

Generalized Polarizabilities of the proton

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Outline

Introduction to the GPs

Overview: Status & Challenges

Recent results (Jlab / MAMI)

Prospects (VCS-II @ Jlab, measuring w positrons, ...)

Proton Polarizabilities

Fundamental structure constants
(such as mass, size, shape, ...)

Response of the nucleon to external EM field

Sensitive to the full excitation spectrum

Accessed experimentally through Compton Scattering

Virtual Compton Scattering:

Virtuality of photon gives access to the GPs : $\alpha_E(Q^2)$ & $\beta_M(Q^2)$ (+ 4 spin GPs)

→ mapping out the spatial distribution of the polarization densities

Fourier transform of densities of electric charges and magnetization of a nucleon deformed by an applied EM field

PDG

150 Baryon Summary Table

| |
|--|
| N BARYONS ($S = 0, I = 1/2$) $p, N^+ = uud; \quad n, N^0 = udd$ |
|--|

p

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

Mass $m = 1.00727646681 \pm 0.00000000009$ u

Mass $m = 938.272046 \pm 0.000021$ MeV [a]

$|m_p - m_{\bar{p}}|/m_p < 7 \times 10^{-10}$, CL = 90% [b]

$\frac{q_p}{m_p} / (\frac{q_e}{m_e}) = 0.99999999991 \pm 0.00000000009$

$|q_p + q_{\bar{p}}|/e < 7 \times 10^{-10}$, CL = 90% [b]

$|q_p + q_e|/e < 1 \times 10^{-21}$ [c]

Magnetic moment $\mu = 2.792847356 \pm 0.000000023$ μ_N

$(\mu_p + \mu_{\bar{p}}) / \mu_p = (0 \pm 5) \times 10^{-6}$

Electric dipole moment $d < 0.54 \times 10^{-23}$ e cm

Electric polarizability $\alpha = (11.2 \pm 0.4) \times 10^{-4}$ fm³

Magnetic polarizability $\beta = (2.5 \pm 0.4) \times 10^{-4}$ fm³ (S = 1.2)

Charge radius, μp Lamb shift = 0.84087 ± 0.00039 fm [d]

Charge radius, $e p$ CODATA value = 0.8775 ± 0.0051 fm [d]

Magnetic radius = 0.777 ± 0.016 fm

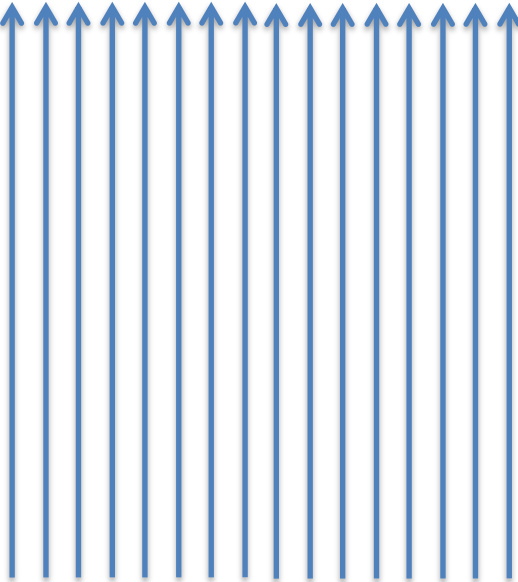
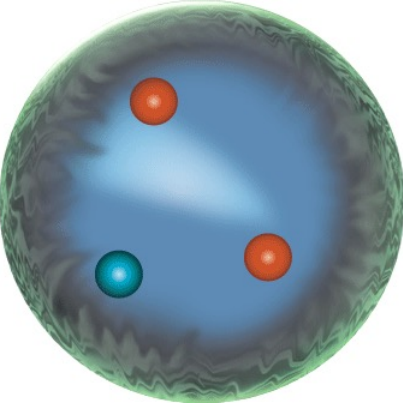
Mean life $\tau > 2.1 \times 10^{29}$ years, CL = 90% [e] ($p \rightarrow$ invisible mode)

Mean life $\tau > 10^{31}$ to 10^{33} years [e] (mode dependent)

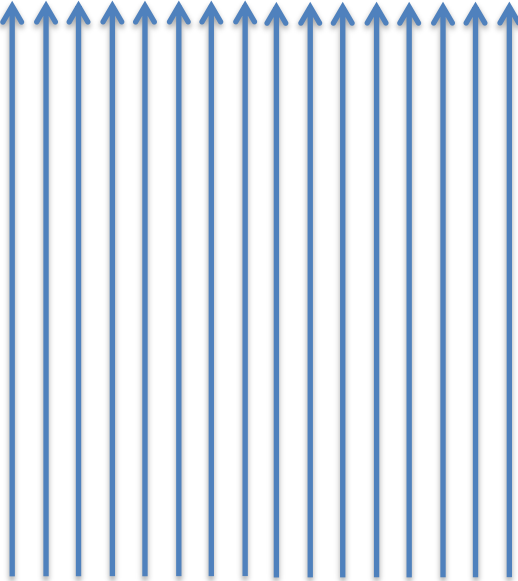
Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon



\vec{E}

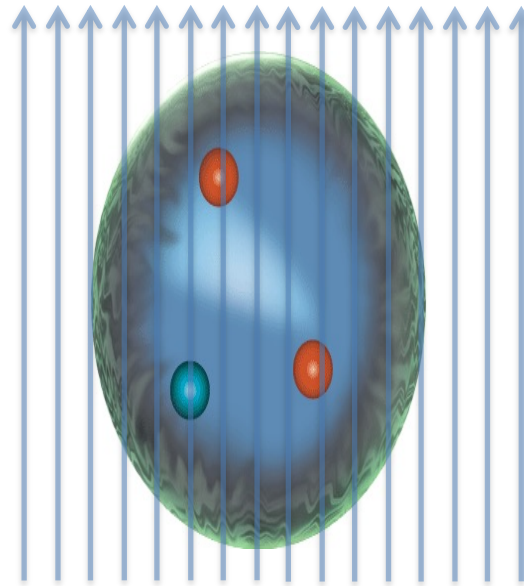
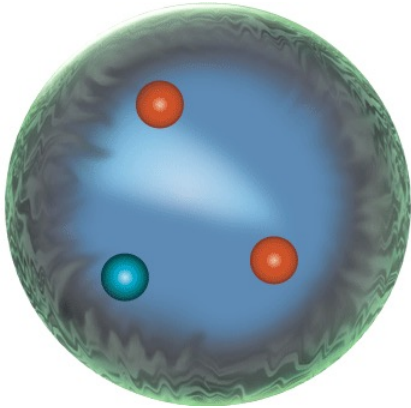


\vec{B}

Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon

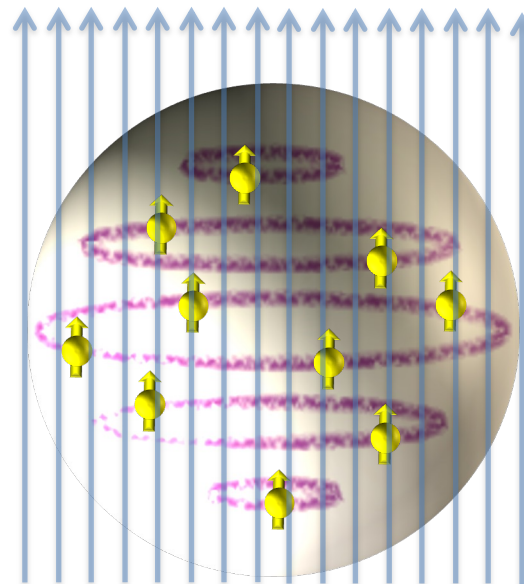


\vec{E}

“stretchability”

$$\vec{d}_{E \text{ induced}} \sim \alpha \vec{E}$$

External field deforms the charge distribution



\vec{B}

“alignability”

$$\vec{d}_{M \text{ induced}} \sim \beta \vec{B}$$

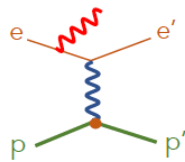
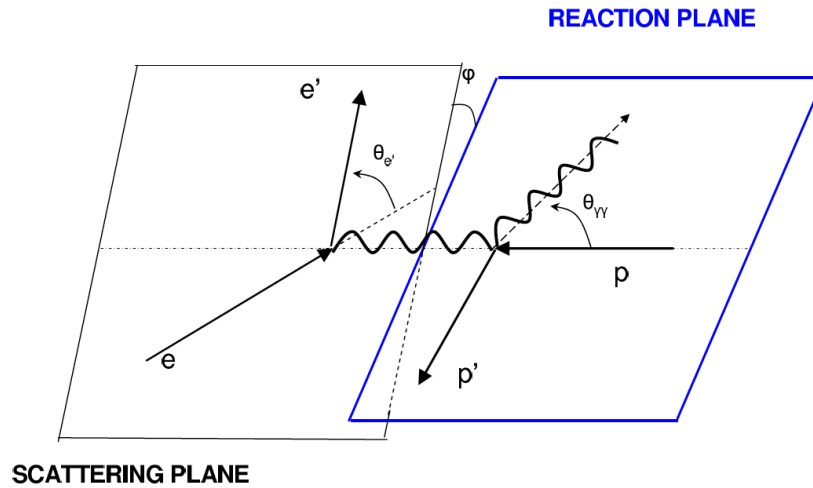
$$\beta_{\text{para}} > 0$$

$$\beta_{\text{diam}} < 0$$

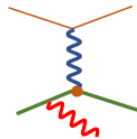
Paramagnetic: proton spin aligns with the external magnetic field

Diamagnetic: π -cloud induction produces field counter to the external perturbation

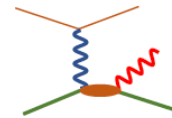
Virtual Compton Scattering



Bethe-Heitler



Born VCS



non-Born VCS

Elastic FFs

GPs

Virtual Compton Scattering

DR

valid below & above
Pion threshold

Dispersive integrals
for Non Born amplitudes

Spin GPs are fixed

Scalar GPs have
an unconstrained part

Fit to the experimental
cross sections at each Q^2

LEX

valid only below
Pion threshold

Response functions

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

$$\Psi_0 = v_1 \cdot \left(P_{LL} - \frac{1}{\epsilon} P_{TT} \right) + v_2 \cdot P_{LT}$$

Subtract the spin part

$$P_{TT} = [P_{TT} \text{ spin}]$$

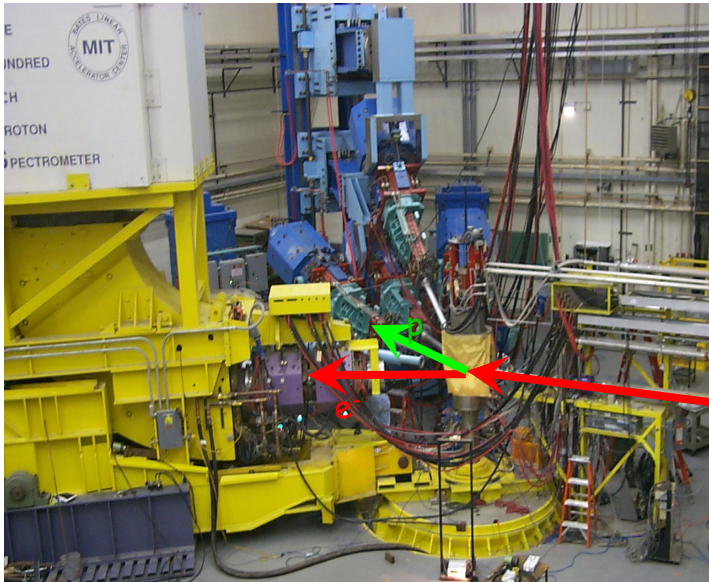
$$P_{LT} = -\frac{2M}{\alpha_{em}} \sqrt{\frac{q'^2_{cm}}{Q^2}} \cdot G_E^p(Q^2) \cdot \beta_M(Q^2) + [P_{LT} \text{ spin}]$$

utilize DR

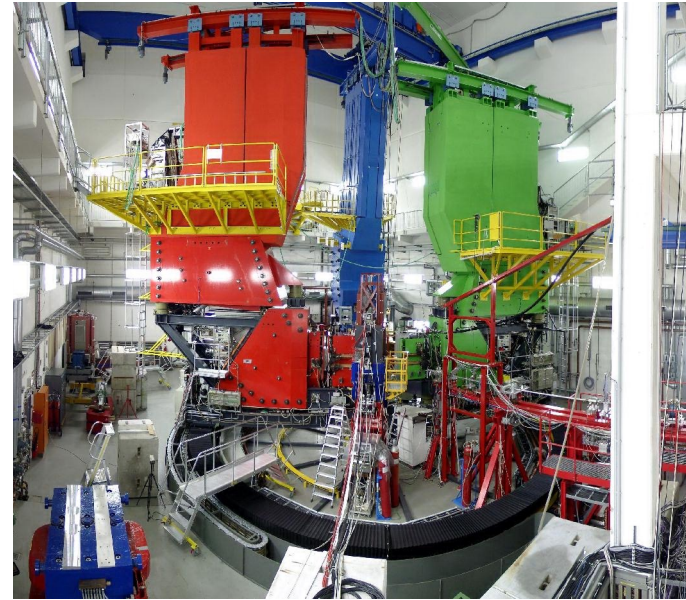
scalar GPs α_E and β_M

Early Experiments

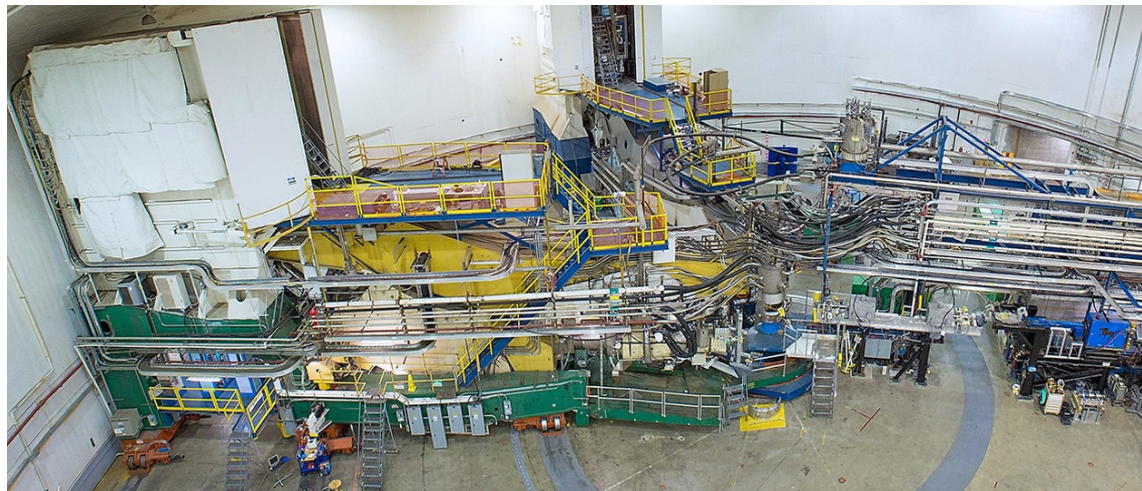
MIT-Bates @ $Q^2=0.06 \text{ GeV}^2$



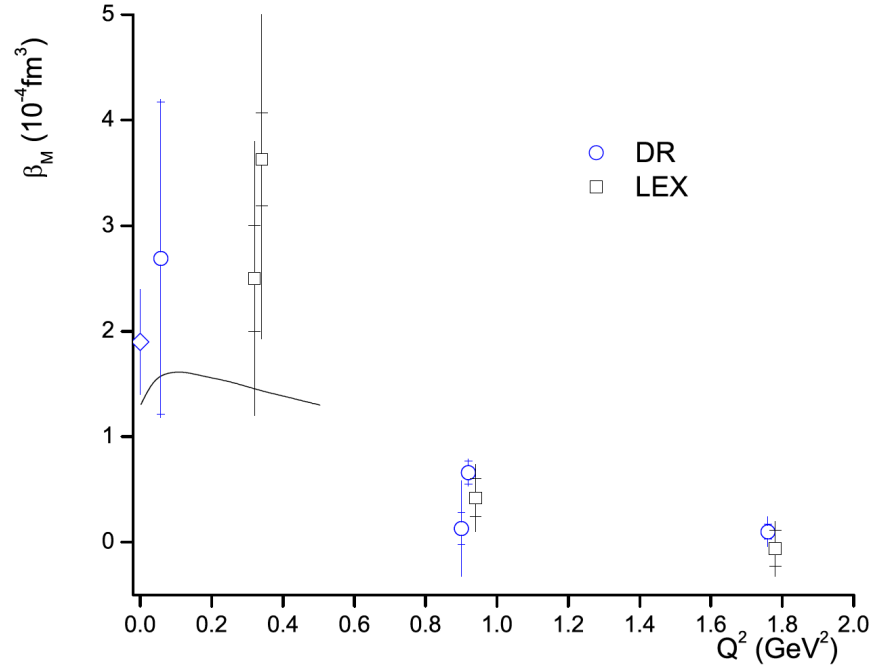
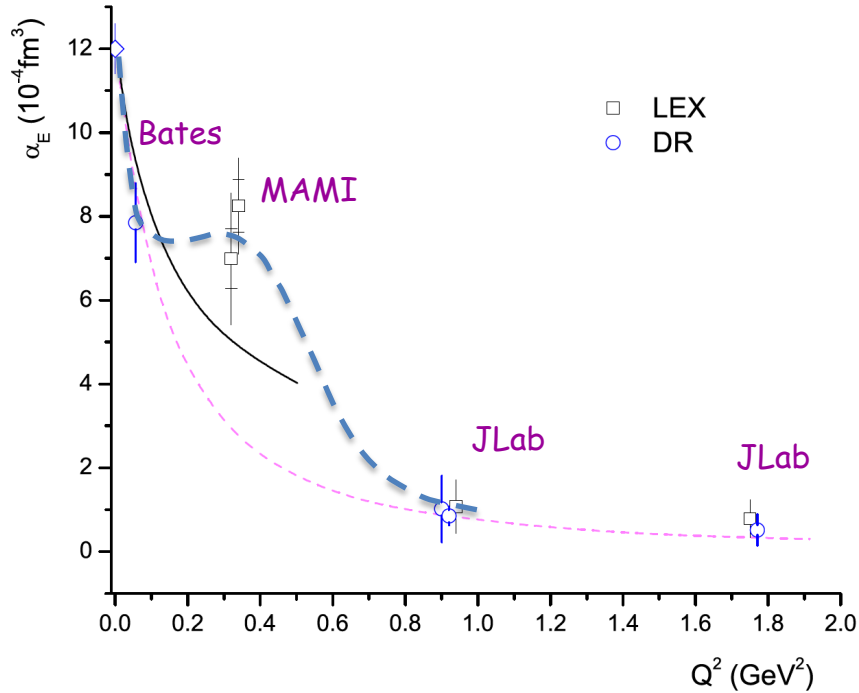
MAMI-A1 @ $Q^2=0.33 \text{ GeV}^2$



Jlab-Hall A @ $Q^2=0.9 \text{ \& } 1.8 \text{ GeV}^2$



Early Experiments



$\alpha_E \approx 10^{-3} V_N$ (stiffness / relativistic character)

Data: non-trivial Q^2 dependence of α_E (?)

