







Cabibbo unitarity status



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DFG

Neutron beta decay review: MG, Seng, Universe 2023, 9(9), 422, arXiv:2307.01145 Nuclear beta decay review: MG, Seng (for Annual Reviews Part. Nucl. Sci. - deadline Nov 2)

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Outline

Status of Cabibbo Angle Anomaly

Status of V_{ud} and V_{us}

RC to β -decays: overall setup, scale separation

Dispersion theory of γW -box

Possible BSM scenarios

Summary & Outlook

Quark Mixing & CKM Unitarity

Cabibbo: mass and flavor eigenstates connected by Cabibbo angle θ_C Strength is distributed among two channels

$$|G_V^{\Delta S=0}| = \cos \theta_C G_\mu$$
$$|G_V^{\Delta S=1}| = \sin \theta_C G_\mu$$



Kobayashi & Maskawa: 3 flavors, CP-violation

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} = V_{CKM} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



CKM unitarity - measure of completeness of the SM: $VV^{\dagger} = \mathbf{1}$ Top-row unitarity constraint $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$

CKM unitarity is among our best precision tools to test the SM!

Status of Cabibbo unitarity

$$|V_{ud}|^{2} + |V_{us}|^{2} + |V_{ub}|^{2} = 0.9985(6)_{V_{ud}}(4)_{V_{us}}$$

$$\sim 0.95 \sim 0.05 \sim 10^{-5}$$
Inconsistencies: between measurements of V_{ud} and V_{us} and SM predictions
Main reason for Cabibbo-angle anomaly: shift in V_{ud} (and small uncertainties?)
$$0.228$$

$$0.226$$

$$K\ell^{3}$$

$$0.222$$

$$K\ell^{3}$$

$$0.222$$

$$V_{ud}$$

Status of V_{us}

$V_{us}\,/\,V_{ud}$ from KI2 decays

$$\frac{|V_{us}|}{|V_{ud}|}\frac{f_K}{f_{\pi}} = \left(\frac{\Gamma_{K_{\mu 2(\gamma)}} m_{\pi^{\pm}}}{\Gamma_{\pi_{\mu 2(\gamma)}} m_{K^{\pm}}}\right)^{1/2} \frac{1 - m_{\mu}^2 / m_{\pi^{\pm}}^2}{1 - m_{\mu}^2 / m_{K^{\pm}}^2} \left(1 - \frac{1}{2} \delta_{\text{EM}} - \frac{1}{2} \delta_{SU(2)}\right)$$

Inputs from experiment:

Inputs from theory:

From K^{\pm} BR fit: **BR** $(K^{\pm}_{\mu 2(\gamma)}) = 0.6358(11)$ $\tau_{K\pm} = 12.384(15)$ ns

From PDG:

BR($\pi^{\pm}_{\mu^{2}(\gamma)}$) = 0.9999 $\tau_{\pi^{\pm}}$ = 26.033(5) ns $\delta_{\rm EM}$ Long-distance EM corrections

 $\frac{\delta_{SU(2)}}{f_K / f_\pi \to f_{K\pm} / f_{\pi\pm}}$ Strong isospin breaking

 $f_K f_{\pi}$ Ratio of decay constants Cancellation of lattice-scale uncertainties from ratio NB: Most lattice results already corrected for SU(2)-breaking: $f_{K\pm} f_{\pi\pm}$

 $|V_{us}/V_{ud}| = 0.23108(23)_{exp}(42)_{lat}(16)_{IB}$ (51)_{tot} = 0.22 %

 $\delta_{SU(2)} + \delta_{EM}$

 f_K/f_{π}

ChPT: $\delta_{SU(2)} + \delta_{EM} = -0.0112(21)$ Cirigliano et al, 2011 LQCD+EM ($N_f = 2 + 1 + 1$): $\delta_{SU(2)} + \delta_{EM} = -0.0122(16)$ Giusti et al, 2018 LQCD+EM ($N_f = 2 + 1 + 1$): $\delta_{SU(2)} + \delta_{EM} = -0.0126(14)$ Di Carlo et al, 2019 LQCD ($N_f = 2 + 1$): $f_K/f_{\pi} = 1.1946(34)$ LQCD ($N_f = 2 + 1 + 1$): $f_K/f_{\pi} = 1.1978(22)$ FLAG 2021 averages

