

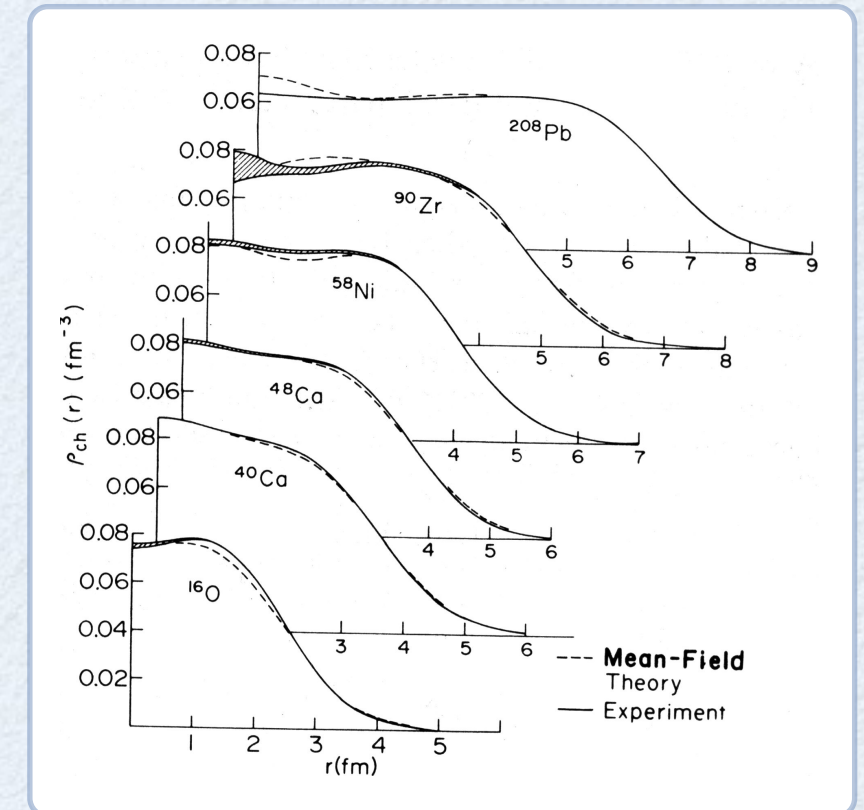
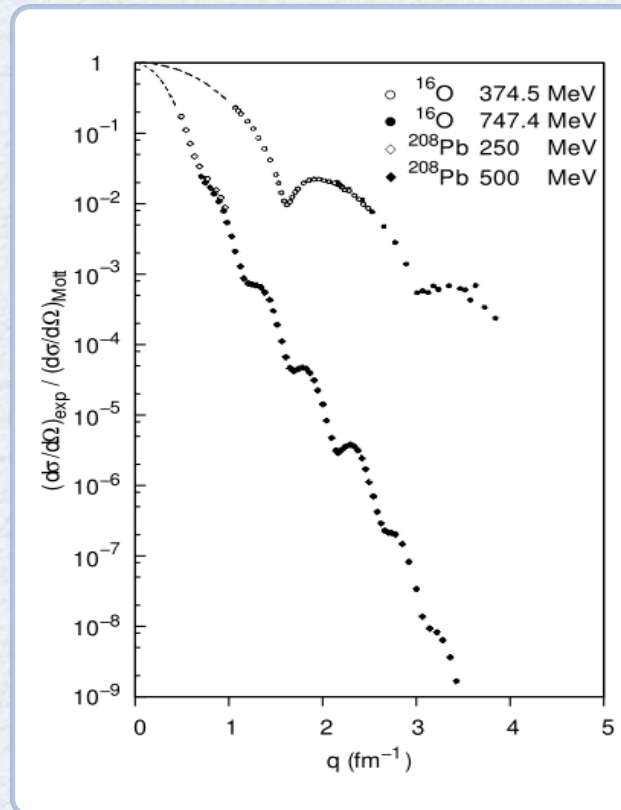
Proton radius: theoretical overview and perspectives

M. Vanderhaeghen

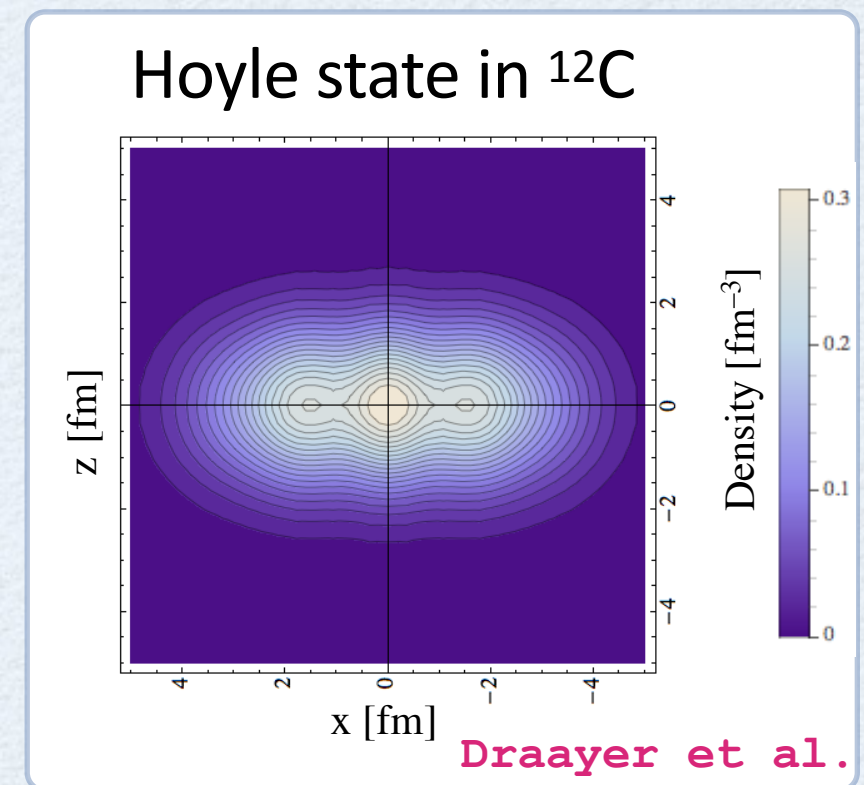
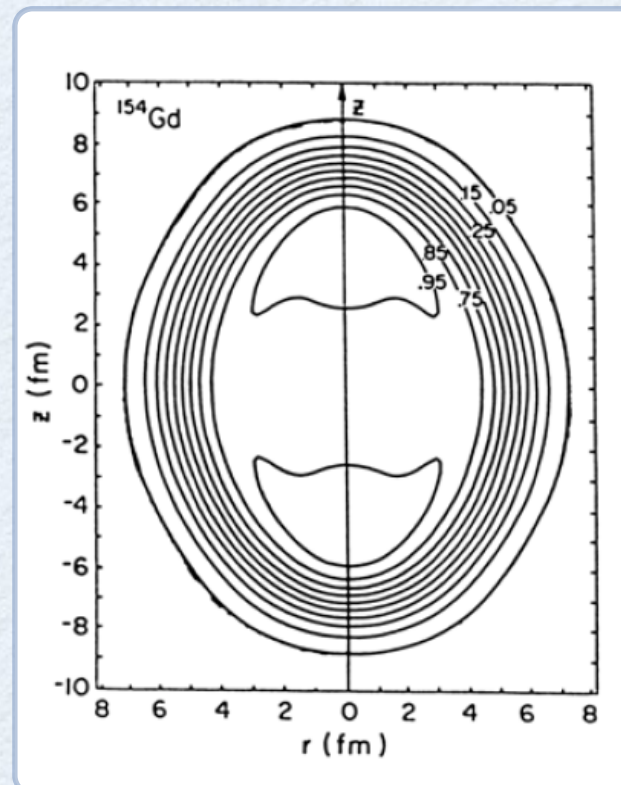
PRES Meeting, March 30-31, 2021, Univ. Mainz

Imaging of atomic nuclei

sizes of nuclei:
as revealed through
elastic electron scattering



shapes of nuclei:
as revealed through
inelastic electron scattering
deformations, coherent states



Atomic nuclei

➔ To define and reconstruct a 3dim charge distribution of system of mass M from electron scattering: one needs to localize the system and fix its c.m.

Condition:

Jaffe (2021)

$$1/M \ll R \ll R_E$$

Compton
wavelength

localization provided
by wave packet

Size of
system

➔ For **atomic nuclei**: $1/M \sim 0.2/A$ fm for $^{12}\text{C} \rightarrow 0.02$ fm
 $R_E \sim 1.2 A^{1/3}$ fm for $^{12}\text{C} \rightarrow 2.5$ fm

$$1/M \lll R_E$$



For static (non-relativistic) systems: the 3D **Fourier transform** of **form factors** gives the distribution of electric charge and magnetization

$$G_E(Q^2) = 1 - \frac{1}{6} R_E^2 Q^2 + \mathcal{O}(Q^4) \quad \rightarrow \quad \text{Charge radius: } R_E^2 \equiv \frac{\int d^3\vec{r} r^2 \rho_{3d}(r)}{\int d^3\vec{r} \rho_{3d}(r)}$$

From nuclei to nucleons

➔ For **proton**:

Compton wavelength: $1/M \sim 0.2 \text{ fm}$

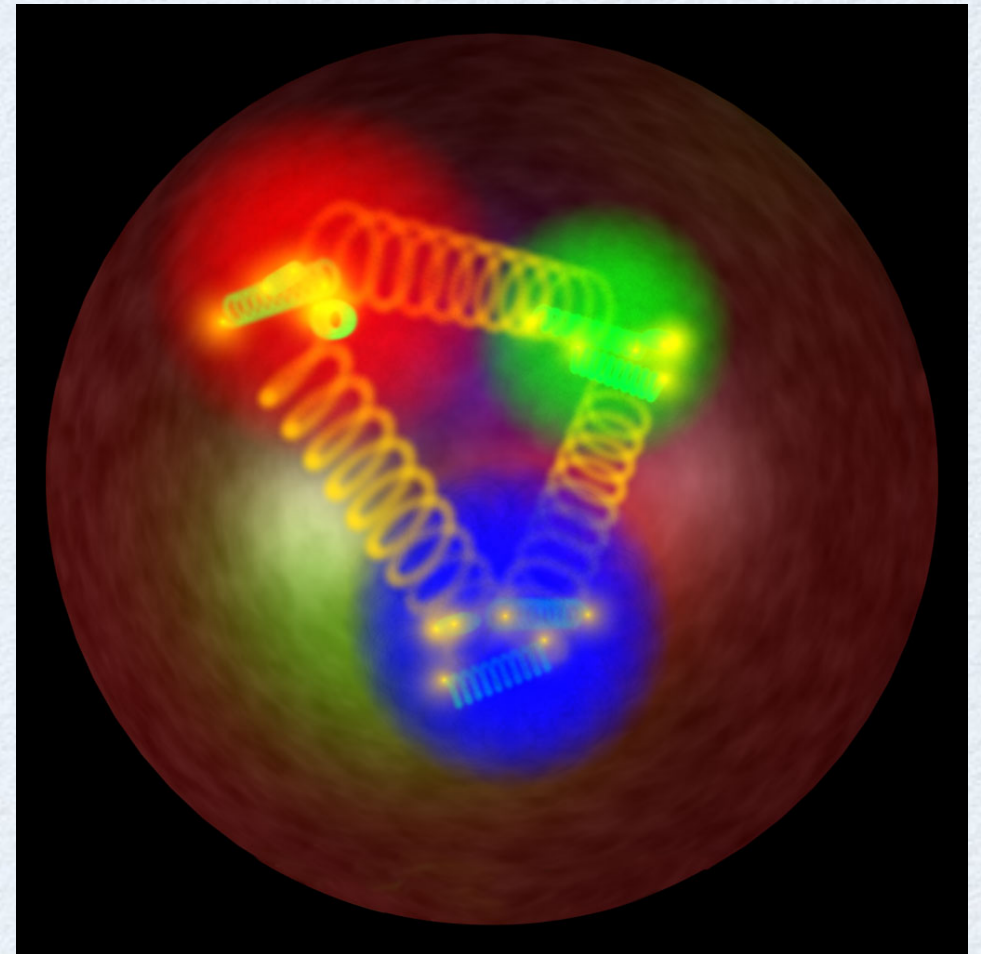
Size of proton: $R_E \sim 0.85 \text{ fm}$