Theory: monotonic fall-off

$Q^2 = 0.33 (\text{GeV}/c)^2$ measured twice at MAMI:

- Phys. Rev. Lett 85, 708 (2000)
- Eur. Phys. J. A37, 1-8 (2008)

β_M small \leftrightarrow cancellation of competing mechanisms

Large uncertainties

Higher precision measurements needed

\rightarrow Quantify the balance between diamagnetism and paramagnetism

Theory

HBChPT

NRQCM

Effective Lagrangian Model

Linear Sigma Model

T.R. Hemmert et al

B. Pasquini et al

A. Yu. Korchin and O. Scholten

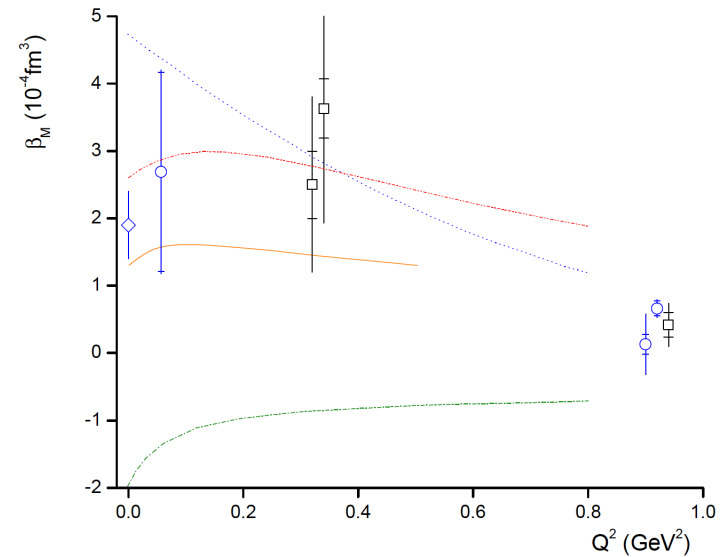
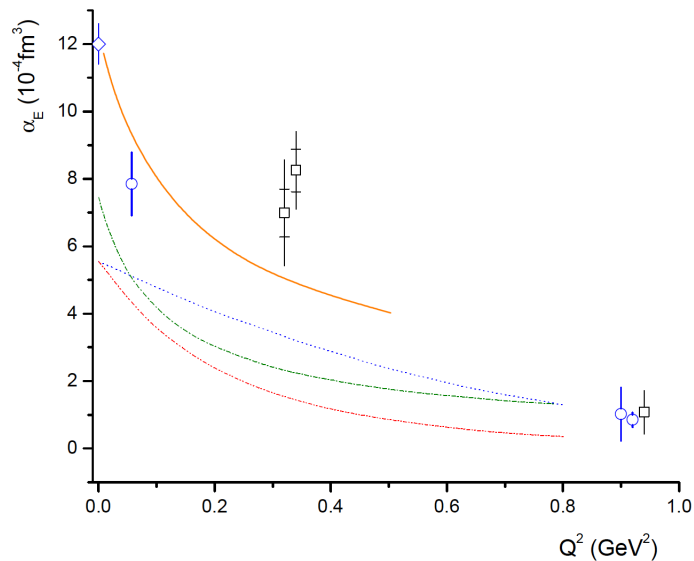
A. Metz and D. Drechsel

Phys. Rev. D 62, 014013 (2000)

Phys. Rev. C 63, 025205 (2001)

Phys. Rev. C 58, 1098 (1998)

Z. Phys. A 356, 351 (1996)



Smooth fall-off for α_E

A non-trivial structure for α_E is not supported by theory

Large spread in the predictions of the magnetic GP

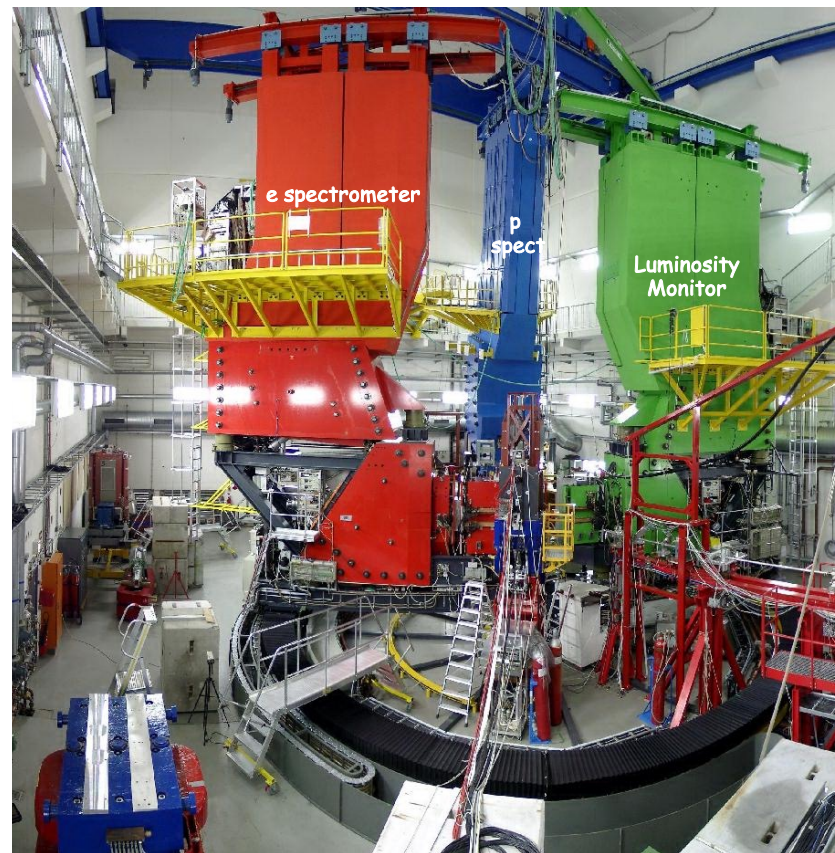
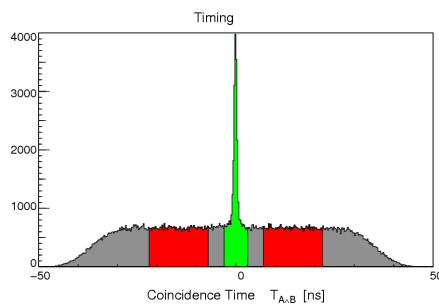
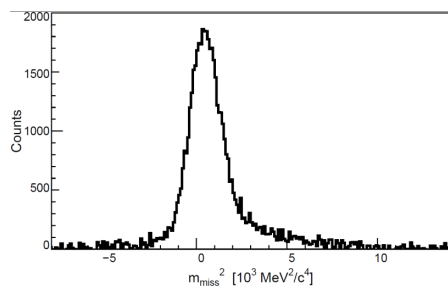
Recent Measurements

Recent Measurements: MAMI

MAMI A1/1-09 (vcsq2) below threshold

MAMI A1/3-12 (vcsdelta) above threshold

Both experiments utilized the A1 setup at MAMI



A1/1-09 @ MAMI

For LEX the higher order terms have to be kept small / under control

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

Refined analysis procedure / phase space masking to keep these terms smaller than $\sim 2\%$ - 3% level

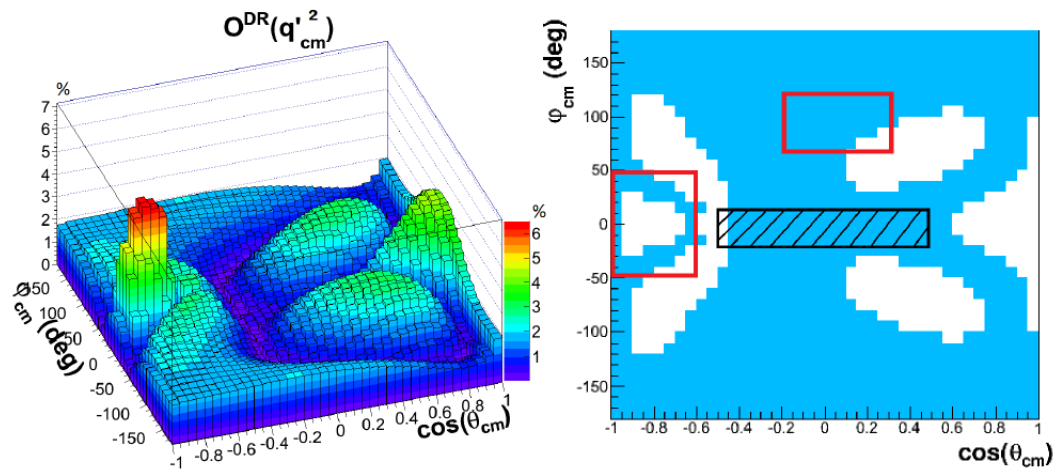


Figure 3.13: (Left) behavior of $\mathcal{O}^{DR}(q'^2_{cm})$ in the $(\cos(\theta_{cm}), \varphi_{cm})$ -plane at $q'_{cm} = 87.5 \text{ MeV}/c$ and (right) two-dimensional representation of the angular region where $\mathcal{O}^{DR}(q'^2_{cm}) < 2\%$ (blue), the red squares correspond to the two areas of interest to perform the GP extraction.

Figure from PhD thesis of L. Correa, Mainz / Cl. Ferrand

MAMI Results

Phys. Rev. Lett 123, 192302

Phys. Rev. C 103, 025205

Eur. Phys. J. A55, 182

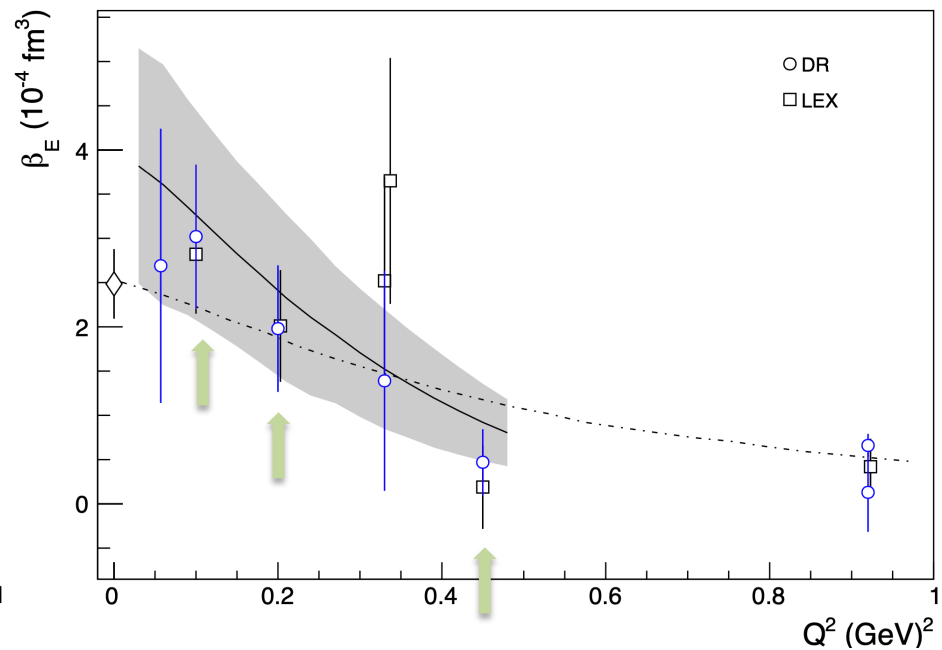
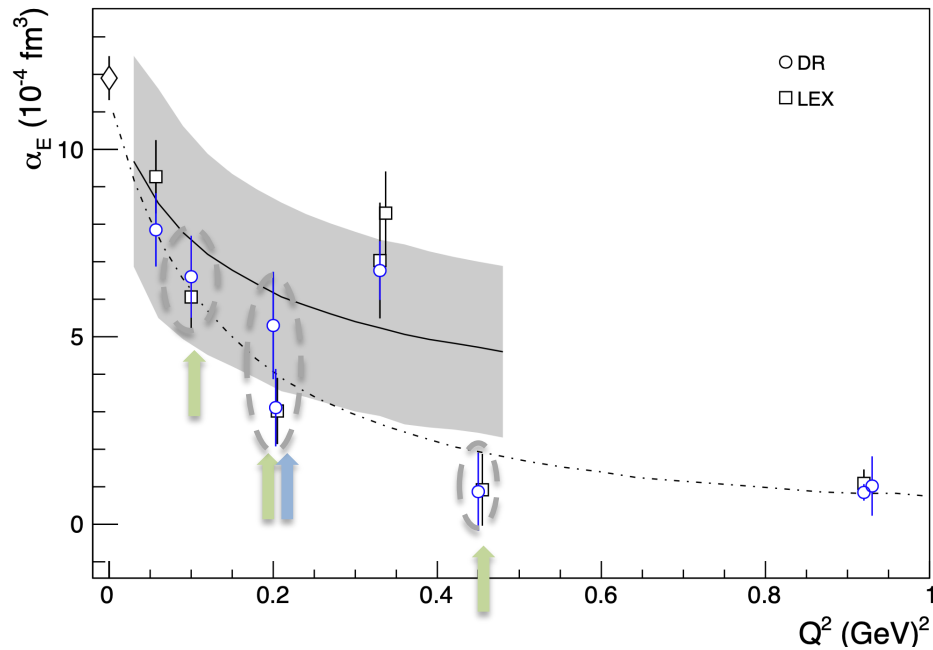
PhD students:

Jure Bericic (Ljubljana Univ.)

Loup Correa (Clermont-Fd Univ.)

Meriem BenAli (Clermont-Fd Univ.)

Adam Blomberg (Temple Univ.)



A1/1-09 @ MAMI

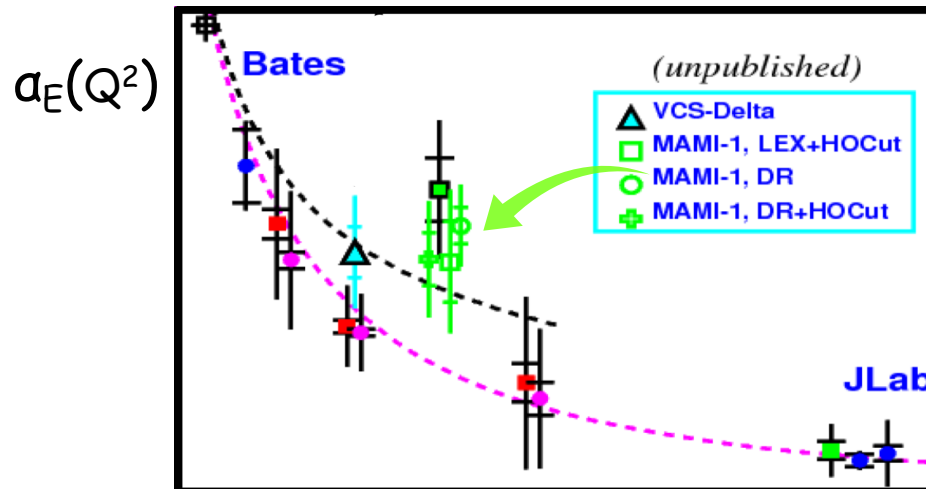
A1/3-12 @ MAMI

Revisiting the $Q^2=0.33 \text{ GeV}^2$ data

$Q^2 = 0.33 \text{ (GeV/c)}^2$ measured twice at MAMI - two different experiments

- Phys. Rev. Lett 85, 708 (2000)
- Eur. Phys. J. A37, 1-8 (2008)

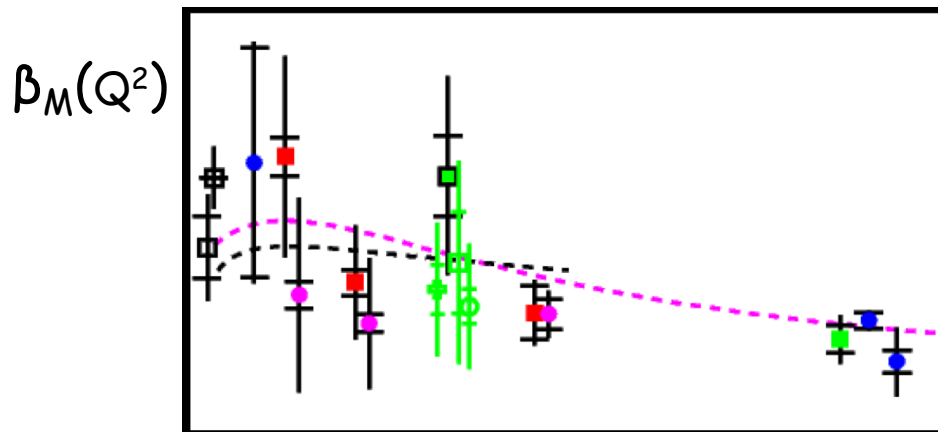
Analysis revisited (unpublished):



Re-fits at
 $Q^2=0.33$
 GeV^2
(H.F.)

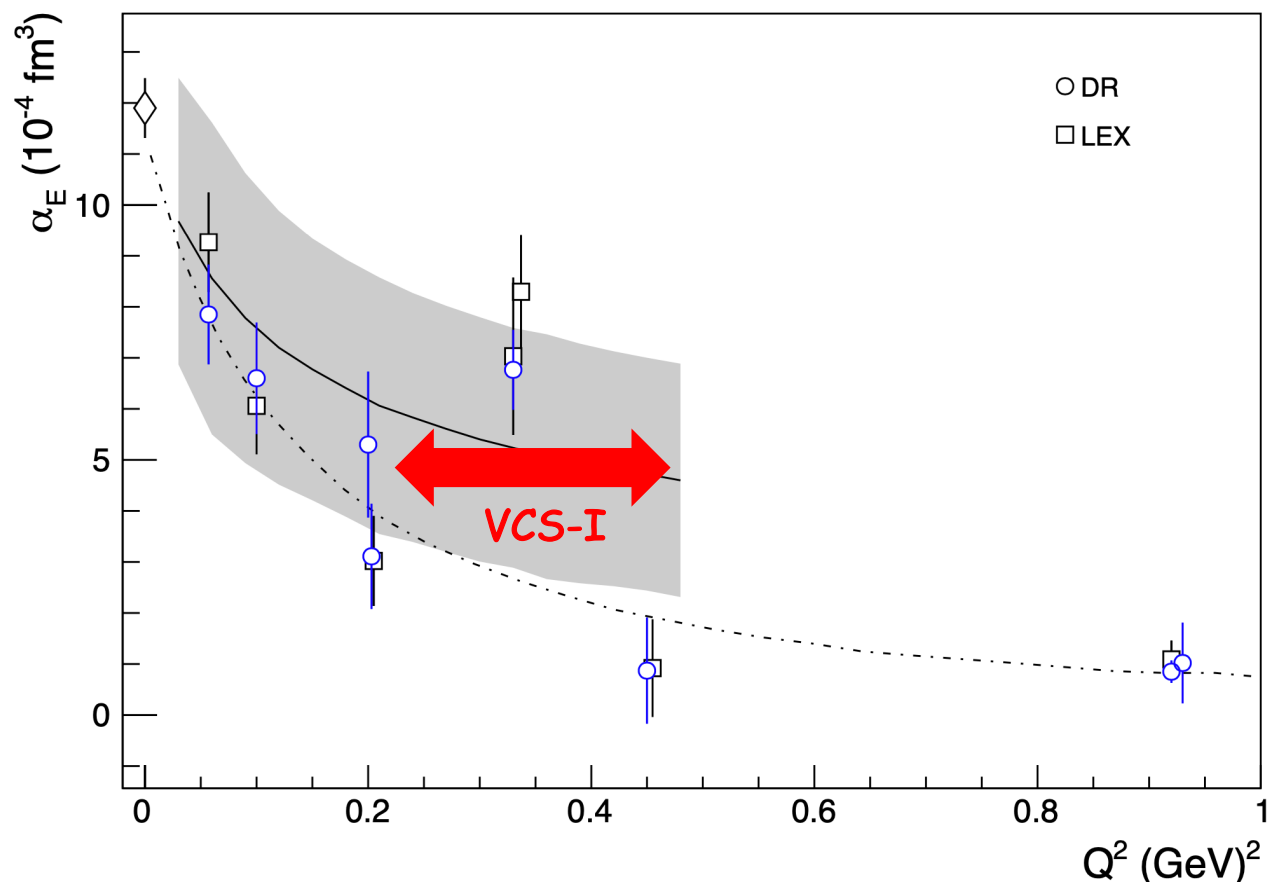
LEX and DR
Updated HO-cut

The α_E puzzle still holds



Jlab : Experiment E12-15-001 (VCS-I) in Hall C

High precision measurements targeting explicitly the kinematics of interest for a_E



Hall C HMS and SHMS

Slide Courtesy
of H. Feneker

SHMS:

- 11-GeV Spectrometer
- Partner of existing 6-GeV HMS

MAGNETIC OPTICS:

- Point-to Point QQQD for easy calibration and wide acceptance.
- Horizontal bend magnet allows acceptance at forward angles (5.5°)

Detector Package:

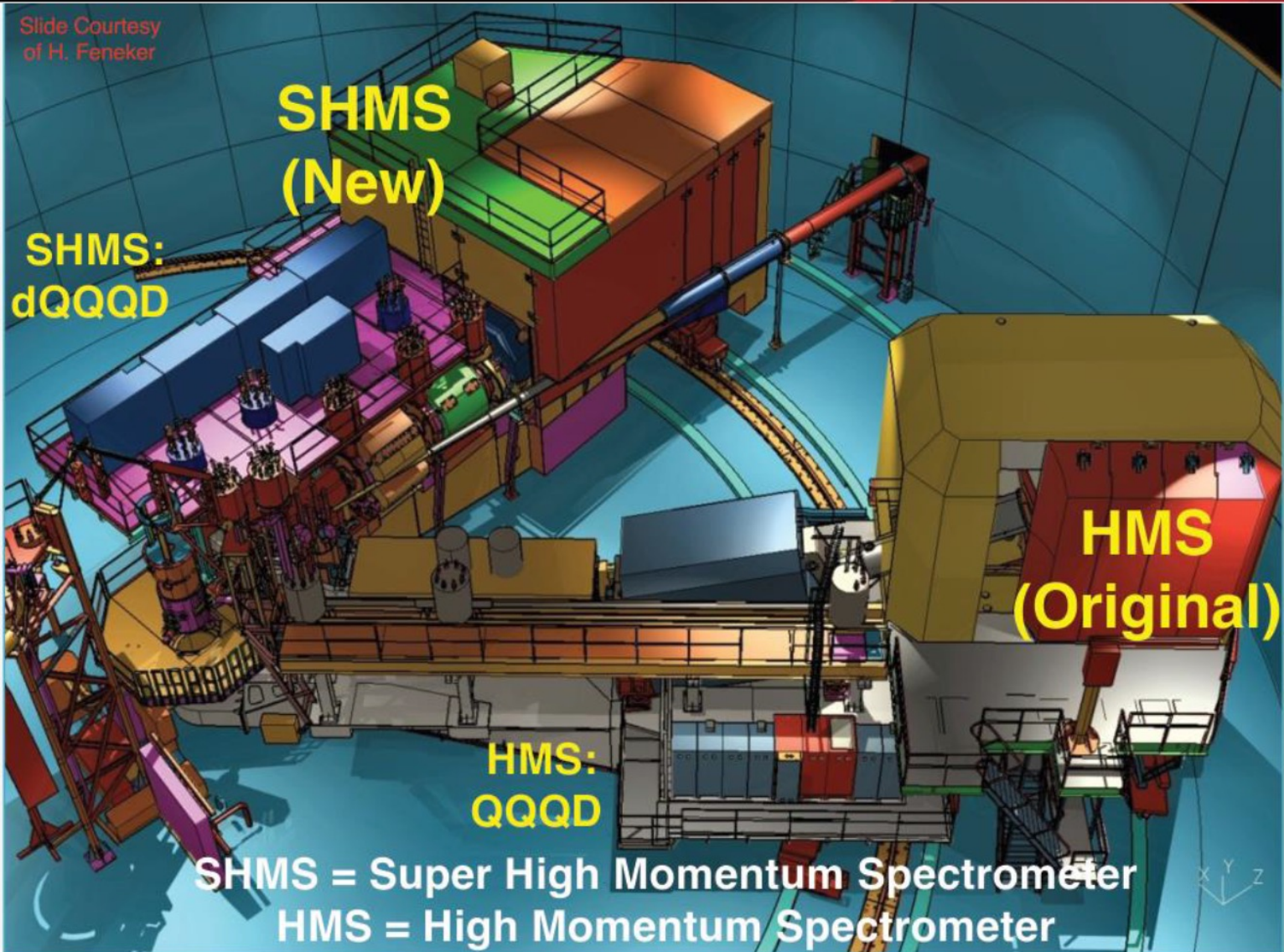
- Drift Chambers
- Hodoscopes
- Cerenkovs
- Calorimeter
- All derived from existing HMS/SOS detector designs

