

Generalized Polarizabilities of the proton

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16th International Conference on Meson-Nucleon Physics
and the Structure of the Nucleon

October 2023

Outline

Introduction to the GPs

Overview / Status & Challenges

Recent results / Jlab & MAMI

Prospects / VCS-II @ Jlab, measuring w positrons, ...

Proton Polarizabilities

Fundamental structure constants
(such as mass, size, shape, ...)

Response of the nucleon to external EM field

Sensitive to the full excitation spectrum

Accessed experimentally through Compton Scattering

RCS: static polarizabilities \rightarrow net effect on the nucleon

Virtual Compton Scattering:

Virtuality of photon gives access to the GPs : $\alpha_E(Q^2)$ & $\beta_M(Q^2)$ (+ 4 spin GPs)

\rightarrow mapping out the spatial distribution of the polarization densities

Fourier transform of densities of electric charges and magnetization of a nucleon deformed by an applied EM field

PDG

150 Baryon Summary Table

N BARYONS
($S = 0, I = 1/2$)

$p, N^+ = uud; \quad n, N^0 = udd$

p

$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$

Mass $m = 1.00727646681 \pm 0.0000000009 \text{ u}$

Mass $m = 938.272046 \pm 0.000021 \text{ MeV} \text{ }^{[a]}$

$|m_p - m_{\bar{p}}|/m_p < 7 \times 10^{-10}, \text{ CL} = 90\% \text{ }^{[b]}$

$|q_p/m_p|/(q_e/m_e) = 0.9999999991 \pm 0.0000000009$

$|q_p + q_{\bar{p}}|/e < 7 \times 10^{-10}, \text{ CL} = 90\% \text{ }^{[b]}$

$|q_p + q_e|/e < 1 \times 10^{-21} \text{ }^{[c]}$

Magnetic moment $\mu = 2.792847356 \pm 0.000000023 \mu_N$

$(\mu_p + \mu_{\bar{p}}) / \mu_p = (0 \pm 5) \times 10^{-6}$

Electric dipole moment $d < 0.54 \times 10^{-23} \text{ e cm}$

Electric polarizability $\alpha = (11.2 \pm 0.4) \times 10^{-4} \text{ fm}^3$

Magnetic polarizability $\beta = (2.5 \pm 0.4) \times 10^{-4} \text{ fm}^3 \text{ }^{(S = 1.2)}$

Charge radius, μp Lamb shift $= 0.84087 \pm 0.00039 \text{ fm} \text{ }^{[d]}$

Charge radius, $e p$ CODATA value $= 0.8775 \pm 0.0051 \text{ fm} \text{ }^{[d]}$

Magnetic radius $= 0.777 \pm 0.016 \text{ fm}$

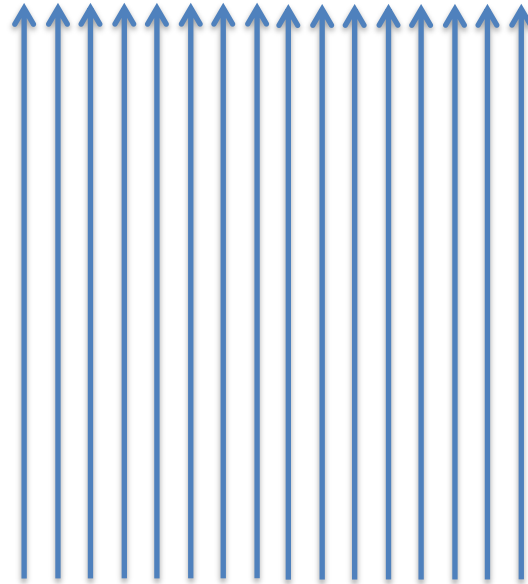
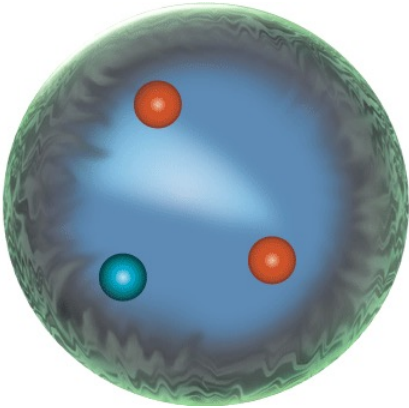
Mean life $\tau > 2.1 \times 10^{29} \text{ years}, \text{ CL} = 90\% \text{ }^{[e]}$ ($p \rightarrow$ invisible mode)

Mean life $\tau > 10^{31} \text{ to } 10^{33} \text{ years} \text{ }^{[e]}$ (mode dependent)

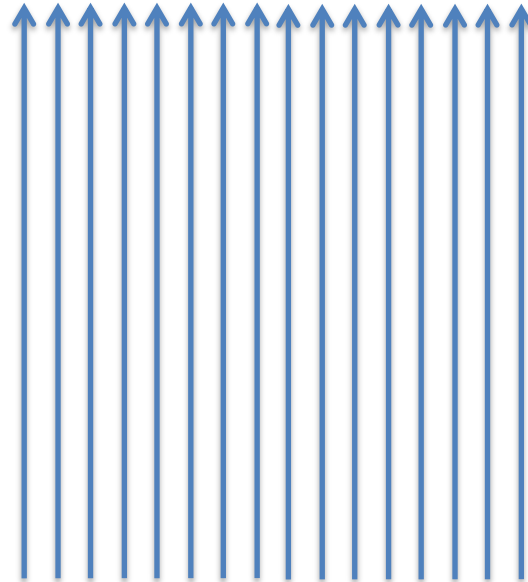
Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon



\vec{E}

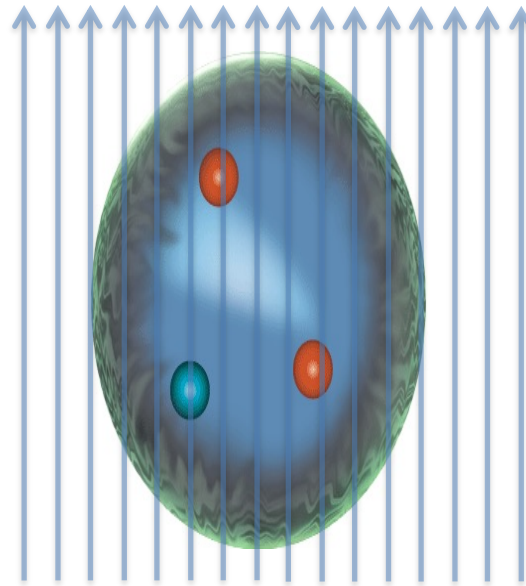
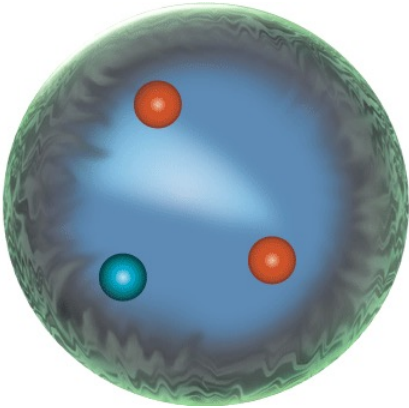


\vec{B}

Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon

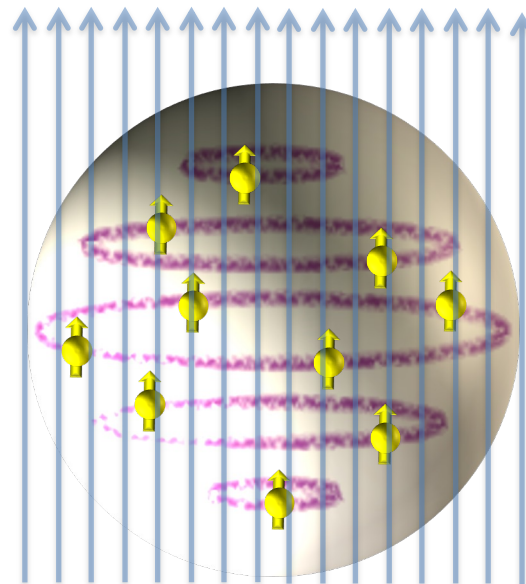


\vec{E}

“stretchability”

$$\vec{d}_{E \text{ induced}} \sim \alpha \vec{E}$$

External field deforms the charge distribution



\vec{B}

“alignability”

$$\vec{d}_{M \text{ induced}} \sim \beta \vec{B}$$

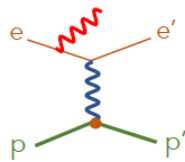
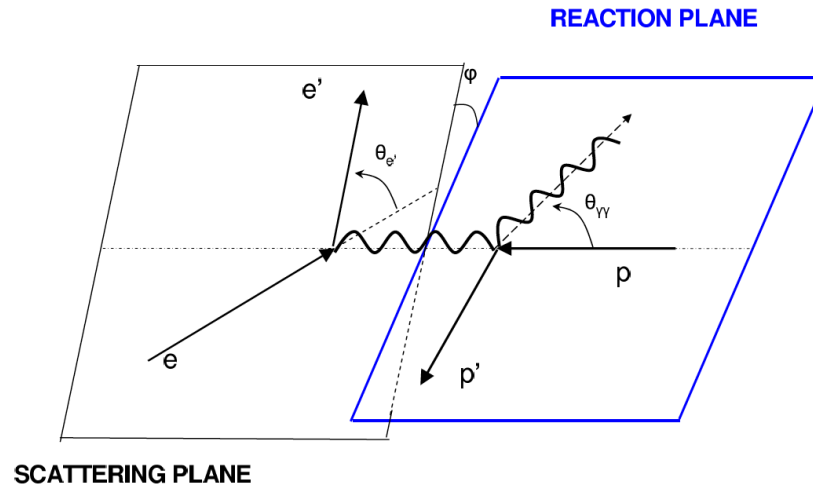
$$\beta_{\text{para}} > 0$$

$$\beta_{\text{diam}} < 0$$

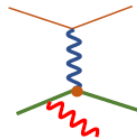
Paramagnetic: proton spin aligns with the external magnetic field

Diamagnetic: π -cloud induction produces field counter to the external perturbation

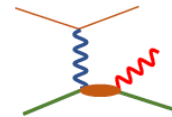
Virtual Compton Scattering



Bethe-Heitler



Born VCS



non-Born VCS

Elastic FFs

GPs

Virtual Compton Scattering

DR

valid below & above
Pion threshold

Dispersive integrals
for Non Born amplitudes

Spin GPs are fixed

Scalar GPs have
an unconstrained part

Fit to the experimental
cross sections at each Q^2

LEX

valid only below
Pion threshold

Response functions

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

$$\Psi_0 = v_1 \cdot \left(P_{LL} - \frac{1}{\epsilon} P_{TT}\right) + v_2 \cdot P_{LT}$$

Subtract the spin part

$$P_{TT} = [P_{TT \text{ spin}}]$$

$$P_{LT} = -\frac{2M}{\alpha_{em}} \sqrt{\frac{q'^2_{cm}}{Q^2}} \cdot G_E^p(Q^2) \cdot \beta_M(Q^2) + [P_{LT \text{ spin}}]$$