Condition $1/M \lll R_E$



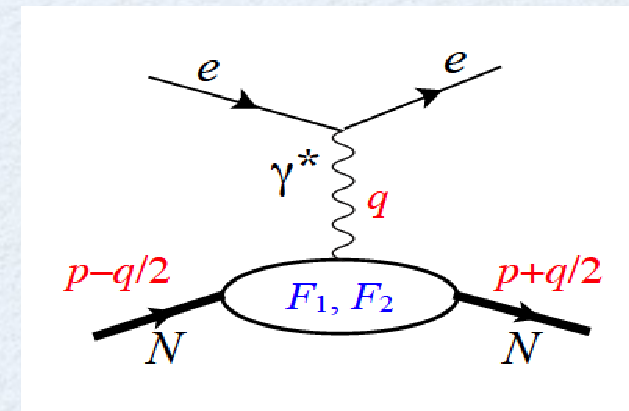
➔ For relativistic systems as proton:
interpretation of **Fourier transform**
of **form factors** as a 3-dim charge or
magnetization distribution cannot
be defined precisely

Jaffe (2021)



Proton electromagnetic form factors

- **elastic electron-proton scattering**
in the *1-photon exchange approximation*:
parameterized by 2 **form factors (FFs)**



$$\langle p + \frac{q}{2}, \lambda' | J^\mu(0) | p - \frac{q}{2}, \lambda \rangle = \bar{u}(p + \frac{q}{2}, \lambda') \left[F_1(Q^2) \gamma^\mu + F_2(Q^2) \frac{i}{2M} \sigma^{\mu\nu} q_\nu \right] u(p - \frac{q}{2}, \lambda)$$

↑ ↑
Dirac FF **Pauli FF**

for proton: $F_1(Q^2 = 0) = 1$ $F_2(Q^2 = 0) = \kappa_p = 1.79$

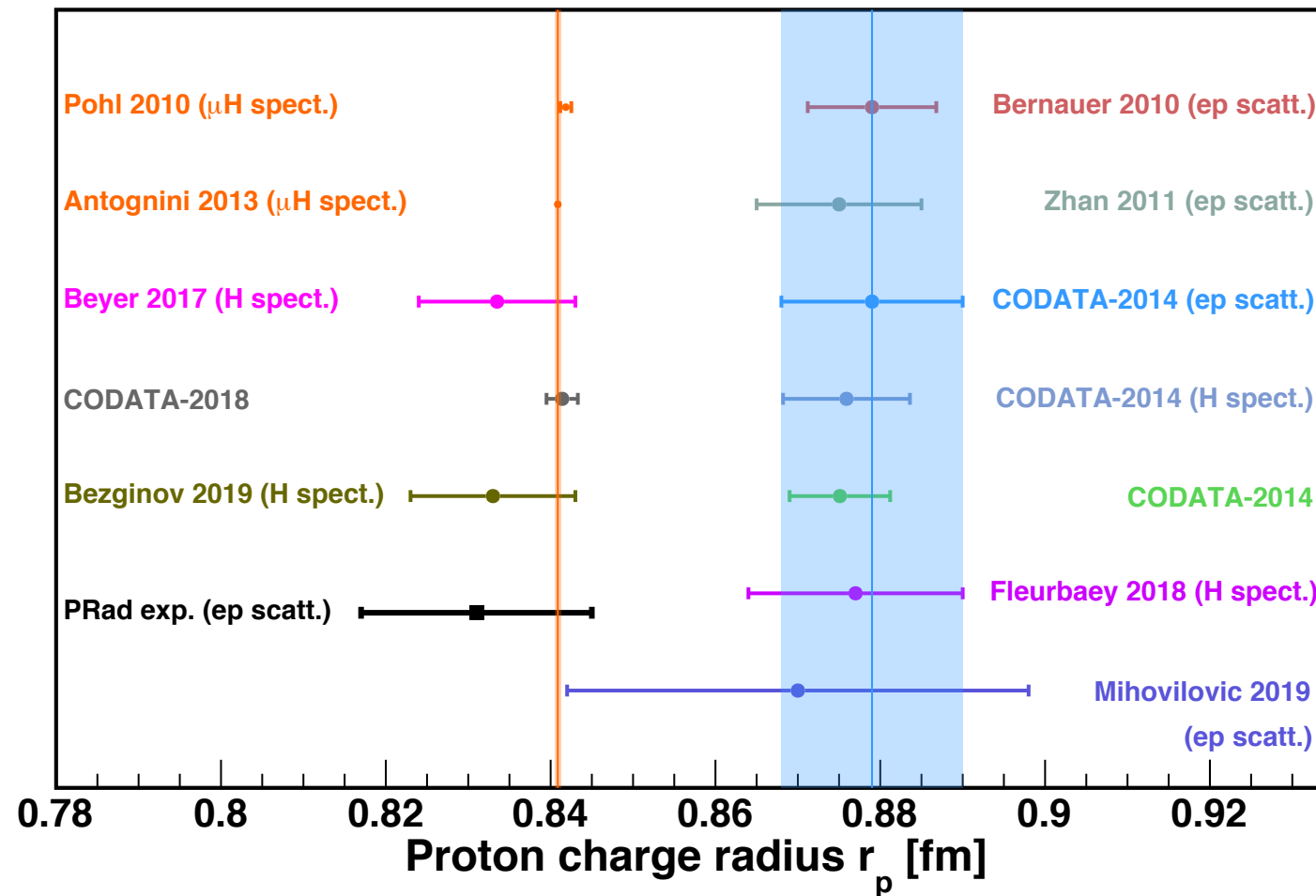
- equivalently: in experiment one uses **Sachs FFs** with $\tau \equiv \frac{Q^2}{4M^2}$

$G_M(Q^2) = F_1(Q^2) + F_2(Q^2)$	→	magnetic FF	$G_E(Q^2) = 1 - \frac{1}{6} \langle r_E^2 \rangle Q^2 + \mathcal{O}(Q^4)$
$G_E(Q^2) = F_1(Q^2) - \tau F_2(Q^2)$	→	electric FF	

charge radius is definition for FF slope!

- In **Hydrogen spectroscopy (Lamb shift)**: quantity which enters is also slope of G_E !

Experimental r_{Ep} determinations



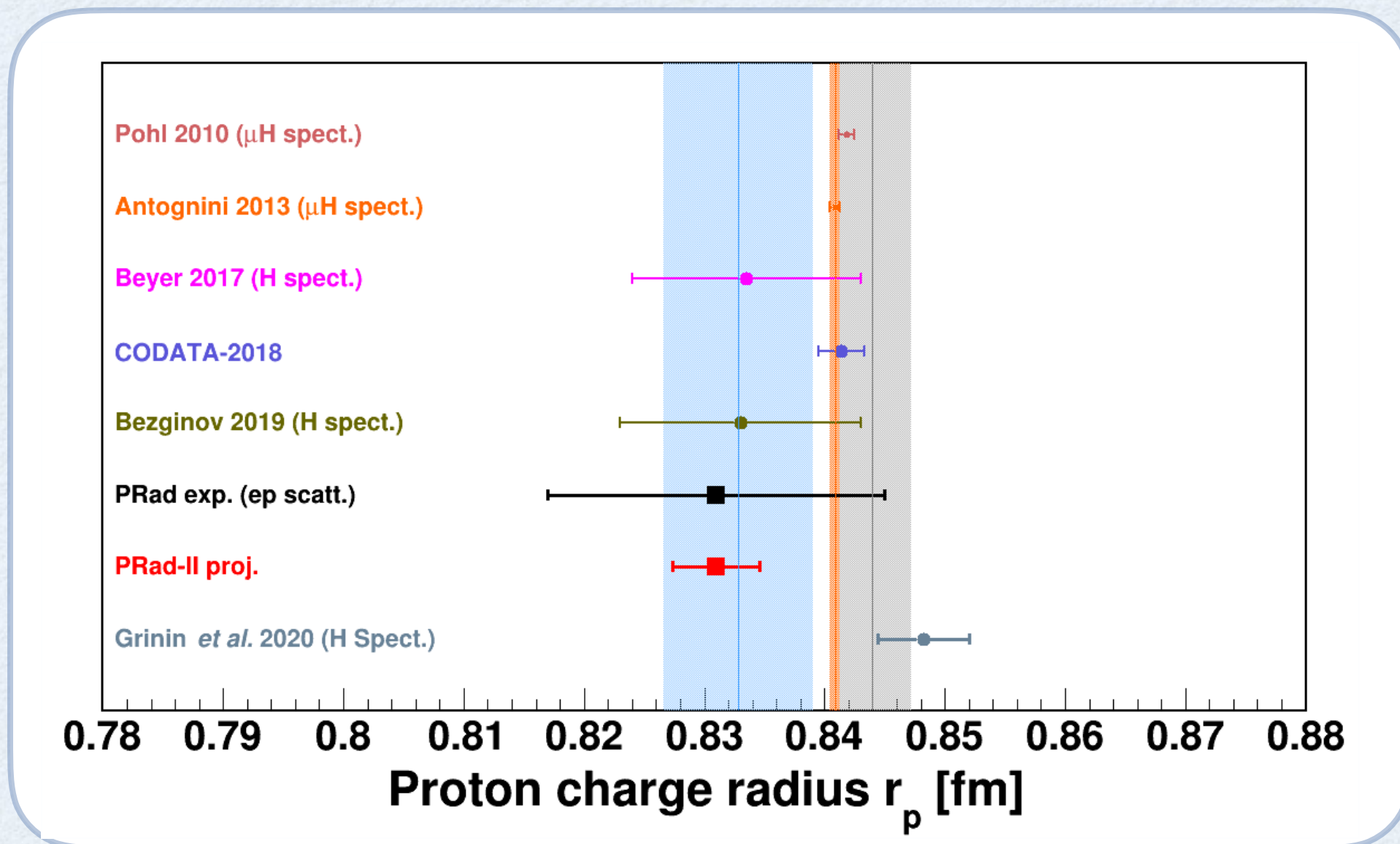
from recent compilation
prepared for Rev. Mod. Phys.

H. Gao, M. Vdh (2021)

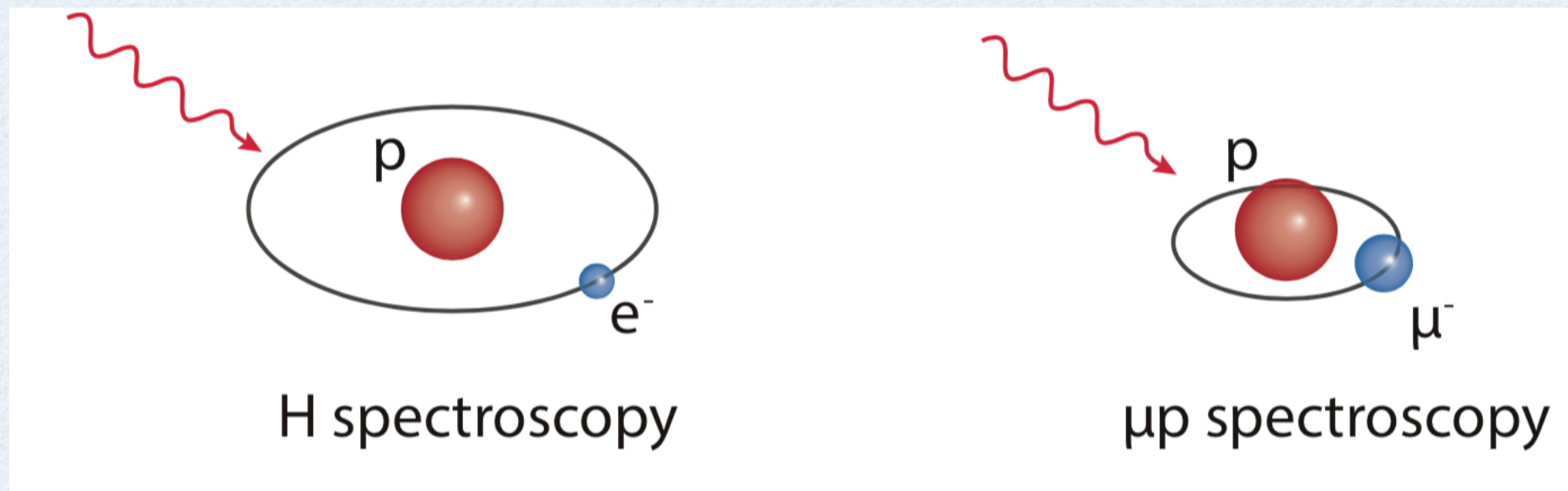
PRad-II plans

PRad-II: new JLab approved experiment Dutta et al. (2020)