• Super High Momentum Spectrometer

- HB, 3 Quads, Dipole
- $P \rightarrow 2 - 11 \text{ GeV}$
- Resolution: $\delta < 0.1\%$
- Acceptance: $\delta \rightarrow 30\%$, 4 msr
- $5.5^\circ < \theta < 40^\circ$
- Good $e/\pi/K/p$ PID

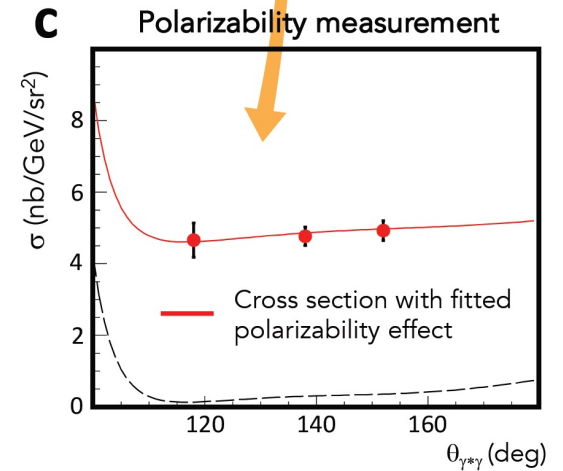
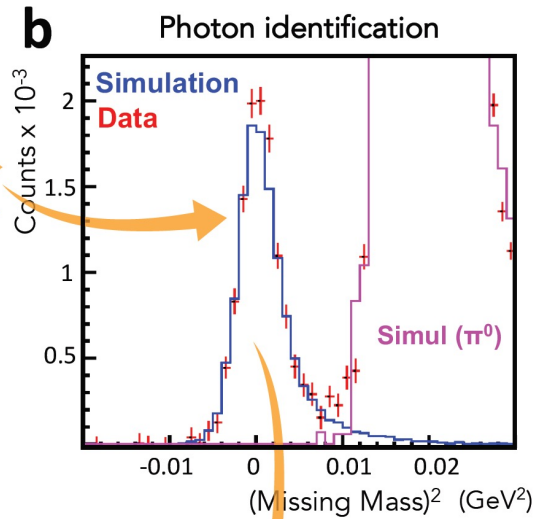
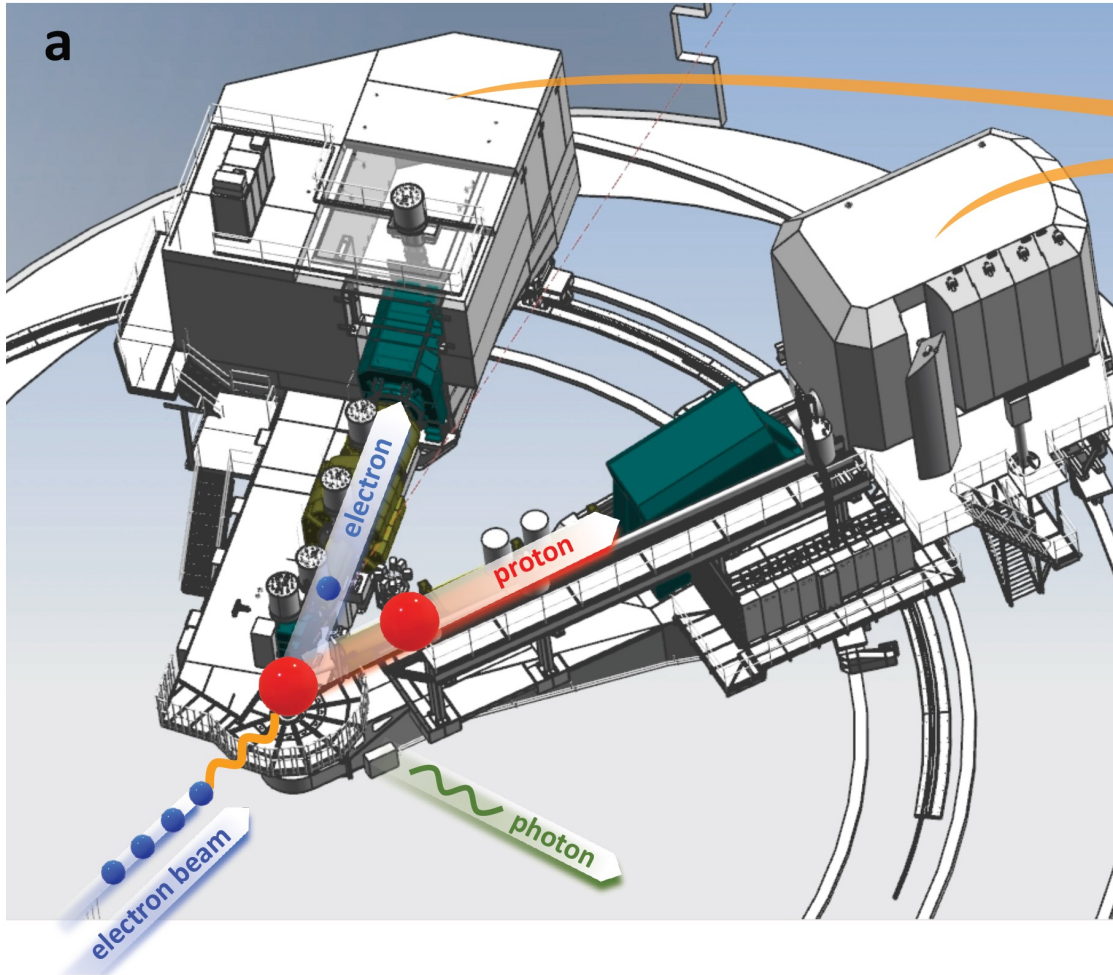
• High Momentum Spectrometer

- 3 Quads, Dipole
- $P \rightarrow 7.5 \text{ GeV}$
- Resolution: $\delta < 0.1\%$
- Acceptance: $\delta \rightarrow 18\%$, 6.5 msr
- $10.5^\circ < \theta < 90^\circ$
- Good $e/\pi/K/p$ PID



SHMS = Super High Momentum Spectrometer
HMS = High Momentum Spectrometer

The experiment

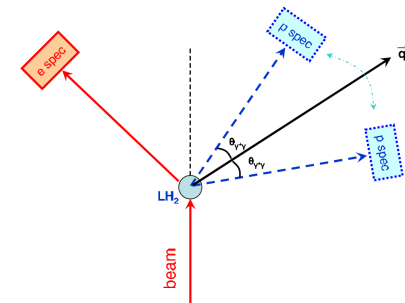


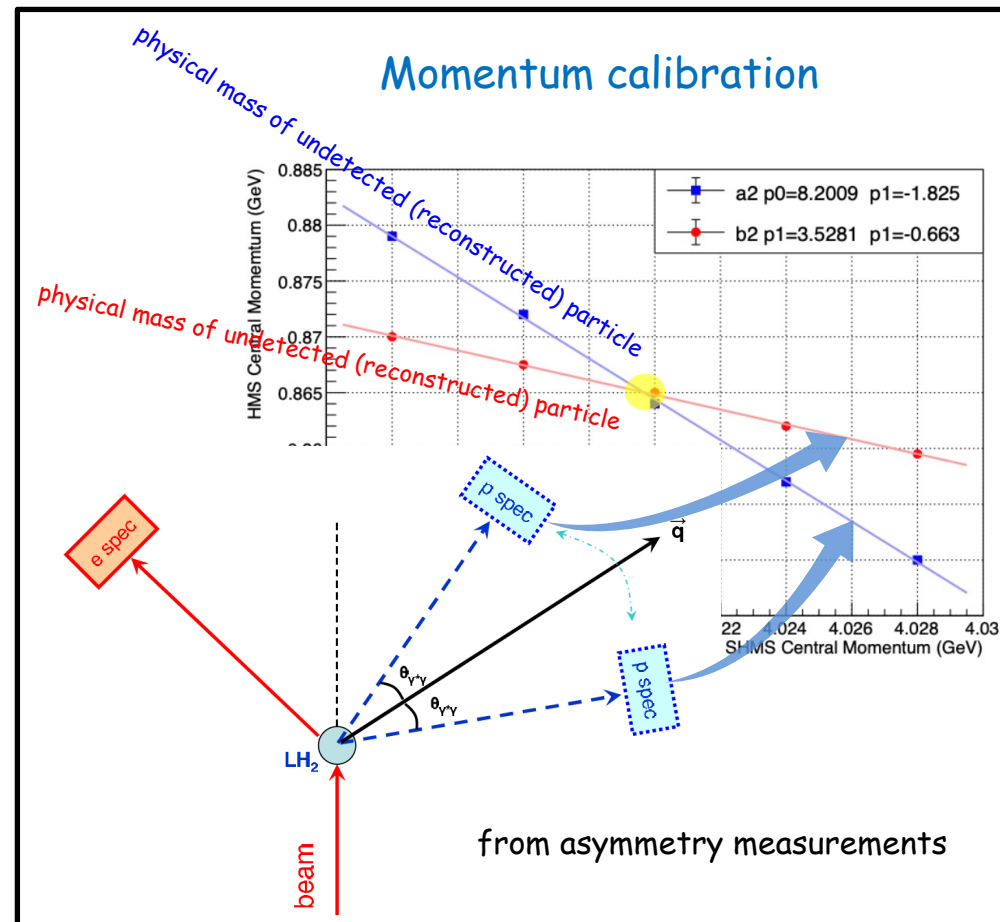
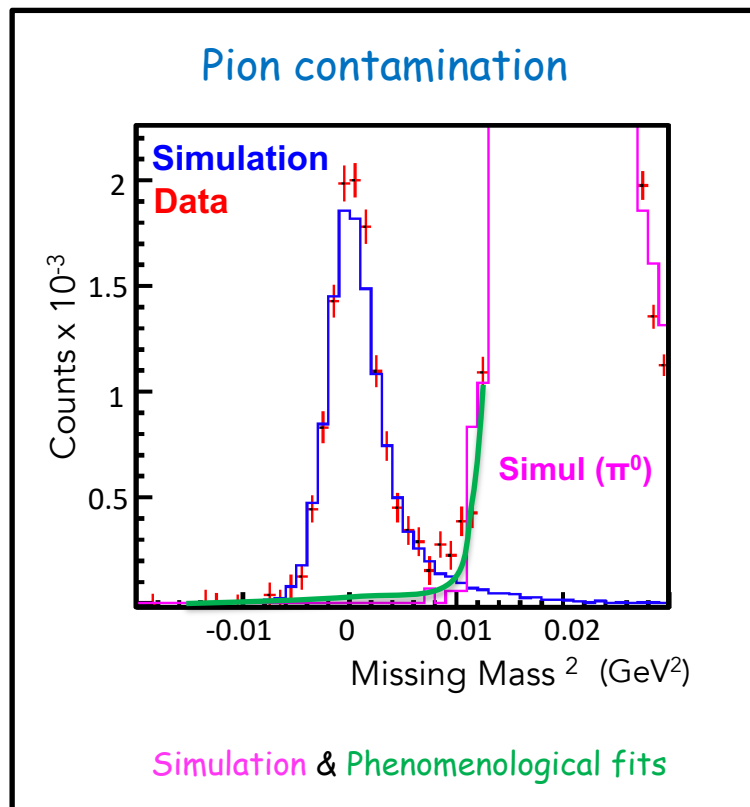
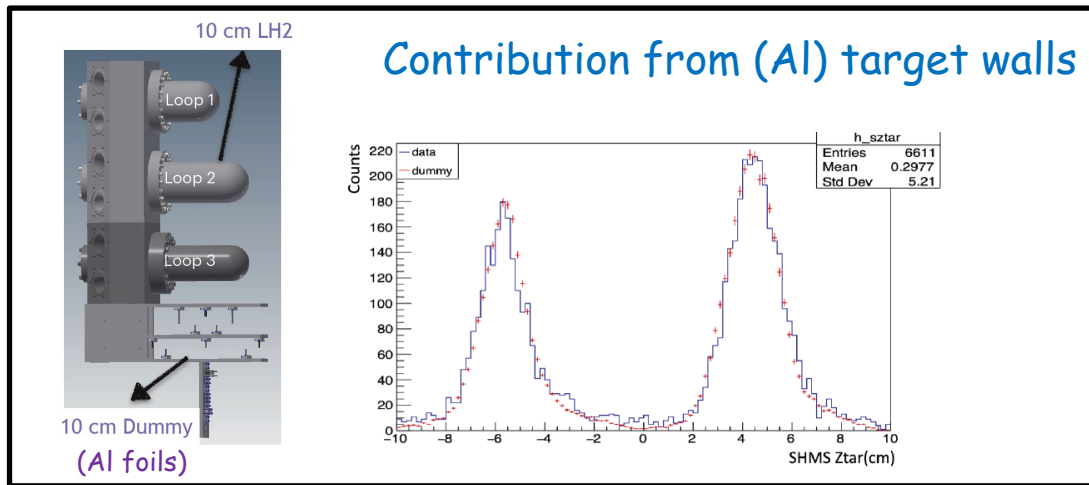
Hall C: SHMS, HMS
 4.56 GeV
 20 μ A
 Liquid hydrogen 10 cm

cross sections & azimuthal asymmetries

$$A_{(\phi_{\gamma^*\gamma}=0,\pi)} = \frac{\sigma_{\phi_{\gamma^*\gamma}=0} - \sigma_{\phi_{\gamma^*\gamma}=180}}{\sigma_{\phi_{\gamma^*\gamma}=0} + \sigma_{\phi_{\gamma^*\gamma}=180}}$$

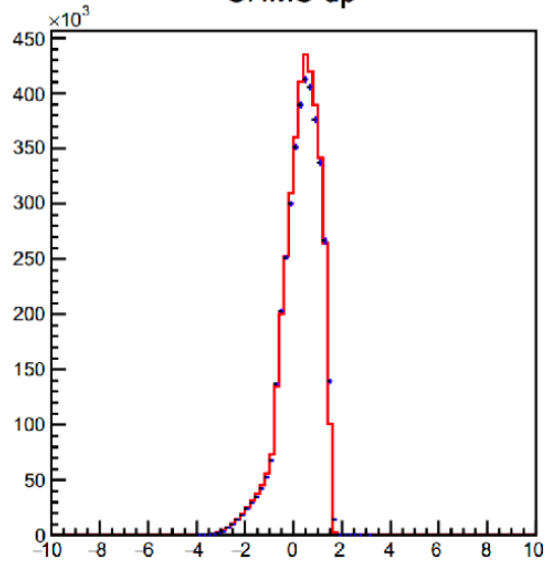
sensitivity to GPs



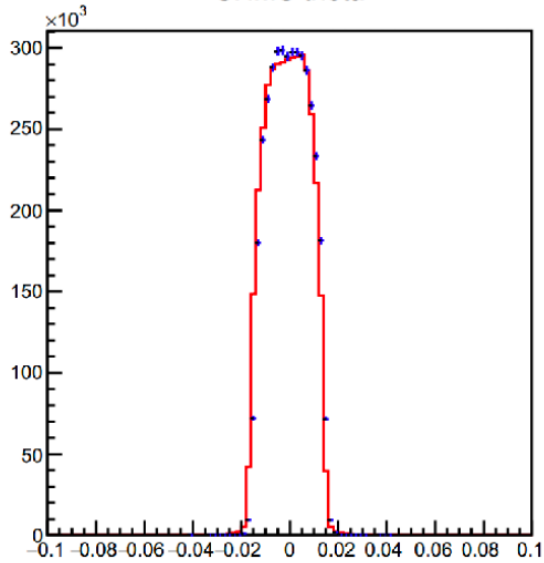


Elastic data

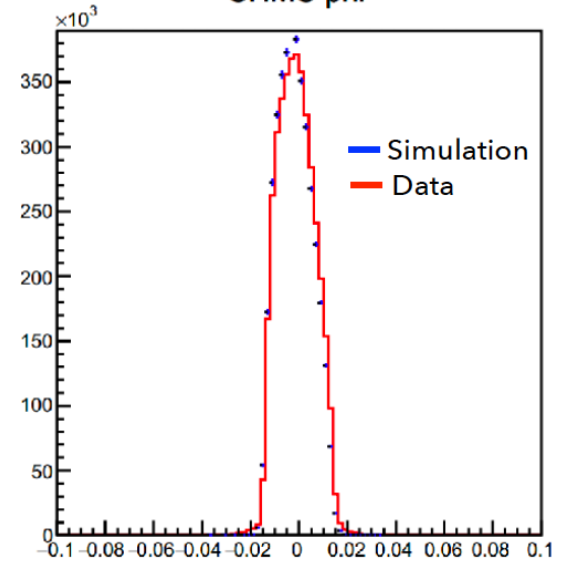
SHMS d_p



SHMS theta

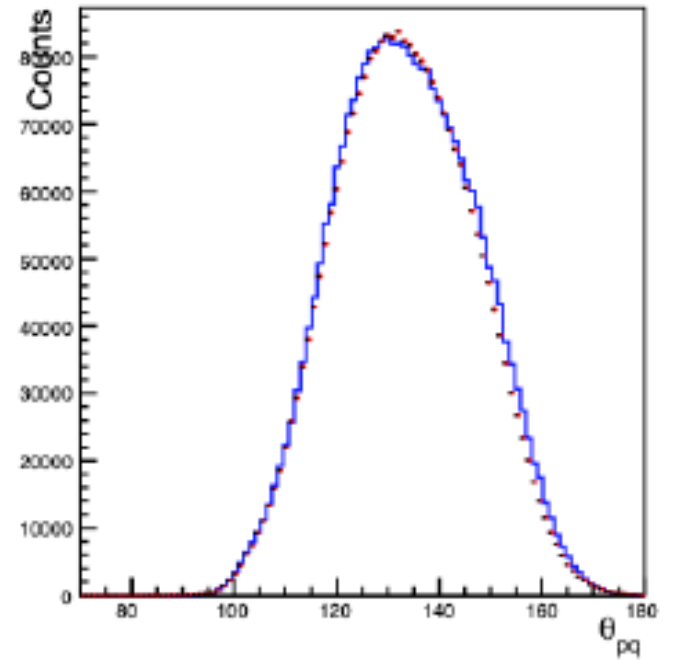
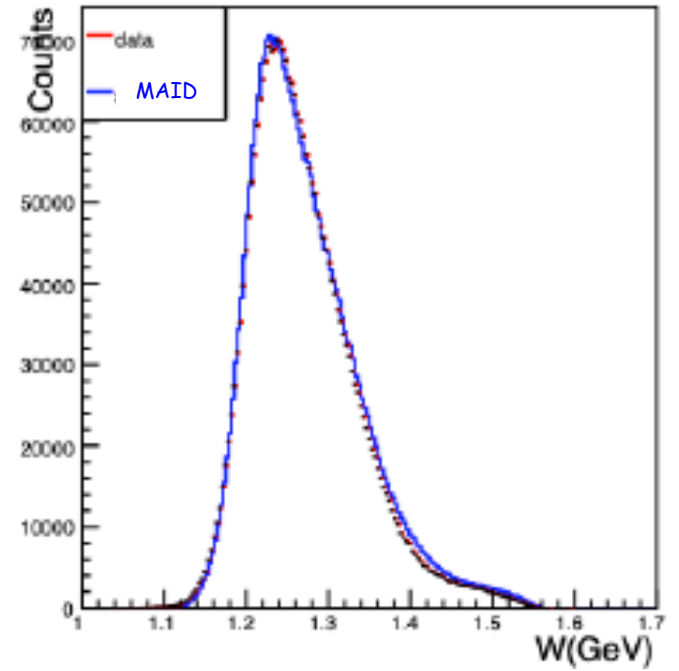
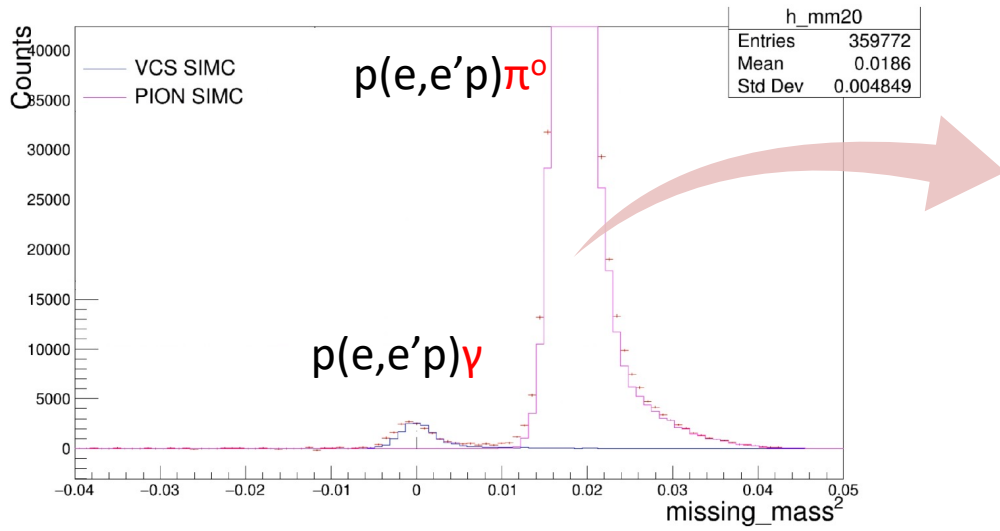


