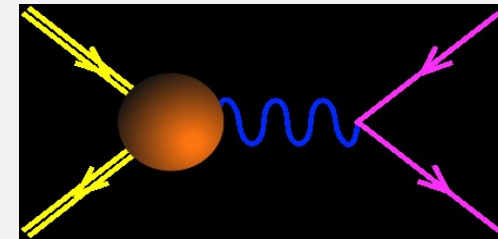
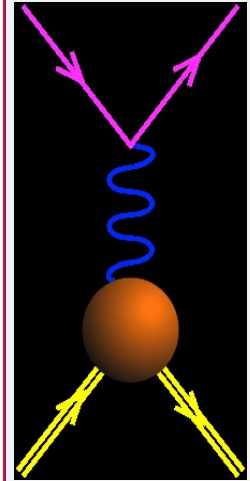


The Proton Heart

Egle Tomasi-Gustafsson

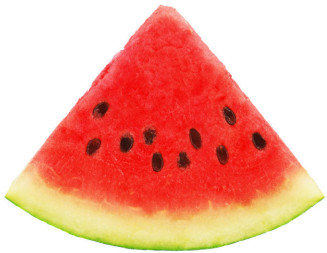
*CEA, IRFU, DPhN and
Université Paris-Saclay
France*



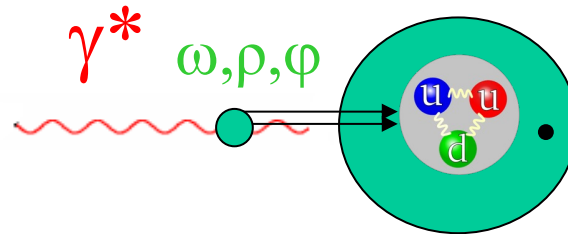
Mainz, June 19-23, 2023



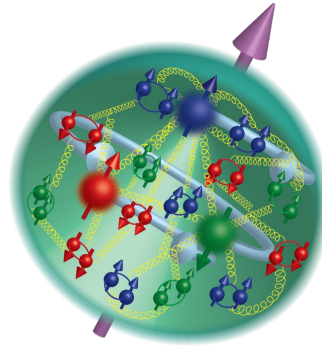
If the proton were a fruit...



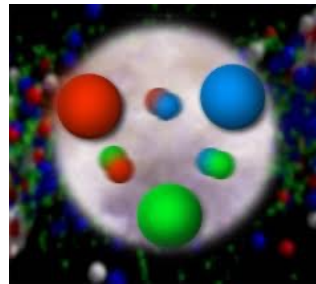
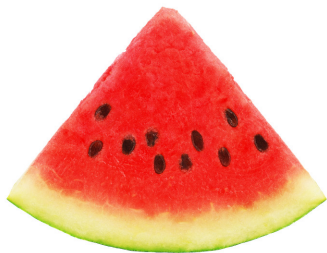
If the proton were a fruit...



VDM :
vector meson dominance



- TMD, GPD...
Parton structure functions



- Instantons:
Mostly Vacuum

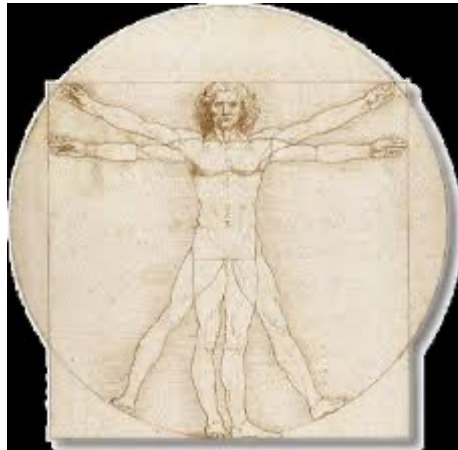


The proton

Proton is the the most common constituent of visible matter...



...BUT...



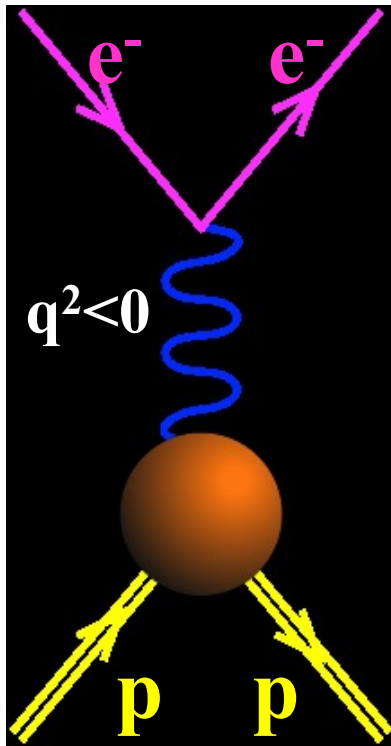
....its fundamental properties as

- Mass
- Spin
- Size

are still object of controversy



Electromagnetic Interaction



The electron vertex is known, γ_μ

The interaction is carried by a virtual photon of 4-mom q^2

The proton vertex is parametrized in terms of FFs: Pauli and Dirac F_1, F_2

$$\Gamma_\mu = \gamma_\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}q^\nu}{2M} F_2(q^2)$$

$$q^2 = -4E_1 E_2 \sin^2 \theta / 2$$

or in terms of Sachs FFs:

$$G_E = F_1 + \tau F_2, \quad G_M = F_1 + F_2, \quad \tau = q^2 / 4M^2$$

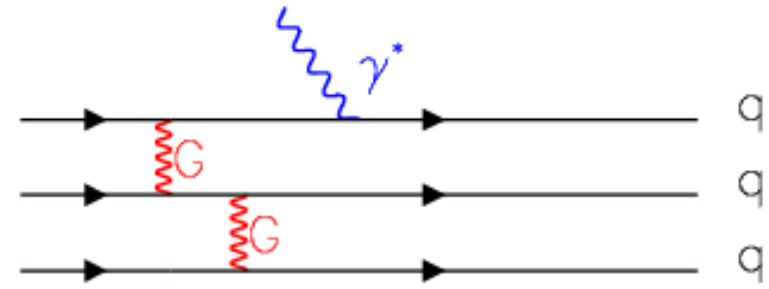
$$G_E(0) = 1(e) \quad G_M(0) = \mu_N$$

What about high order radiative corrections?



Dipole Approximation & pQCD

Dimensional scaling



- $F_n(Q^2) = C_n [1/(1+Q^2/m_n)^{n-1}]$,
 - $m_n = n\beta^2$, $\langle \text{quark momentum squared} \rangle$
 - n is the number of constituent quarks
- Setting $\beta^2 = (0.471 \pm 0.010) \text{ GeV}^2$ (fitting pion data)
 - pion: $F_\pi(Q^2) = C_\pi [1/(1+Q^2/0.471 \text{ GeV}^2)^1]$,
 - nucleon: $F_N(Q^2) = C_N [1/(1+Q^2/0.71 \text{ GeV}^2)^2]$,
 - deuteron: $F_d(Q^2) = C_d [1/(1+Q^2/1.41 \text{ GeV}^2)^5]$

V. A. Matveev, R. M. Muradian, and A. N. Tavkhelidze (1973), Brodsky and Farrar (1973), Politzer (1974), Chernyak & Zhitnisky (1984), Efremov & Radyuskin (1980)...



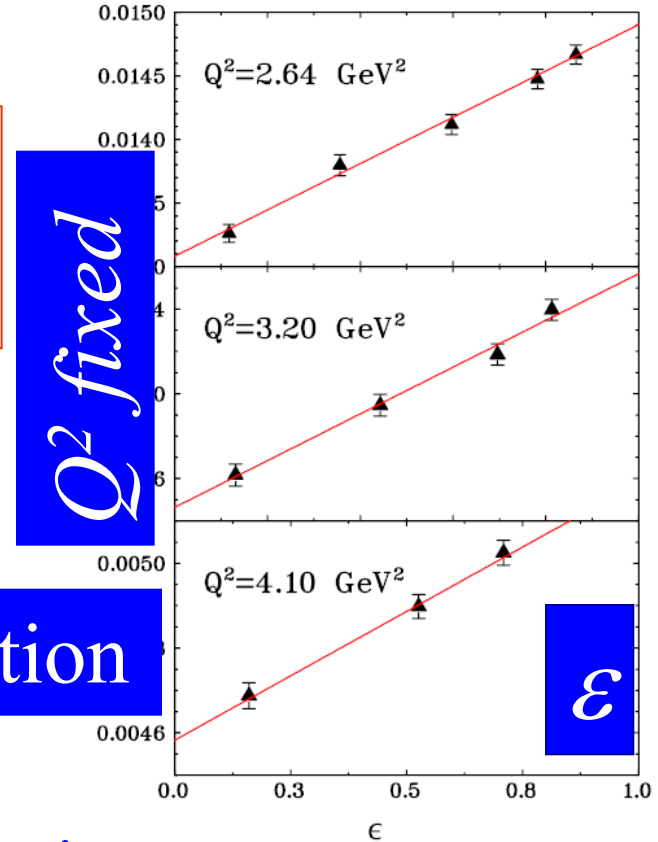
ep-elastic scattering : Rosenbluth separation

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \frac{1}{(1+\tau)} \left(G_E^2(Q^2) + \frac{\tau}{\varepsilon} G_M^2(Q^2) \right)$$

1950

$$\varepsilon = \left(1 + 2(1+\tau) \tan^2 \left(\frac{\theta_e}{2} \right) \right)^{-1}, \quad \tau = \frac{Q^2}{4M^2}$$

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2$$



PRL 94, 142301 (2005)

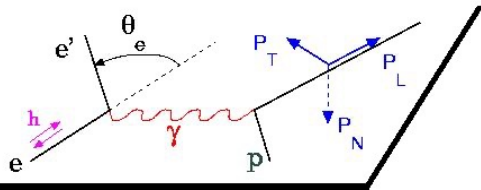
Linearity of the reduced cross section

→ $\tan^2 \theta_e$ dependence

→ Holds for 1γ exchange only

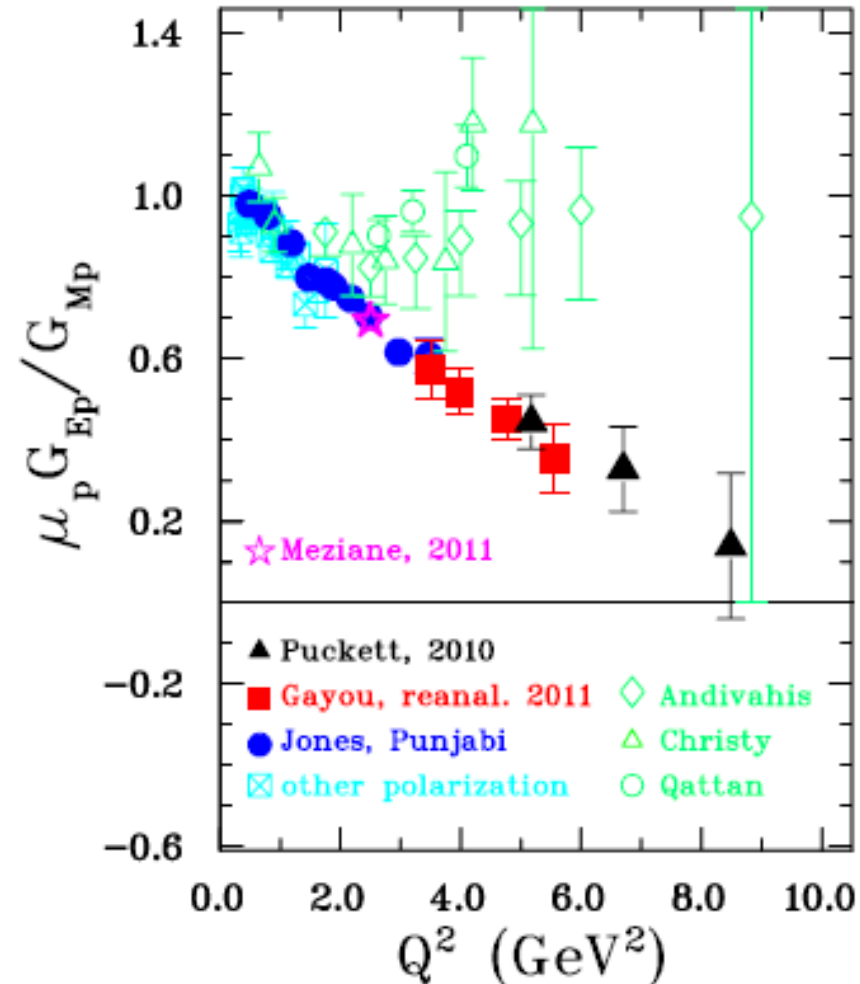


The Akhiezer-Rekalo recoil proton polarization- Method (1967) GEp Experiments (>2000)