- improvement of stat. uncertainty by factor 4 in comparison to PRad
- various improvements to reduce syst. uncertainties
- Factor 3.8 reduction in overall exp. uncertainty on proton radius compared to PRad



Hadronic input to muonic atom spectroscopy program



Muonic atom spectroscopy needs nucleon/nuclear input

2S-2P Lamb Shift:

THEORY

EXPERIMENT

	$\Delta E_{TPE} \pm \delta_{theo} (\Delta E_{TPE})$	Ref.	$\delta_{exp}(\Delta_{LS})$	Ref.
μH	$33 \mu\text{eV} \pm 2 \mu\text{eV}$	Antognini et al. (2013)	$2.3 \mu\text{eV}$	Antognini et al. (2013)
μD	$1710 \mu\text{eV} \pm 15 \mu\text{eV}$	Krauth et al. (2015)	$3.4 \mu\text{eV}$	Pohl et al. (2016)
$\mu^3\text{He}^+$	$15.30 \text{ meV} \pm 0.52 \text{ meV}$	Franke et al. (2017)	0.05 meV	
$\mu^4\text{He}^+$	$9.34 \text{ meV} \pm 0.25 \text{ meV}$ $-0.15 \text{ meV} \pm 0.15 \text{ meV} (3\text{PE})$	Diepold et al. (2018) Pachucki et al. (2018)	0.05 meV	Krauth et al. (2020)

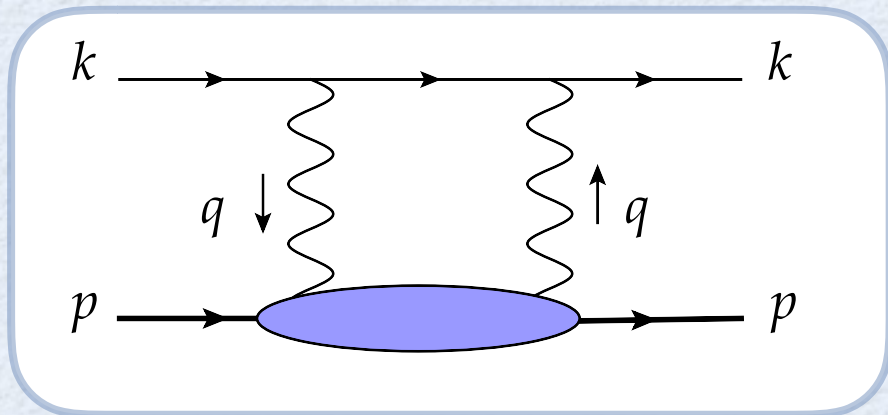
μH :

present accuracy comparable with experimental precision
Future: factor 5 improvement planned

$\mu\text{D}, \mu^3\text{He}^+, \mu^4\text{He}^+$:

present accuracy factor 5-10 worse than experimental precision

Two-photon exchange: hadronic corrections



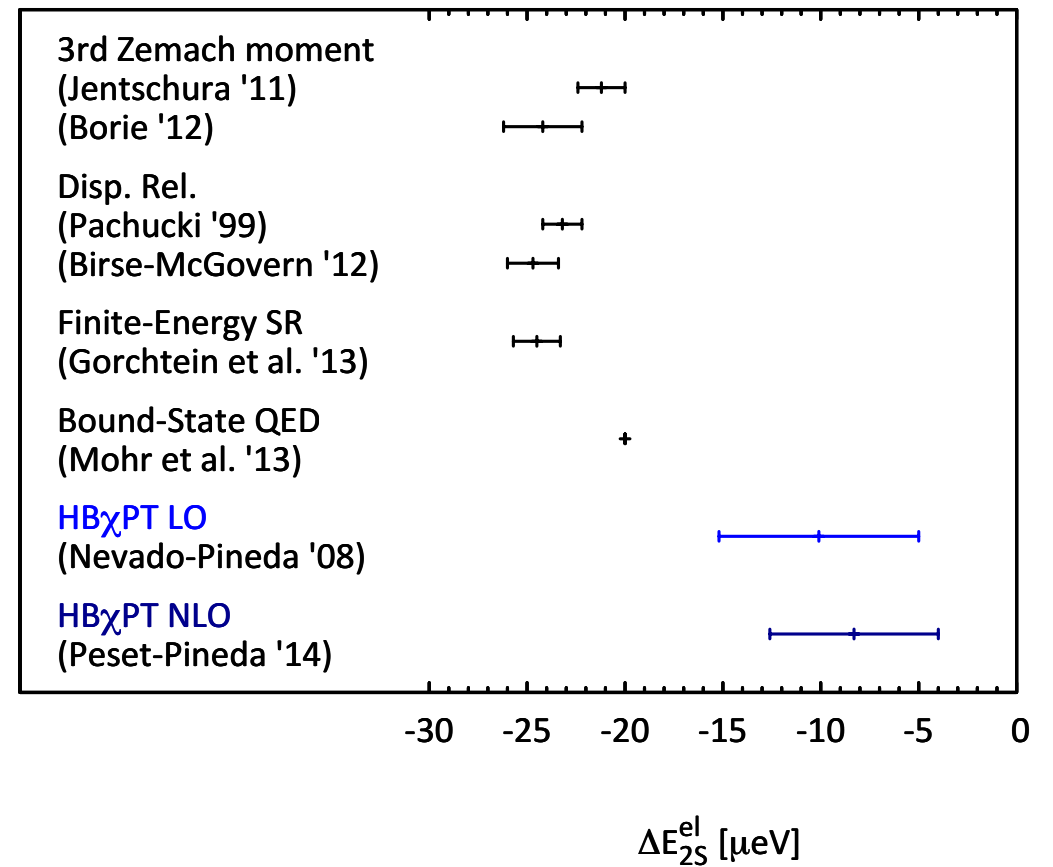
$$T^{\mu\nu}(p, q) = \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) T_1(\nu, Q^2) + \frac{1}{M^2} \left(p^\mu - \frac{p \cdot q}{q^2} q^\mu \right) \left(p^\nu - \frac{p \cdot q}{q^2} q^\nu \right) T_2(\nu, Q^2)$$

- **Two-photon exchange (TPE)**: lower blob contains both elastic (nucleon) and inelastic states
- **Lamb shift**: described by unpolarized amplitudes T_1 , T_2 : functions of energy ν and Q^2
- **Hyperfine splitting**: described by polarized amplitudes S_1 , S_2
- **Imaginary parts**: directly proportional to nucleon structure functions F_1 , F_2 resp. g_1 , g_2
- **Real parts**: obtained as dispersion integral over the imaginary parts modulo a subtraction function in case of T_1

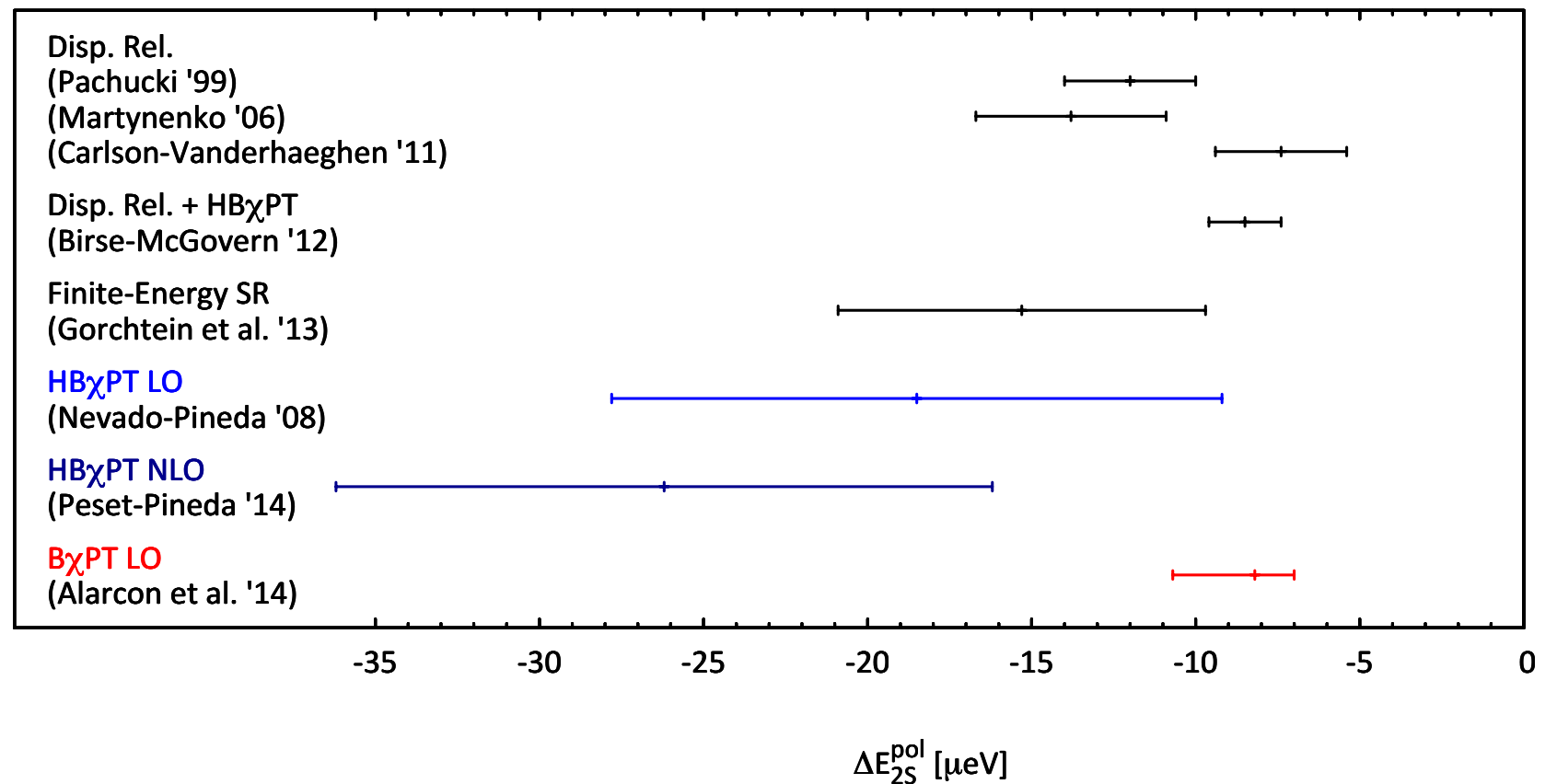
$\Delta E = \Delta E^{el}$	→	Elastic state: involves nucleon form factors
$+ \Delta E^{subtr}$	→	Subtraction: involves nucleon polarizabilities
$+ \Delta E^{inel}$	→	Inelastic state: involves nucleon structure functions

Hadron/Nuclear physics input needed !