utilize DR

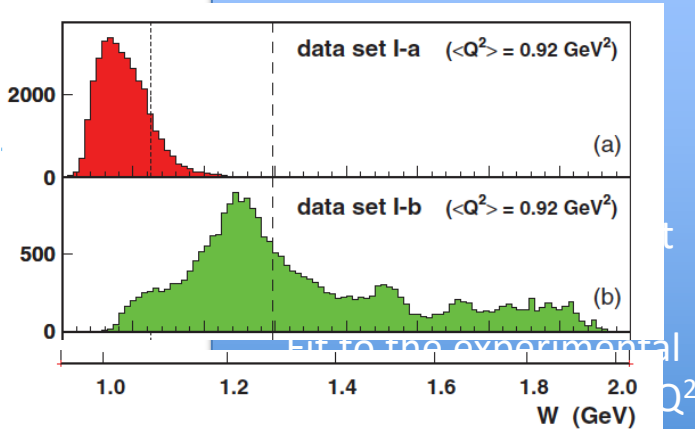
scalar GPs α_E and β_M

Virtual Compton Scattering

DR

valid below & above
Pion threshold

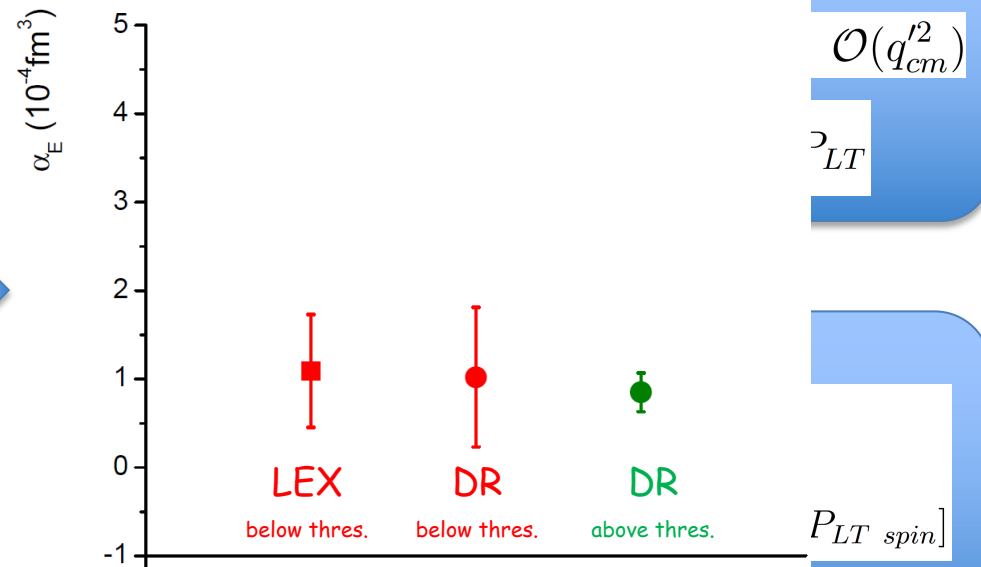
Dispersive integrals
for Non Born amplitudes