Jlab-GEp collaboration (>2000)

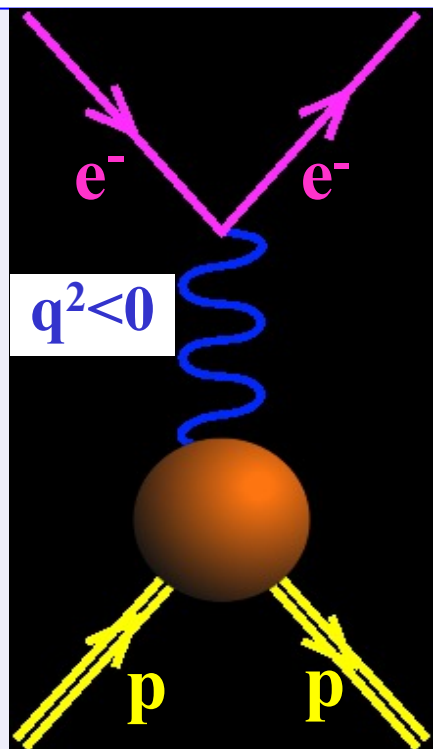
- 1) "standard" **dipole function** for the nucleon magnetic FFs **GMp** and **GMn**
- 2) **linear deviation** from the dipole function for the electric proton FF **Gep**
- 3) **QCD scaling** not reached
- 3) **Zero crossing** of Gep?
- 4) **contradiction between polarized and unpolarized measurements**



A.J.R. Puckett et al, Phys. Rev. C96, 055203 (2017).



Proton Charge and Magnetic Distributions



$$G_E(0) = 1$$

$$G_M(0) = \mu_p$$

*Space-like
FFs are real*

$$e + p \rightarrow e + p$$

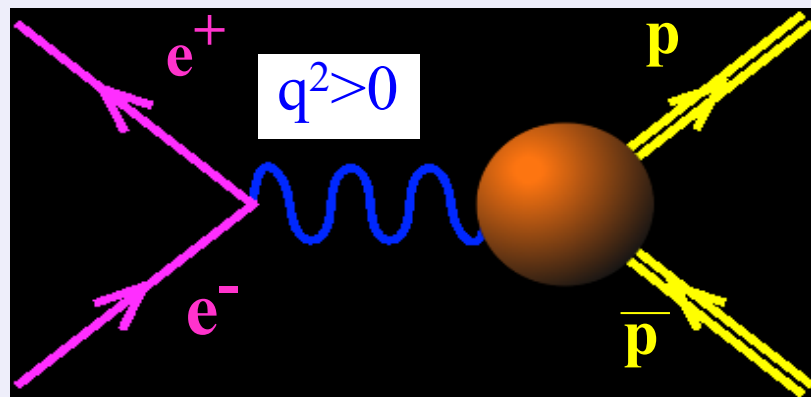
0

$$4m_p^2$$

$$G_E = G_M$$

Unphysical region
 $p + \bar{p} \leftrightarrow e^+ e^- + \pi^0$

Crossing symmetry
Asymptotics
- QCD
- analyticity



*Time-Like
FFs are complex*

$$p + \bar{p} \leftrightarrow e^+ + e^-$$

q^2



Time-like observables: $|G_E|^2$ and $|G_M|^2$.

-The cross section for $\bar{p} + p \rightarrow e^+ + e^-$ (1 γ -exchange):

$$\frac{d\sigma}{d(\cos\theta)} = \frac{\pi\alpha^2}{8m^2\sqrt{\tau-1}} [\tau|G_M|^2(1 + \cos^2\theta) + |G_E|^2\sin^2\theta]$$

θ : angle between e^- and \bar{p} in cms.

A. Zichichi, S. M. Berman, N. Cabibbo, R. Gatto, Il Nuovo Cimento XXIV, 170 (1962)

B. Bilenkii, C. Giunti, V. Wataghin, Z. Phys. C 59, 475 (1993)

G. Gakh, E.T-G., Nucl. Phys. A761,120 (2005)

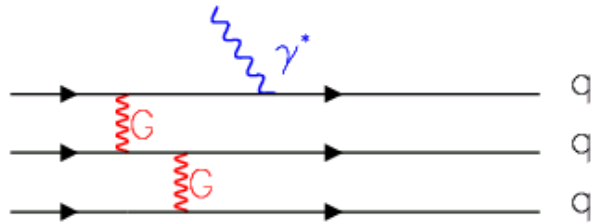
As in SL region:

- Dependence on q^2 contained in FFs
- Even dependence on $\cos^2\theta$ (1 γ exchange)
- No dependence on sign of FFs
- Enhancement of the magnetic term

but TL form factors are complex!



The Time-like Region

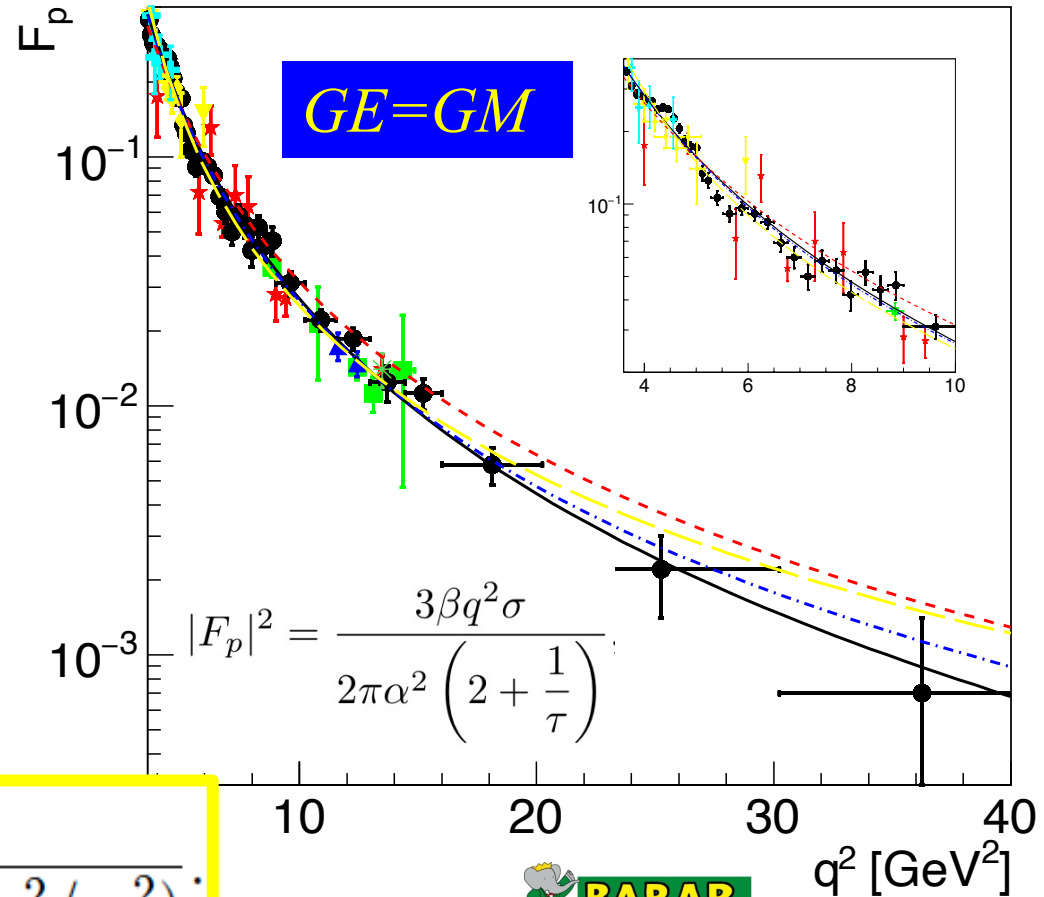
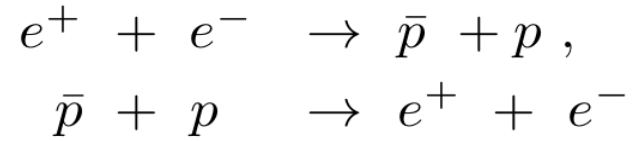


Expected QCD scaling $(q^2)^2$

$$\frac{A}{(q^2)^2 [\log^2(q^2/\Lambda^2) + \pi^2]}$$

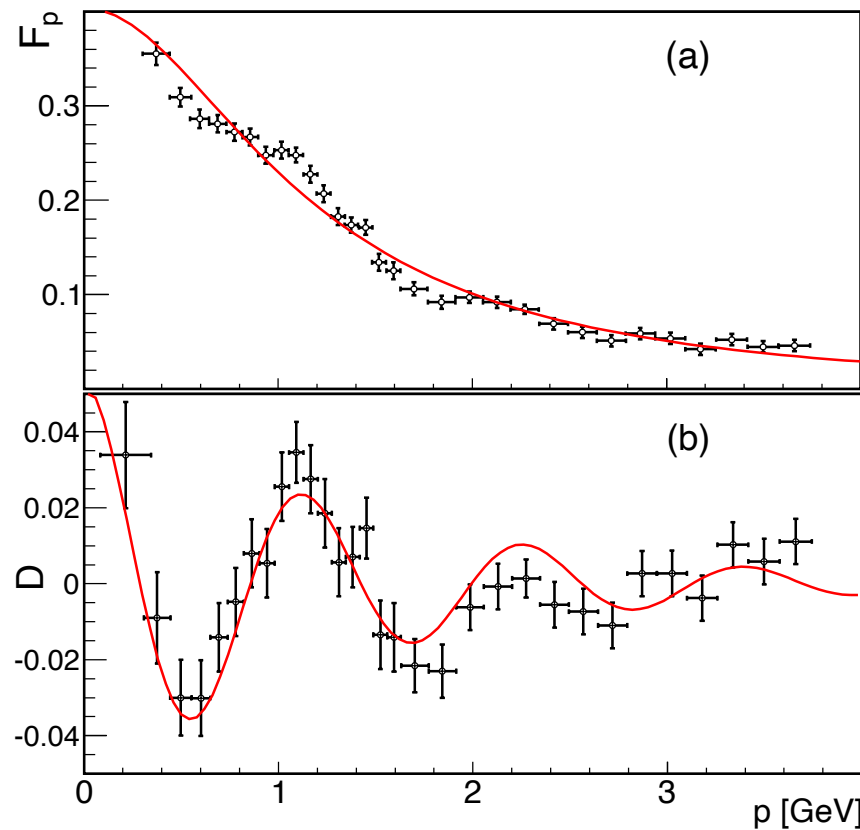
$$\frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2}$$

$$|F_{T3}(q^2)| = \frac{A}{(1 - q^2/m_1^2)(2 - q^2/m_2^2)}$$



Oscillations : regular pattern in p_{Lab}

The relevant variable is p_{Lab} associated to the relative motion of the final hadrons