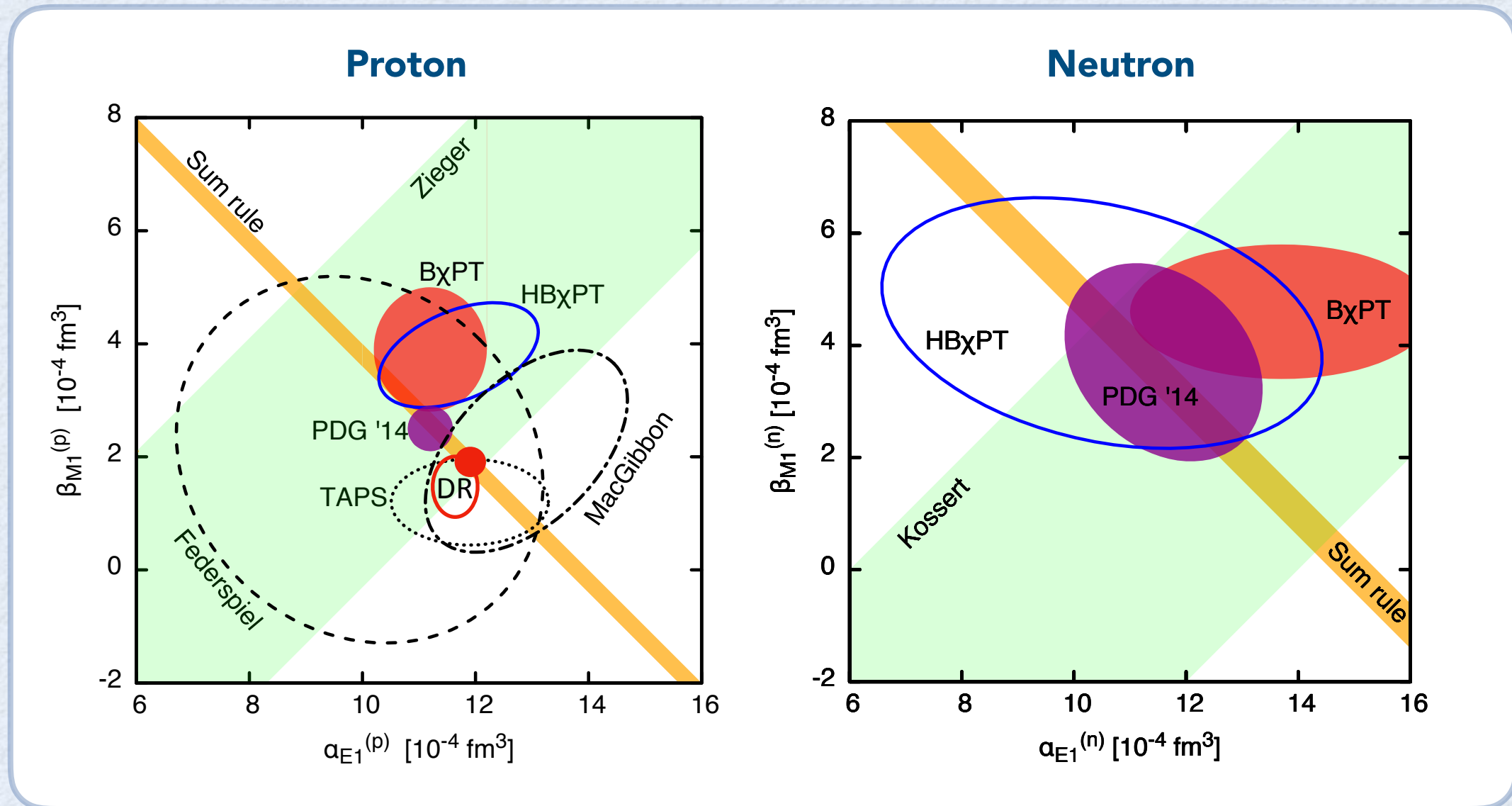
TPE elastic correction:



TPE polarizability correction:



TPC program A2@MAMI: neutron polarizabilities



$$\alpha_E = (11.2 \pm 0.4) \times 10^{-4} \text{ fm}^3$$

PDG '14 values

$$\beta_M = (2.5 \pm 0.4) \times 10^{-4} \text{ fm}^3$$

$$\alpha_E = (11.8 \pm 1.1) \times 10^{-4} \text{ fm}^3$$

$$\beta_M = (3.7 \pm 1.2) \times 10^{-4} \text{ fm}^3$$

New A2@MAMI consisted data set (factor 5 increase of world data) in final stage of analysis

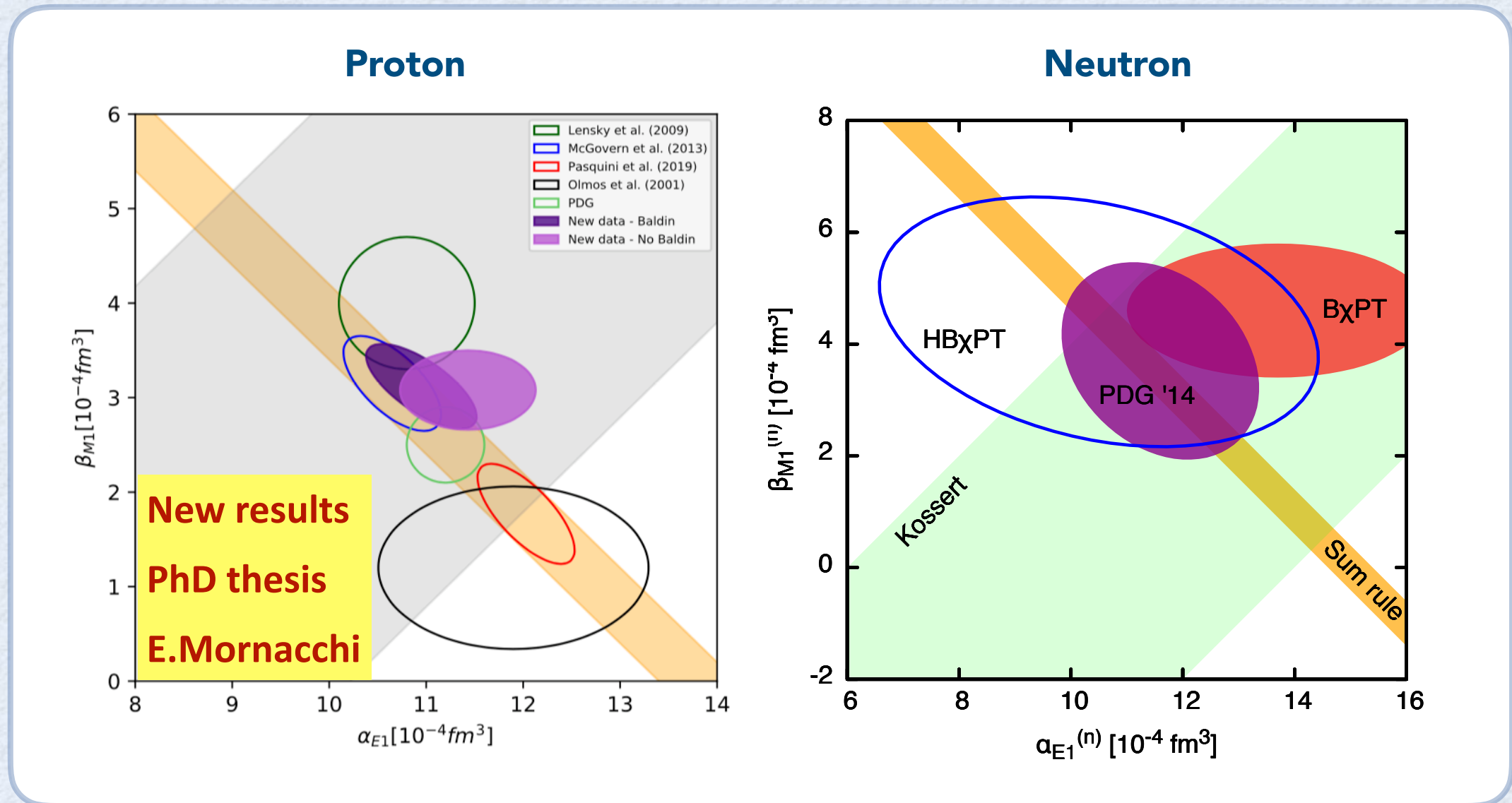
→ will consolidate PDG errors

A reduction of error by factor of 2

→ precision comparable with proton

→ impact light muonic atom program

TPC program A2@MAMI: neutron polarizabilities



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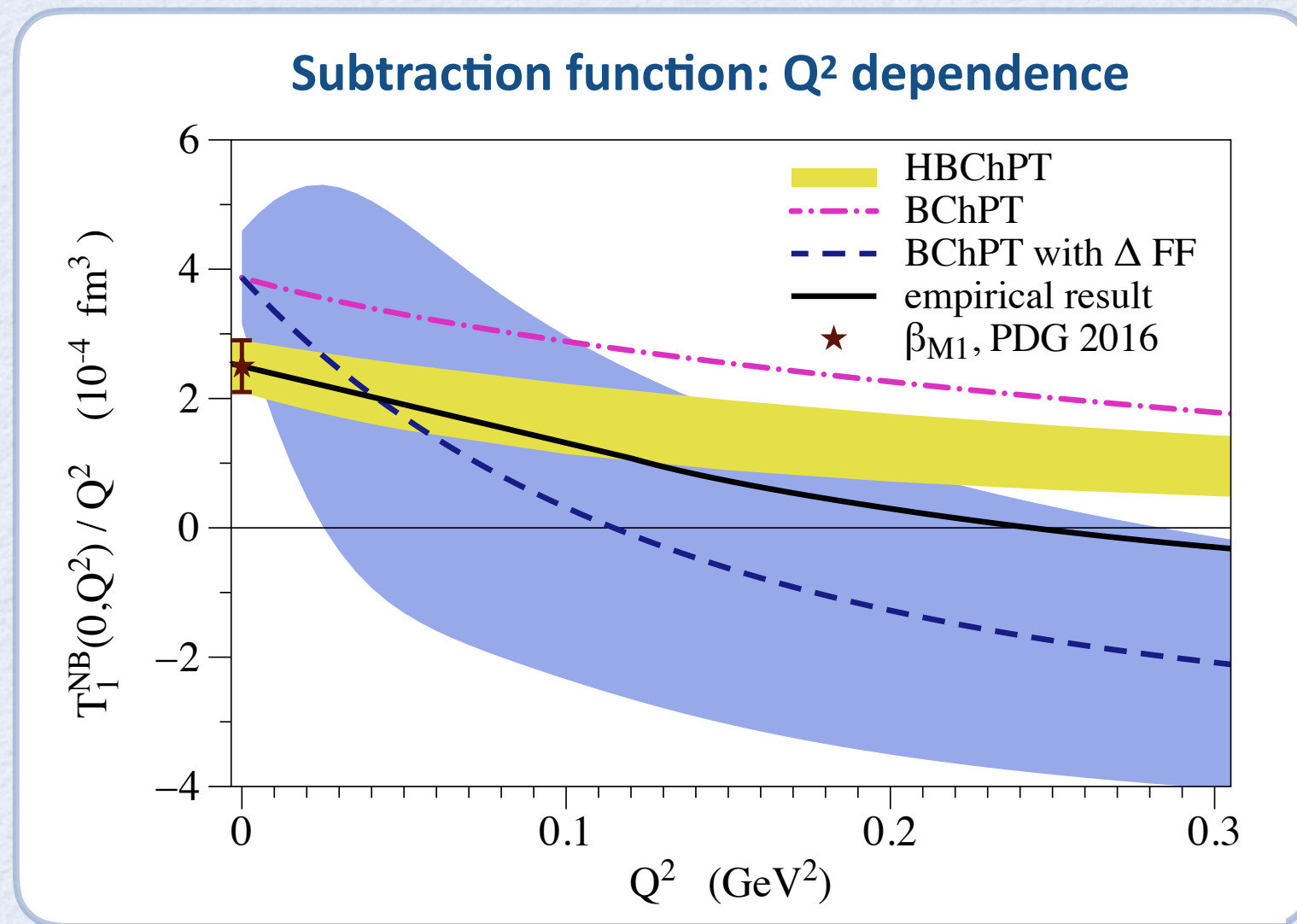
A reduction of error by factor of 2

→ precision comparable with proton

→ impact light muonic atom program

LS: Improved determinations of subtraction function

Lensky, Hagelstein,
Pascalutsa, Vdh
(2018)



To improve on uncertainty due to subtraction function: **3 avenues**

- Full NLO calculation in Baryon ChPT Pascalutsa et al.
- New formalism for lattice determination of subtraction function Hagelstein, Pascalutsa (2020)
- Empirical determination of Q^4 term using dilepton production process