Vus from Ke3

$$\Gamma(K_{\ell 3(\gamma)}) = \frac{C_K^2 G_F^2 m_K^5}{192\pi^3} S_{\rm EW} |V_{us}|^2 |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}(\lambda_{K\ell}) \left(1 + 2\Delta_K^{SU(2)} + 2\Delta_{K\ell}^{\rm EM}\right)$$

with $K \in \{K^+, K^0\}$; $\ell \in \{e, \mu\}$, and: C_K^2 1/2 for K^+ , 1 for K^0 $S_{\rm EW}$ Universal SD EW correction (1.0232)

Inputs from experiment:

 $\Gamma(K_{\ell 3(\gamma)})$

Rates with well-determined treatment of radiative decays:

- Branching ratios
- Kaon lifetimes

 $I_{K\ell}(\{\lambda\}_{K\ell})$

Integral of form factor over phase space: λ s parameterize evolution in *t*

 $\Delta_{K\ell}^{\rm EM}$

 $\Delta_{K}^{SU(2)}$

 $f_{+}^{K^{0}\pi^{-}}(0)$

 $K_{\ell 3}$: $|V_{us}| = 0.22330(35)_{exp}(39)_{lat}(8)_{IB}$ (53)_{tot} = 0.24 %

Inputs from theory:

Hadronic matrix element (form factor) at zero momentum transfer (t = 0)

Form-factor correction for *SU*(2) breaking

Form-factor correction for long-distance EM effects

 $f_{+}(0)$

LQCD $(N_f = 2 + 1)$: $f_+(0) = 0.9677(27)$ LQCD $(N_f = 2 + 1 + 1)$: $f_+(0) = 0.9698(17)$ ChPT: $f_+(0) = 0.970(8)$ Ecker et al 2015; Bijnens, Ecker 2014

RC to Kl3

Until 2021: best way to compute long-distance EM RC was with ChPT

	$I_{K\ell}^{(0)}(\lambda_i)$	$\delta^{K\ell}_{ m EM}(\mathcal{D}_3)(\%)$	$\delta_{\mathrm{EM}}^{K\ell}(\mathcal{D}_{4-3})(\%)$	$\delta^{K\ell}_{ m EM}(\%)$
K_{e3}^{0}	0.103070	0.50	0.49	0.99 ± 0.30
K_{e3}^{\pm}	0.105972	-0.35	0.45	0.10 ± 0.30
$K_{\mu 3}^{0}$	0.068467	1.38	0.02	1.40 ± 0.30
$K_{\mu 3}^{\mu 3}$	0.070324	0.007	0.009	0.016 ± 0.30

Cirigliano, Gianotti, Neufeld 0807.4507

A series of works reformulated the problem as a hybrid of current algebra and ChPT, plus input from lattice QCD calculations of γW -box for $\pi e3$ and $K\ell3$

	$\delta_{\rm EM}^{K\ell}$ [10 ⁻³]	ChPT
$K^0 e$	$11.6(2)_{\text{inel}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{e^2p^4}$	$9.9(1.9)_{e^2p^4}(1.1)_{\text{LEC}}$
K^+e	$2.1(2)_{\text{inel}}(1)_{\text{lat}}(4)_{\text{NF}}(1)_{e^2p^4}$	$1.0(1.9)_{e^2p^4}(1.6)_{\rm LEC}$
$K^0\mu$	$15.4(2)_{\text{inel}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2p^4}$	$14.0(1.9)_{e^2p^4}(1.1)_{\rm LEC}$
$K^+\mu$	$0.5(2)_{\text{inel}}(1)_{\text{lat}}(4)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2p^4}$	$0.2(1.9)_{e^2p^4}(1.6)_{\text{LEC}}$

Seng, Galviz, Meißner 1910.13208 Seng, Galviz, MG, Meißner 2103.04843 Seng, Galviz, MG, Meißner 2203.05217

Feng, MG, Jin, Ma, Seng 2003.09798 Ma, Feng, MG, Jin, Seng 2102.12048

Uncertainties reduced by an o.o.m. Long-distance EM RC not responsible for the $K\ell^2 - K\ell^3$ discrepancy!

