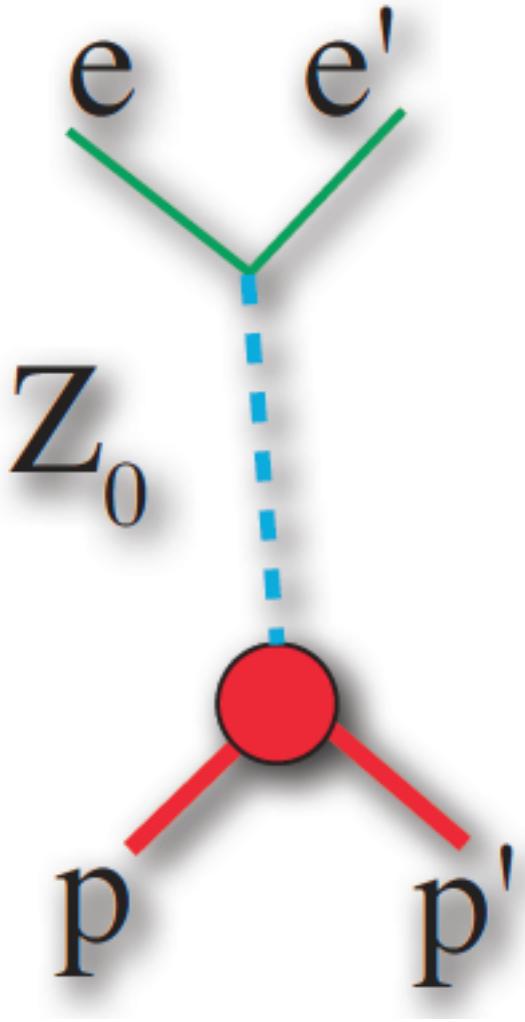






$$\begin{aligned}\sigma &\sim \mathcal{M} \mathcal{M}^* \text{ Phasespace} \\ &\sim (\mathbf{j}_\mu \frac{1}{Q^2} J^\mu)(\mathbf{j}_\mu \frac{1}{Q^2} J^\mu)^*\end{aligned}$$
$$\mathbf{j}_\mu \sim \bar{e} \gamma_\mu e \text{ Vector Current}$$

$$\begin{aligned}J_\gamma^\mu &\sim \left\langle N | q^{\textcolor{red}{u}} \bar{u} \gamma_\mu u + q^{\textcolor{blue}{d}} \bar{d} \gamma_\mu d + q^{\textcolor{green}{s}} \bar{s} \gamma_\mu s | N' \right\rangle \\ &= \overline{\mathcal{P}} [\gamma^\mu \mathbf{F}_1 - i \sigma^{\mu\nu} q_\nu \frac{\kappa_p}{2M_N} \mathbf{F}_2] \mathcal{P}\end{aligned}$$



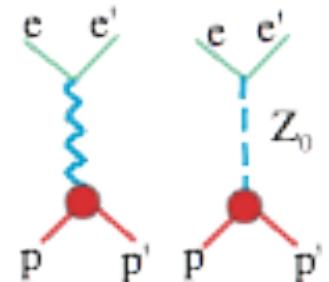
$$\tilde{q}^d V = \tau_3 - 2q^d \sin^2(\theta_W)$$

$$\begin{aligned}\tilde{J}_Z^\mu &\sim \left\langle N | \tilde{q}^u \bar{u} \gamma_\mu u + \tilde{q}^d \bar{d} \gamma_\mu d + \tilde{q}^s \bar{s} \gamma_\mu s | N' \right\rangle \\ &= \overline{\mathcal{P}} [\gamma^\mu \tilde{F}_1 - i \sigma^{\mu\nu} q_\nu \frac{\kappa_p}{2M_N} \tilde{F}_2] \mathcal{P}\end{aligned}$$



Parity violating cross section asymmetry

$$A_{ep} = \left[\frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \right] \frac{\epsilon G_E^\gamma G_E^Z + \tau G_M^\gamma G_M^Z - (1 - 4 \sin^2 \theta_W) \epsilon' G_M^\gamma G_A^Z}{\epsilon (G_E^\gamma)^2 + \tau (G_M^\gamma)^2}$$



$$A_{RL} = \underbrace{A_V + A_A}_{= A_0} + A_S \begin{cases} A_V = -a\rho'_{eq} \left[(1 - 4 \sin^2 \theta_W) - \frac{\epsilon G_E^p G_E^n + \tau G_M^p G_M^n}{\epsilon (G_E^p)^2 + \tau (G_M^p)^2} \right] \\ A_A = a \frac{(1 - 4 \sin^2 \theta_W) \sqrt{1 - \epsilon^2} \sqrt{\tau(1 + \tau)} G_M^p \tilde{G}_A^p}{\epsilon (G_E^p)^2 + \tau (G_M^p)^2} \\ A_S = a\rho'_{eq} \frac{\epsilon G_E^p G_E^s + \tau G_M^p G_M^s}{\epsilon (G_E^p)^2 + \tau (G_M^p)^2} \end{cases}$$

$$a = -G_F q^2 / 4\pi\alpha\sqrt{2}, \quad \tau = -q^2 / 4M_p^2, \quad \epsilon = [1 + 2(1 + \tau) \tan^2 \theta_W / 2]^{-1}$$