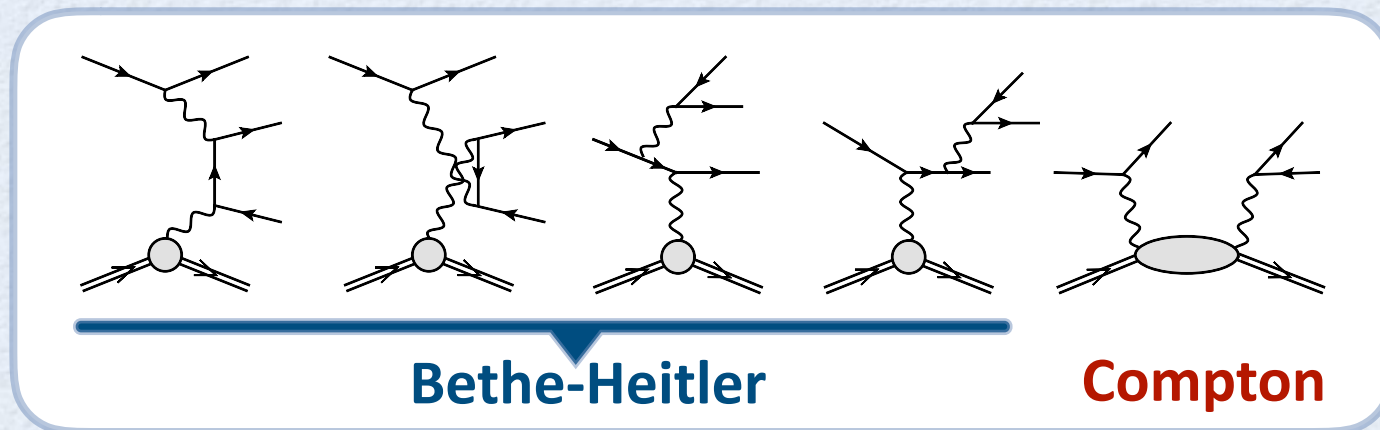
Empirical determination of Q^4 term through $e^- p \rightarrow e^- p l^- l^+$

$$T_1^{NB}(0, Q^2) = \beta_{M1} Q^2 + \left(\frac{1}{6} \beta_{M2} + 2\beta'_{M1} + \alpha_{em} b_{3,0} + \frac{1}{(2M)^2} \beta_{M1} \right) Q^4 + \mathcal{O}(Q^6)$$

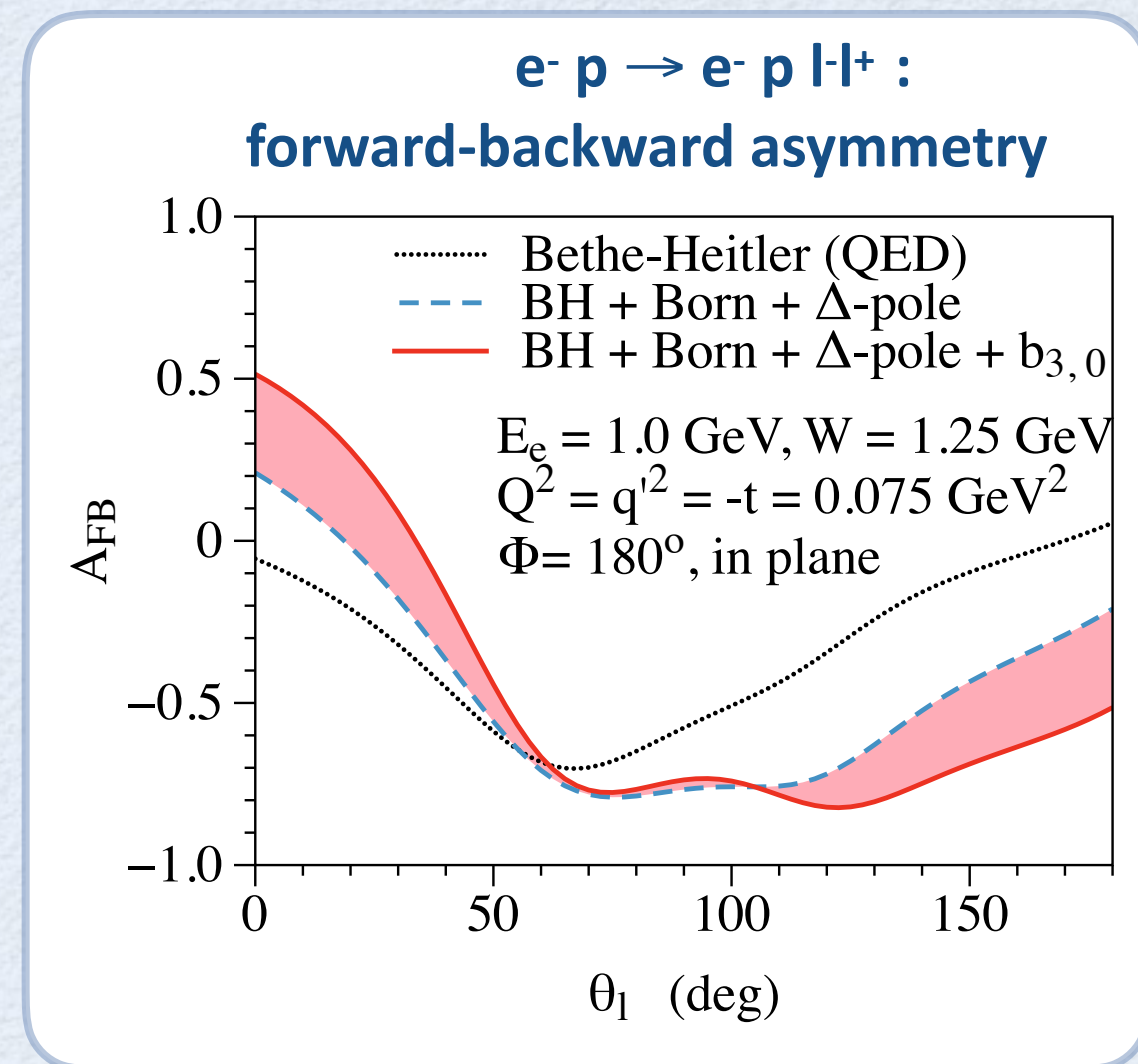
Subtraction
function T_1

only empirical
unknown

- Forward-backward asymmetry of $e^- p \rightarrow e^- p l^- l^+$:
allows access to low-energy structure constant $b_{3,0}$

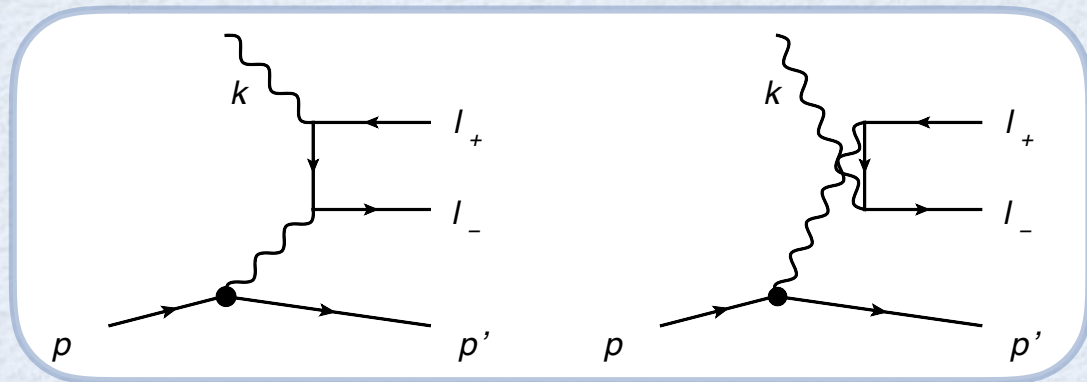


- First CLAS12@JLab data on high-energy timelike Compton (2020): FB asymmetry in 10 - 30% range

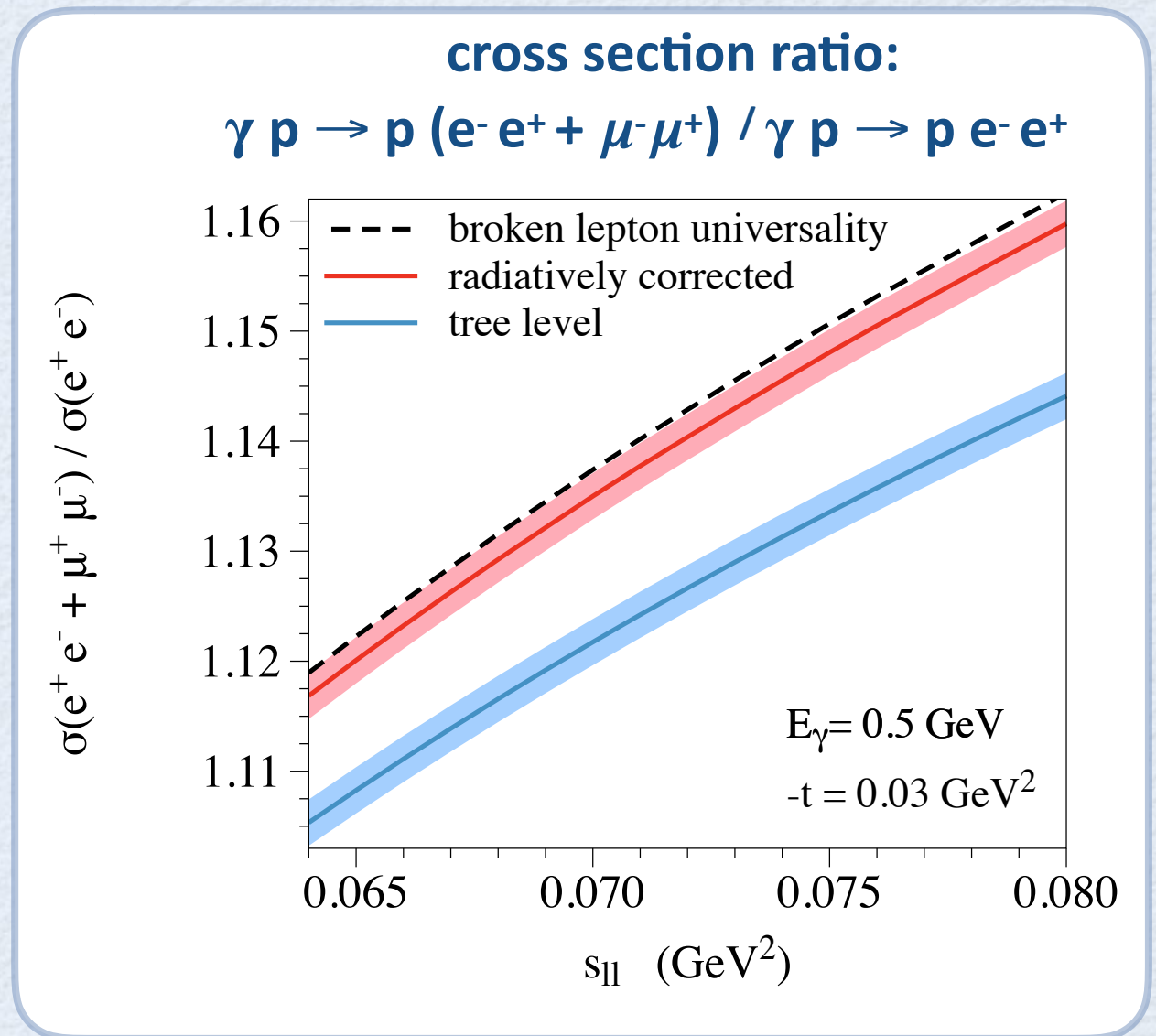


Pauk, Carlson, Vdh (2020)

Dilepton photoproduction $\gamma p \rightarrow p l^- l^+$: proton form factor, lepton universality