$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$

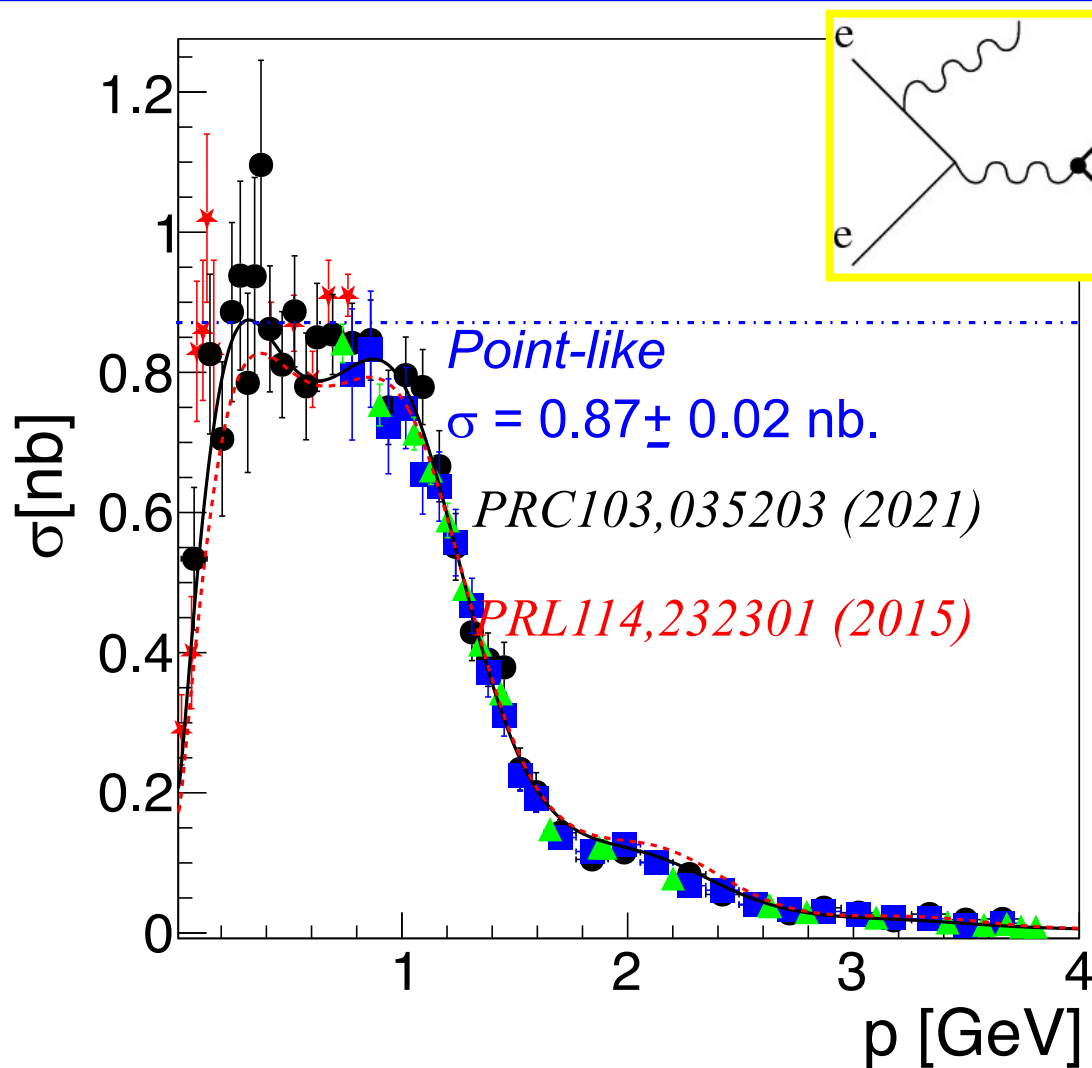
$A \pm \Delta A$	$B \pm \Delta B$	$C \pm \Delta C$	$D \pm \Delta D$	$\chi^2/n.d.f$
	$[GeV]^{-1}$	$[GeV]^{-1}$		
0.05 ± 0.01	0.7 ± 0.2	5.5 ± 0.2	0.03 ± 0.3	1.2

A: Small perturbation B: damping
C: $r < 1\text{fm}$ D=0: maximum at $p=0$

Simple oscillatory behaviour
Small number of coherent sources

A. Bianconi, E. T-G. Phys. Rev. Lett. 114,232301 (2015)

Cross section from $e^+e^- \rightarrow p\bar{p} (\gamma)$



Novosibirsk 38pt

$1.9 < 2E < 4.5$

PLB794,64 (2019)

BaBar 85pt

$1.9 < 2E < 4.5$

PRD87,092005 (2013)

ISR-ISR-SA 30pt

$2 < 2E < 3.6$

PRD99,092002 (2019)

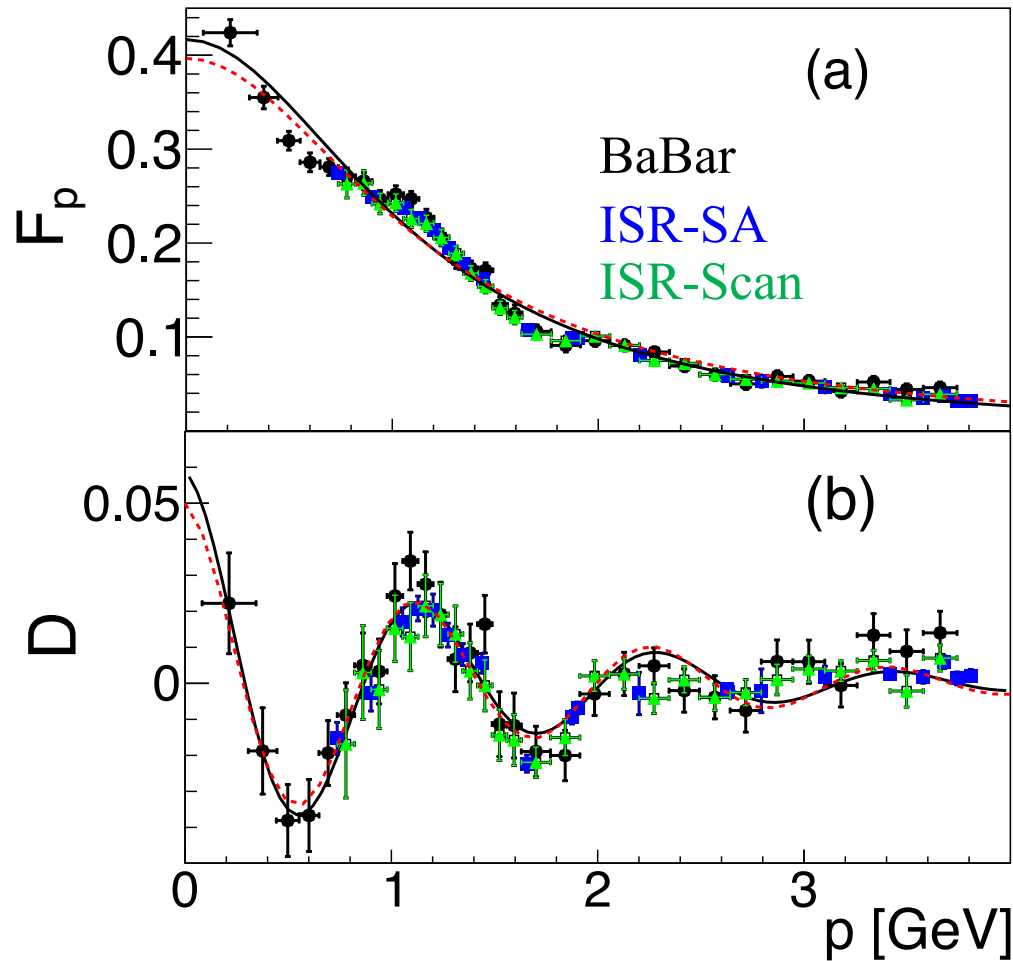
ISR-Scan 22pt

$2 < 2E < 3.1$

PRL124,042001 (2020)

E.T.-G., A. Bianconi, S. Pacetti, Phys.Rev.C 103 (2021) 3, 035203

Confirmation of regular oscillations



$$F_p^{\text{fit}}(s) = F_{3p}(s) + F_{\text{osc}}(p(s))$$

$$F_{3p}(s) = \frac{F_0}{\left(1 + \frac{s}{m_a^2}\right) \left(1 - \frac{s}{m_0^2}\right)^2},$$

$$F_{\text{osc}}(p(s)) = Ae^{-Bp} \cos(Cp + D).$$

$$s = 2m_p \left(m_p + \sqrt{p^2 + m_p^2} \right),$$

$$p = \sqrt{s \left(\frac{s}{4m_p^2} - 1 \right)}.$$

E.T.-G., A. Bianconi, S. Pacetti, *Phys.Rev.C* 103 (2021) 3, 035203

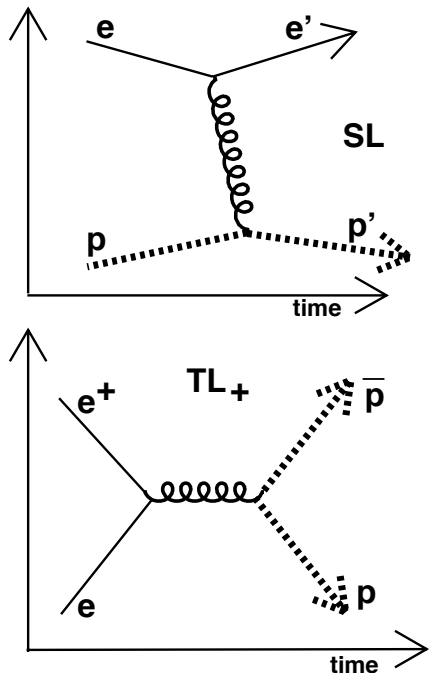


Time- and Space-Like Form Factors

Unified definition

$$F(q^2) = \int_{\mathcal{D}} d^4x e^{iq_\mu x^\mu} \rho(x), \quad q_\mu x^\mu = q_0 t - \vec{q} \cdot \vec{x}$$

$\rho(x) = \rho(\vec{x}, t)$ space-time distribution of the electric charge in the space-time volume \mathcal{D} .



SL photon 'sees' a charge density
TL photon can not test a space distribution but sees the time evolution from the annihilation point to the formed hadron

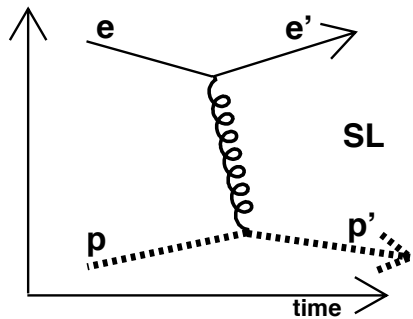


Time- and Space-Like Form Factors

Unified definition

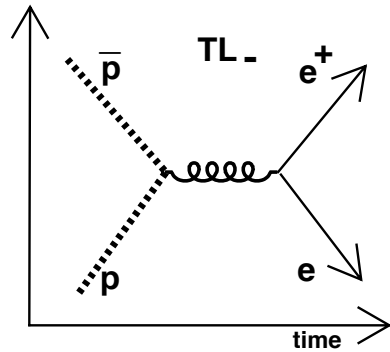
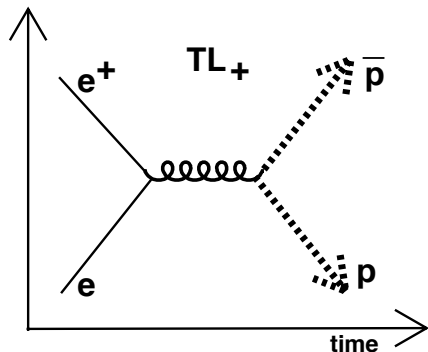
$$F(q^2) = \int_{\mathcal{D}} d^4x e^{iq_\mu x^\mu} \rho(x), \quad q_\mu x^\mu = q_0 t - \vec{q} \cdot \vec{x}$$

$\rho(x) = \rho(\vec{x}, t)$ space-time distribution of the electric charge in the space-time volume \mathcal{D} .



What do we measure? Projections!