Status of V_{ud}

Vud from neutron decay

Neutron decay: 2 measurements needed

$$|V_{ud}|^2 = \frac{5024.7 \text{ s}}{\tau_n (1 + 3g_A^2)(1 + \Delta_R^V)}$$

RC Δ_R^V : bottleneck since 40 years

Pre-2018: $\Delta_R^V = 0.02361(38)$ Marciano, Sirlin PRL 2006 Post-2018: $\Delta_R^V = 0.02479(21)$ MG, Seng Universe 2023

Since 2018: DR+data+pQCD+EFT+LQCD Δ_R^V uncertainty: factor 2 reduction

C-Y Seng et al., PRL 2018; PRD 2019 A. Czarnecki, B. Marciano, A. Sirlin, PRD 2018 K. Shiells et al, PRD 2021; L. Hayen PRD 2021 P-X Ma, X. Feng, MG, L-C Jin, et al 2308.16755

Experiment: factor 3-5 uncertainties improvement; discrepancies in τ_n and g_A

3.4 σ $g_A = -1.27641(56)$ $g_A = -1.2677(28)$ 4 σ $\tau_n = 877.75(28)^{+16}_{-12}$ $\tau_n = 887.7(2.3)$

PERKEO-III B. Märkisch et al, Phys.Rev.Lett. 122 (2019) 24, 242501 **aSPECT** M. Beck et al, Phys. Rev. C101 (2020) 5, 055506; 2308.16170

UCNτ F. M. Gonzalez et al. Phys. Rev. Lett. 127 (2021) 162501 BL1 (NIST) Yue et al, PRL 111 (2013) 222501

PDG average $|V_{ud}^{\text{free n}}| = 0.9743 (3)_{\tau_n} (8)_{g_A} (1)_{RC} [9]_{total}$ Single best measurements only $|V_{ud}^{\text{free n}}| = 0.9740 (2)_{\tau_n} (3)_{g_A} (1)_{RC} [4]_{total}$

V_{ud} from superallowed decays

Advantages:

- 1. Only conserved vector current
- 2. 15 measured to better than 0.2%
- 3. 5 measured better than τ_n
- 4. Internal consistency as a check
- 5. SU(2) good —> corrections ~small
- 6. We know a lot about nuclei
- 7. Only scalar (or vector) BSM accessible

Exp.: **f** - phase space (Q value)

 \boldsymbol{t} - partial half-life (t_{1/2}, branching ratio)







ft values: same within ~2% but not exactly! Reason: SU(2) slightly broken

- a. RC (e.m. interaction does not conserve isospin)
- b. Nuclear WF are not SU(2) symmetric(proton and neutron distribution not the same)

Vud extraction: Universal RC and Universal Ft

To obtain Vud —> absorb all decay-specific corrections into universal Ft



$$|V_{ud}^{0^+-0^+}| = 0.9737 \,(1)_{exp,\,nucl} (3)_{NS} \,(1)_{RC} [3]_{total}$$

Vud from semileptonic pion decay

Pion decay $\pi^+ \rightarrow \pi^0 e^+ \nu_e$: theoretically cleanest, experimentally tough

$$|V_{ud}|^2 = \frac{0.9799}{(1+\delta)} \frac{\Gamma_{\pi\ell3}}{0.3988(23)\,\mathrm{s}^{-1}} \qquad |V_{ud}^{\pi\ell3}| = 0.9739\,(27)_{exp}\,(1)_{RC}$$

RC to semileptonic pion decay δ uncertainty: factor 3 reduction

ChPT: $\delta = -0.0334(10)_{\text{LEC}}(3)_{\text{HO}}$ Cirigliano et al, 2003; Passera et al, 2011 DR + LQCD + ChPT: $\delta = 0.0332(1)_{\gamma W}(3)_{\text{HO}}$ Feng et al, 2020; Yoo et al, 2023

Future exp: 1 o.o.m. (PIONEER @ PSI)

RC to beta decay: overall setup

RC to beta decay: overall setup $\nu_e(\bar{\nu}_e)$ Tree-level amplitude $i = n, A(0^+)$ Radiative corrections to tree-level amplitude $\sim \alpha/2\pi \approx 10^{-3}$ 1×10^{-4} Precision goal for V_{ud} extraction Weak boson scale Electron carries away energy E < Q-value of a decay $M_7, M_W \sim 90 \,\mathrm{GeV}$ E-dep RC: $\frac{\alpha}{2\pi} \left(\frac{E}{\Lambda}, \ln \frac{E}{\Lambda}, \dots \right)$ Hadronic scale Universal $\Lambda_{\rm had} = 300 \,{\rm MeV}$ Energy scales Λ Nuclear scale Nuclear structure dependent $\Lambda_{\rm nuc} = 10 - 30 \,{\rm MeV}$ (QCD) Decay Q-value (endpoint energy) $Q_{if} = M_i - M_f = 1 - 10 \text{ MeV}$ Nucleus-specific **Electron mass** Nuclear structure independent $m_{\rho} \approx 0.5 \,\mathrm{MeV}$ (QED)