SHMS phi

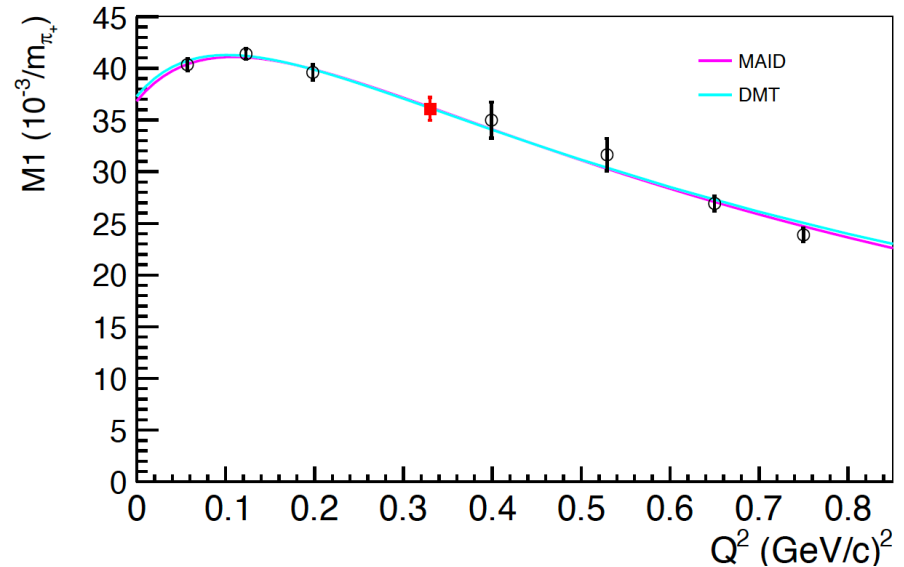
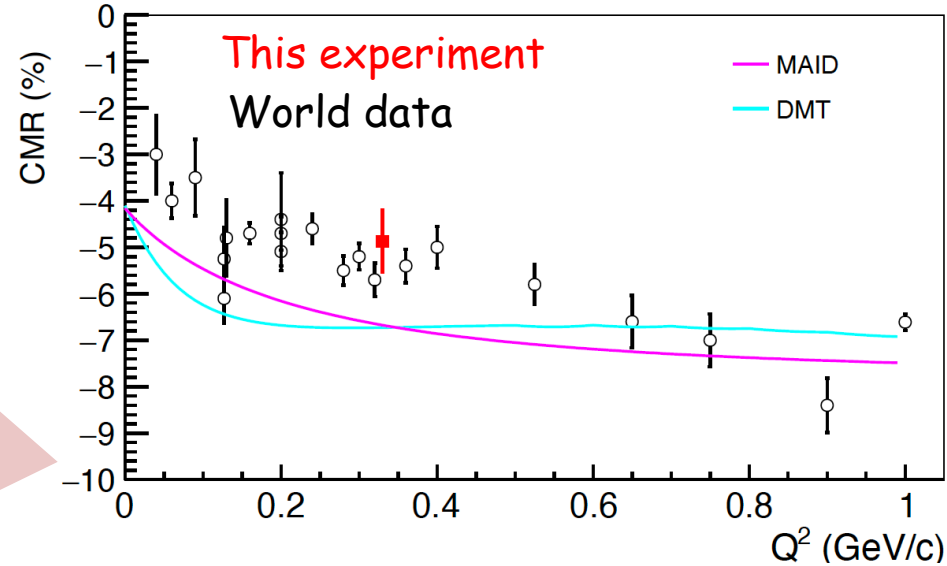
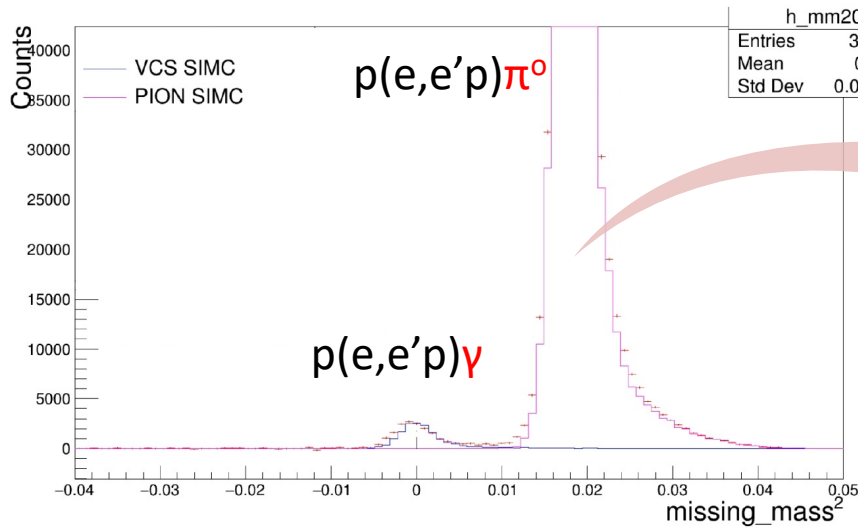


| Kinematic | θ_e° | $P_e(\text{GeV}/c)$ | θ_p° | $P_p(\text{GeV}/c)$ |
|-------------|------------------|---------------------|------------------|---------------------|
| Elastic I | 10.76 | 4.193 | 61.16 | 0.893 |
| Elastic II | 10.41 | 4.214 | 61.95 | 0.863 |
| Elastic III | 9.64 | 4.259 | 63.76 | 0.795 |

$\rho(e,e'p)\pi^0$



$N \rightarrow \Delta$ TFFs



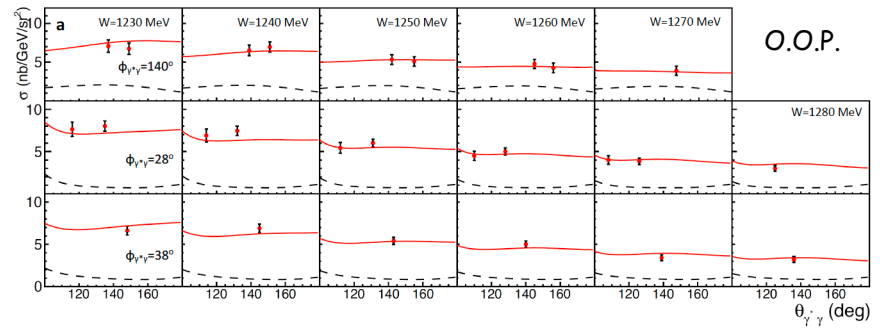
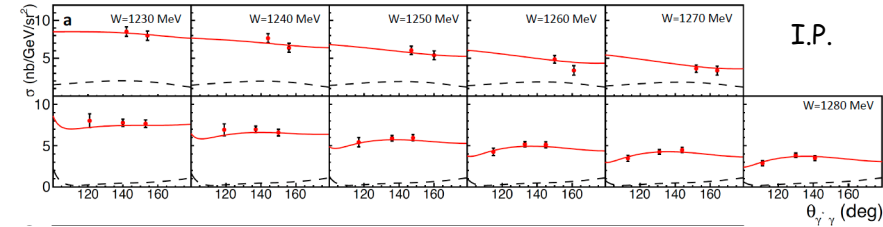
Simultaneous measurement
of the $N \rightarrow \Delta$ TFFs

TFFs well known
→ Real time normalization control

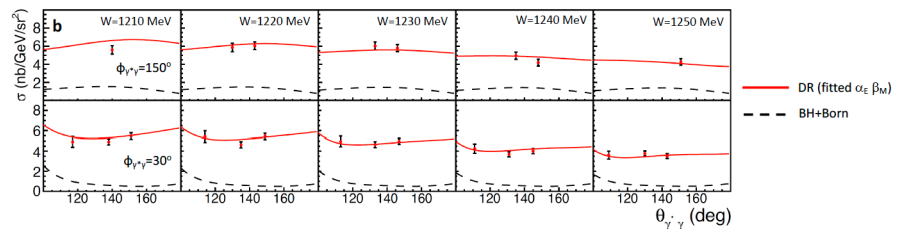
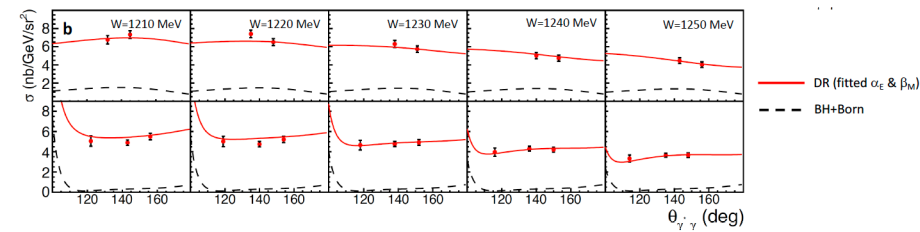
Good understanding of
spectrometer acceptance

New results: VCS cross sections

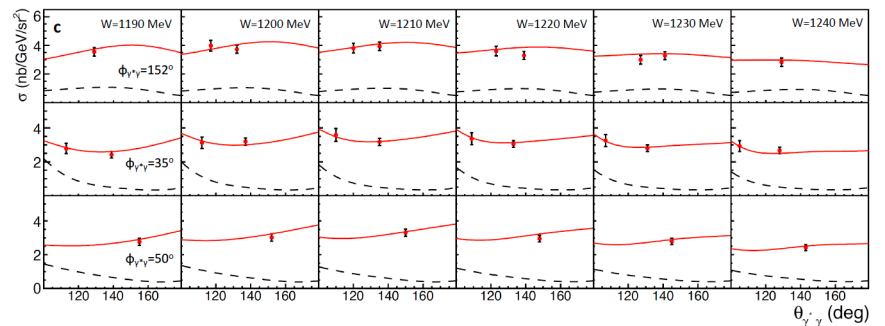
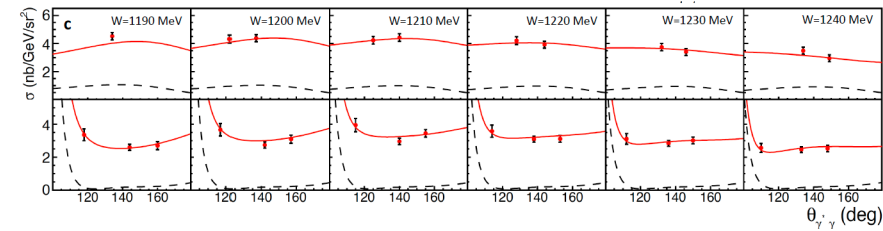
$Q^2=0.27 \text{ GeV}^2$



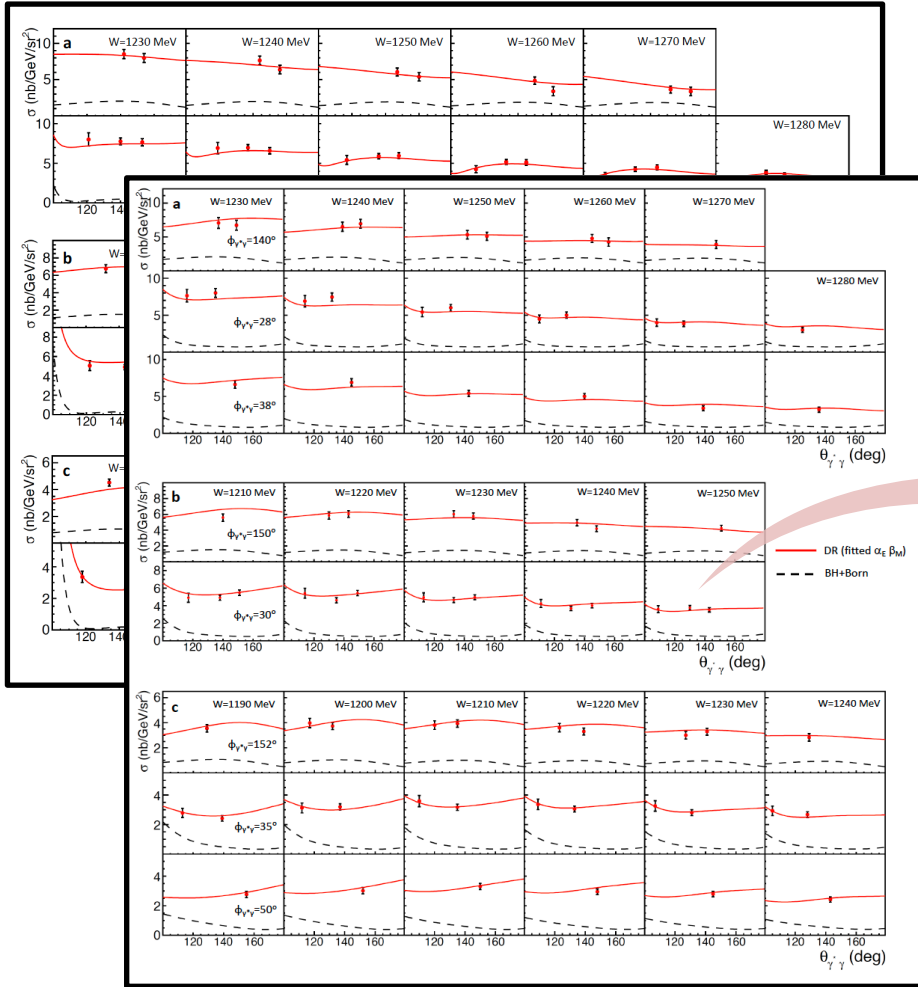
$Q^2=0.33 \text{ GeV}^2$



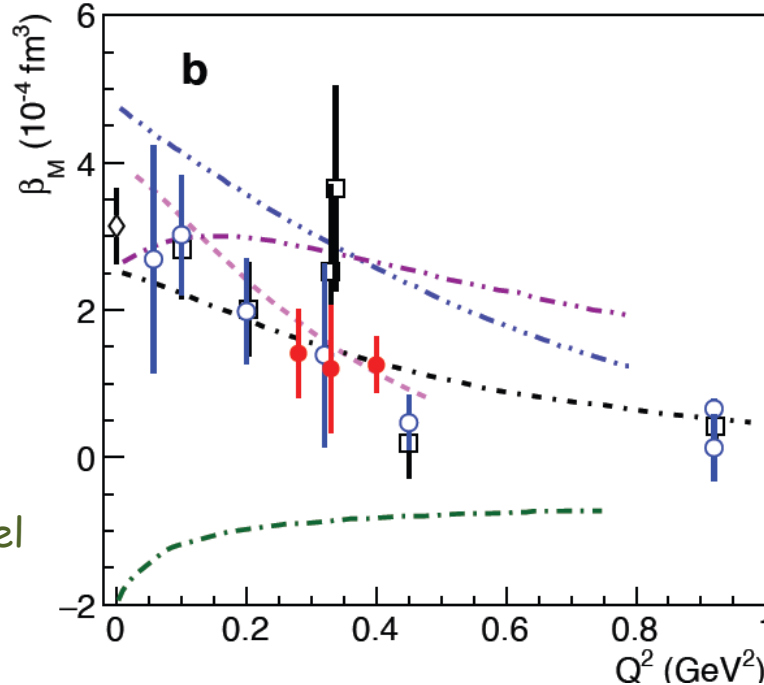
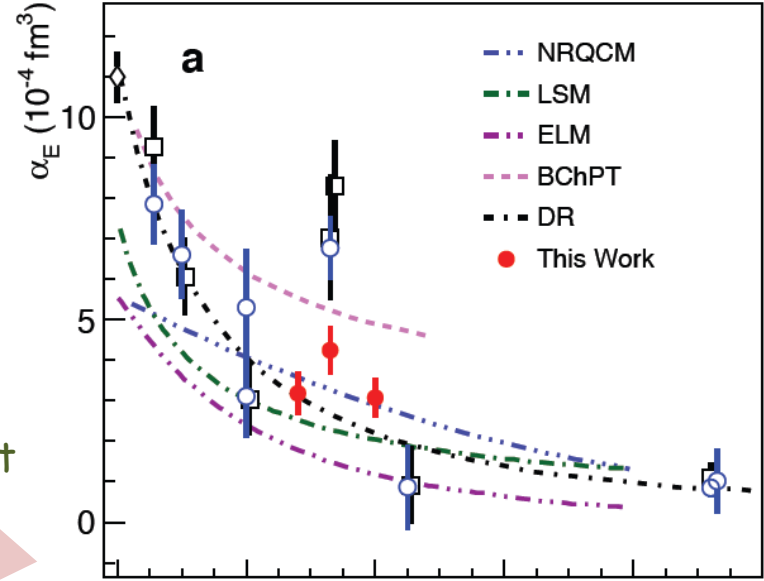
$Q^2=0.40 \text{ GeV}^2$



New results: GPs



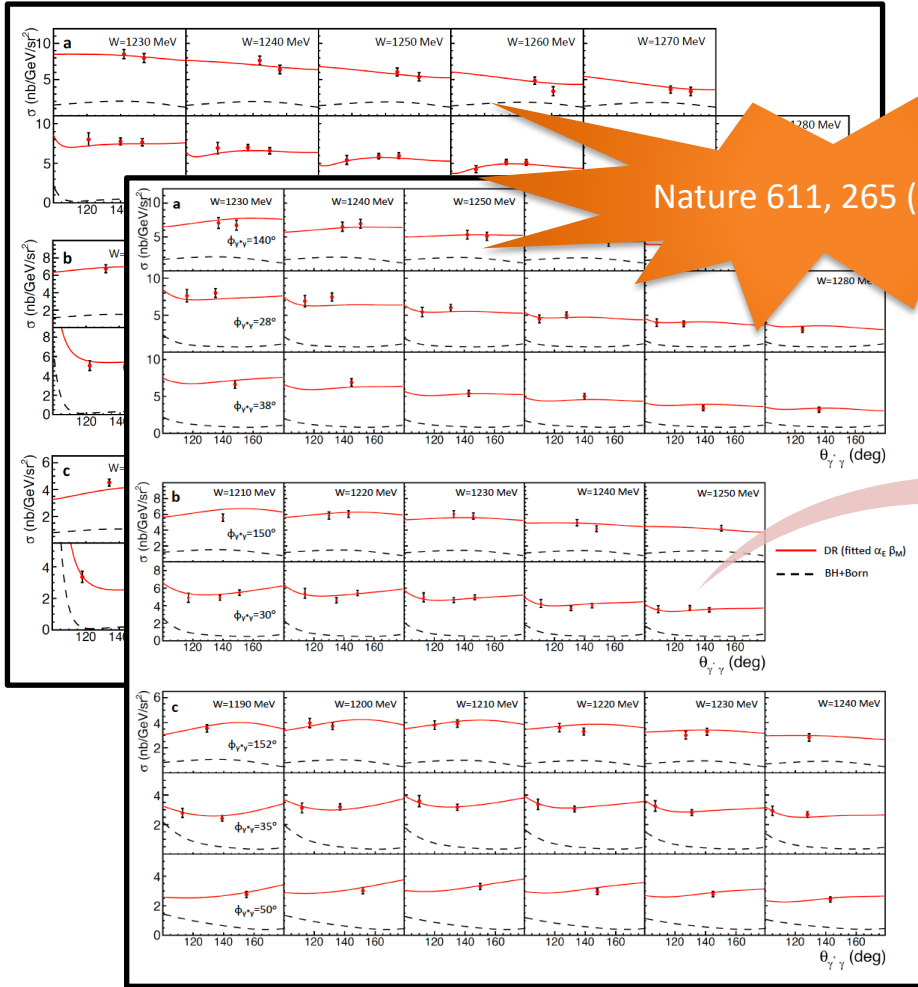
DR fit



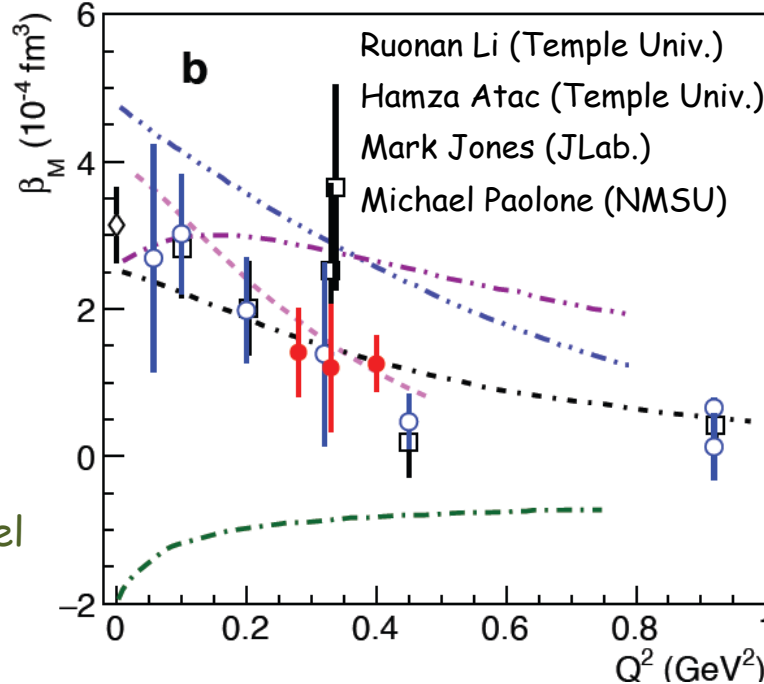
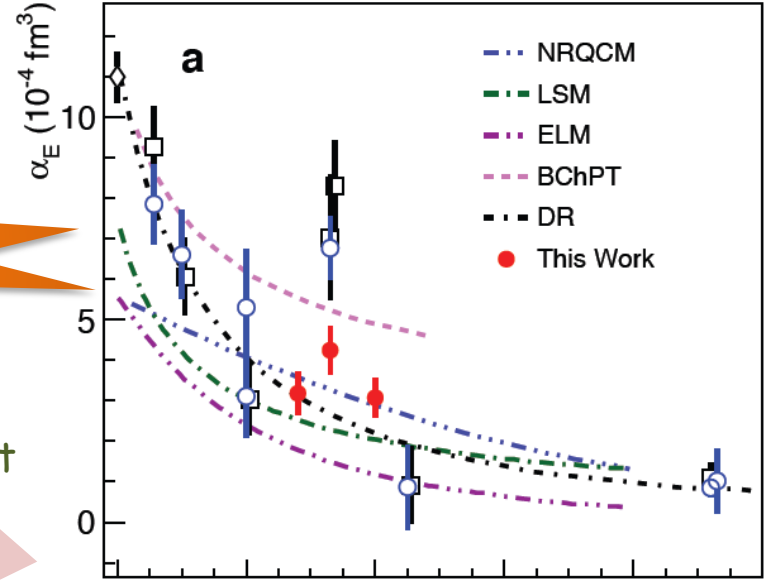
Experimental cross sections are compared to the DR model predictions for all possible values for the GPs