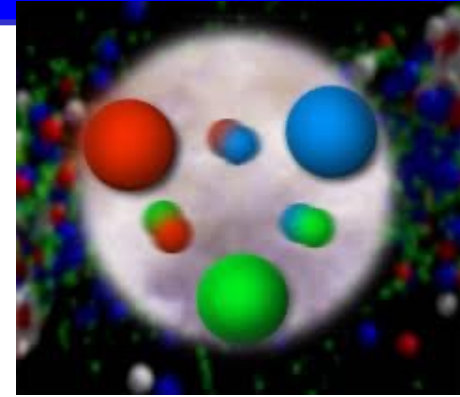
-on the space-axis in Breit system
-on the time-axis in CMS



How to connect and understand the amplitudes?



The nucleon



3 valence quarks and a neutral sea of $q\bar{q}$ pairs

Antisymmetric state of colored quarks:

$$\begin{aligned} |p\rangle &\sim \epsilon_{ijk} |u^i u^j d^k\rangle \\ |n\rangle &\sim \epsilon_{ijk} |u^i d^j d^k\rangle \end{aligned}$$

Assumption:

Does not hold in the spatial center of the nucleon: the center of the nucleon *is electrically neutral*, due to strong gluonic field
Inner region: gluonic condensate of clusters with randomly oriented chromo-magnetic field (Vainshtein, 1982): no color degree of freedom!

Charge screening as in a plasma

E.A. Kuraev, E. T-G, A. Dbeyssi, Phys.Lett. B712 (2012) 240



Predictions for SL and TL

Quark counting rules apply to the vector part of the potential

$$G_M^{(p,n)}(Q^2) = \mu G_E(Q^2);$$

$$G_E^{(p,n)}(Q^2) = G_D(Q^2) = \left[1 + Q^2/(0.71 \text{ GeV}^2)\right]^{-2}$$

The neutral plasma acts on the distribution of the electric charge (not magnetic).

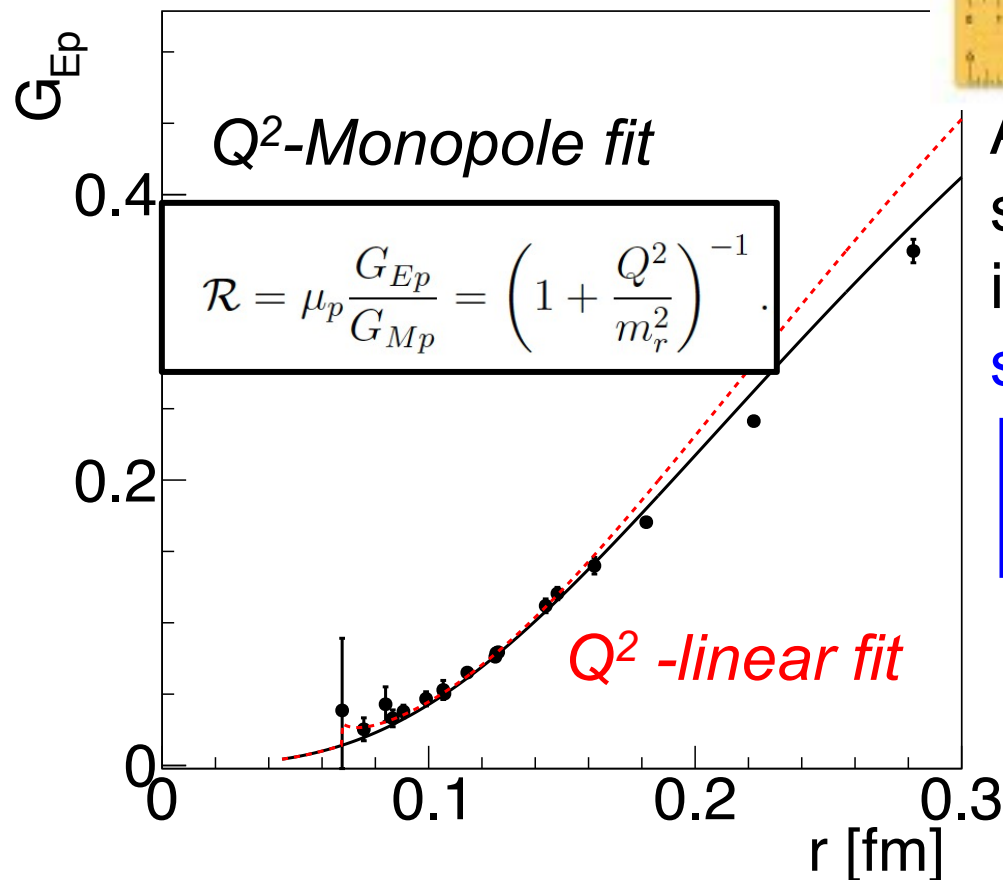
Additional suppression due to the **neutral plasma**

$$G_E(Q^2) = \frac{G_M(Q^2)}{\mu} \left(1 + Q^2/q_1^2\right)^{-1}$$

Similar behavior in SL and TL regions



SL- the most precise ruler



Additional suppression for the scalar part due to colorless internal region: “charge screening as in a plasma”:

$$G_E(Q^2) = \frac{G_M(Q^2)}{\mu} \left(1 + Q^2/q_1^2\right)^{-1}$$

Zero crossing?

Prediction: NO

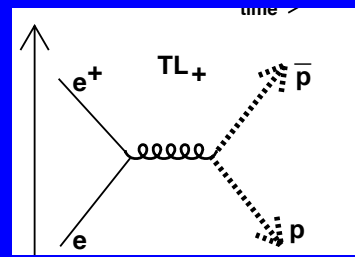
The photon ‘sees’ the neutral, screened region

$G_{Ep} \approx 0$ for $r < 0.06$ fm

$$r \text{ [fm]} = \lambda = \hbar c / \sqrt{Q^2} = 0.197 \text{ [GeV fm]} / \sqrt{Q^2 \text{ [GeV]}},$$

E. T-G., S. Pacetti, Phys. Rev. C 106 (2022) 3, 035203





Time-like region



Antisymmetric state
of colored quarks



Colorless quarks:
Pauli principle

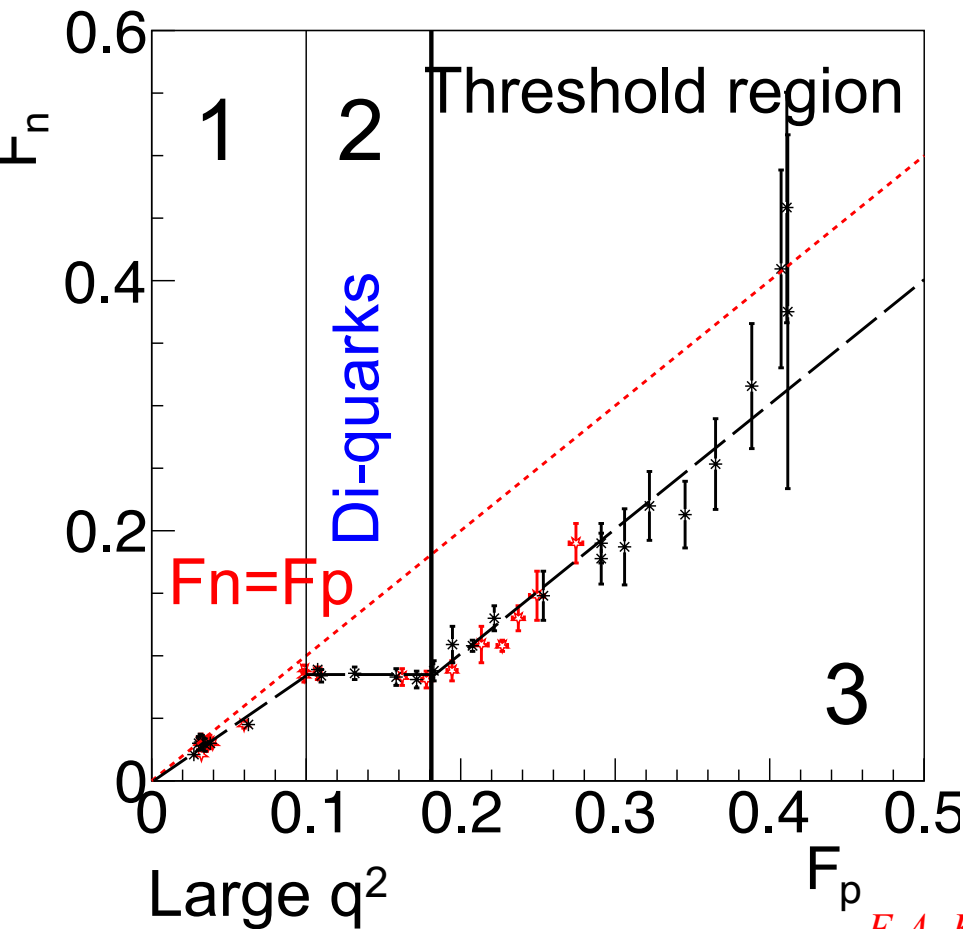
The vacuum state transfers all the released energy to a state of matter consisting at least of 6 massless valence quarks, a set of gluons, sea of $\bar{q}q$ with $q_0 > 2M_p$, $J=1$, dimensions $\hbar/(2M_p) \sim 0.1 \text{ fm}$.

- 1) uu (dd) quarks are repulsed from the inner region
- 2) The 3rd quark u (p) or d (n) is attracted by one of the identical quarks, forming *a compact di-quark: competition between attraction force and stochastic force of the gluon field*
- 3) The color state is restored: the 'point-like' hadron expands and cools down: *the current quarks and antiquarks absorb gluons and transform into constituent quarks*

E.A. Kuraev, A. Dbeyssi, E. T-G. Phys. Lett. 712, 240 (2012)



TL - np-correlation : 3 steps



Experimental points at
the same P_L

Proton values calculated
from the 6-parameter fit

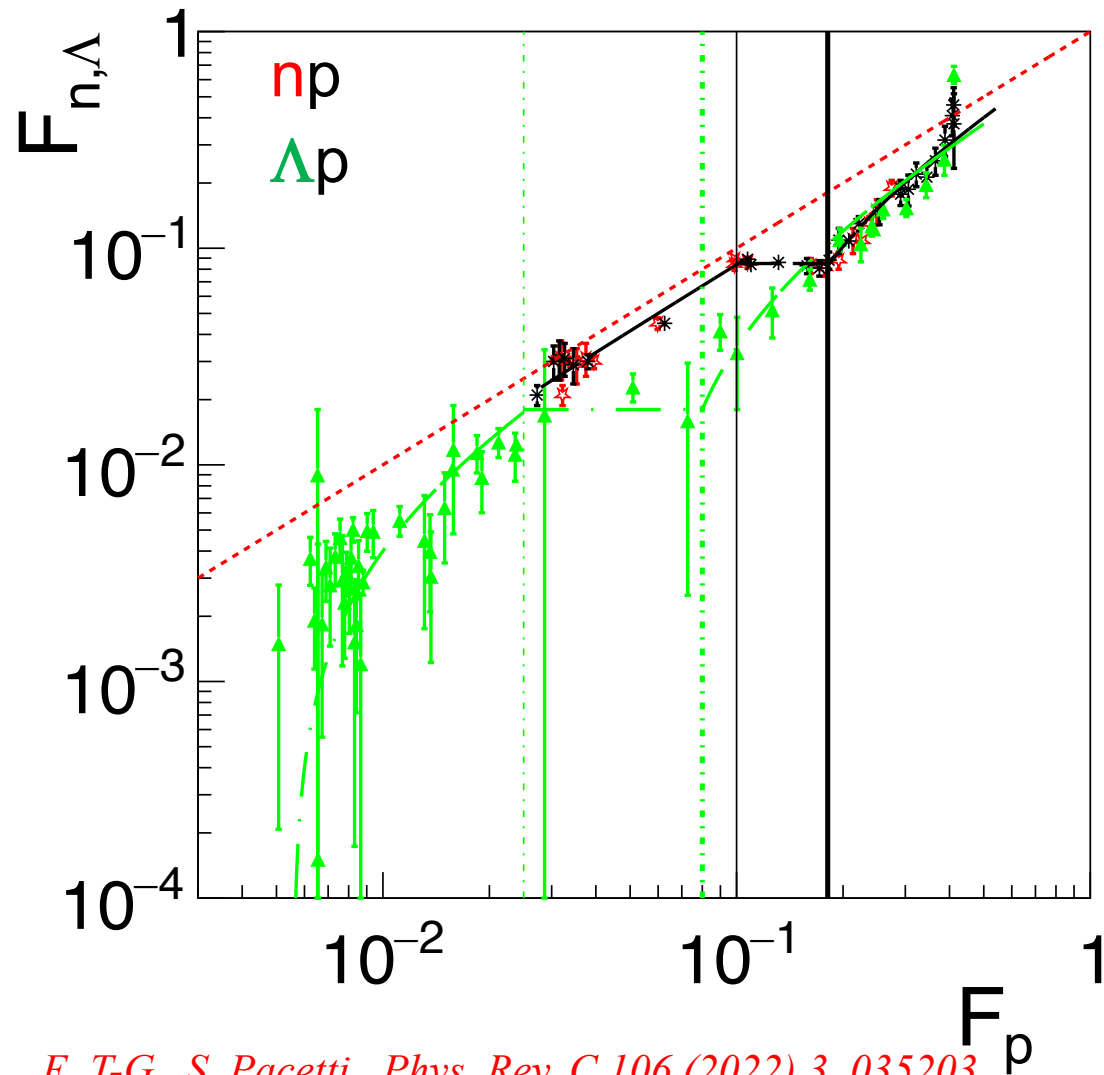
- 1) pQCD applies
- 2) di-quark phase
charge redistributed
- 3) The hadron is formed