LEX

valid only below
Pion threshold

Response functions

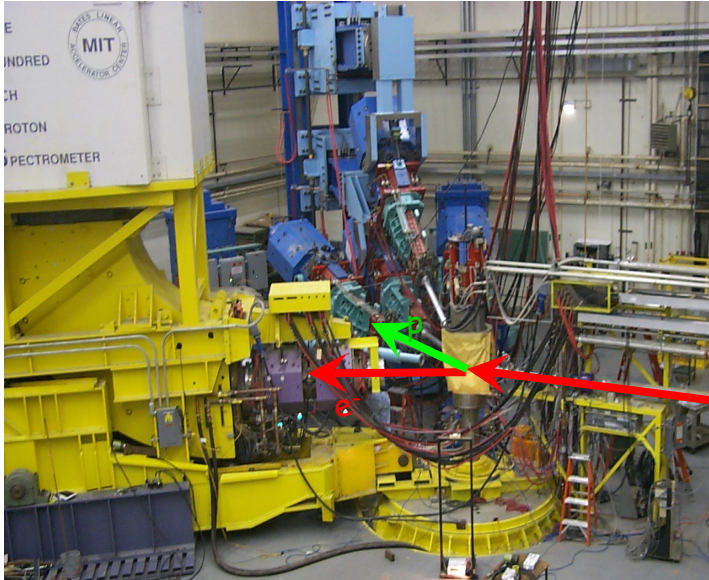


utilize DR

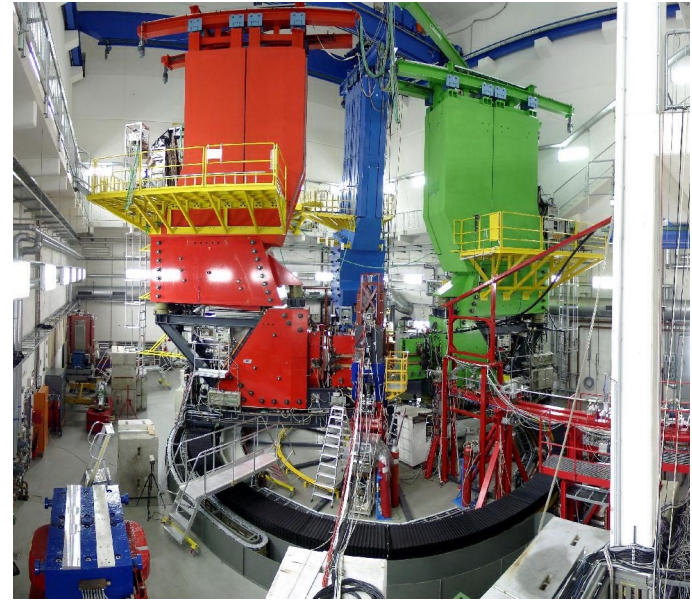
scalar GPs α_E and β_M

Early Experiments

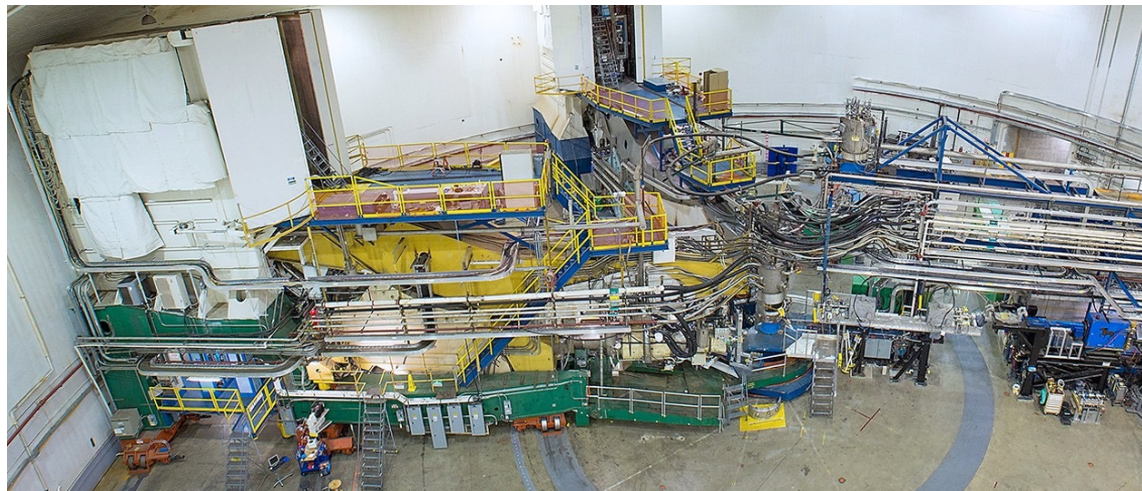
MIT-Bates @ $Q^2=0.06 \text{ GeV}^2$



MAMI-A1 @ $Q^2=0.33 \text{ GeV}^2$



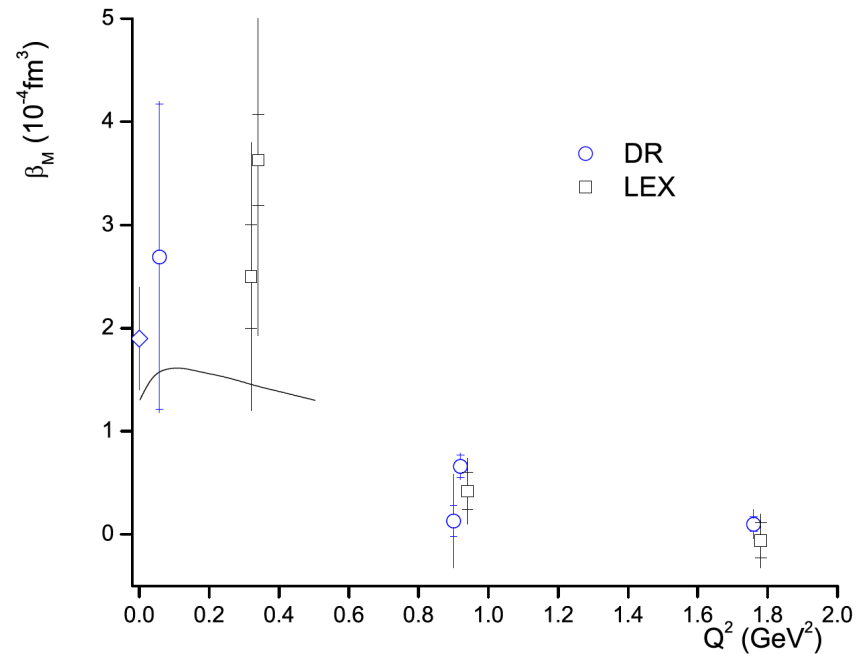
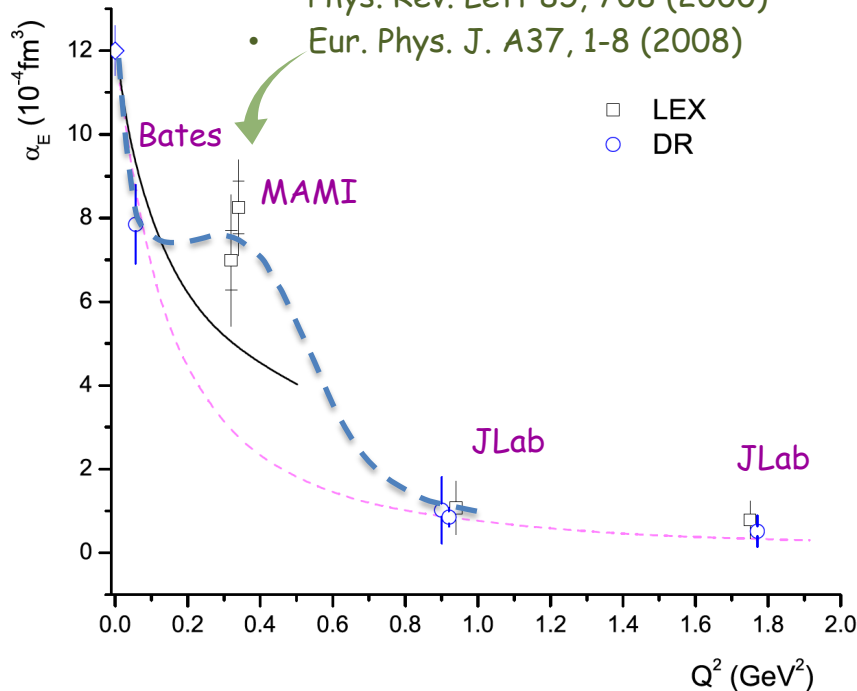
Jlab-Hall A @ $Q^2=0.9 \text{ \& } 1.8 \text{ GeV}^2$



Early Experiments

$Q^2 = 0.33 \text{ (GeV/c)}^2$ measured twice at MAMI:

- Phys. Rev. Lett 85, 708 (2000)
- Eur. Phys. J. A37, 1-8 (2008)



$\alpha_E \approx 10^{-3} V_N$ (stiffness / relativistic character)

Data: non-trivial Q^2 dependence of α_E (?)