RC to beta decay: overall setup

Generically: only IR and UV extremes feature large logarithms! Works by Sirlin (1930-2022) and collaborators: all large logs under control

IR: Fermi function + Sirlin function

Fermi function: resummation of $(Z\alpha)^n \longrightarrow$ Dirac - Coulomb problem

UV: large EW logs + pQCD corrections

Inner RC: energy- and model-independent

W,Z - loops UV structure of SM



γW -box: sensitive to all scales

New method for computing EW boxes: dispersion theory Combine exp. data with pQCD, lattice, EFT, ab-initio nuclear



 $(\operatorname{Re} c)_{\mathrm{m.d}} = 8\pi^2 \operatorname{Re} \int \frac{d^2 q}{(2\pi)^4}$

RC to β decay - scale separation

Fermi function (pure Coulomb + nuclear size & recoil + atomic) —> phase-space **f** Fermi, Behrens-Bühring, Wilkinson...

Soft Bremsstrahlung: universal Sirlin's function + nucleus specific corrections —> δ'_R All IR-sensitive pieces: recent review Hayen et al RMP 2018

UV-sensitive RC on free neutron Δ_R^V : Sirlin, Marciano, Czarnecki 1967 - 2006

$$g_V^2 = |V_{ud}|^2 \left[1 + \frac{\alpha}{2\pi} \left\{ 3\ln\frac{M_Z}{M_p} + \ln\frac{M_Z}{M_W} + \tilde{a}_g \right\} + \delta_{\text{QED}}^{HO} + 2\Box_{\gamma W} \right]$$

RC on nuclei: extract the universal (free n) Δ_R^V : Nuclear structure correction: $\delta_{\rm NS} = 2(\Box_{\gamma W}^{\rm Nucl} - \Box_{\gamma W}^{\rm free n})$ All non-enhanced terms $\sim \alpha/2\pi \sim 10^{-3}$ — only need to ~10% Approximations are legitimate (isospin limit, ...)

γW -box from DR

γW -box from dispersion relations

Model-dependent part or RC: γW -box





$$\int dx e^{iqx} \langle H_f(p) | T\{J_{em}^{\mu}(x)J_W^{\nu,\pm}(0)\} | H_i(p) \rangle$$

Generalized (non-diagonal) Compton amplitudes



Commutator (Im part) - only on-shell hadronic states — related to data

$$\int dx e^{iqx} \langle H_f(p) | [J^{\mu}_{em}(x), J^{\nu,\pm}_W(0)] | H_i(p) \rangle$$

Interference structure functions

Physics of taming model dependence with dispersion relations:

virtual photon polarizes the nucleon/nucleus;

Long- and intermediate-range part of the box sensitive to hadronic **polarizabilities** Polarizabilities related to the excitation spectrum via dispersion relation (Cf. Kramers-Kronig)

Universal RC from dispersion relations

Interference γW structure functions

$$\mathrm{Im}T^{\mu\nu}_{\gamma W} = \dots + \frac{i\varepsilon^{\mu\nu\alpha\beta}p_{\alpha}q_{\beta}}{2(pq)}F^{\gamma W}_{3}(x,Q^{2})$$

After some algebra (isospin decomposition, loop integration)

$$\Box_{\gamma W}^{b,e}(E_e) = \frac{\alpha}{\pi} \int_0^\infty dQ^2 \frac{M_W^2}{M_W^2 + Q^2} \int_{\nu_{thr}}^\infty \frac{d\nu'}{\nu'} \frac{\nu' + 2\sqrt{\nu'^2 + Q^2}}{(\nu' + \sqrt{\nu'^2 + Q^2})^2} \frac{F_{3,-}(\nu',Q^2)}{Mf_+(0)} + \mathcal{O}(E_e^2)$$

$$\Box_{\gamma W}^{b,o}(E_e) = \frac{2\alpha E_e}{3\pi} \int_0^\infty dQ^2 \int_{\nu_{thr}}^\infty \frac{d\nu'}{\nu'} \frac{\nu' + 3\sqrt{\nu'^2 + Q^2}}{(\nu' + \sqrt{\nu'^2 + Q^2})^3} \frac{F_{3,+}(\nu',Q^2)}{Mf_+(0)} + \mathcal{O}(E_e^3)$$