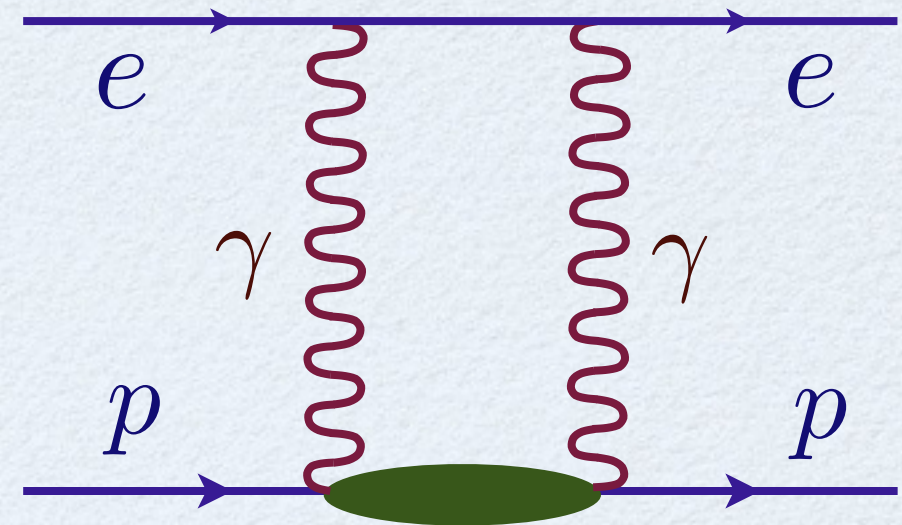
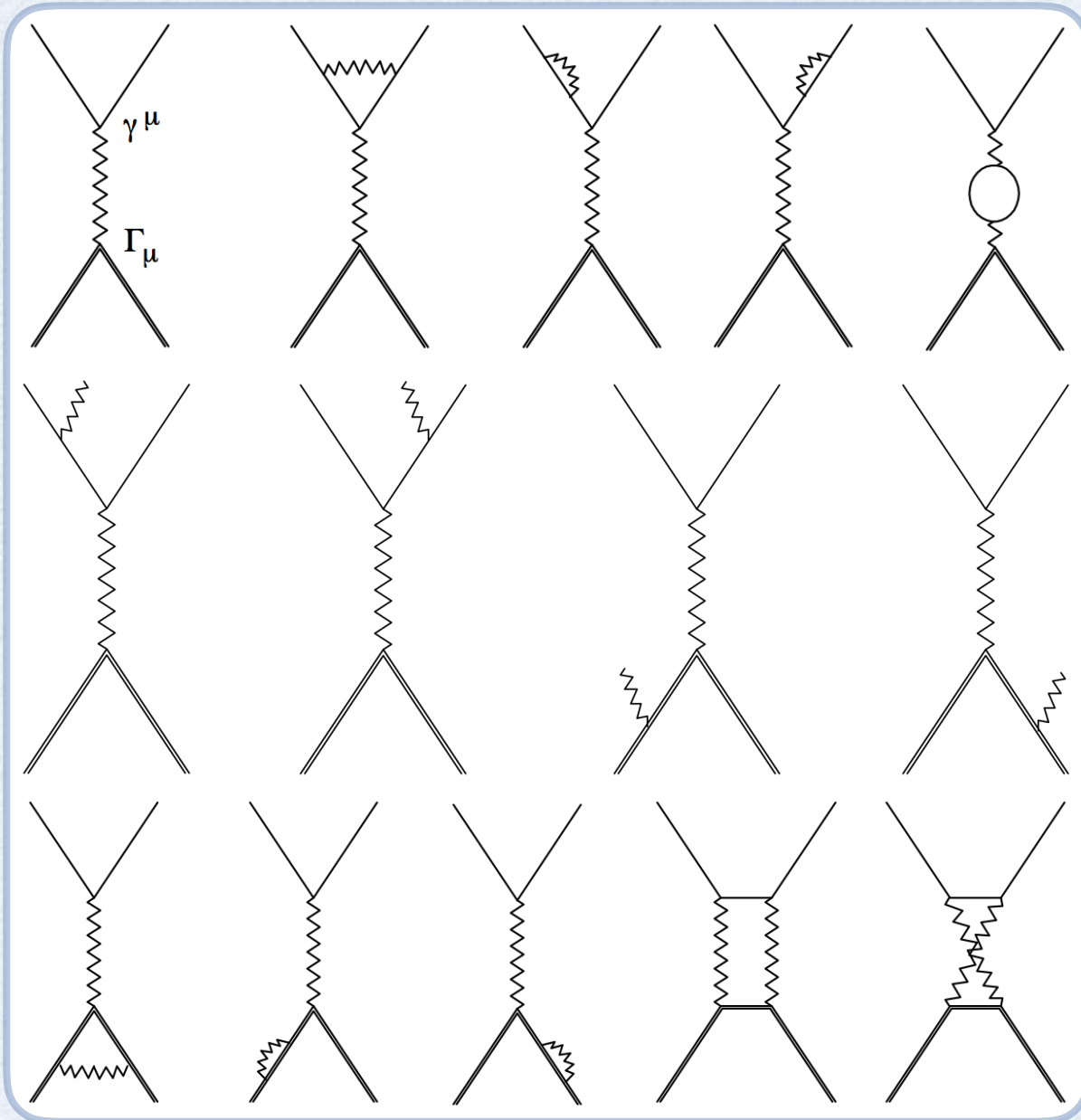
- e^-e^+ photoproduction using **active TPC + Crystal Ball @MAMI**
- full one-loop radiatively corrected amplitude needed
- e^-e^+ vs $\mu^-\mu^+$ photoproduction: cross section ratio above/below dimuon threshold : different proton radius values correspond to **0.2% effect** in ratio
- Complements ongoing muon scattering program (**MUSE, COMPASS**)



Pauk, Vdh (2015)

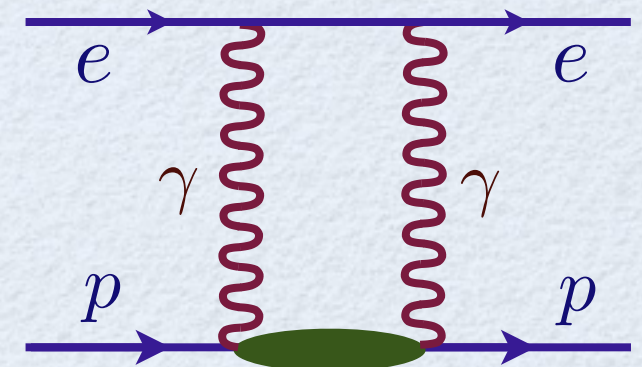
Heller, Tomalak, Wu, Vdh (2018, 2019)

Radiative corrections in lepton-nucleon scattering



status of radiative corrections

- ➔ Accuracy goal PRES: 0.2% absolute precision on $d\sigma/dt$
- ➔ For experiment with recoil proton detection: **Fadin & Gerasimov (2019)**
to 1-loop: real and virtual radiative corrections on electron side cancel
both for $\log(Q^2/m^2)$ terms and constant terms
beyond 1-loop: cancellation for log terms
- ➔ Vacuum polarization is dominant contribution: **A. Arbuzov**
Can be accurately calculated: 1 - 1.6 % in Q^2 range: 0.001 - 0.02 GeV^2
- ➔ Largest uncertainty from subdominant two-photon exchange between e and p



2γ-exchange at low Q²

$$\delta_{2\gamma} \sim a \sqrt{Q^2} + b Q^2 \ln Q^2 + c Q^2 \ln^2 Q^2$$

Feshbach
inelastic
elastic

McKinley, Feshbach (1948)

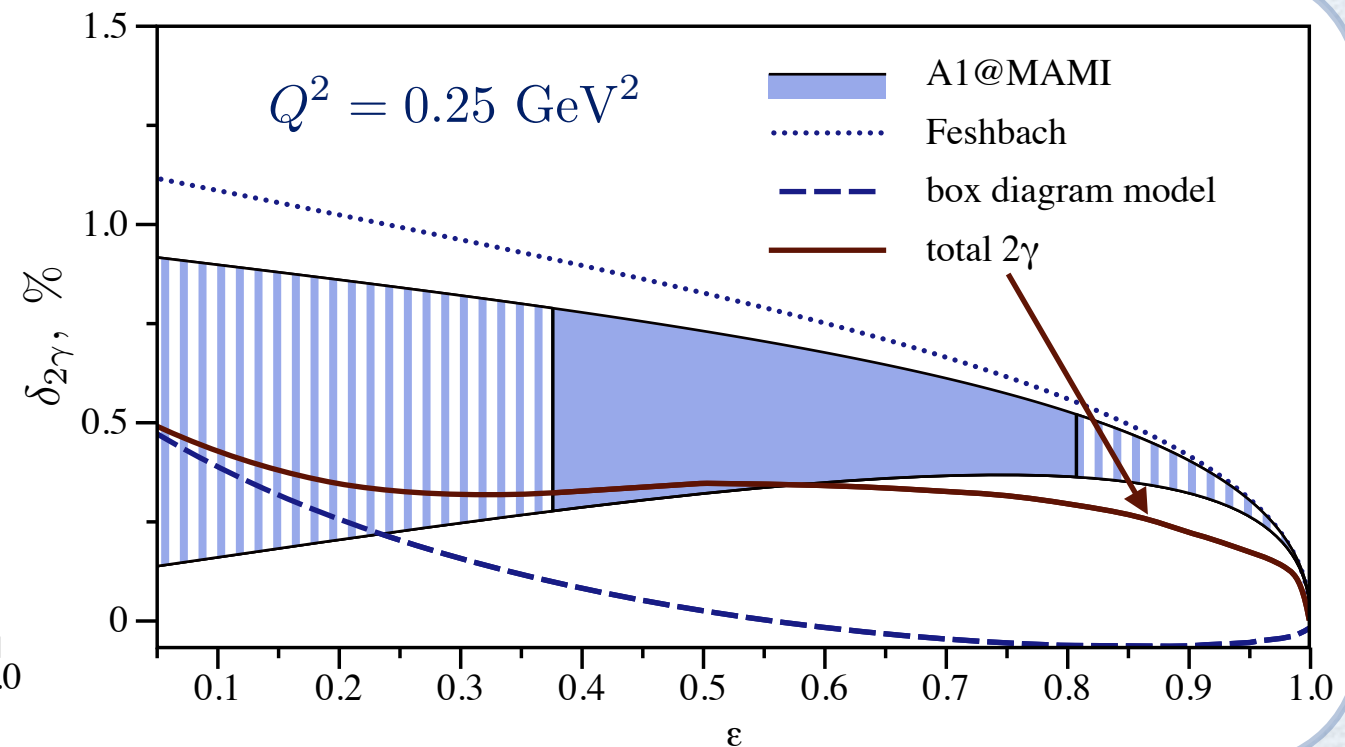
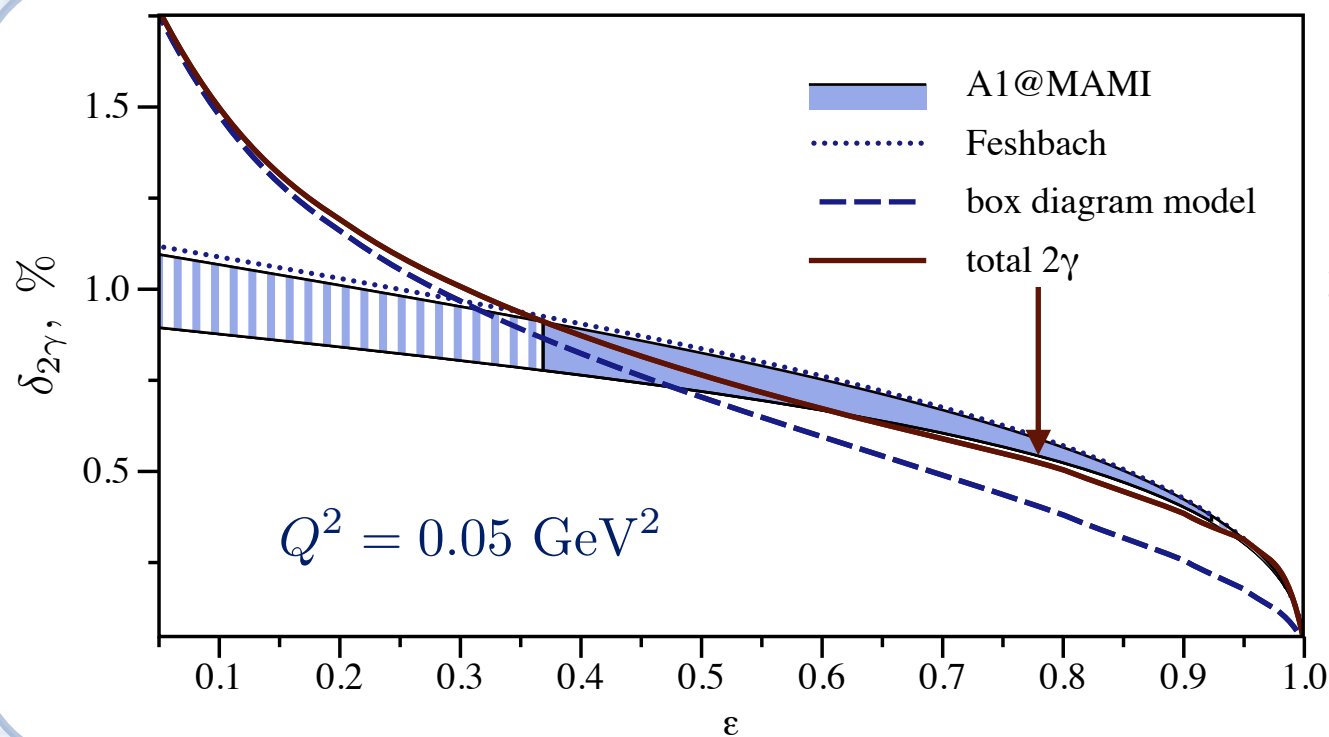
R.W. Brown (1970)