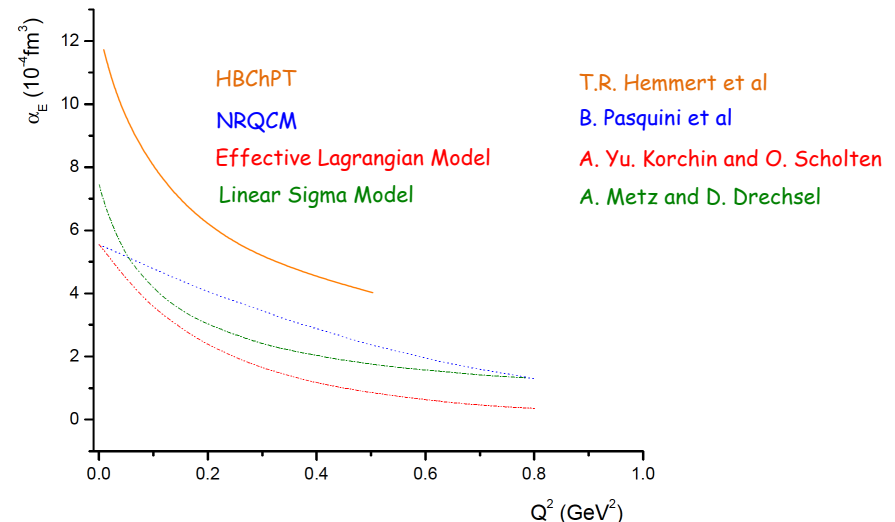
Theory: monotonic fall-off

β_M small \leftrightarrow cancellation of competing mechanisms

Large uncertainties

Higher precision measurements needed

→ Quantify balance between dia/para-magnetism



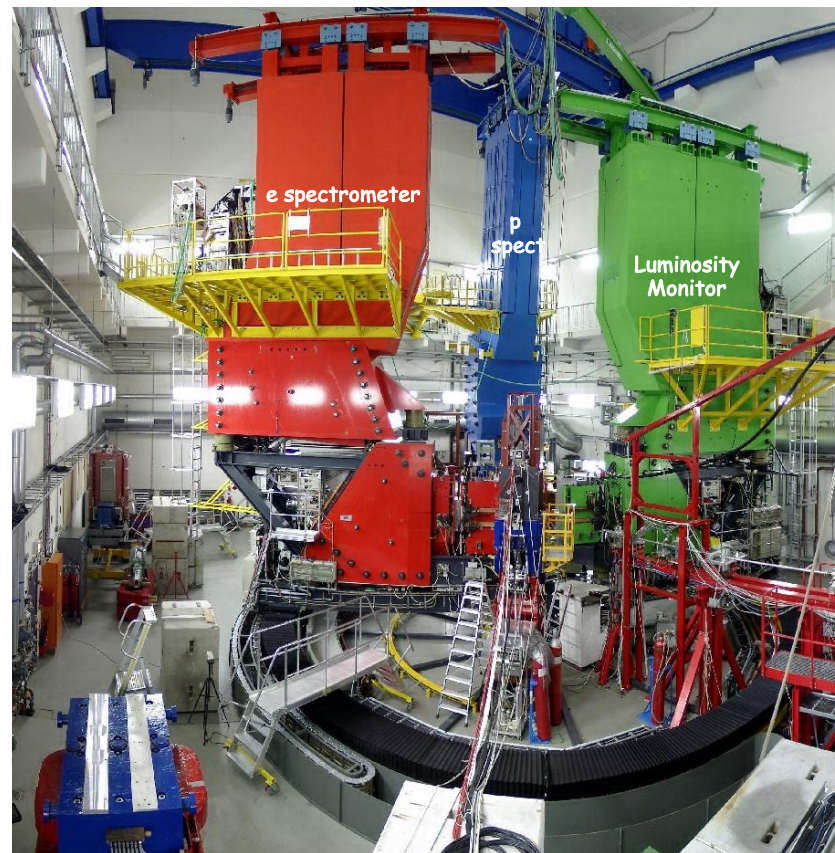
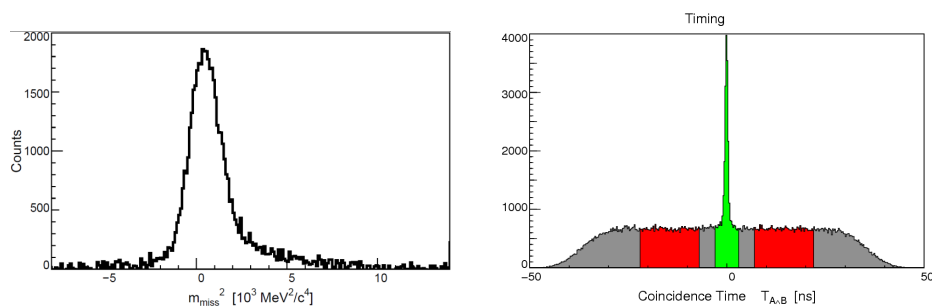
Recent Experiments

Recent Measurements: MAMI

MAMI A1/1-09 (vcsq2) below threshold

MAMI A1/3-12 (vcsdelta) above threshold

Both experiments utilized
the A1 setup at MAMI



For LEX the higher order terms have to be kept small / under control

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

Refined analysis procedure / phase space masking to keep these terms smaller than $\sim 2\%$ - 3% level