→ $\alpha_E(Q^2)$ and $\beta_M(Q^2)$ are fitted by a χ^2 minimization

New results: GPs



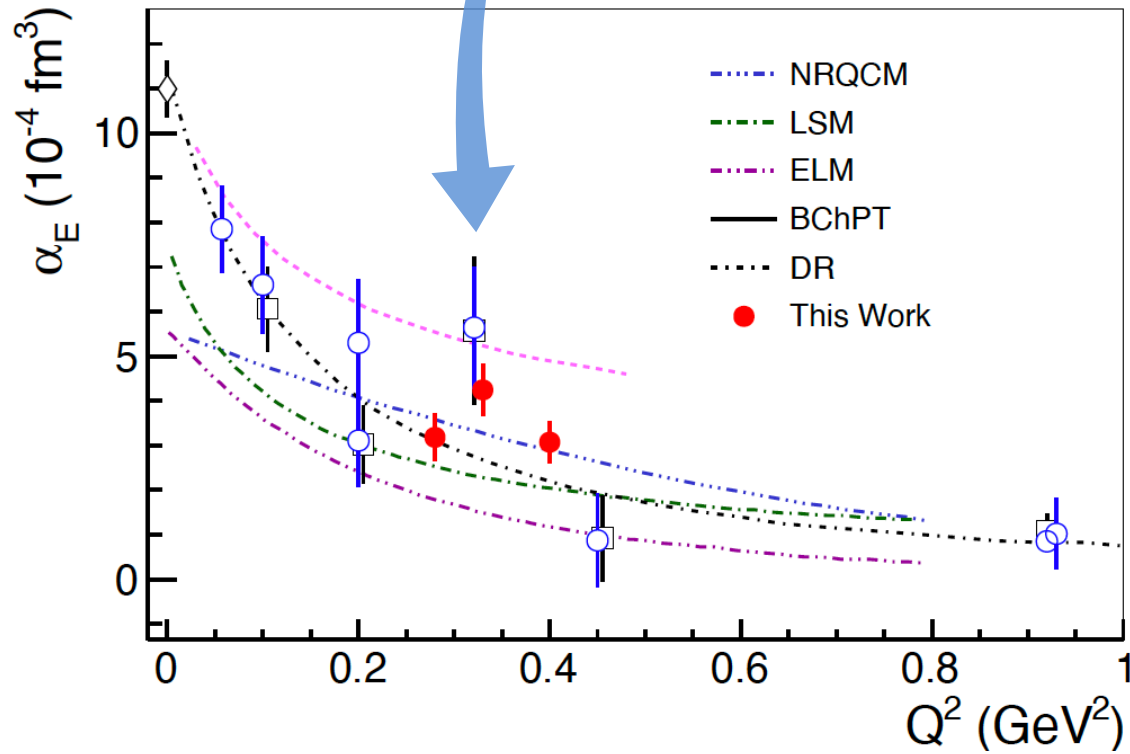
DR fit



Experimental cross sections are compared to the DR model predictions for all possible values for the GPs

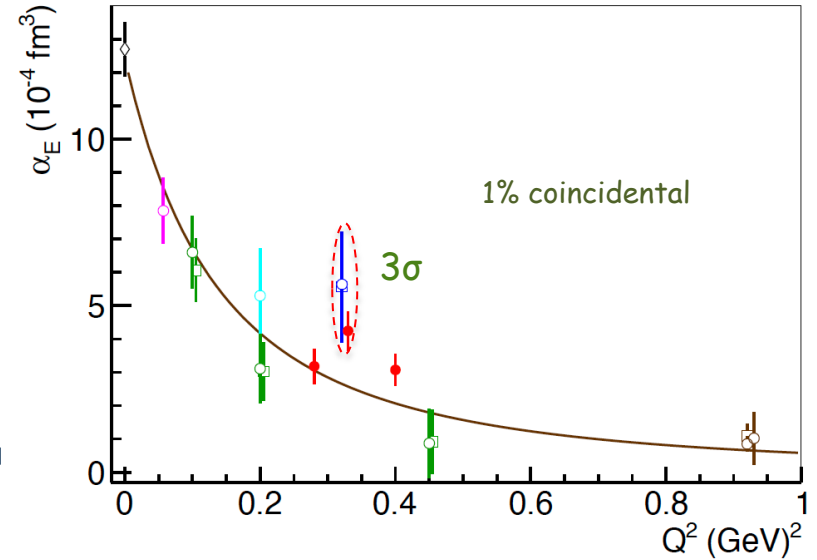
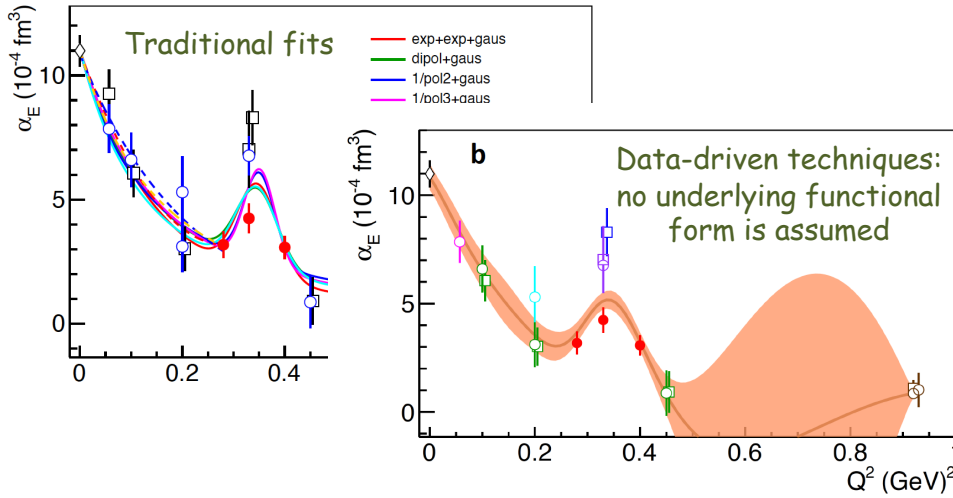
→ $\alpha_E(Q^2)$ and $\beta_M(Q^2)$ are fitted by a χ^2 minimization

MAMI-I re-analysis (unpublished)



Is there a non-trivial structure?

Electric GP



Is the observed a_E structure coincidental or not?

If true: Measure the shape precisely \rightarrow input to theory

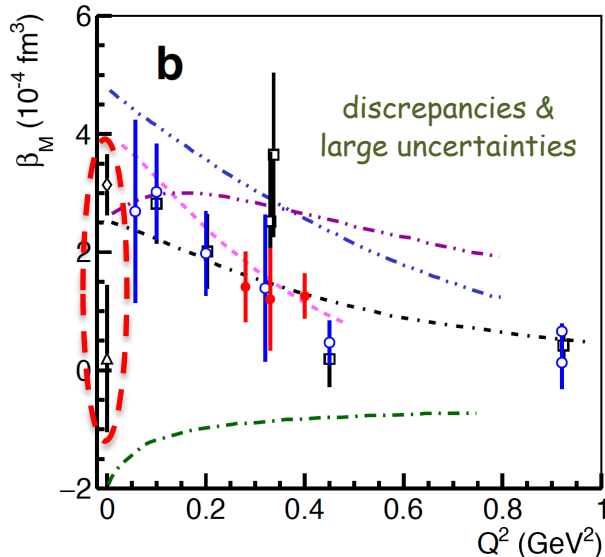
If not: We are able to show it with more measurements

Strong tension between world data (?)

Things we do not yet understand well?

Underestimated uncertainties? ...

Magnetic GP



Magnetic GP: Large uncertainties & discrepancies

Instrumental to disentangle diamagnetism and paramagnetism in the proton

Unique opportunity to measure a_E and β_M with superb precision and with consistent systematics across Q^2

Theory: B χ PT

Generalized polarizabilities of the nucleon in baryon chiral perturbation theory

Vadim Lensky^{1,2,3,a}, Vladimir Pascalutsa¹, Marc Vanderhaeghen¹

¹ Institut für Kernphysik, Cluster of Excellence PRISMA, Johannes Gutenberg Universität Mainz, 55128 Mainz, Germany

² Institute for Theoretical and Experimental Physics, Moscow 117218, Russia

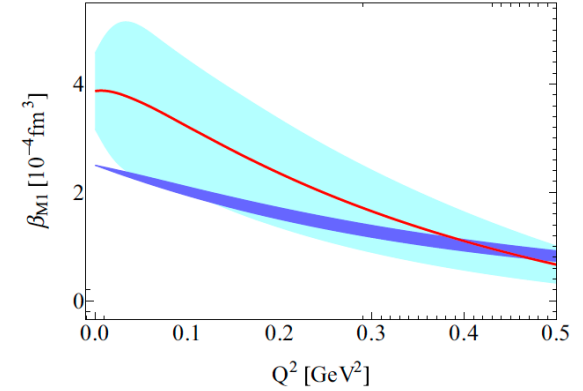
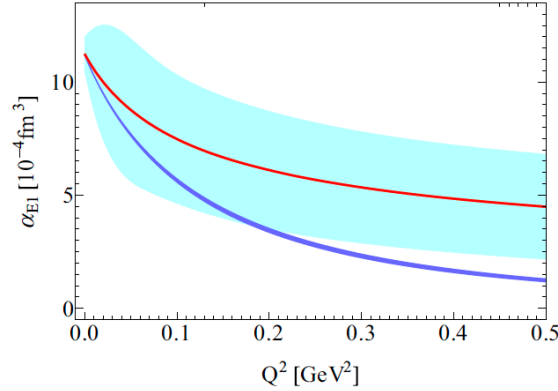
³ National Research Nuclear University MEPhI (Moscow Engineering Physics Institute), Moscow 115409, Russia



B χ PT calculation to NLO
in the δ -counting scheme



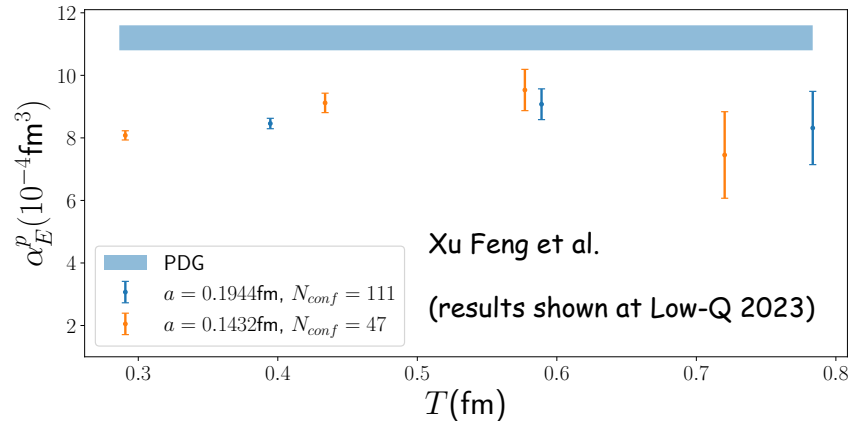
DR calculation
D. Drechsel, B. Pasquini, M. Vanderhaeghen,
Phys. Rep. 378,99 (2003)



Theory: Lattice QCD

Lattice QCD results for
the static polarizabilities

Next step: Lattice QCD
calculations for the GPs



Spatial dependence of induced polarizations

Nucleon form factor data → light-front quark charge densities

Formalism extended to the deformation of these quark densities when applying an external e.m. field:

GPs → spatial deformation of charge & magnetization densities under an applied e.m. field

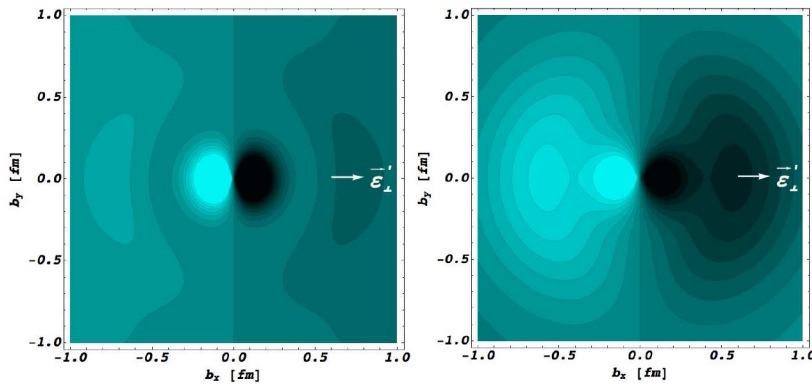
Induced polarization in a proton when submitted to an e.m. field

Phys. Rev. Lett. 104, 112001 (2010)

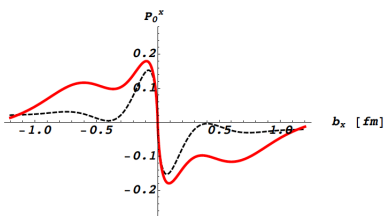
M. Gorchtein, C. Lorce, B. Pasquini, M. Vanderhaeghen

GP I

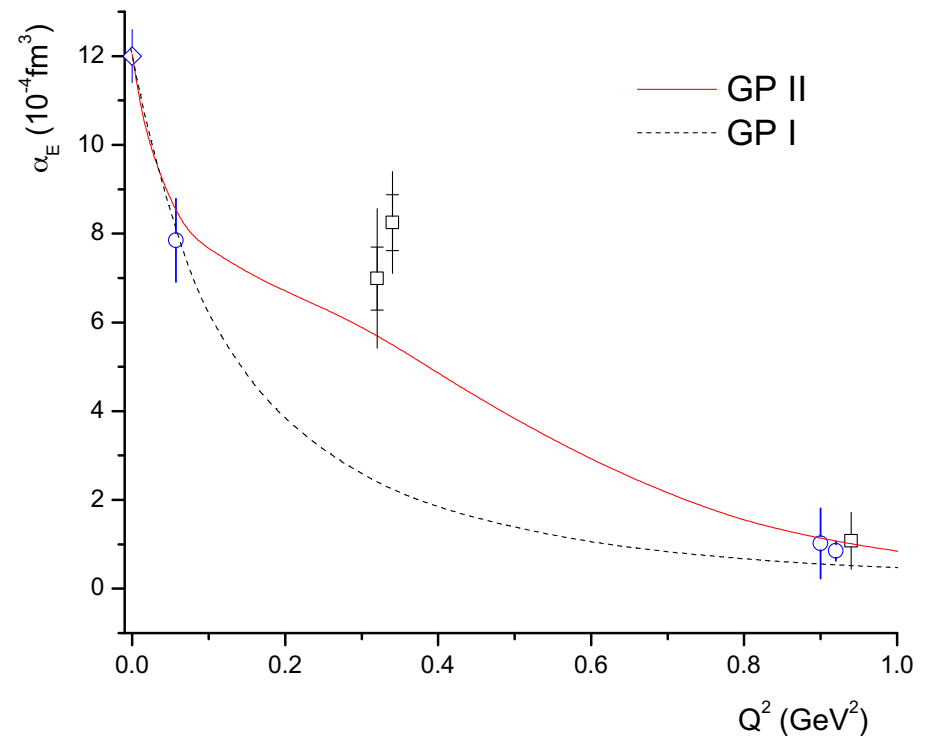
GP II



Light (dark) regions → largest (smaller) values
(photon polarization along x-axis, as indicated)



Induced polarization along $b_y=0$

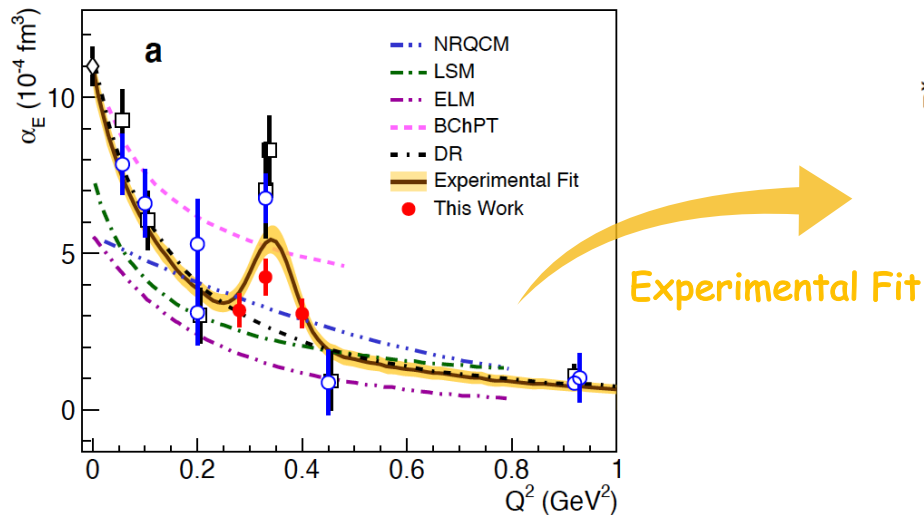


Spatial dependence of induced polarizations

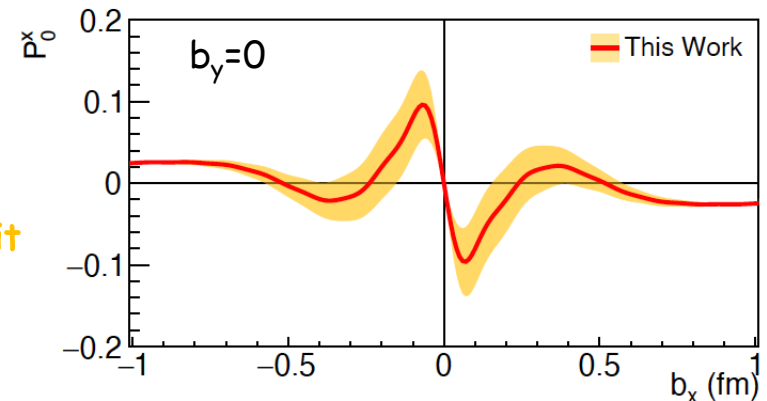
Nucleon form factor data → light-front quark charge densities

Formalism extended to the deformation of these quark densities when applying an external e.m. field:

GPs → spatial deformation of charge & magnetization densities under an applied e.m. field



Induced polarization in a proton when submitted to an e.m. field

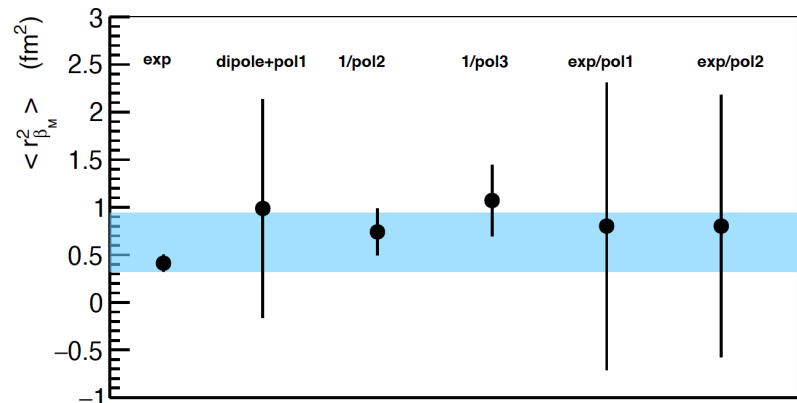
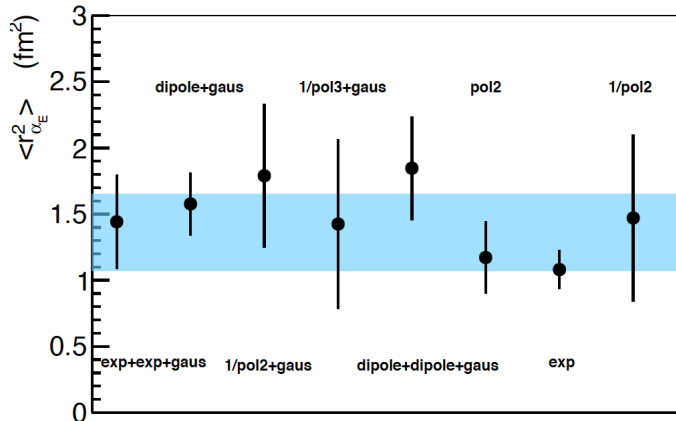
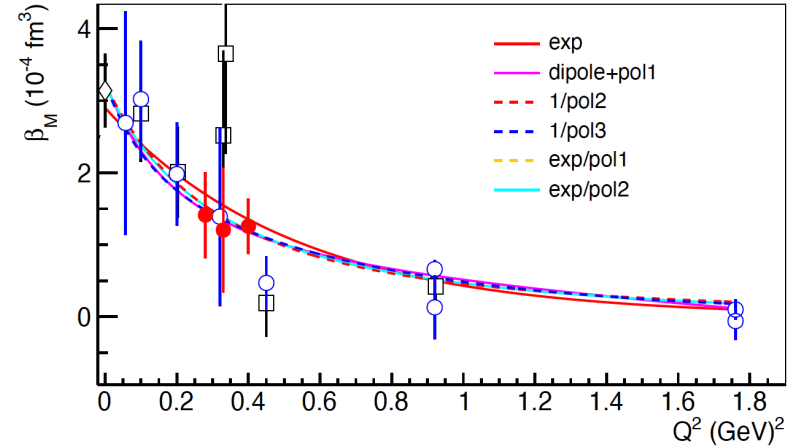
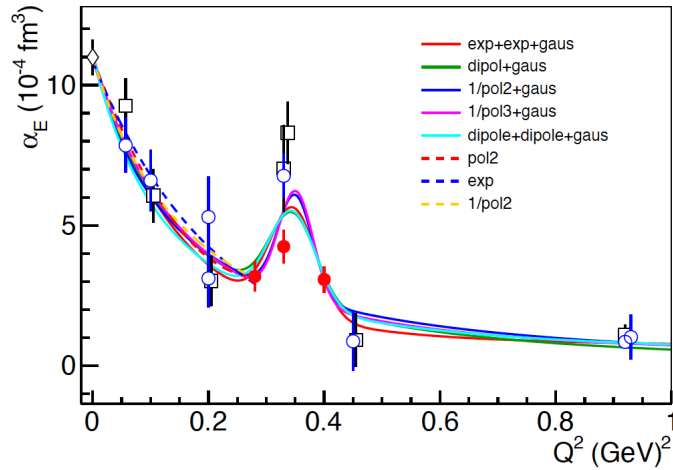


x - y defines the transverse plane with the z -axis being the direction of the fast-moving proton

Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$



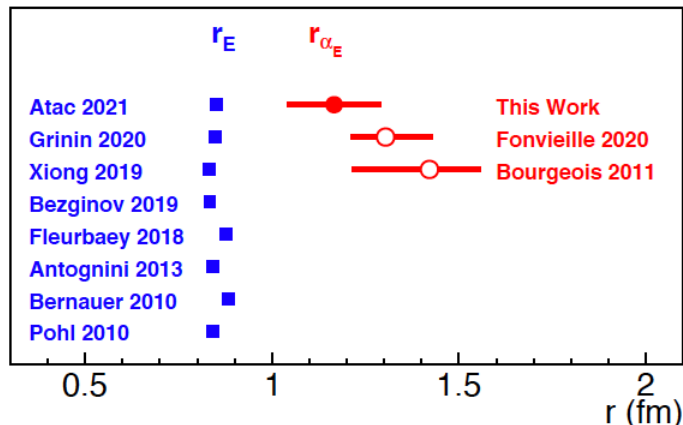
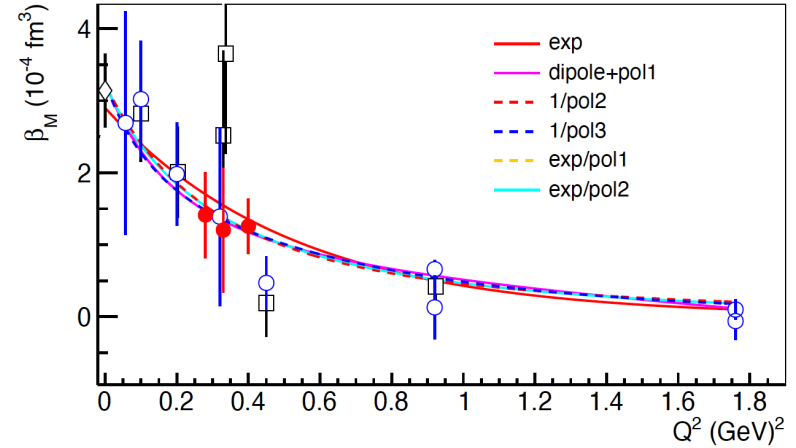
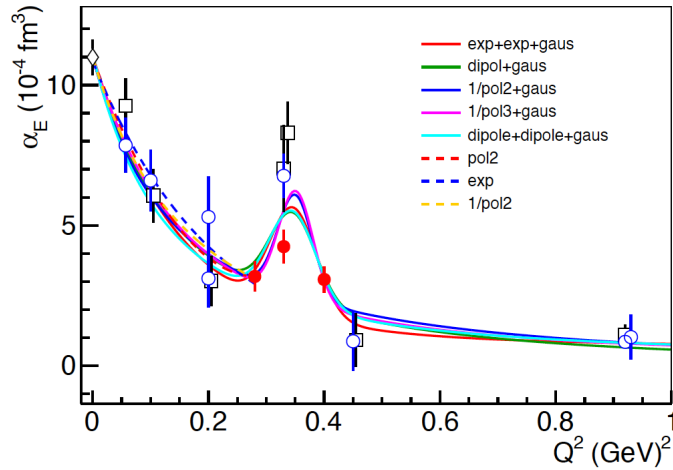
$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$

$$\langle r_{\beta_M}^2 \rangle = 0.63 \pm 0.31 \text{ fm}^2$$

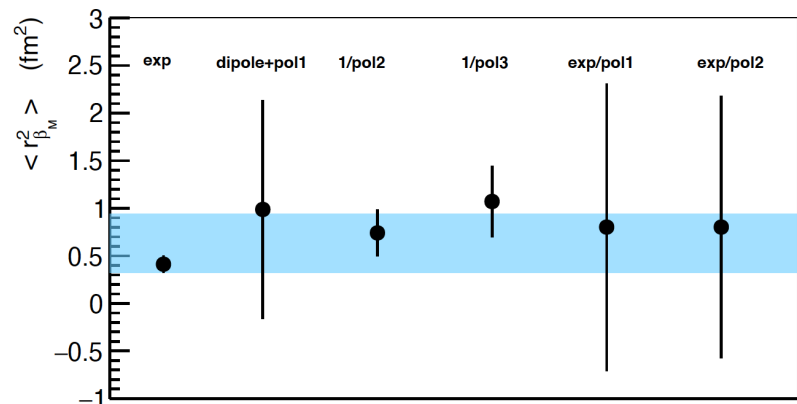
Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$



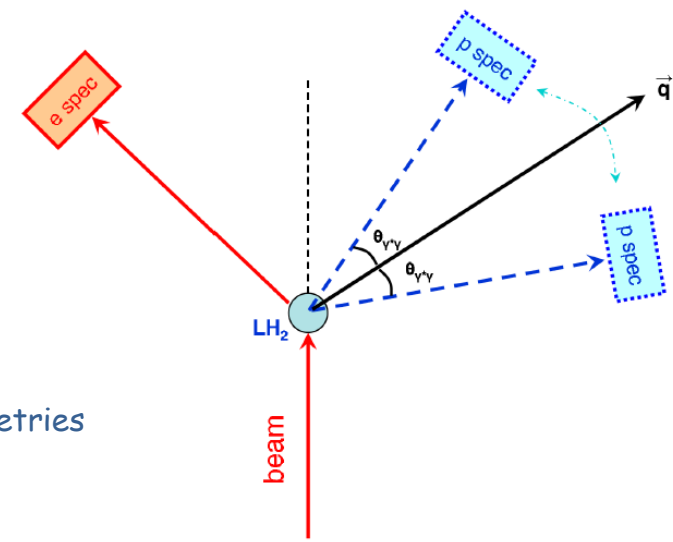
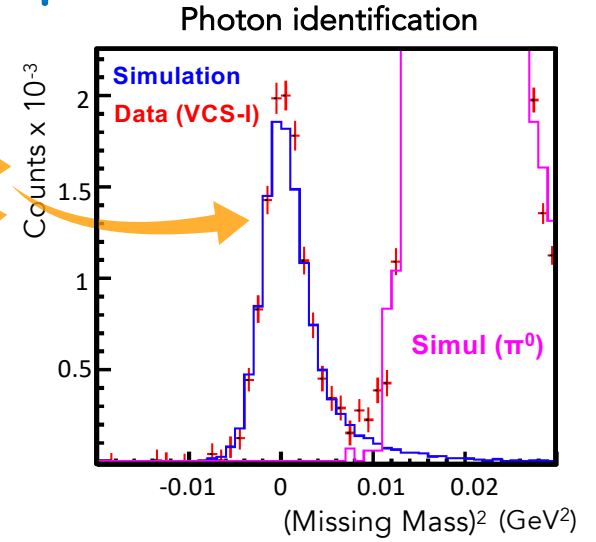
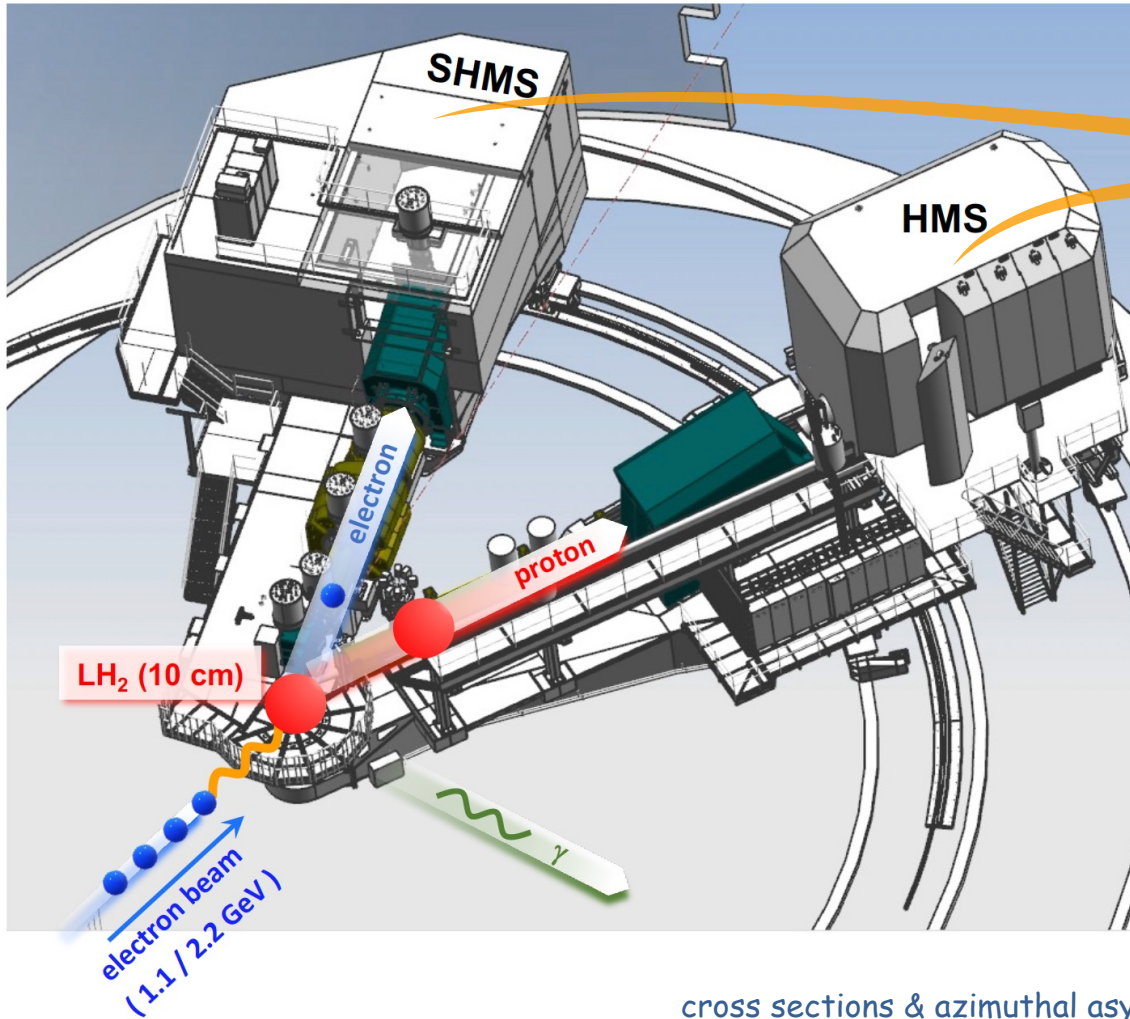
$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$



$$\langle r_{\beta_M}^2 \rangle = 0.63 \pm 0.31 \text{ fm}^2$$

Moving Forward

VCS-II Experimental Setup



cross sections & azimuthal asymmetries

$$A_{(\phi_{\gamma^*\gamma}=0,\pi)} = \frac{\sigma_{\phi_{\gamma^*\gamma}=0} - \sigma_{\phi_{\gamma^*\gamma}=180}}{\sigma_{\phi_{\gamma^*\gamma}=0} + \sigma_{\phi_{\gamma^*\gamma}=180}}$$

suppression of systematic uncertainties

Extra handle to spectrometer momentum calibration

Hall C:

SHMS, HMS

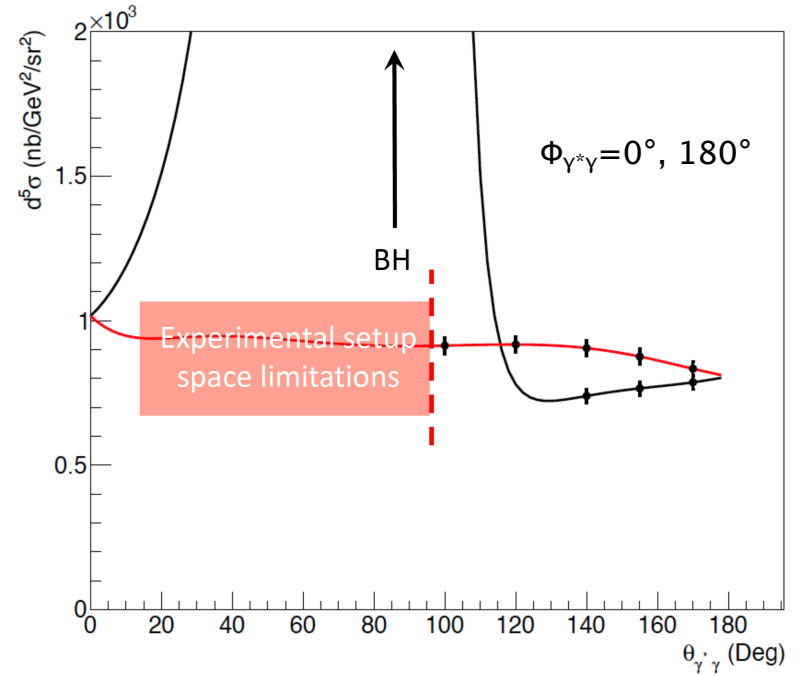
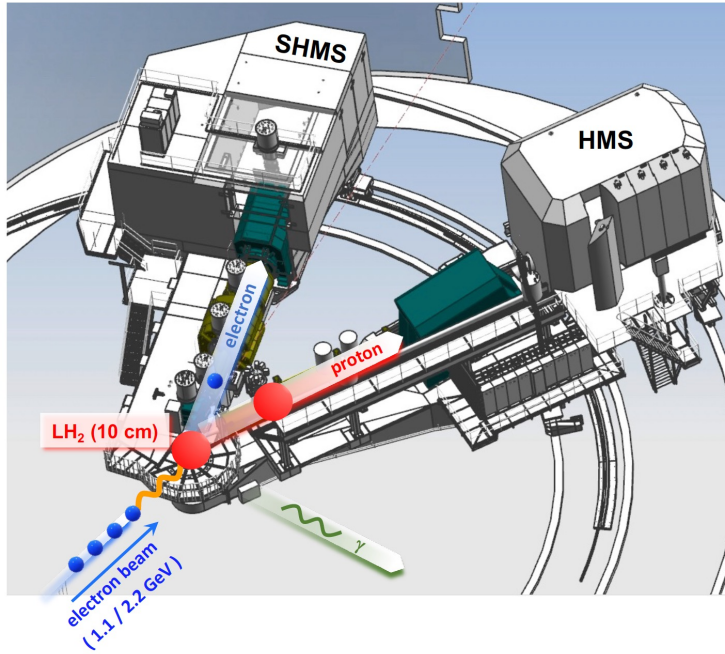
E= 1.1 GeV (lowest Q²)

E=2.2 GeV (all other settings)

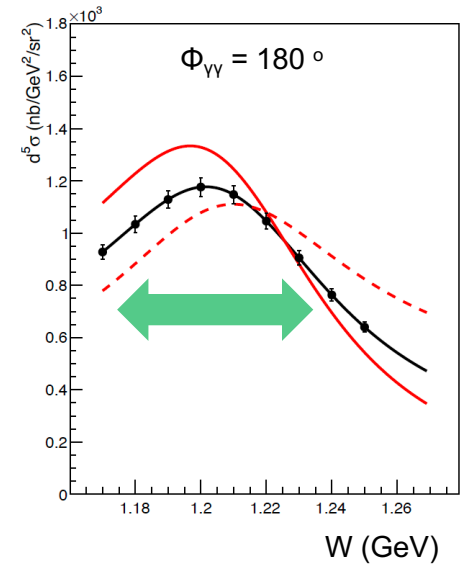
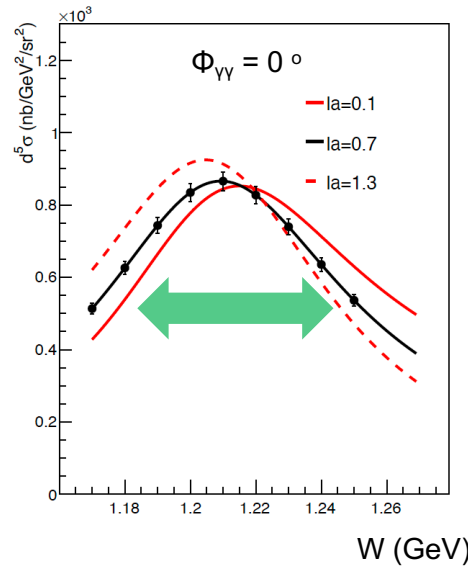
Liquid hydrogen 10 cm

VCS-II Experimental Setup

Selection of kinematics

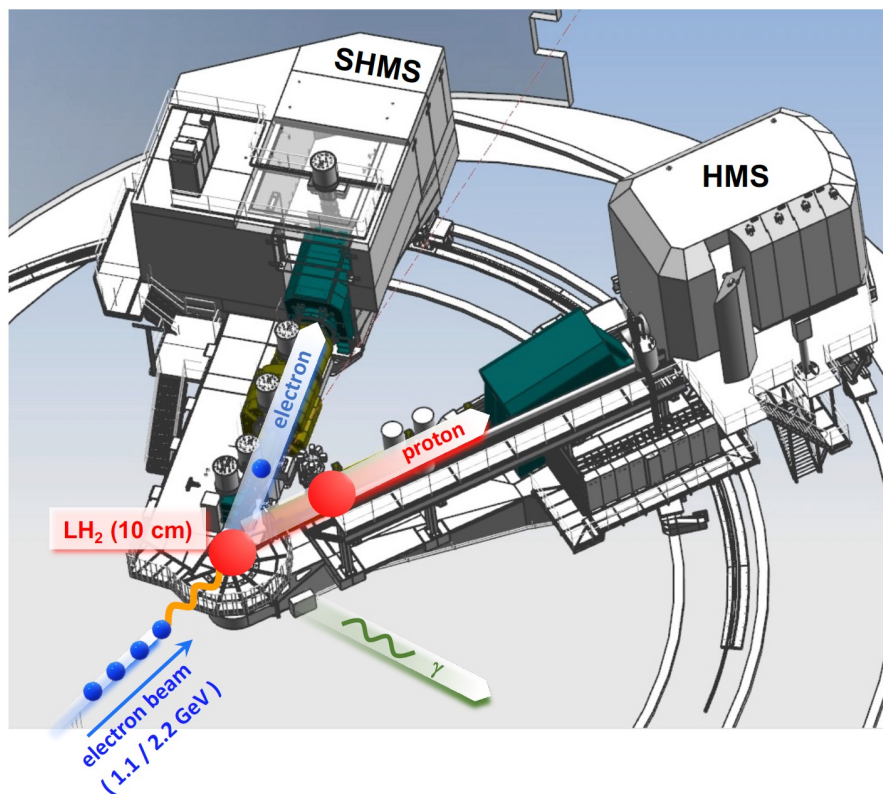


Targeted measurements to fully exploit the sensitivity to the GPs



Kinematics

Q^2 spans 0.05 GeV^2 to 0.5 GeV^2



Production ($E_0 = 1.1 \text{ GeV}$): 6 days

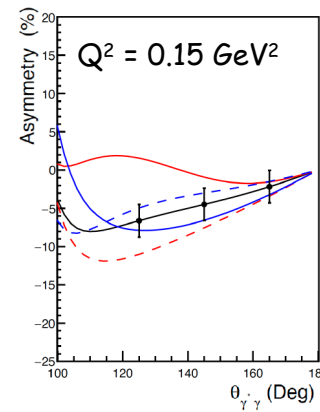
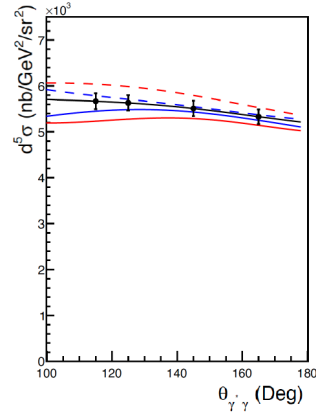
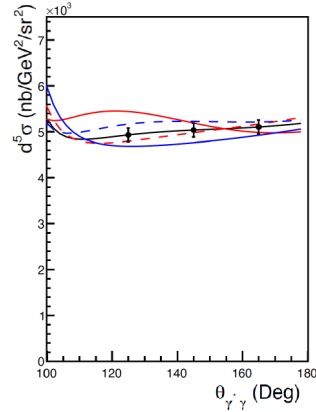
Production ($E_0 = 2.2 \text{ GeV}$): 53 days