M. Gorchtein (2013)

$$\delta_{2\gamma} = \int d\nu_\gamma d\tilde{Q}^2 (\omega_1(\nu_\gamma, \tilde{Q}^2) \cdot F_1(\nu_\gamma, \tilde{Q}^2) + \omega_2(\nu_\gamma, \tilde{Q}^2) \cdot F_2(\nu_\gamma, \tilde{Q}^2))$$

unpolarized proton structure

Tomalak, Vdh (2016)

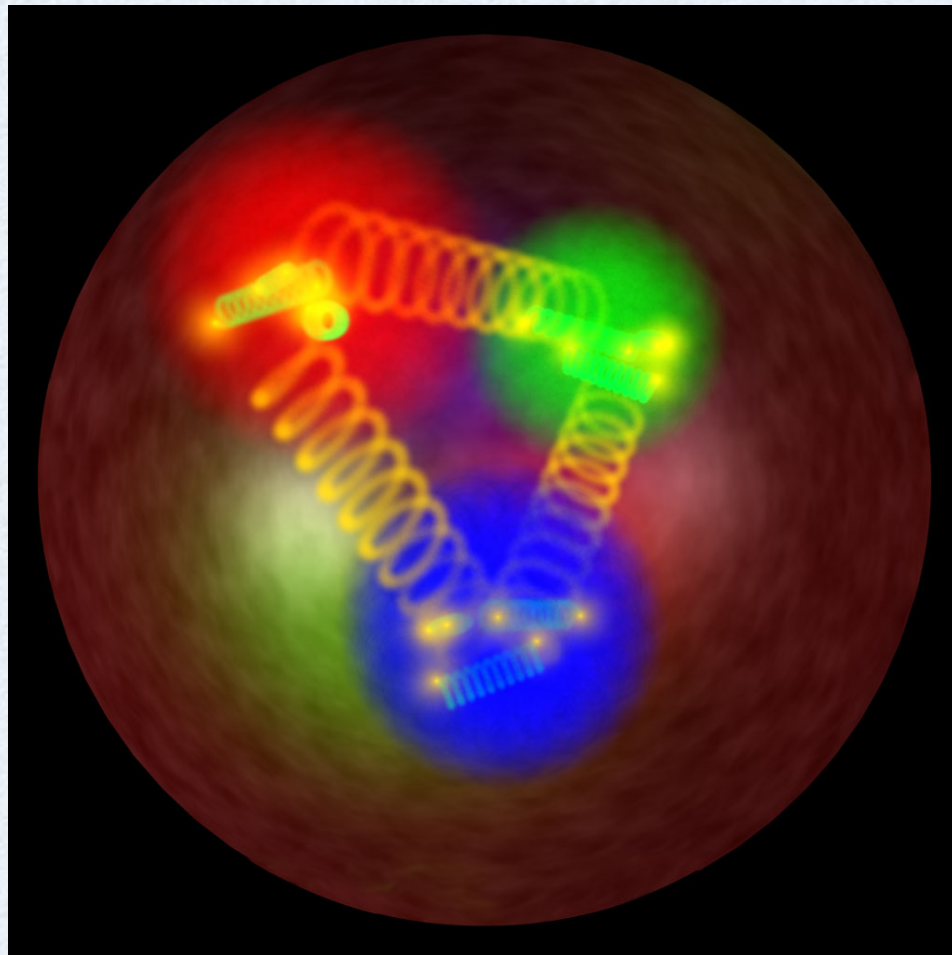


2γ at large ε agrees with empirical fit

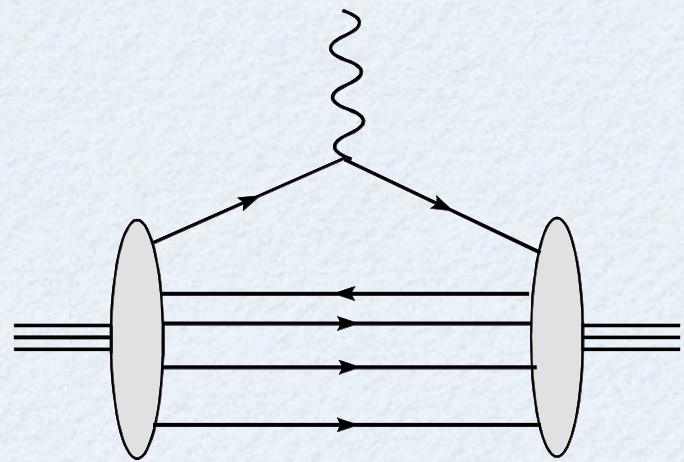
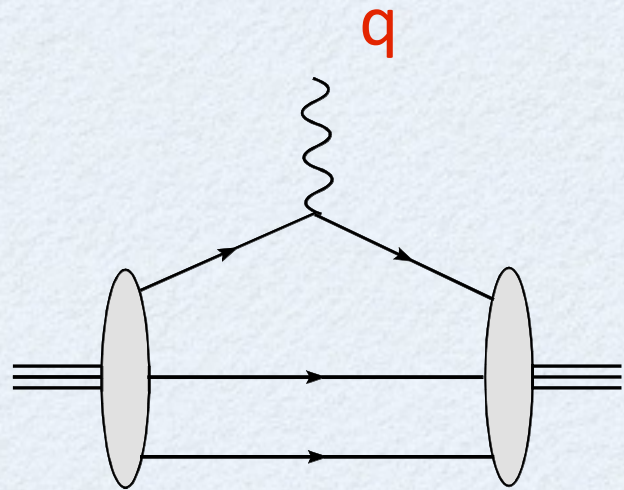
r_E extraction ✓

TPE varies between 0.05% to 0.2% in Q² range: 0.001 - 0.02 GeV²

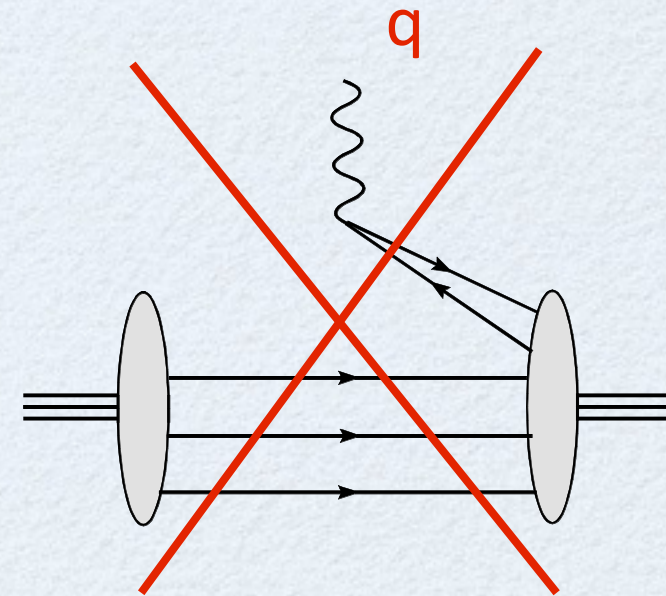
How to properly define a charge density for nucleons ?



Interpretation of nucleon form factor as density



overlap of wave function
Fock components
with **same** number of quarks



overlap of wave function
Fock components
with **different** number of quarks
NO probability / charge density
interpretation

absent in a light-front frame!

$$q^+ = q^0 + q^3 = 0$$

quark transverse charge densities in nucleon

transverse c.m. can be fixed in a light-front frame !

→ longitudinally polarized nucleon

$$\begin{aligned} \rho_0^N(\vec{b}) &\equiv \int \frac{d^2\vec{q}_\perp}{(2\pi)^2} e^{-i\vec{q}_\perp \cdot \vec{b}} \frac{1}{2P^+} \langle P^+, \frac{\vec{q}_\perp}{2}, \lambda | J^+(0) | P^+, -\frac{\vec{q}_\perp}{2}, \lambda \rangle \\ &= \int_0^\infty \frac{dQ}{2\pi} Q J_0(bQ) F_1(Q^2) \end{aligned}$$

Soper (1997)

Burkardt (2000)

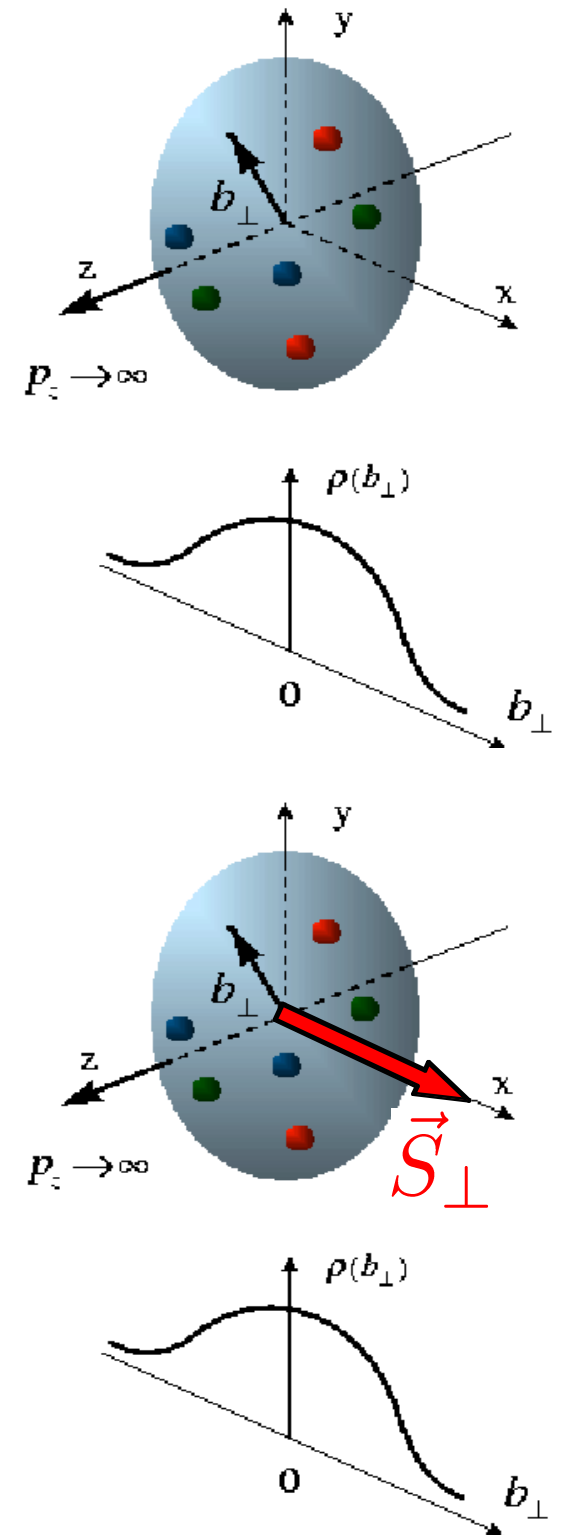
Miller (2007)