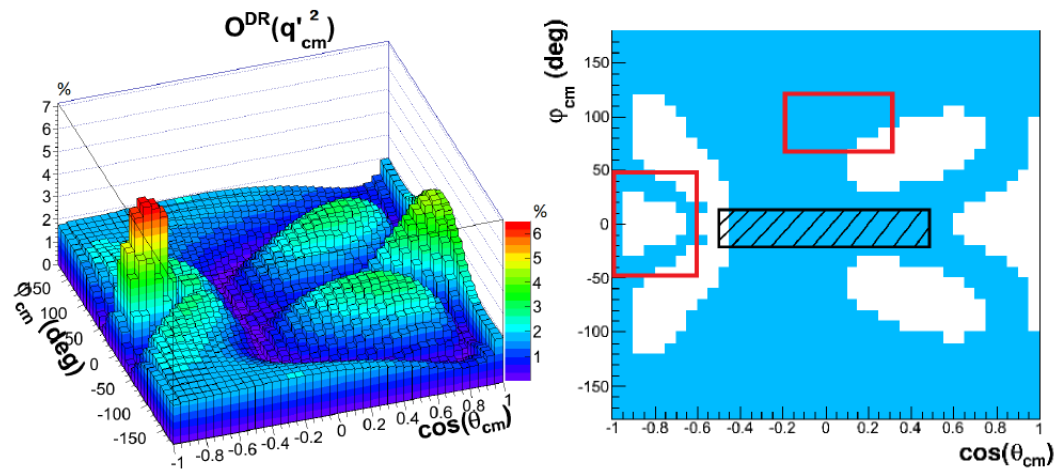
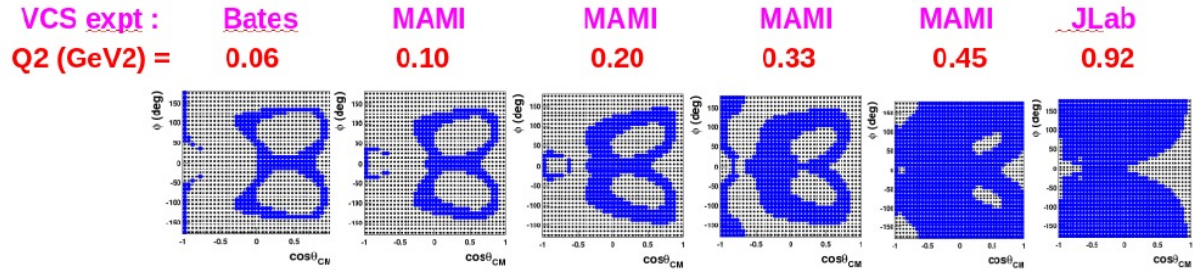
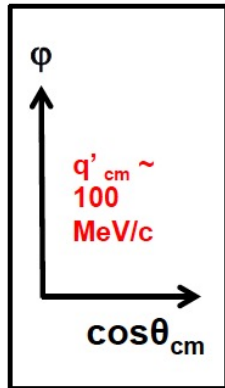


Figure 3.13: (Left) behavior of $\mathcal{O}^{DR}(q'^2_{cm})$ in the $(\cos(\theta_{cm}), \varphi_{cm})$ -plane at $q'_{cm} = 87.5 \text{ MeV}/c$ and (right) two-dimensional representation of the angular region where $\mathcal{O}^{DR}(q'^2_{cm}) < 2\%$ (blue), the red squares correspond to the two areas of interest to perform the GP extraction.

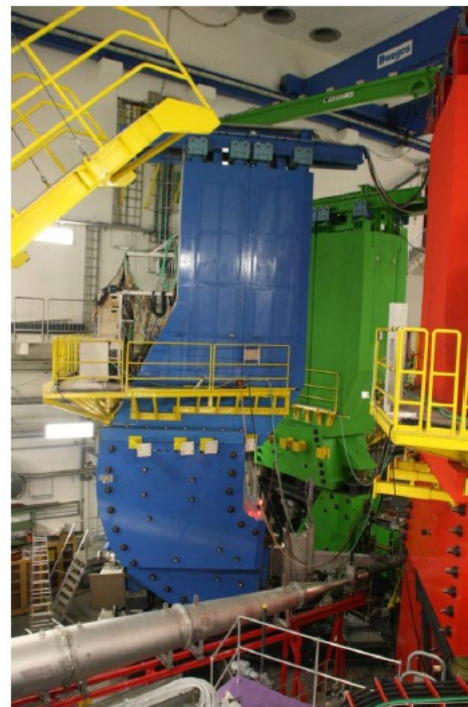
Figure from PhD thesis of L. Correa, Mainz / Cl. Ferrand

Blue bins = where the higher-order estimator is $< 3\%$
(LEX truncation « valid »)

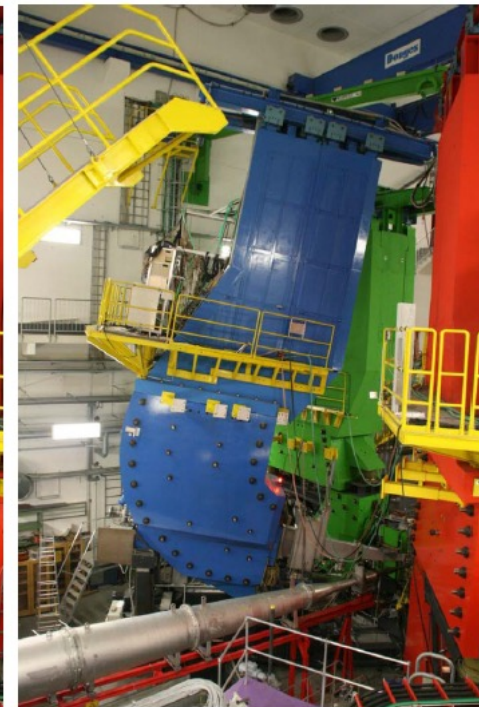


New « vcsq2 » data:

- OOP kinematics (to access the blue region)
- LEX Fit done with bin selection at $Q^2 = 0.1$ and 0.2 GeV².
- was found not necessary at $Q^2 = 0.45$ GeV².