Studies (optics/dummy/calibrations): 3 days

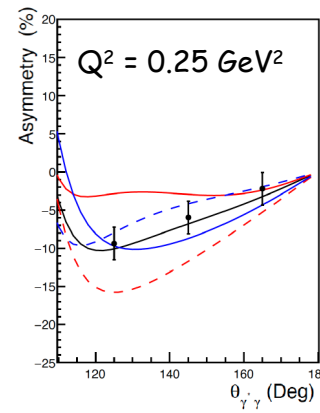
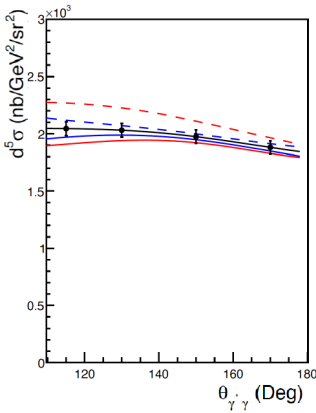
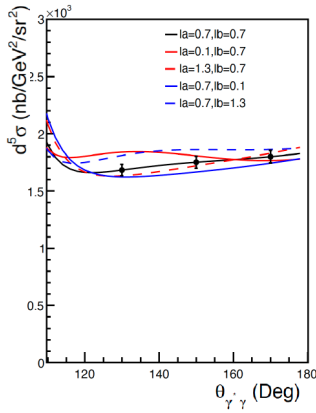
Total: 62 days

| Kinematic Group | Kinematic Setting | $\theta_{\gamma^* \gamma}$ | θ_e° | $P'_e(\text{MeV}/c)$ | θ_p° | $P'_p(\text{MeV}/c)$ | $I (\mu\text{A})$ | beam time (days) | |
|-----------------|-------------------|----------------------------|------------------|----------------------|------------------|----------------------|-------------------|------------------|------|
| GI | Kin I | 110 | 14.3 | 736.3 | 54.45 | 493.93 | 15 | 1.00 | |
| | Kin II | 133 | 14.3 | 736.3 | 44.93 | 556.10 | 15 | 1.00 | |
| | Kin IIIa | 147 | 14.3 | 736.3 | 11.26 | 583.05 | 15 | 1.00 | |
| | Kin IIIb | 147 | 14.3 | 736.3 | 39.06 | 583.05 | 15 | 1.00 | |
| | Kin IVa | 160 | 14.3 | 736.3 | 16.73 | 599.95 | 15 | 1.00 | |
| | Kin IVb | 160 | 14.3 | 736.3 | 33.59 | 599.95 | 15 | 1.00 | |
| GII | Kin I | 115 | 11.22 | 1783.0 | 15.33 | 615.69 | 10 | 1.50 | |
| | Kin IIa | 125 | 11.22 | 1783.0 | 56.54 | 647.85 | 10 | 2.50 | |
| | Kin IIb | 125 | 11.22 | 1783.0 | 18.60 | 647.85 | 10 | 1.50 | |
| | Kin IIIa | 145 | 11.22 | 1783.0 | 49.77 | 697.99 | 10 | 1.50 | |
| | Kin IIIb | 145 | 11.22 | 1783.0 | 25.37 | 697.99 | 10 | 1.00 | |
| | Kin IVa | 165 | 11.22 | 1783.0 | 42.82 | 726.87 | 10 | 1.00 | |
| | Kin IVb | 165 | 11.22 | 1783.0 | 32.32 | 726.87 | 10 | 1.00 | |
| | GIII | Kin I | 115 | 14.73 | 1729.7 | 20.58 | 706.89 | 30 | 1.75 |
| Kin IIa | | 130 | 14.73 | 1729.7 | 54.89 | 758.24 | 30 | 2.00 | |
| Kin IIb | | 130 | 14.73 | 1729.7 | 24.78 | 758.24 | 30 | 1.75 | |
| Kin IIIa | | 150 | 14.73 | 1729.7 | 48.99 | 808.24 | 30 | 1.75 | |
| Kin IIIb | | 150 | 14.73 | 1729.7 | 30.68 | 808.24 | 30 | 1.75 | |
| Kin IVa | | 170 | 14.73 | 1729.7 | 42.90 | 834.12 | 30 | 1.00 | |
| Kin IVb | | 170 | 14.73 | 1729.7 | 36.76 | 834.12 | 30 | 1.00 | |
| GIV | | Kin I | 100 | 16.32 | 1749.3 | 23.83 | 664.52 | 35 | 1.75 |
| | Kin II | 120 | 16.32 | 1749.3 | 28.01 | 738.39 | 50 | 1.25 | |
| | Kin IIIa | 140 | 16.32 | 1749.3 | 32.84 | 795.37 | 70 | 1.00 | |
| | Kin IIIb | 140 | 16.32 | 1749.3 | 53.80 | 795.37 | 70 | 2.00 | |
| | Kin IVa | 155 | 16.32 | 1749.3 | 36.69 | 824.46 | 70 | 1.50 | |
| | Kin IVb | 155 | 16.32 | 1749.3 | 49.95 | 824.46 | 70 | 2.50 | |
| | Kin Va | 170 | 16.32 | 1749.3 | 40.66 | 840.48 | 70 | 1.00 | |
| | Kin Vb | 170 | 16.32 | 1749.3 | 45.99 | 840.48 | 70 | 1.00 | |
| | GV | Kin I | 100 | 17.72 | 1676.41 | 19.75 | 723.69 | 35 | 2.00 |
| | | Kin II | 120 | 17.72 | 1676.41 | 24.25 | 808.93 | 50 | 1.50 |
| Kin IIIa | | 140 | 17.72 | 1676.41 | 29.34 | 874.74 | 70 | 1.50 | |
| Kin IIIb | | 140 | 17.72 | 1676.41 | 51.12 | 874.74 | 70 | 2.00 | |
| Kin IVa | | 155 | 17.72 | 1676.41 | 33.36 | 908.37 | 70 | 2.00 | |
| Kin IVb | | 155 | 17.72 | 1676.41 | 47.10 | 908.37 | 70 | 2.00 | |
| Kin Va | | 170 | 17.72 | 1676.41 | 37.47 | 926.91 | 70 | 1.00 | |
| Kin Vb | | 170 | 17.72 | 1676.41 | 42.99 | 926.91 | 70 | 1.00 | |
| GVI | Kin I | 120 | 20.45 | 1623.1 | 25.31 | 886.59 | 75 | 1.00 | |
| | Kin IIa | 140 | 20.45 | 1623.1 | 29.91 | 956.82 | 75 | 1.00 | |
| | Kin IIb | 140 | 20.45 | 1623.1 | 49.81 | 956.82 | 75 | 1.50 | |
| | Kin IIIa | 155 | 20.45 | 1623.1 | 33.58 | 992.83 | 75 | 1.50 | |
| | Kin IIIb | 155 | 20.45 | 1623.1 | 46.14 | 992.83 | 75 | 2.00 | |

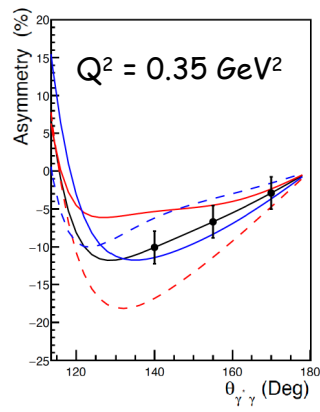
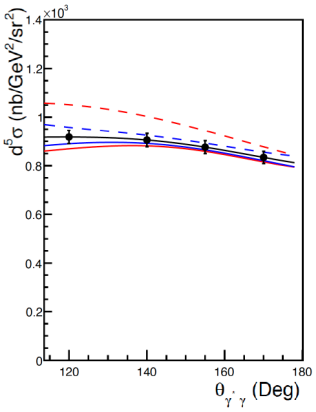
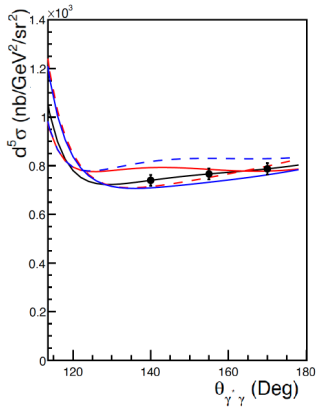
Projected Cross sections



VCS-II Projected

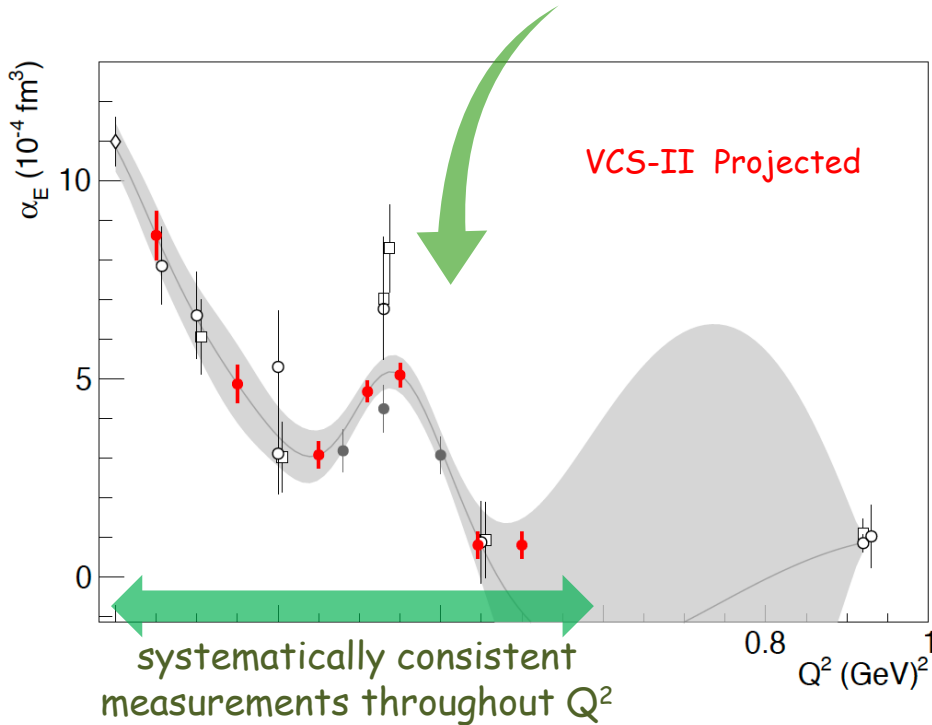


$\delta(\text{stat}) = 1\% - 2\%$
 $\delta(\text{syst}) = 3.5\% - 4\%$

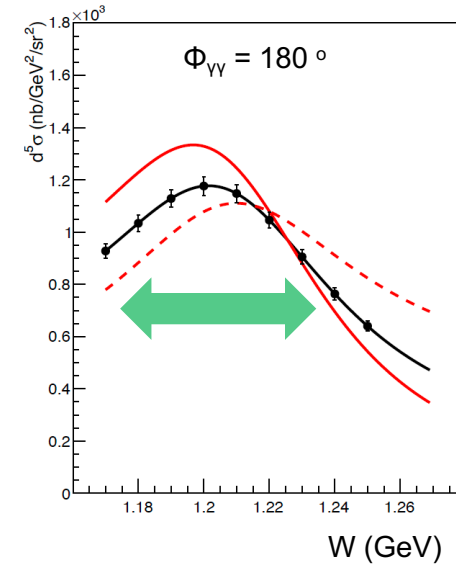
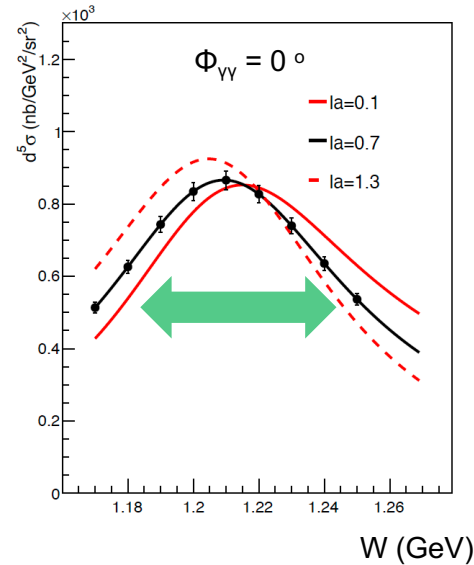


VCS-II Projected Measurements

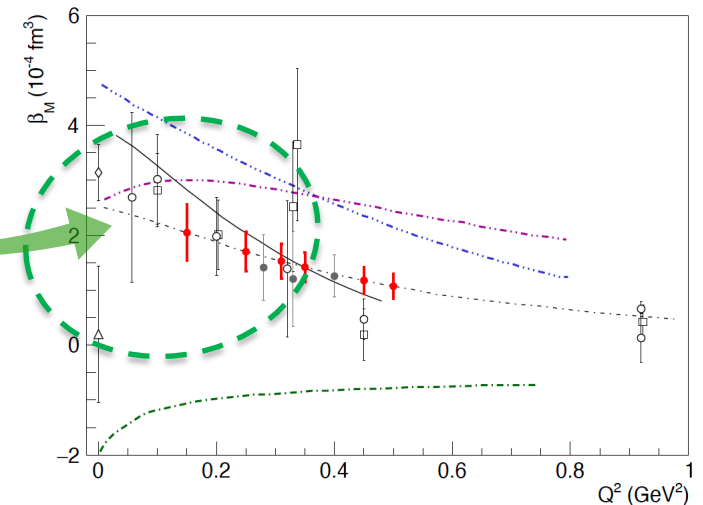
High precision measurements combined with a fine mapping in Q^2



Targeted measurements to fully exploit the sensitivity to the GPs



Improve upon β_M :
Pin down the competing para/dia-magnetic contributions in the nucleon



Other prospects: Measuring with positrons or with polarized e⁻ beam

Positrons allow for an independent path to access experimentally the GPs

Eur. Phys. J. A 57 (2021) 11, 316

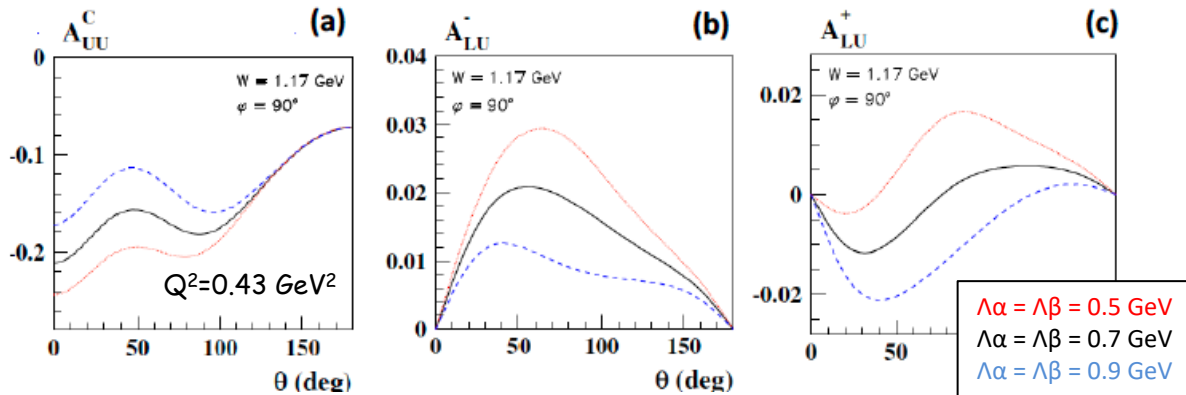
Virtual Compton scattering at low energies with a positron beam

Barbara Pasquini^{a,1,2}, Marc Vanderhaeghen^{b,3}

¹Dipartimento di Fisica, Università degli Studi di Pavia, 27100 Pavia, Italy

²Istituto Nazionale di Fisica Nucleare, Sezione di Pavia, 27100 Pavia, Italy

³Institut für Kernphysik and PRISMA⁺ Cluster of Excellence, Johannes Gutenberg Universität, D-55099 Mainz, Germany



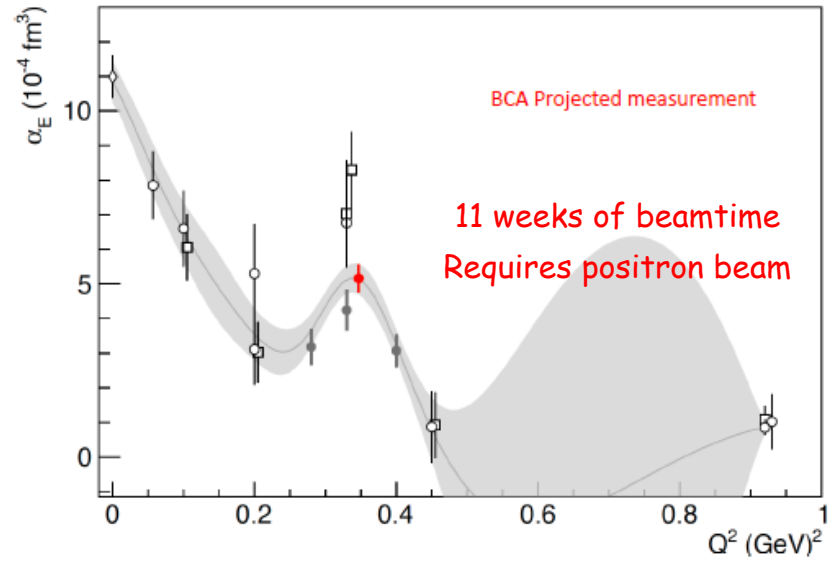
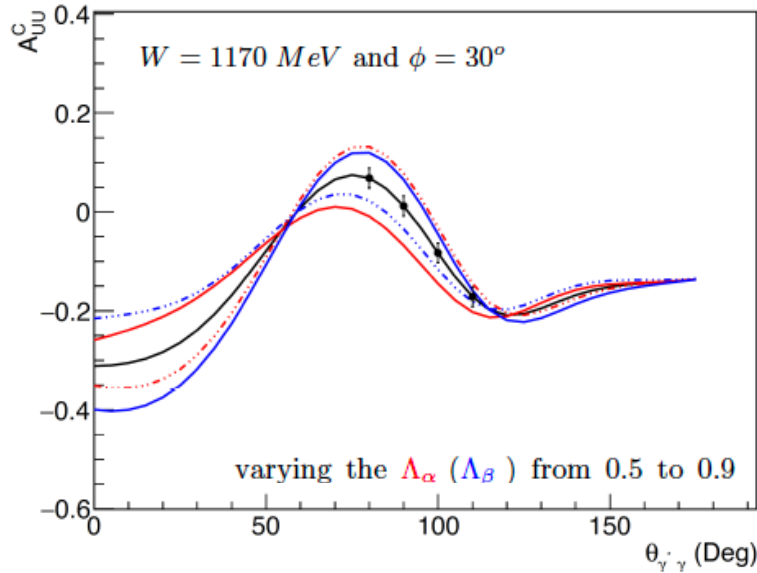
(a): The beam-charge asymmetry as a function of the photon scattering angle at $Q^2 = 0.43 \text{ GeV}^2$.

(b) & (c): The electron and positron beam-spin asymmetry as a function of the photon scattering angle for out-of-plane kinematics.

Unpolarized beam charge asymmetry (BCA):
$$A_{UU}^C = \frac{(d\sigma_+^+ + d\sigma_-^+) - (d\sigma_+^- + d\sigma_-^-)}{d\sigma_+^+ + d\sigma_-^+ + d\sigma_+^- + d\sigma_-^-}$$

Lepton beam spin asymmetry (BSA):
$$A_{LU}^e = \frac{d\sigma_+^e - d\sigma_-^e}{d\sigma_+^e + d\sigma_-^e}$$

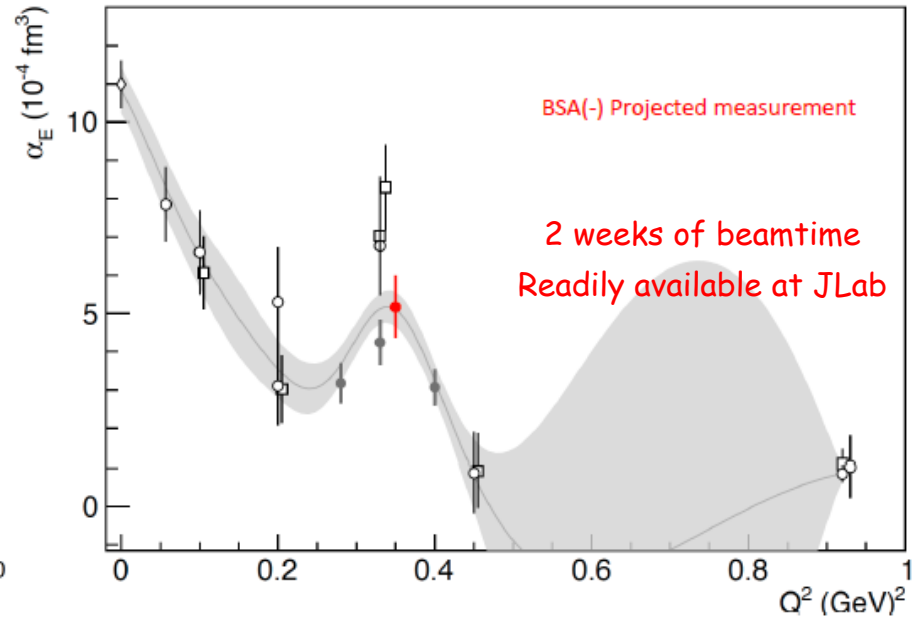
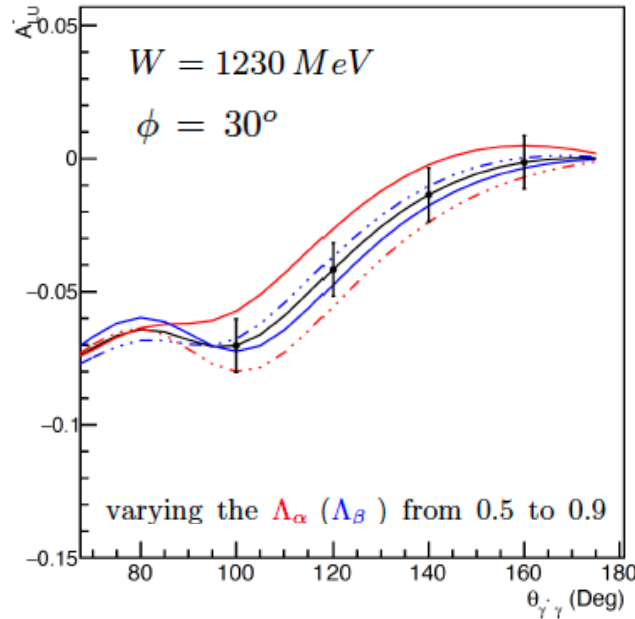
BCA



Hall C (SHMS / HMS)

- e^- : ~ 1 week of beamtime @ $50 \mu\text{A}$
- and
- e^+ : ~ 10 weeks of beamtime @ $5 \mu\text{A}$

BSA



Hall C (SHMS / HMS)

- e^- (pol. 85% @ 70 μA) : ~ 2 weeks of beamtime
- or
- e^+ (pol. 60% @ 50 nA) : ~ 3 orders of magnitude more beamtime

**Measurement of the Generalized Polarizabilities of the Proton
 with positron and polarized electron beams**

Letter of Intent to Jefferson Lab PAC-51

Summary

Progress in the study of fundamental system properties that characterize the proton's response to an EM field

Insight to spatial deformation of the nucleon densities under an applied EM field, interplay of para/dia-magnetic mechanisms in the proton, polarizability radii, ...