E.A. Kuraev, A. Dbeyssi, E. T-G. Phys. Lett. 712, 240 (2012)

E. T-G., S. Pacetti, Phys. Rev. C 106 (2022) 3, 035203



np Λ -correlation

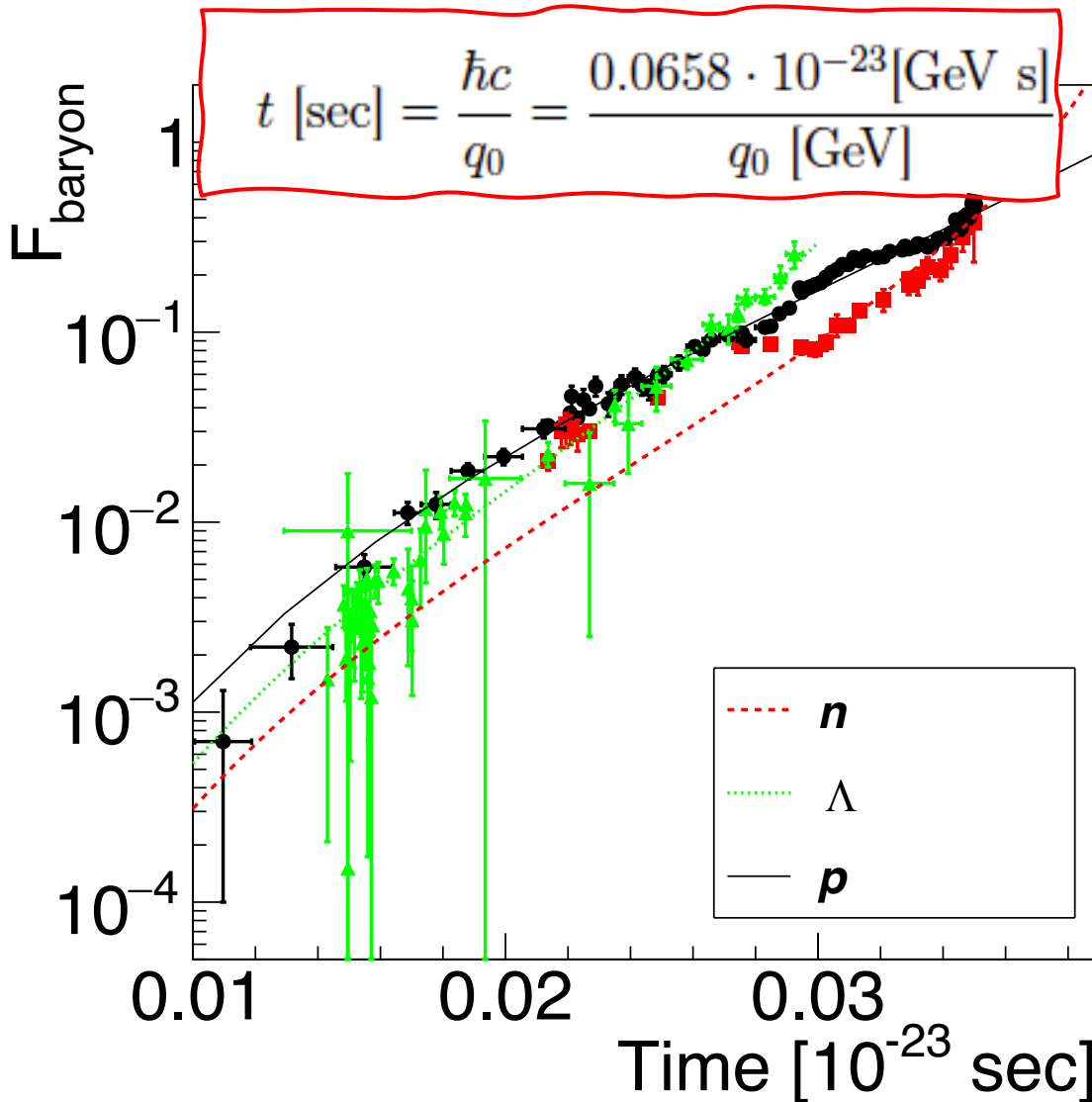


Quark pairs created by quantum vacuum fluctuations: all quark flavors are equally probable, but, due to Heisenberg principle, the associated time depends on the energy (baryon mass)

E. T-G., S. Pacetti, Phys. Rev. C 106 (2022) 3, 035203



TL- the most precise clock



10^{-23} s is the time for
light to cross a proton

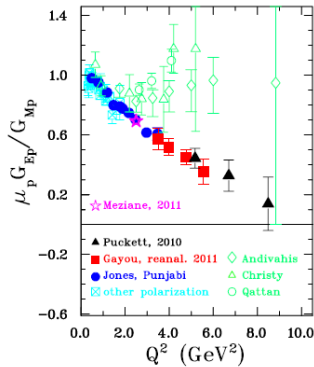
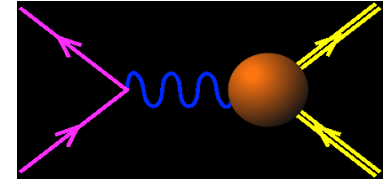
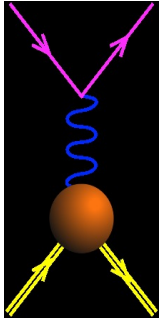
Di-quark phase dominant
at $t \sim 0.02-0.03$ [10^{-23} s]

E. T-G., S. Pacetti, Phys. Rev. C 106 (2022) 3, 035203

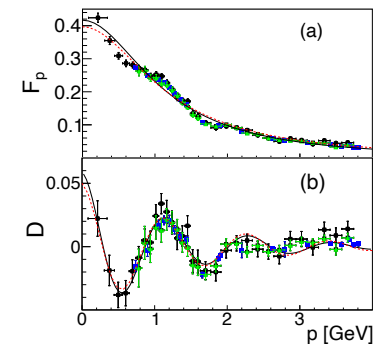


Conclusions

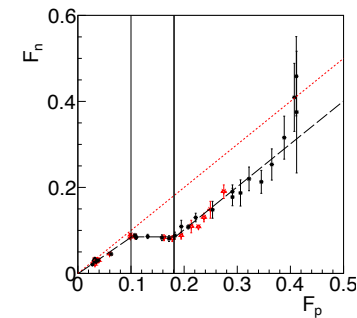
- Disentangle *structure* and *reaction* mechanism:
1 γ exchange mechanism



- Global understanding of *scattering* and *annihilation* reactions

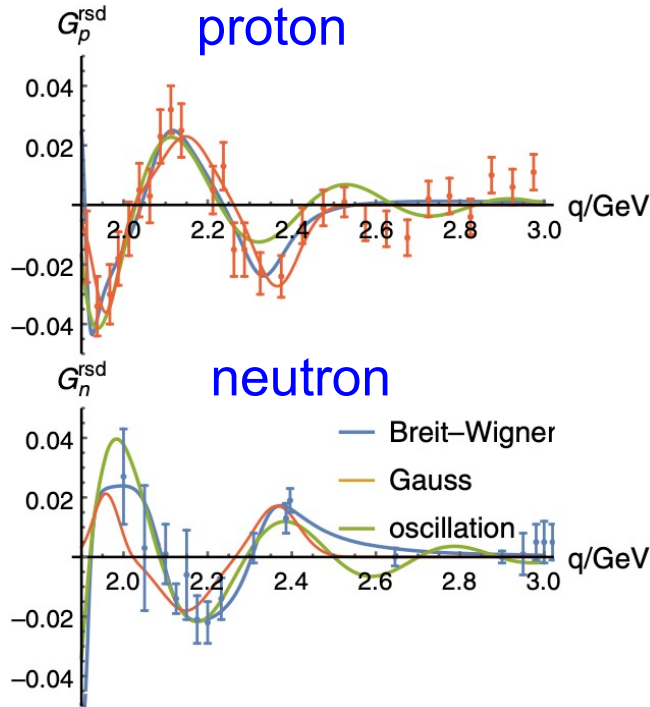


- Dynamical structure in *four* dimensions: space and time



Timelike nucleon electromagnetic form factors: All about interference of isospin amplitudes

Xu Cao^{id,1,2,*} Jian-Ping Dai^{id,3,†} and Horst Lenske^{4,‡}



- No multimeson rescattering processes
- **Competing isospin-clean vector meson intermediate states**
 $\phi(2170)$ ($I=0$) and $\rho(2150)$ ($I=1$)
- Sinusoidal modulation
- Energy dependent relative phase
- Related to the imaginary part of FFs

$$\frac{|I_1^{\text{rsd}} + I_0^{\text{rsd}}|}{|I_1^{\text{rsd}} - I_0^{\text{rsd}}|} = \frac{A_p}{A_n} = 0.88 \pm 0.35$$

Balanced isospin content
Not depending on energy
Limited range for the fit

Dipole Approximation & charge density

$$G_D = (1 + Q^2 / 0.71 \text{ GeV}^2)^{-2}$$

- Classical approach

- Nucleon FF (in non relativistic approximation or in the Breit system) are Fourier transform of the charge or magnetic distribution.

The diagram illustrates the Breit system for a nucleon. It shows two horizontal arrows representing the nucleons. The left arrow is labeled $\gamma^*(\mathbf{q}_B)$ and the right arrow is labeled $P_2(\mathbf{q}_B/2)$. Above the right arrow, there is a label $P_1(\mathbf{q}_B/2)$ with a double-headed arrow pointing towards the center. An equals sign is placed between the two main arrows, indicating the exchange of a photon between the two nucleons.

Breit system

- The dipole approximation corresponds to an exponential density distribution.