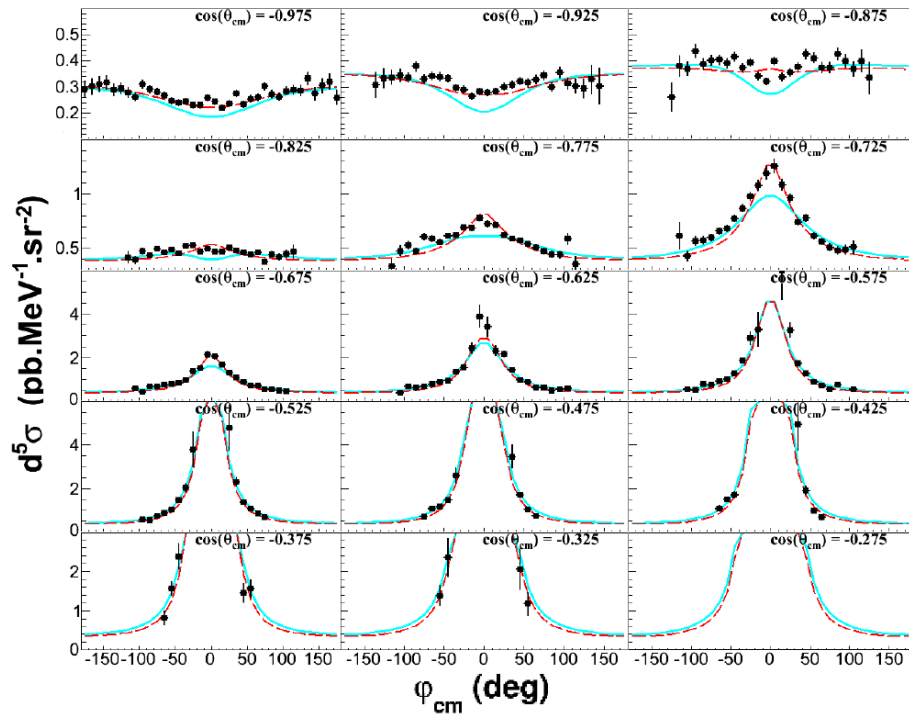
In-plane



8.5 deg OOP

~ 1.0 GeV beam

$Q^2 = 0.1 \text{ (GeV/c)}^2, 0.2 \text{ (GeV/c)}^2, \text{ and } 0.45 \text{ (GeV/c)}^2$



BH+B

Polarizability
effect

GP effect typically 5% - 15%
of the cross section

Polarizability fits:

DR fit:

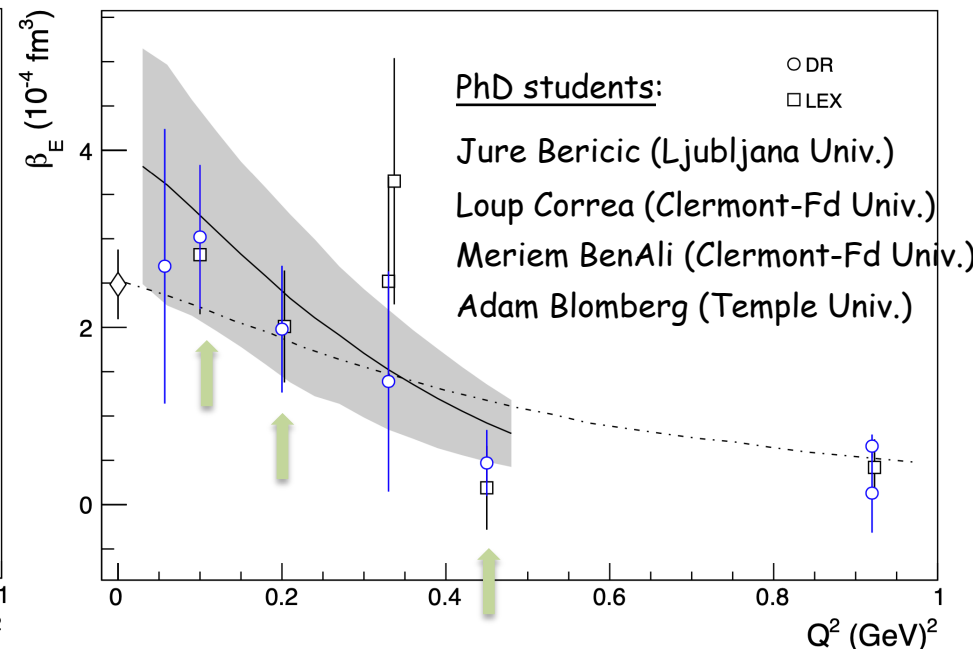
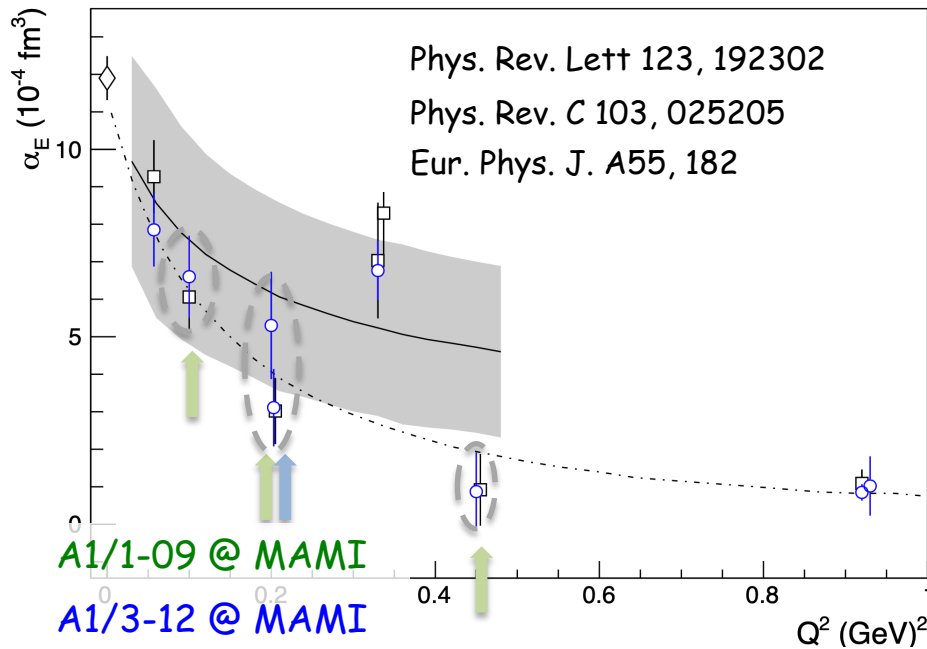
DR calculation includes full dependency in q'_{cm}

LEX fit:

truncated in q'_{cm} . Suppress contribution
from higher order terms

Figure 5.8: Setting INP: measured $ep \rightarrow ep\gamma$ cross section at fixed $q'_{cm} = 112.5 \text{ MeV/c}$ with respect to φ_{cm} for all the $\cos(\theta_{cm})$ -bins. The curves follow the convention of figure 5.6.