Advantage to previous approach (Marciano & Sirlin):

- Explicit 2-fold integral, isospin decomposition and energy dependence

Nachtmann moments
play a role in DIS
$$M_3(n, Q^2) = \frac{n+1}{n+2} \int_0^1 \frac{dx\xi^n}{x^2} \frac{2x(n+1) - n\xi}{n+1} F_3(x, Q^2), \qquad \xi = \frac{2x}{1 + \sqrt{1 + 4M^2 x^2/Q^2}}$$

Hiding the nu-integration in the Nachtmann moments:

$$\Box_{\gamma W}^{b}(E_{e}) = \frac{3\alpha}{2\pi} \int_{0}^{\infty} \frac{dQ^{2}}{Q^{2}} \frac{M_{W}^{2}}{M_{W}^{2} + Q^{2}} \left[M_{3,-}(1,Q^{2}) + \frac{8E_{e}M}{9Q^{2}} M_{3,+}(2,Q^{2}) \right] + \mathcal{O}(E_{e}^{2})$$

Input into dispersion integral

W

q

p

Rev'



Input into dispersion integral - $\nu/\bar{\nu}$ data

Mixed CC-NC γW SF (no data) <—> Purely CC WW SF (inclusive neutrino data) Isospin symmetry: vector-isoscalar current related to vector-isovector current Only useful if we know the physical mechanism (Born, DIS, Regge, Resonance,...) Were able to identify the missing part with Regge (multiparticle continuum)



Marciano, Sirlin 2006: $\Delta_R^V = 0.02361(38) \longrightarrow |V_{ud}| = 0.97420(10)_{Ft}(18)_{RC}$ DR (Seng et al. 2018): $\Delta_R^V = 0.02467(22) \longrightarrow |V_{ud}| = 0.97370(10)_{Ft}(10)_{RC}$

γW -box from DR + Lattice QCD input

Currently available neutrino data at low Q^2 - low quality; Look for alternative input — compute Compton amplitude on the lattice



$$\mathcal{H}_{\mu\nu}^{VA}(x) = \left\langle \pi^{0}(p) \left| T[J_{\mu}^{\text{em}}(x)J_{\nu}^{W,A}(0)] \right| \pi^{-}(p) \right\rangle$$
$$M_{\pi}(Q^{2}) = -\frac{1}{6\sqrt{2}} \frac{\sqrt{Q^{2}}}{m_{\pi}} \int d^{4}x \omega(Q,x) \epsilon_{\mu\nu\alpha0} x_{\alpha} \mathcal{H}_{\mu\nu}^{VA}(x)$$

Direct LQCD computation for $\pi^- \rightarrow \pi^0 e^- \nu_e$

Feng, MG, Jin, Ma, Seng 2003.09798

5 LQCD gauge ensembles at physical pion mass Generated by RBC/UKQCD collaboration w. 2+1 flavor domain wall fermion

Ensemble	m_{π} [MeV]	L	Т	a^{-1} [GeV]	$N_{\rm conf}$	N _r	$\Delta t/a$
24D	141.2(4)	24	64	1.015	46	1024	8
32D	141.4(3)	32	64	1.015	32	2048	8
32D-fine	143.0(3)	32	64	1.378	71	1024	10
48I	135.5(4)	48	96	1.730	28	1024	12
64I	135.3(2)	64	128	2.359	62	1024	18

Blue: DSDR Red : Iwasaki



Quark contraction diagrams

Implications for the free nucleon γW -box

Seng, MG, Feng, Jin, 2003.11264

Indirectly constrains the free neutron γW -box

Independent confirmation of the empirical DR result AND uncertainty

 $\Delta_R^V = 0.02467(22)_{\rm DR} \rightarrow 0.02477(24)_{\rm LQCD+DR}$



Free-n RC in agreement by several groups & methods

Method	Δ_R^V		
DR with neutrino data (1)	0.02467(22)		
DR with neutrino data (2)	0.02471(18)		
DR with indirect lattice data	0.02477(24)		
Non-DR (1)	0.02426(32)		
Non-DR (2)	0.02473(27)		