- $\rho = \rho_0 \exp(-r/r_0)$,
- $r_0^2 = (0.24 \text{ fm})^2$, $\langle r^2 \rangle \sim (0.81 \text{ fm})^2 \leftrightarrow m_D^2 = 0.71 \text{ GeV}^2$



Fourier Transform

A. Bianconi, E. T-G., Phys. Rev. Lett. 114, 232301 (2015)

$$F_0(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M_0(r)$$

$$F(p) = F_0(p) + F_{osc}(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M(r).$$

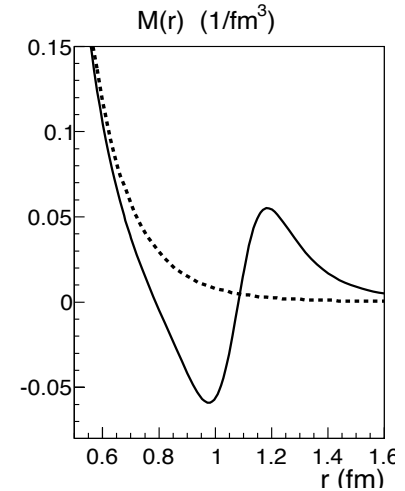
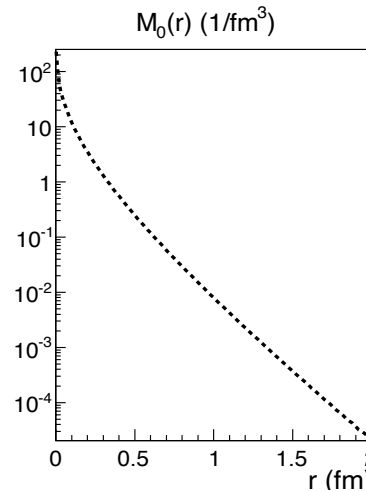
p: relative momentum

r: distance between the center of the forming hadrons

(p,r) conjugate variables, $r \leftrightarrow t$

$$F_0 = \frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2},$$

$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$



- Rescattering processes
- Large imaginary part
- Related to the time evolution of the charge density?
(E.A. Kuraev, E. T-G., A. Dbeyssi, PLB712 (2012) 240)
- Consequences for the SL region?
- Data from BESIII, expected from PANDA



Crossing symmetry

Scattering and annihilation channels:

- Described by the same amplitude :

$$|\overline{\mathcal{M}}(e^\pm h \rightarrow e^\pm h)|^2 = f(s, t) = |\overline{\mathcal{M}}(e^+ e^- \rightarrow \bar{h} h)|^2,$$

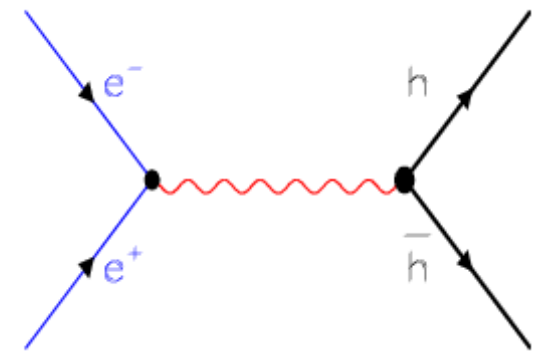
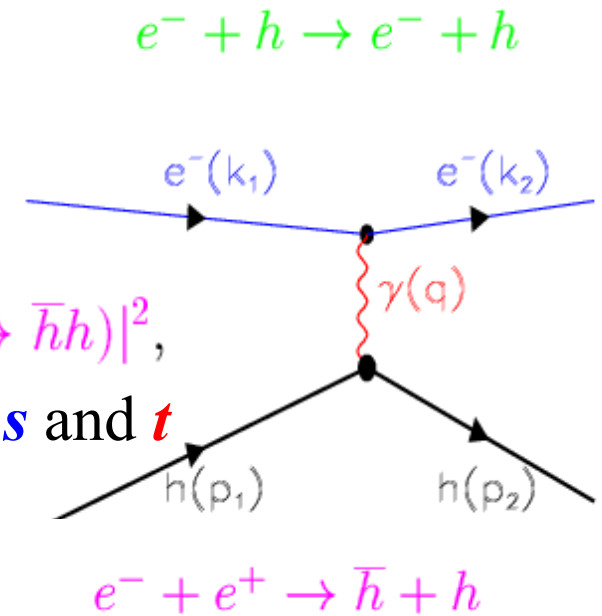
- function of two kinematical variables, s and t

$$s = (k_1 + p_1)^2$$

$$t = (k_1 - k_2)^2$$

- which scan different kinematical regions

$$\cos^2 \tilde{\theta} = 1 + \frac{st + (s - M^2)^2}{t\left(\frac{t}{4} - M^2\right)} \rightarrow 1 + \frac{ctg^2 \frac{\theta}{2}}{1 + \tau}$$



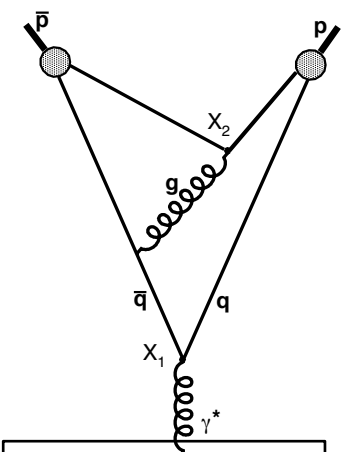
Photon-Charge coupling

$$\rho(\vec{x})$$

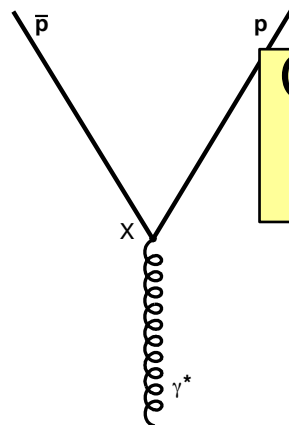
Fourier transform of a stationary charge and current distribution

$$R(t)$$

Amplitude for creating *charge-anticharge pairs* at time t



Resolved



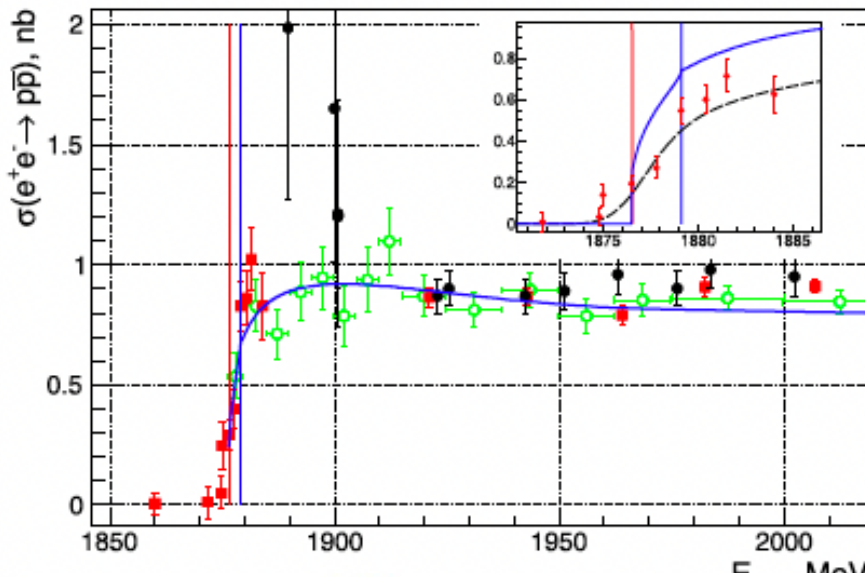
Unresolved

Charge distribution: distribution in time of $\gamma^* \rightarrow$ *charge-anticharge vertices*

The simplest picture: qq pair + compact di-quark

representation

R.R. Akhmetshin et al. / Physics Letters B 794 (2019) 64–68



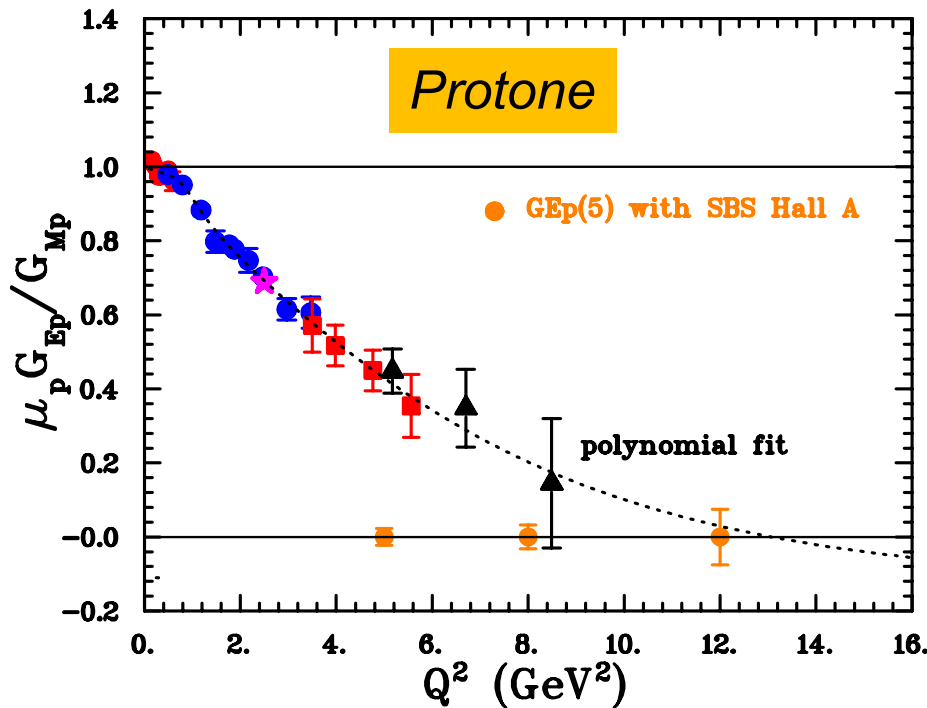
Convolution of the radiative cross section with the cm spread energy function ($\Delta E \sim 0.95$ MeV)

$$\sigma_{\text{vis}}(E_{\text{c.m.}}) = \frac{1}{\sqrt{2\pi}\sigma_{E_{\text{c.m.}}}} \int dE'_{\text{c.m.}} \sigma_{f\gamma}(E'_{\text{c.m.}}) \cdot \exp\left(-\frac{(E_{\text{c.m.}} - E'_{\text{c.m.}})^2}{2\sigma_{E_{\text{c.m.}}}^2}\right),$$

$$\sigma_{f\gamma}(E_{\text{c.m.}}) = \int_0^{E_{\gamma}^{\text{max}}} dE_{\gamma} \cdot \sigma_{\text{Born}}(E_{\text{c.m.}} \sqrt{1 - E_{\gamma}/E_{\text{c.m.}}}) \cdot F(E_{\text{c.m.}}, E_{\gamma}),$$

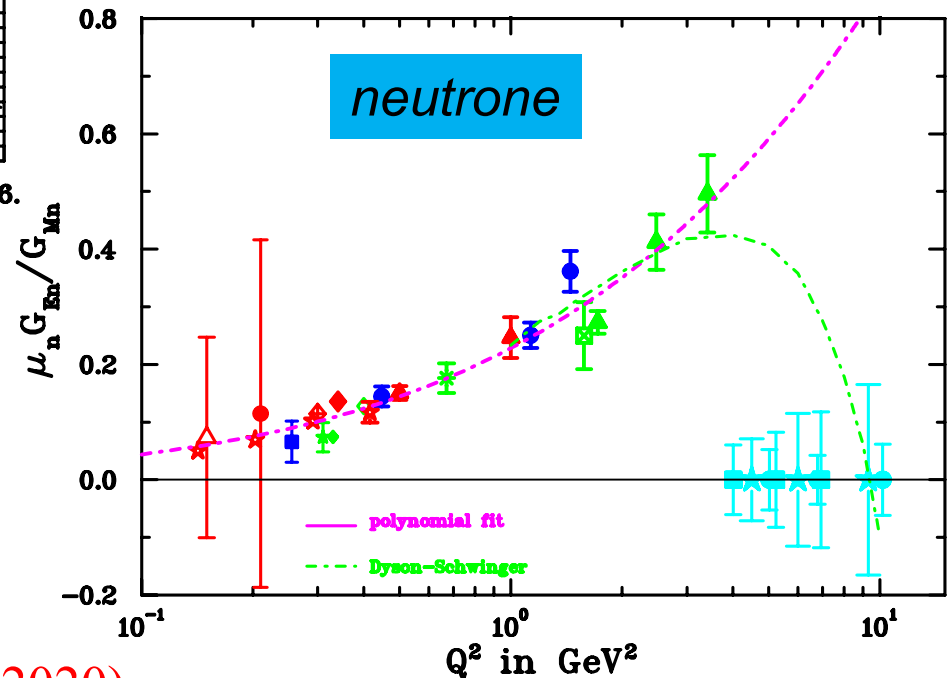
$$\sigma_{\text{Born}}(E_{\text{c.m.}}) = A + B \left[1 - \exp\left(-\frac{(E_{\text{c.m.}} - E_{\text{thr}})}{\sigma_{\text{thr}}}\right) \right],$$