Electric GP: $\left\{ \begin{array}{l} \text{possibility for a non-trivial (non-monotonic) behavior in } a_E(Q^2) \\ \text{(albeit with a smaller magnitude than originally suggested)} \\ \text{or} \\ \text{at minimum: strong tension between world data} \end{array} \right.$

Strong constraints to theoretical predictions (and we can improve further)

High precision benchmark data for theory & upcoming LQCD calculations

Future measurements:

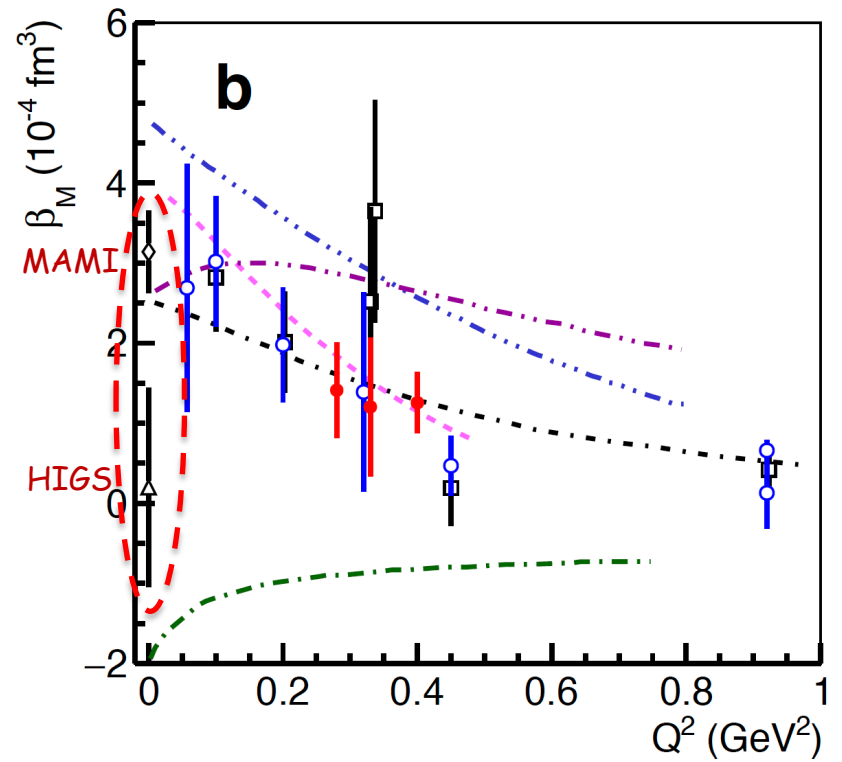
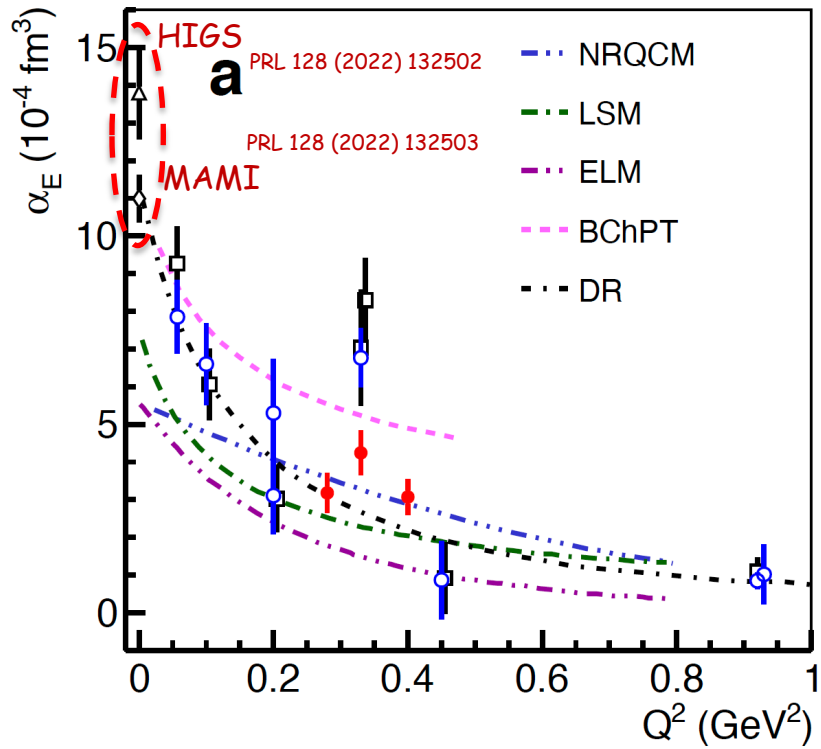
Pinn down with precision the shape of the a_E structure (if it exists) - important input for the theory

Measure via a different channel (positrons)

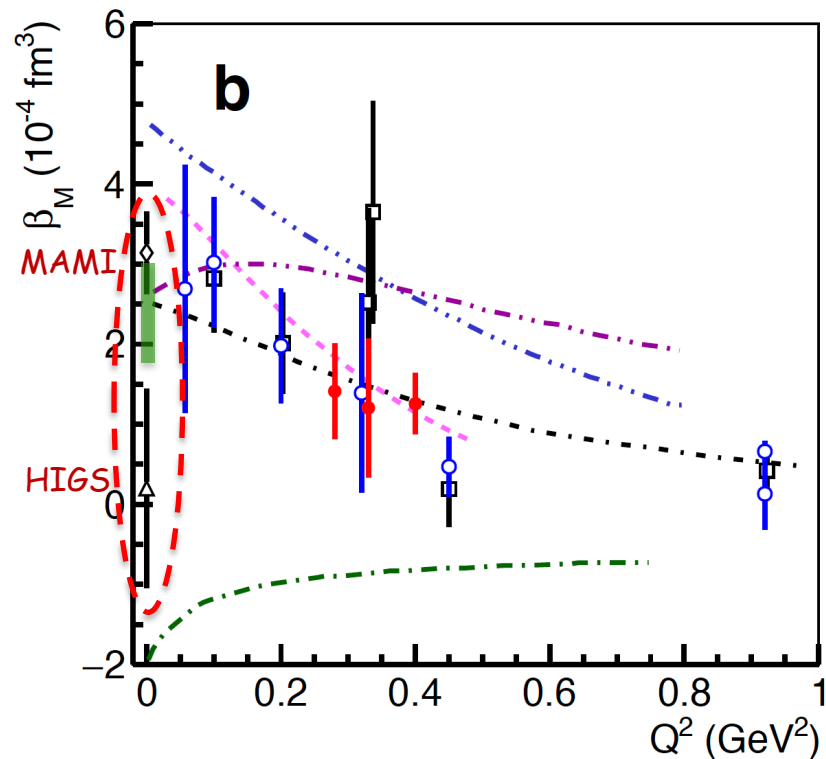
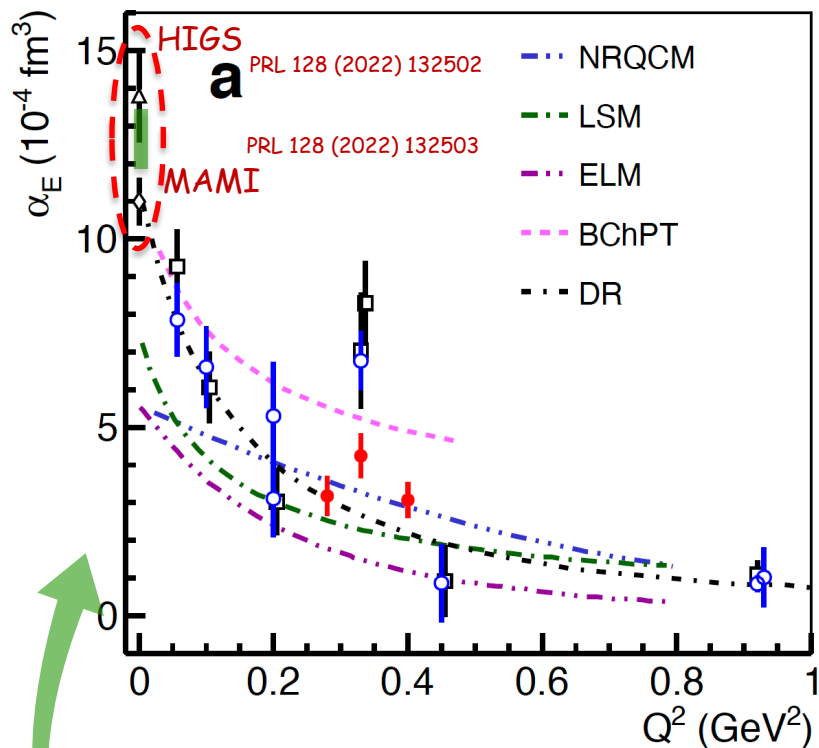
Thank you!

BACK UP

Static Polarizabilities



Static Polarizabilities



PHYSICAL REVIEW LETTERS 129, 102501 (2022)

First Concurrent Extraction of the Leading-Order Scalar and Spin Proton Polarizabilities

E. Mornacchi^{1,*}, S. Rodini², B. Pasquini^{3,4} and P. Pedroni⁴

¹Institut für Kernphysik, Johannes Gutenberg-Universität Mainz, D-55099 Mainz, Germany

²Institut für Theoretische Physik, Universität Regensburg, D-93040 Regensburg, Germany

³Dipartimento di Fisica, Università degli Studi di Pavia, I-27100 Pavia, Italy

⁴INFN Sezione di Pavia, I-27100 Pavia, Italy

(Received 3 May 2022; revised 11 July 2022; accepted 2 August 2022; published 31 August 2022)

We performed the first simultaneous extraction of the six leading-order proton polarizabilities. We reached this milestone thanks to both new high-quality experimental data and an innovative bootstrap-based fitting method. These new results provide a self-consistent and fundamental benchmark for all future theoretical and experimental polarizability estimates.

$$\alpha_{E1} = [12.7 \pm 0.8(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^3,$$

$$\beta_{M1} = [2.4 \pm 0.6(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^3,$$

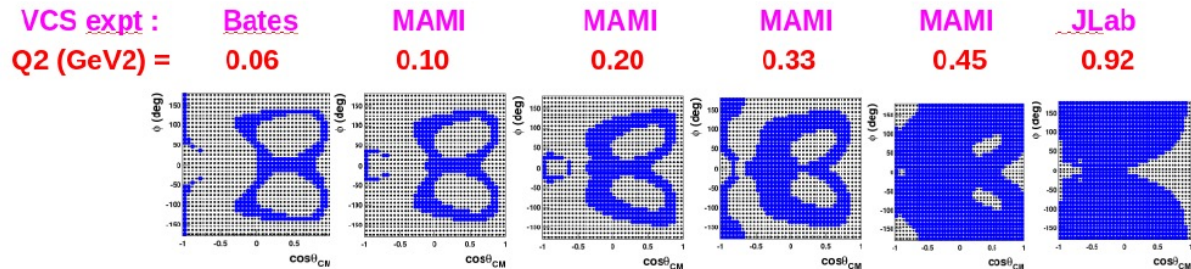
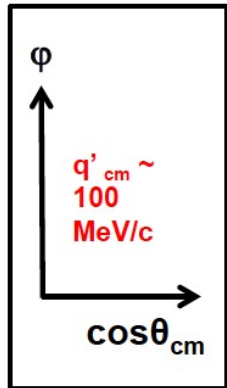
$$\gamma_{E1E1} = [-3.0 \pm 0.6(\text{fit}) \pm 0.4(\text{model})] \times 10^{-4} \text{ fm}^4,$$

$$\gamma_{M1M1} = [3.7 \pm 0.5(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^4,$$

$$\gamma_{E1M2} = [-1.2 \pm 1.0(\text{fit}) \pm 0.3(\text{model})] \times 10^{-4} \text{ fm}^4,$$

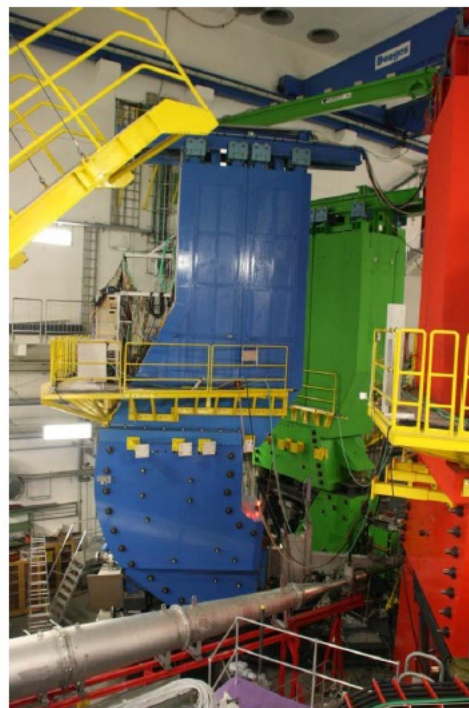
$$\gamma_{M1E2} = [2.0 \pm 0.7(\text{fit}) \pm 0.4(\text{model})] \times 10^{-4} \text{ fm}^4,$$

Blue bins = where the higher-order estimator is $< 3\%$
(LEX truncation « valid »)

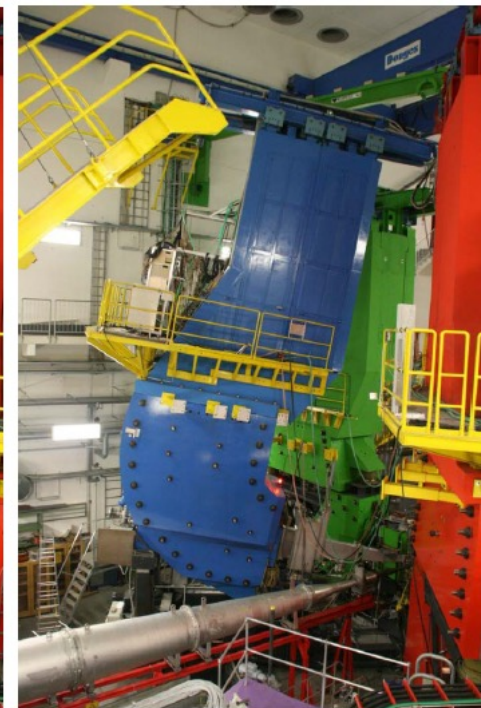


New « vcsq2 » data:

- OOP kinematics (to access the blue region)
- LEX Fit done with bin selection at $Q^2 = 0.1$ and 0.2 GeV^2 .
- was found not necessary at $Q^2 = 0.45 \text{ GeV}^2$.



In-plane

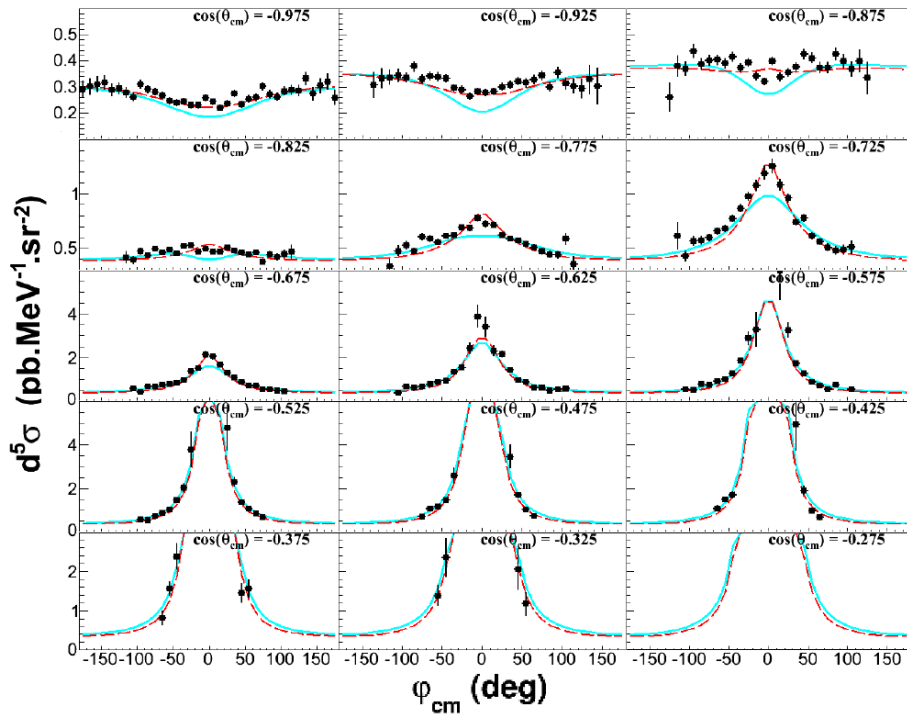


8.5 deg OOP

A1/1-09 @ MAMI

~ 1.0 GeV beam

$Q^2 = 0.1 (GeV/c)^2, 0.2 (GeV/c)^2, \text{ and } 0.45 (GeV/c)^2$



BH+B ---
Polarizability effect ---

GP effect typically 5% - 15% of the cross section

Polarizability fits:

DR fit:
DR calculation includes full dependency in q'_{cm}

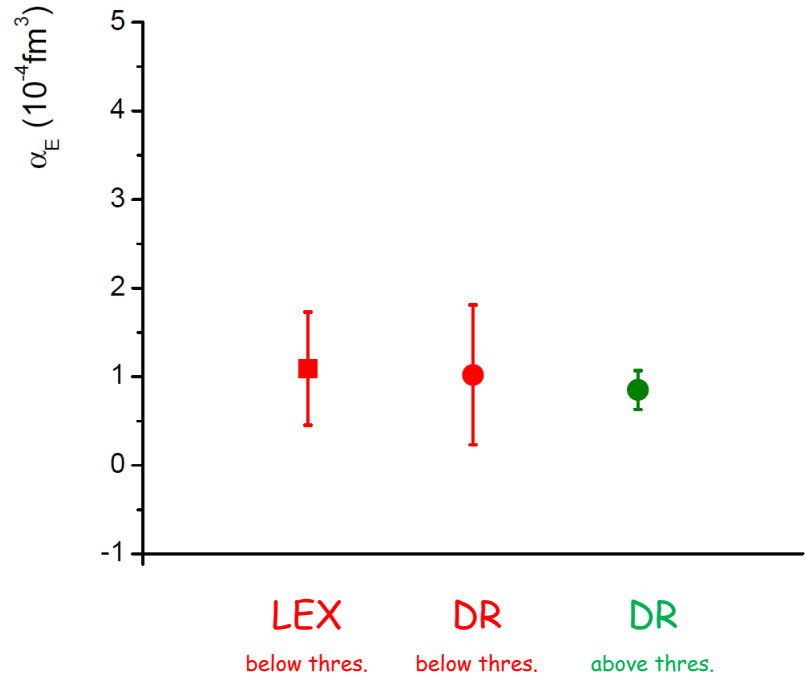
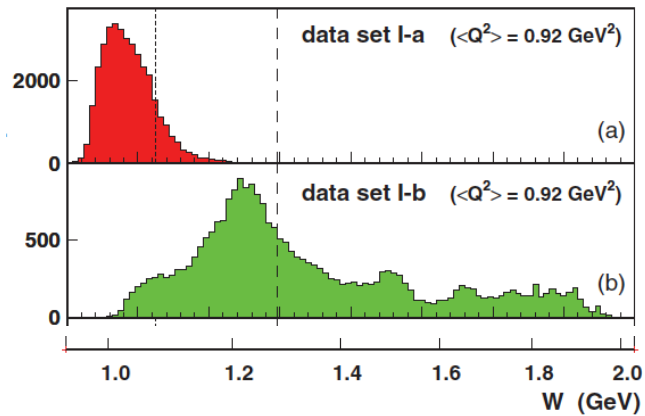
LEX fit:
truncated in q'_{cm} . Suppress contribution from higher order terms

Figure 5.8: Setting INP: measured $ep \rightarrow ep\gamma$ cross section at fixed $q'_{cm} = 112.5 MeV/c$ with respect to φ_{cm} for all the $\cos(\theta_{cm})$ -bins. The curves follow the convention of figure 5.6.

Virtual Compton Scattering

Phys. Rev C 86, 015210 (2012)

Phys. Rev Lett. 93, 122001 (2004)



Sensitivity to the GPs grows as we measure above pion threshold