C-Y Seng et al., Phys.Rev.Lett. 121 (2018) 24, 241804; C-Y Seng, MG, M.J. Ramsey-Musolf, Phys.Rev. D 100 (2019) Shiells, Blunden, Melnitchouk, Phys.Rev.D 104 (2021) 3, 033003 Seng, MG, Feng, Jin, 2003.11264 Czarnecki, Marciano, Sirlin, Phys.Rev. D 100 (2019) 7, 073008 Hayen, Phys.Rev.D 103 (2021) 11, 113001

Recently: first direct LQCD calculation on free neutron *P-XMa, X. Feng, MG, L-C Jin,* et al 2308.16755 $\Delta_R^V = 0.02439(19)_{LQCD}$ Discrepancy and uncertainty to be understood!

RC to neutron decay in EFT

Cirigliano et al, 2306.03138

р

п

Effective Field Theory: in principle, the proper tool to separate scales!

Formal consistency built in, RGE, transparent error estimation...

Precision limited by matching (LEC) and HO — relies on inputs (e.g. γW -box from DR!)

To improve: need to go to higher order — new LECs, still tractable?



Total RC: 1 + $\Delta_{\rm TOT}$ = 1.07761(27) %

Total RC from DR:1 + $\Delta_{TOT} = 1.07735(27)$ %

At present: order $O(\alpha, \alpha \alpha_s, \alpha^2)$ — consistent with matching input? realistic to go beyond? Nuclear structure doable?

Interpretation of Cabibbo Angle Anomaly

CAA summary - 3 anomalies!

3 observables: $|V_{us}|^{K\ell 3}$, $|V_{us}/V_{ud}|^{K\mu 2}$, V_{ud} 2 quantities to determine: V_{us} , V_{ud}

3 ways to test unitarity

$$\Delta_{\text{CKM}}^{(1)} = |V_{ud}|^2 + |V_{us}^{K_{\ell 3}}|^2 - 1 = -0.00176(56) -3.1\sigma$$

$$\Delta_{\text{CKM}}^{(2)} = |V_{ud}|^2 \left[1 + \left(\left| \frac{V_{us}}{V_{ud}} \right|^{K_{\mu 2}} \right)^2 \right] - 1 = -0.00098(58) -1.7\sigma$$

 $K_{\mu 2}$ result shows better agreement with unitarity than $K_{\ell 3}$ result when $|V_{ud}|$ obtained from beta decays:

$$\Delta V_{us}(K_{\ell 3} - K_{\mu 2}) = V_{us}^{K_{\ell 3}} - V_{ud} \left(\frac{V_{us}}{V_{ud}}\right)^{K_{\mu 2}} = -0.0174(73) -2.4\sigma$$

 $\Delta^{(3)}_{CKM}$ uses no information from β decays:

$$\Delta_{\text{CKM}}^{(3)} = |V_{us}^{K_{\ell 3}}|^2 \left[\left(\frac{1}{|V_{us}/V_{ud}|^{K_{\mu 2}}} \right)^2 + 1 \right] - 1 = -0.0164(63) -2.6\sigma$$

CAA in presence of RH currents

- In SM, W couples only to LH chiral fermion states
- New physics with couplings to RH currents could explain both unitarity deficit and K_{ℓ_3} - K_{μ_2} difference
- Define ϵ_R = admixture of RH currents in non-strange sector $\epsilon_R + \Delta \epsilon_R$ = admixture of RH currents in strange sector



From current fit:

 $\epsilon_R = -0.69(27) \times 10^{-3} (2.5\sigma)$ $\Delta \epsilon_R = -3.9(1.6) \times 10^{-3} (2.4\sigma)$ $\epsilon_R = \Delta \epsilon_R = 0$ excluded at 3.1 σ



Cirigliano et al.

PLB 838 (2023)

Summary and Outlook

Summary & Outlook

Cabibbo unitarity deficit at 2-3 σ observed

Great improvement in theory of RC in past 5 years EFT formulation of free-n RC - promising avenue

Great improvement in free-n decay exp., more coming! Barely any improvement in $0^+ - 0^+$ nuclear decays

Nuclear uncertainties under scrutiny

Meson decays: future more precise experiments

Cabibbo anomaly interpretable in terms of BSM (e.g. flavor-dependent RH currents)

Keep tuned!