...and future plans



Ch. Perdrisat, Jlab PR12-07-109

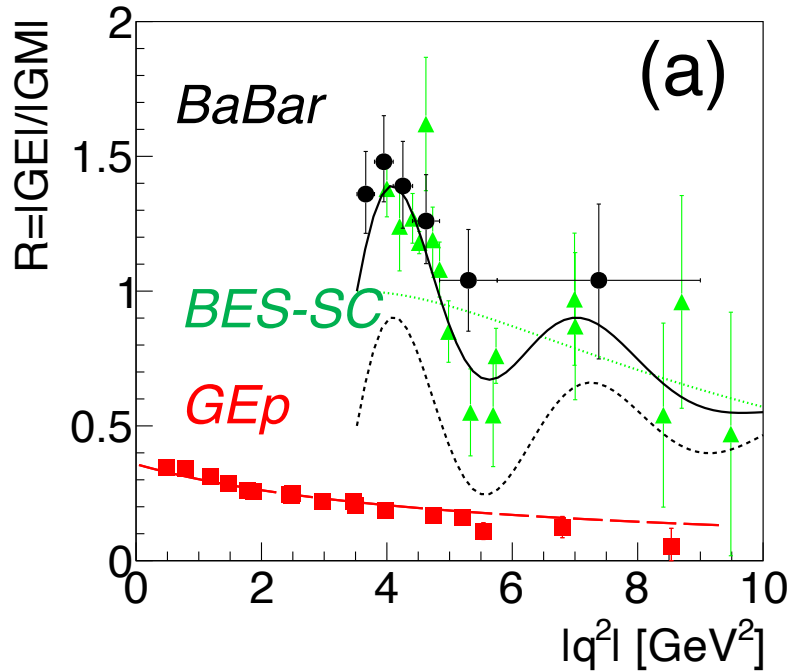
J. Anderson Jlab PR12-09-009
J. Annand Jlab PR12-09-019



S.N. Basilev et al, Eur.Phys.J.A 56 (2020)



Form Factor Ratio $R=|GE|/|GM|$



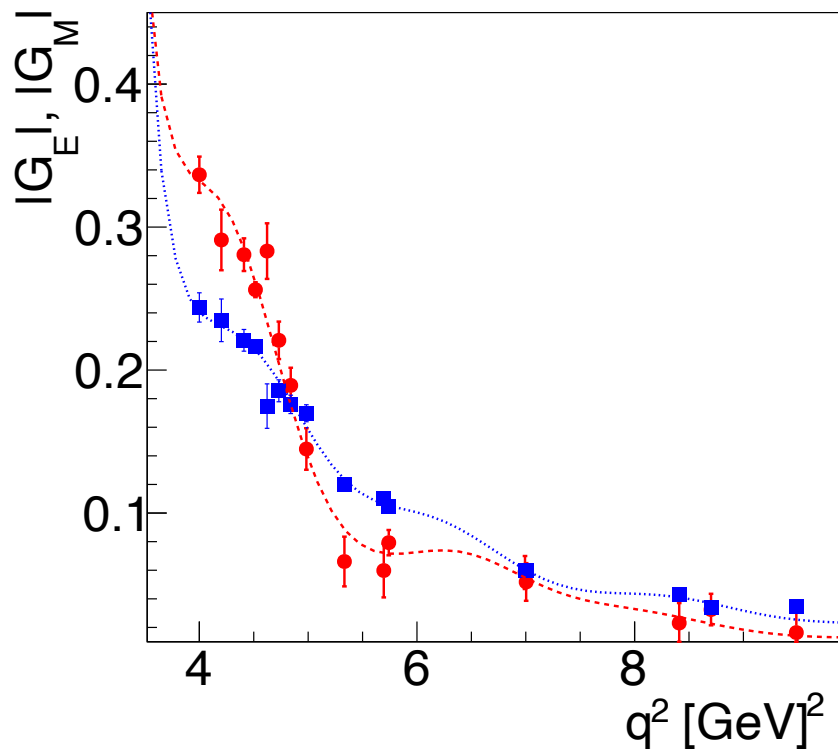
- Precise data from BESIII
- Dip at $|q^2| \sim 5.8$ GeV²
- Comparison with SL (Jlab-GEp data)
- Oscillations on top of a monopole: from GE or GM?

$$F_R(\omega(s)) = \frac{1}{1 + \omega^2/r_0} [1 + r_1 e^{-r_2 \omega} \sin(r_3 \omega)], \quad \omega = \sqrt{s} - 2m_p,$$



Sachs form factors: $|G_E|$, $|G_M|$

From the fit on F_p and the fit on R ,
the Sachs FFs (moduli) can be reconstructed



$$|G_E(s)| = F_p(s) \sqrt{\frac{1 + 2\tau}{R^2(s) + 2\tau/R^2(s)}}$$
$$|G_M(s)| = F_p(s) \sqrt{\frac{1 + 2\tau}{R^2(s) + 2\tau}}$$

Threshold constrain $R=1$ for $\tau=1$

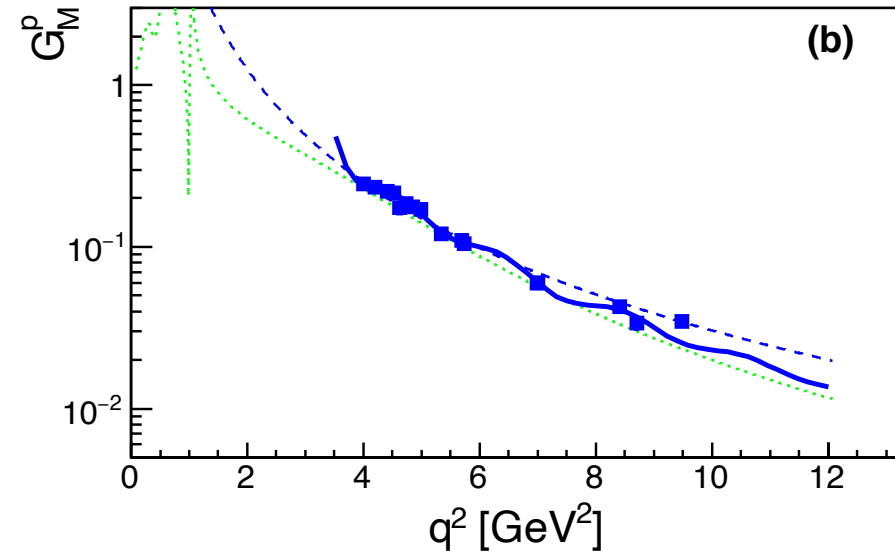
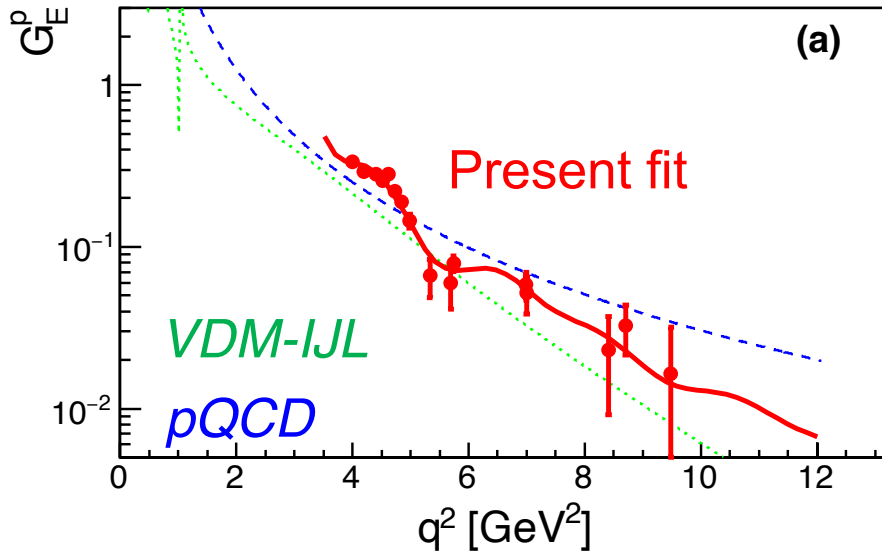
The fit gives :

$$|G_E| = |G_M| = 0.48$$



Models

Parametrizations have been determined by fitting F_p & R



$|G_E|$: more pronounced oscillations
faster q^2 -decrease

Threshold constrain $R=1$ for $\tau=1$

The fit gives : pQCD : 0.34

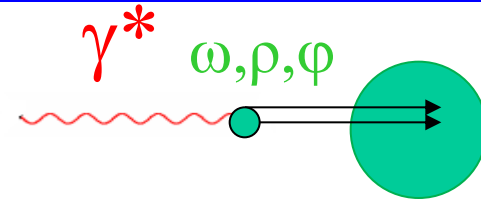
$|G_E| = |G_M| = 0.48$ VDM-IJL : 0.29

E.T.-G., A. Bianconi, S. Pacetti, Phys.Rev.C 103 (2021) 3, 035203



VMD: Iachello, Jackson and Landé (1973)

Isoscalar and isovector FFs



$$F_1^s(Q^2) = \frac{g(Q^2)}{2} \left[(1 - \beta_\omega - \beta_\phi) + \beta_\omega \frac{\mu_\omega^2}{\mu_\omega^2 + Q^2} + \beta_\phi \frac{\mu_\phi^2}{\mu_\phi^2 + Q^2} \right],$$

$$F_1^v(Q^2) = \frac{g(Q^2)}{2} \left[(1 - \beta_\rho) + \beta_\rho \frac{\mu_\rho^2 + 8\Gamma_\rho \mu_\pi / \pi}{(\mu_\rho^2 + Q^2) + (4\mu_\pi^2 + Q^2)\Gamma_\rho \alpha(Q^2) / \mu_\pi} \right],$$

$$F_2^s(Q^2) = \frac{g(Q^2)}{2} \left[(\mu_p + \mu_n - 1 - \alpha_\phi) \frac{\mu_\omega^2}{\mu_\omega^2 + Q^2} + \alpha_\phi \frac{\mu_\phi^2}{\mu_\phi^2 + Q^2} \right],$$

$$F_2^v(Q^2) = \frac{g(Q^2)}{2} \left[(\mu_p - \mu_n - 1) \frac{\mu_\rho^2 + 8\Gamma_\rho \mu_\pi / \pi}{(\mu_\rho^2 + Q^2) + (4\mu_\pi^2 + Q^2)\Gamma_\rho \alpha(Q^2) / \mu_\pi} \right],$$

$$g(Q^2) = \frac{1}{(1 + \gamma e^{i\theta} Q^2)^2}$$

Intrinsic factor

Meson Cloud

$$2F_i^p = F_i^s + F_i^v,$$

$$2F_i^n = F_i^s - F_i^v.$$

$$\alpha(Q^2) = \frac{2}{\pi} \sqrt{\frac{Q^2 + 4\mu_\pi^2}{Q^2}} \ln \left[\frac{\sqrt{(Q^2 + 4\mu_\pi^2)} + \sqrt{Q^2}}{2\mu_\pi} \right]$$

Few # parameters, with physical meaning
Naturally arising TL imaginary part



Total Cross Section from $e^+e^- \rightarrow p\bar{p}$

$$\sigma_{e^+e^- \rightarrow p\bar{p}}(s) = \frac{4\pi\alpha^2\beta\mathcal{C}(\beta)}{3s} \left(|G_M(s)|^2 + \frac{1}{2\tau} |G_E(s)|^2 \right)$$