→ transversely polarized nucleon

$$\begin{aligned} \rho_T^N(\vec{b}) &\equiv \int \frac{d^2\vec{q}_\perp}{(2\pi)^2} e^{-i\vec{q}_\perp \cdot \vec{b}} \frac{1}{2P^+} \langle P^+, \frac{\vec{q}_\perp}{2}, s_\perp = +\frac{1}{2} | J^+(0) | P^+, -\frac{\vec{q}_\perp}{2}, s_\perp = +\frac{1}{2} \rangle \\ &= \rho_0^N(b) + \sin(\phi_b - \phi_S) \int_0^\infty \frac{dQ}{2\pi} \frac{Q^2}{2M} J_1(bQ) F_2(Q^2) \end{aligned}$$

dipole field pattern

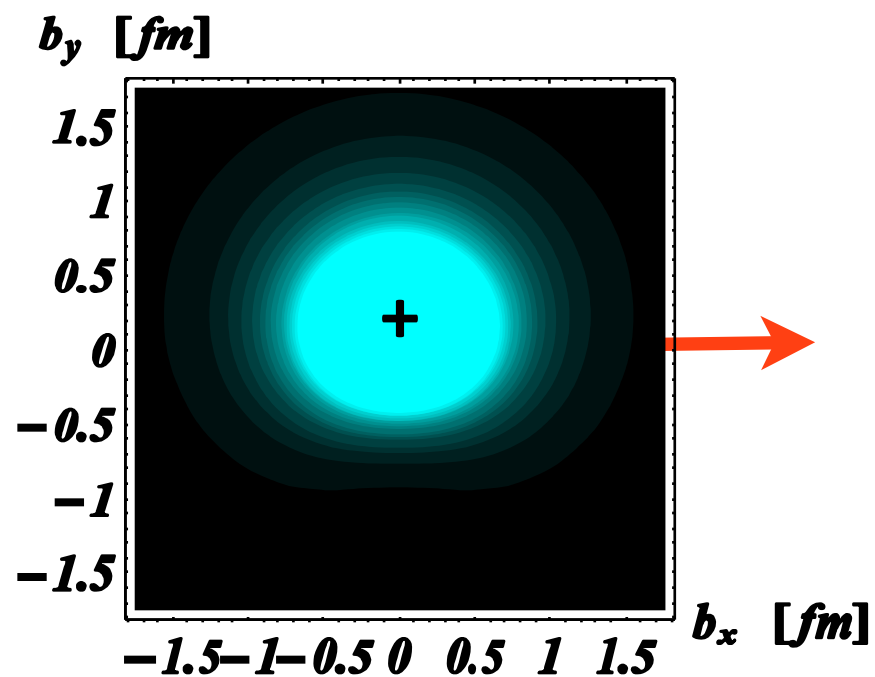
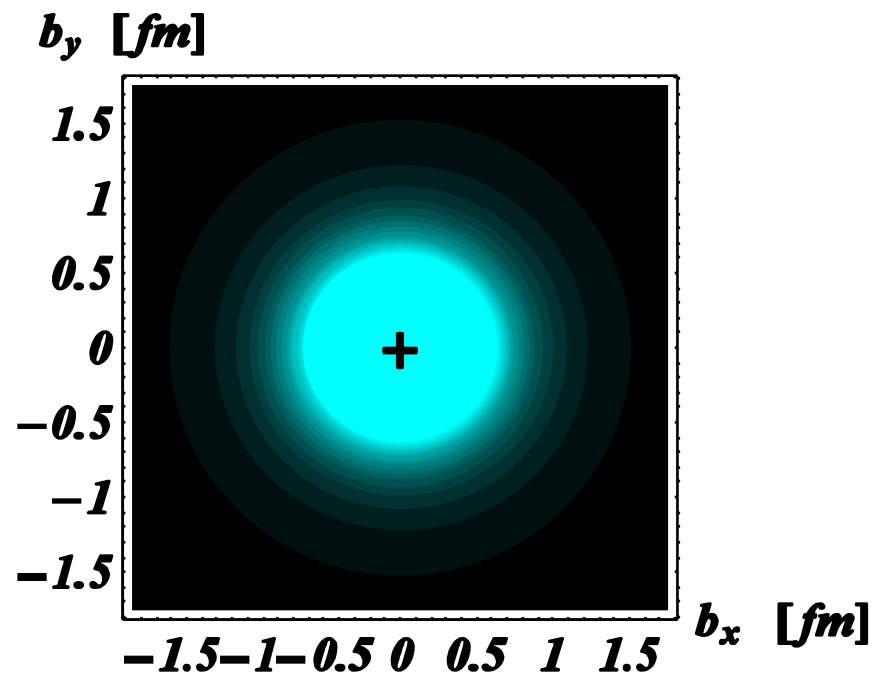
Carlson, Vdh (2007)



spatial imaging of nucleons

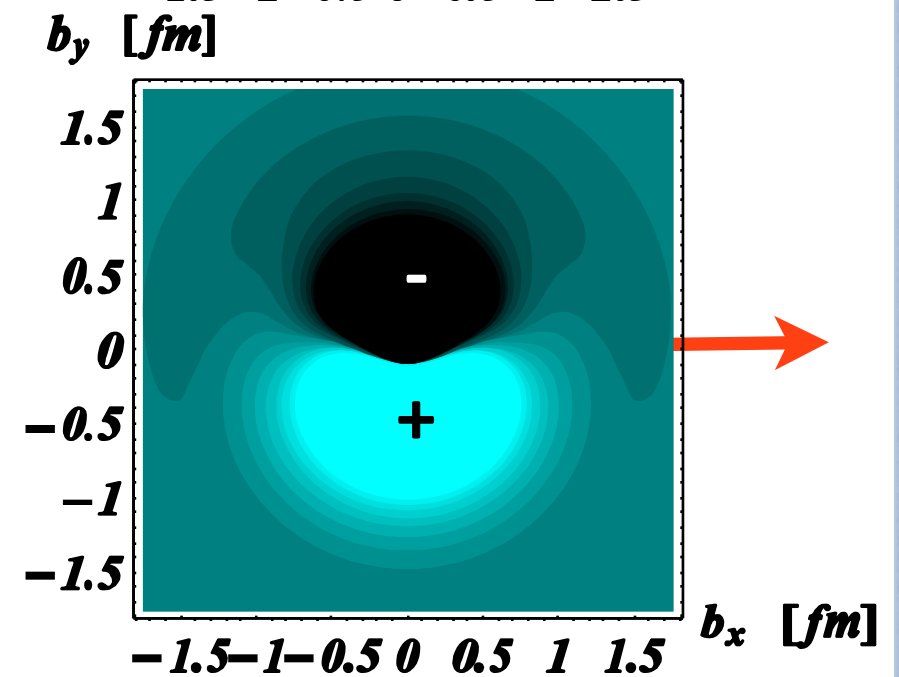
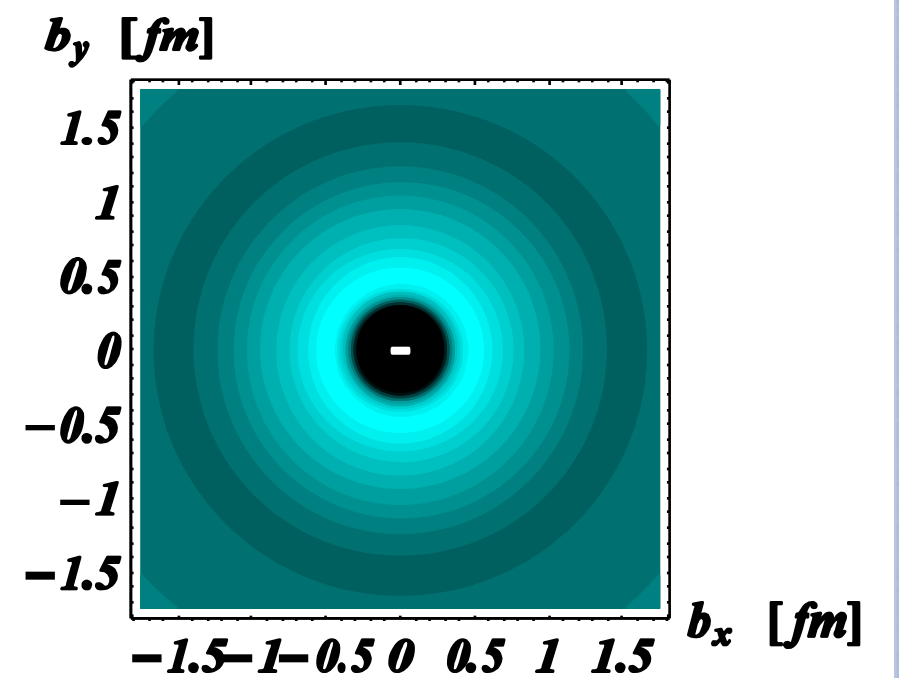
proton

neutron



induced
electric dipole
moment:

$$d_y = \kappa \frac{e}{2M}$$



Miller (2007)

Carlson, Vdh (2007)

Radii of charge distribution for nucleons

➔ Radius of 2-dim transverse distribution of quarks of flavor q in **proton**:

$$\langle b^2 \rangle^q = \frac{\int d^2\mathbf{b} \, b^2 \, \rho^q(b)}{\int d^2\mathbf{b} \, \rho^q(b)} = -4 \frac{F_1'^q(0)}{F_1^q(0)}$$

with: $F_1'^q(0) \equiv \left. \frac{dF_1^q}{dQ^2} \right|_{Q^2=0}$

➔ Isospin symmetry:

$$F_{1p} = e_u F_1^u + e_d F_1^d$$

$$F_{1n} = e_u F_1^d + e_d F_1^u$$

$$\langle b^2 \rangle_p = \frac{4}{3} \langle b^2 \rangle^u - \frac{1}{3} \langle b^2 \rangle^d = -4F_{1p}'(0)$$

$$\langle b^2 \rangle_n = \frac{2}{3} \langle b^2 \rangle^d - \frac{2}{3} \langle b^2 \rangle^u = -4F_{1n}'(0)$$

➔ Charge radii for proton and neutron:

		r_E^2 (fm ²)	$-\frac{3\kappa_N}{2M^2}$ (fm ²)	$-6F_1'(0)$ (fm ²)	$\langle b^2 \rangle$ (fm ²)
Cui et al. (2021)	proton (e-p)	0.717 ± 0.014		0.598 ± 0.014	0.399 ± 0.009
Antognini et al. (2013)	proton (μ H)	0.7071 ± 0.0007	-0.1189	0.5882 ± 0.0007	0.3921 ± 0.0005
	neutron (PDG)	-0.1161 ± 0.0022	0.1266	0.0105 ± 0.0022	0.0070 ± 0.0015

**H. Gao, M. Vdh
(2021)**

Summary and outlook



electron scattering:

- new PRad proton r_E result is consistent with muonic hydrogen spectroscopy value
- PRad-II aims at factor 3.8 improvement on proton charge radius r_E
- Plans in Mainz: PRES, MAGIX@MESA



hadronic corrections to Lamb shift in muonic atoms:

- **μH** : present TPE accuracy (2 μeV) is comparable with present Lamb shift accuracy (2.3 μeV)
Future plan: factor 5 improvement on LS for muonic H
- **μD , $\mu\text{}^3\text{He}^+$, $\mu\text{}^4\text{He}^+$** : present TPE accuracy is 5-10 times worse than Lamb shift accuracy
Largest uncertainty from neutron polarizability correction



muon scattering experiments:

- Ongoing plans: MUSE@PSI, AMBER@COMPASS
- dilepton production complementary: cross section ratio above/below dimuon threshold



Theoretical understanding:

For system as proton composed of relativistic quarks and gluons:

- 3-dim charge distribution cannot be defined precisely: r_E is definition of G_E slope at origin !
- instead we can define a 2-dim transverse charge distribution
as seen by observer which moves with speed of light, and which is obtained as integral over quark longitudinal momentum fraction x of partonic distributions
- we can properly define and extract the proton charge radius in this transverse plane