- Effective FF: $\sigma_{\text{Tot}} \sim F_p^2$

$$F_p(s)^2 = \frac{2\tau |G_M(s)|^2 + |G_E(s)|^2}{2\tau + 1}$$

- Equivalent to:

$$|G_E(s)| = |G_M(s)| \equiv F_p(s)$$

Strictly valid at threshold, where only one amplitude is present



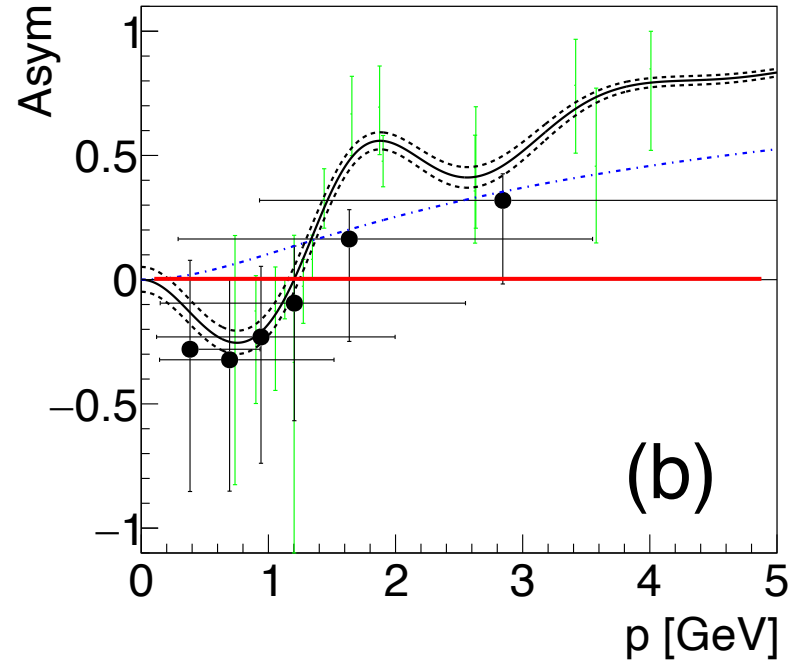
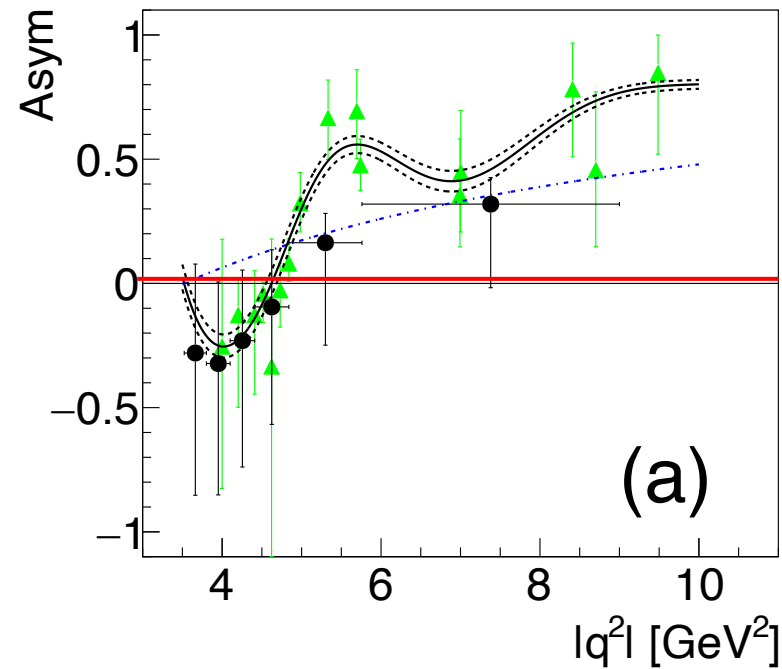
Focus to threshold

- Large activity at all world facilities both in Space and Time-like regions
- Theory: unified models in SL and TL regions:
 - describe all 4 FFs:
proton and neutron, electric and magnetic
- Features in SL and TL regions explained by an inner vacuum:
 - monopole decrease of R in n&p, SL & TL, diquark phase
prediction of non-zero crossing of SL ratio
- $p\bar{p}$ annihilation at threshold : pointlike?:
 - R=1 : but what is the value of FFs? test of models
 - Are the $e^+e^- \leftrightarrow pp$ really equivalent? related by time reversal, but ... FSI, Coulomb interaction, radiative corrections...

Time structure of the nucleon: enter the 4th dimension !



Angular Asymmetry



$$\frac{d\sigma_{e^+e^- \rightarrow \bar{p}p}}{d\Omega}(s, \theta) = \sigma_0(s) |1 + \mathcal{A}(s) \cos^2(\theta)|$$

$$\sigma_0(s) = \frac{\alpha^2 \beta \mathcal{C}(\beta)}{4s} \left(|G_M(s)|^2 + \frac{1}{\tau} |G_E(s)|^2 \right)$$

$$\mathcal{A}(s) = \frac{\tau |G_M(s)|^2 - |G_E(s)|^2}{\tau |G_M(s)|^2 + |G_E(s)|^2} = \frac{\tau - R(s)^2}{\tau + R(s)^2}$$

$$q^2 = (4.60 \pm 0.07) \text{ GeV}^2$$

$$p = (1.20 \pm 0.04) \text{ GeV}$$

Zero of the angular asymmetry

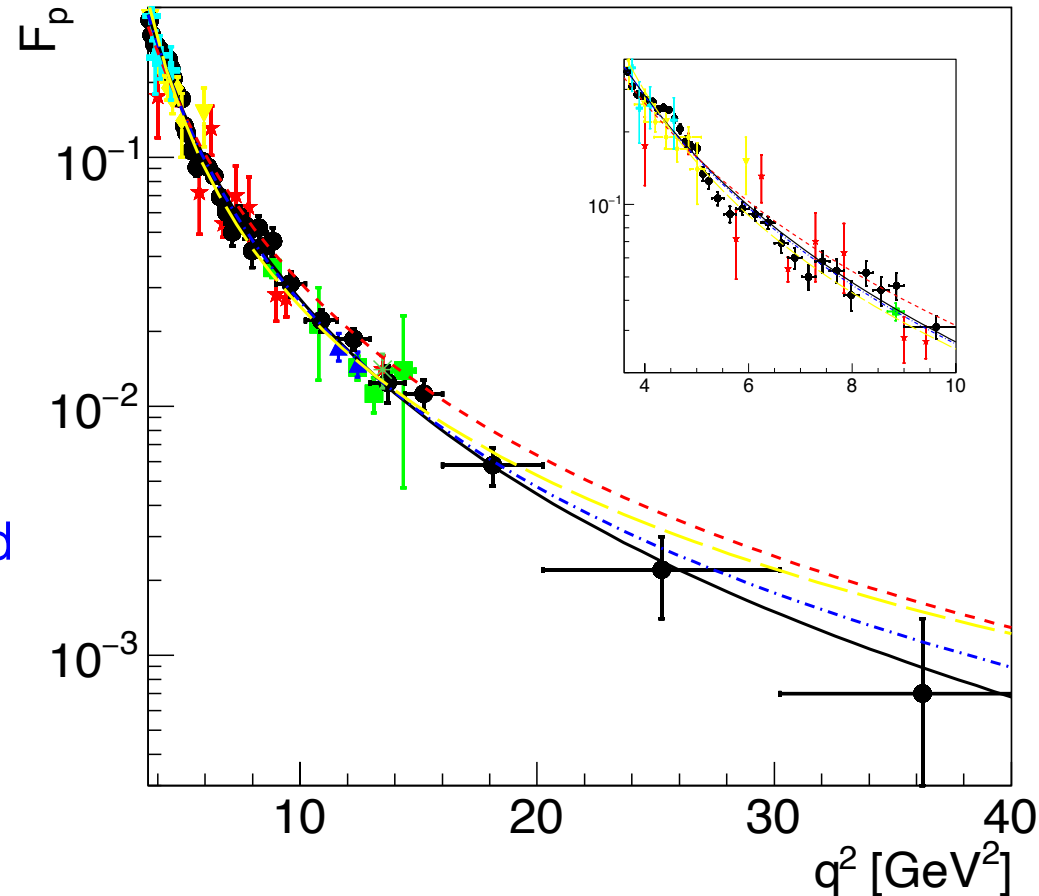
The Time-like Region

$$G_E = G_M$$

The Experimental Status

- Individual determination of G_E and G_M : only recently!
- TL proton FFs twice larger than in SL at the same Q^2
- Steep behaviour at threshold
- BaBAR: Structures?
Resonances?

Confirmed by BES



S. Pacetti, R. Baldini-Ferroli, E.T-G, Physics Reports, 514 (2014) 1

Panda contribution: M.P. Rekalo, E.T-G, DAPNIA-04-01, ArXiv:0810.4245.



ep-elastic scattering : The Akhiezer-Rekalo method

PHYSICS

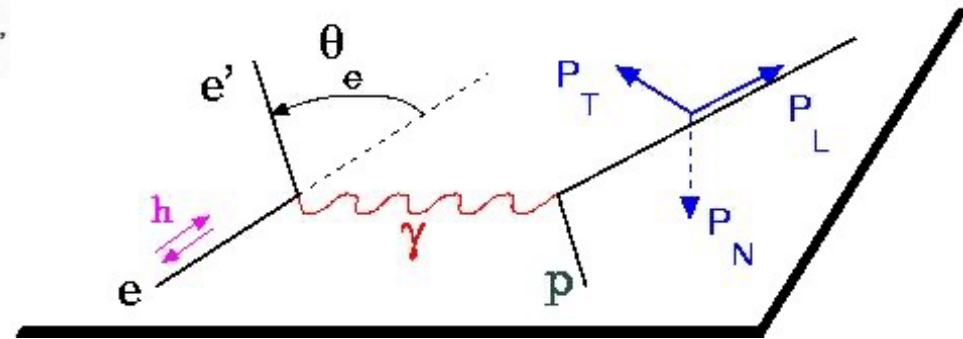
1967

POLARIZATION PHENOMENA IN ELECTRON
SCATTERING BY PROTONS IN THE
HIGH-ENERGY REGION

Academician A. I. Akhiezer* and M. P. Rekalo

Physicotechnical Institute, Academy of Sciences of the Ukrainian SSR
Translated from Doklady Akademii Nauk SSSR, Vol. 180, No. 5,
pp. 1081-1083, June, 1968
Original article submitted February 26,

$$s_2 \frac{d\sigma}{d\Omega_R} = 4p_2 \frac{(\mathbf{s} \cdot \mathbf{q})}{1 + \tau} \Gamma(\theta, \varepsilon_1) \left[\tau G_M (G_M + G_E) - \frac{1}{4\varepsilon_1} G_M (G_E - \tau G_M) \right],$$



The polarization induces a term in the cross section proportional to $G_E G_M$
Polarized beam and target or
polarized beam and recoil proton polarization

The polarization method (exp: 2000)

Transferred polarization is:

*C. Perdrisat, V. Punjabi, et al.,
JLab-GEp collaboration*

$$P_n = 0$$

$$\pm h P_t = \mp h 2\sqrt{\tau(1+\tau)} G_E^p G_M^p \tan\left(\frac{\theta_e}{2}\right) / I_0$$

$$\pm h P_l = \pm h (E_e + E_{e'}) (G_M^p)^2 \sqrt{\tau(1+\tau)} \tan^2\left(\frac{\theta_e}{2}\right) / M / I_0$$

Where, $h = |h|$ is the beam helicity

$$I_0 = (G_E^p(Q^2))^2 + \frac{\tau}{\epsilon} (G_M^p(Q^2))^2$$

$$\Rightarrow \frac{G_E^p}{G_M^p} = -\frac{P_t}{P_l} \frac{E_e + E_{e'}}{2M} \tan\left(\frac{\theta_e}{2}\right)$$

The simultaneous measurement of P_t and P_l reduces the systematic errors