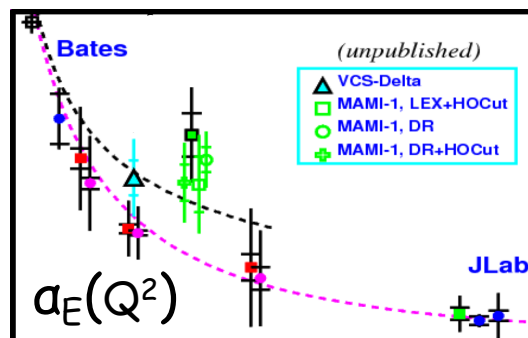
MAMI Results



Revisiting the $Q^2=0.33 \text{ GeV}^2$ data

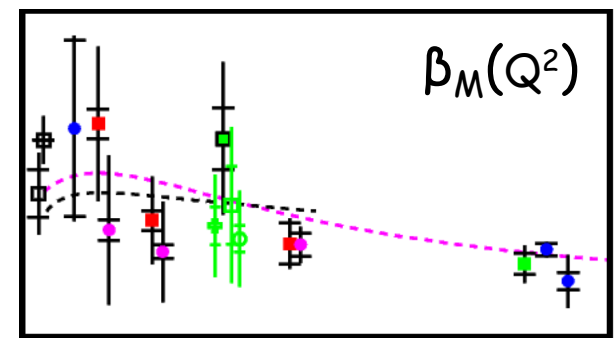
Analysis revisited (unpublished):

The α_E puzzle still holds



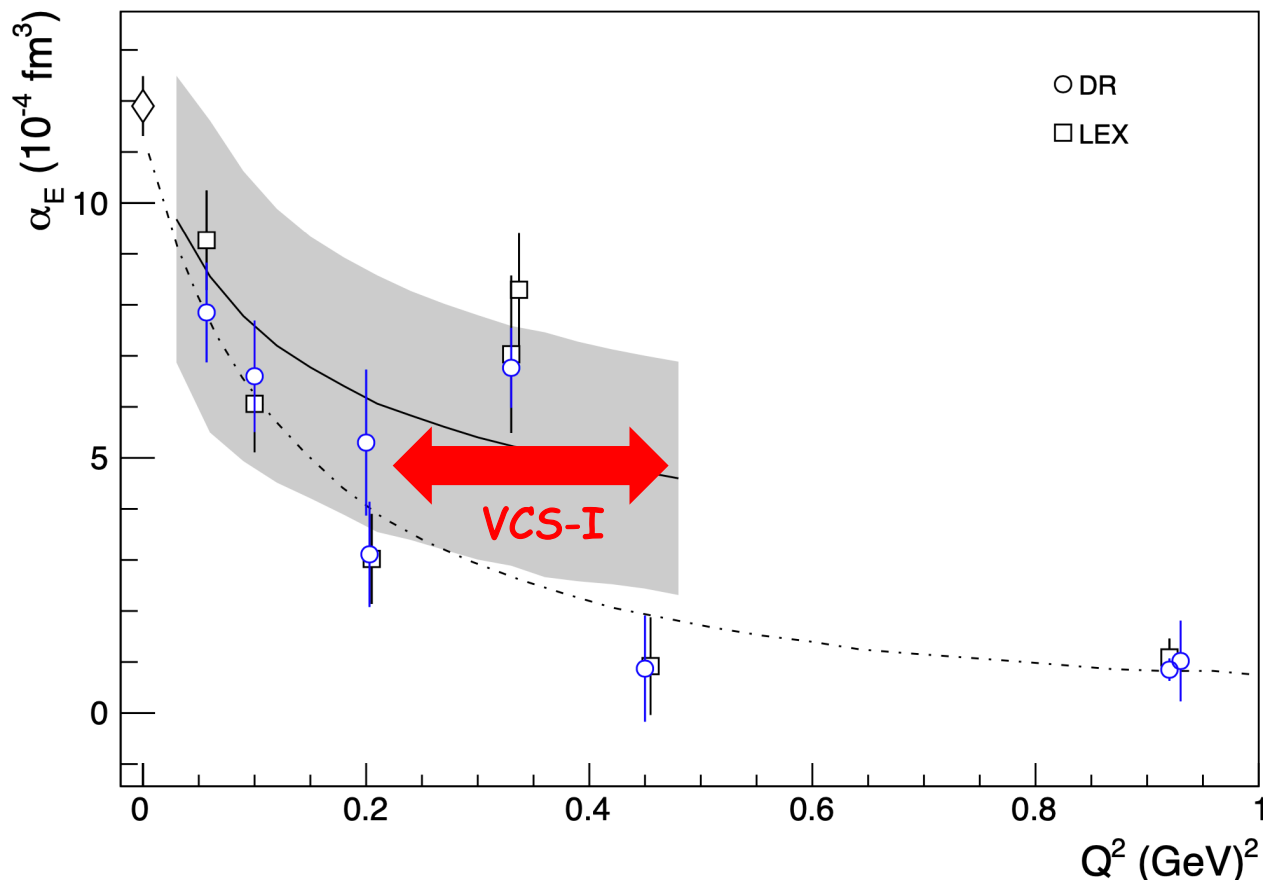
Re-fits at
 $Q^2=0.33$
 GeV^2
(H.F.)

LEX and DR
Updated HO-cut

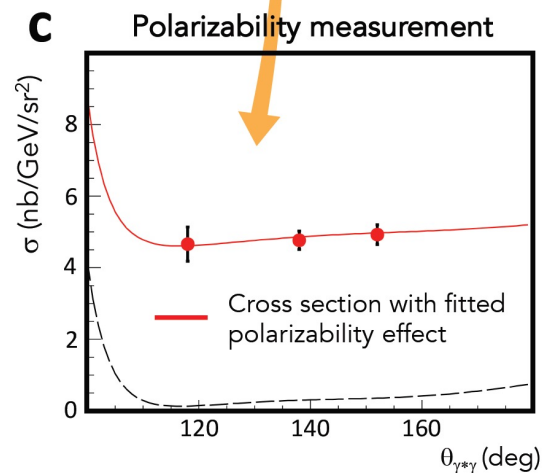
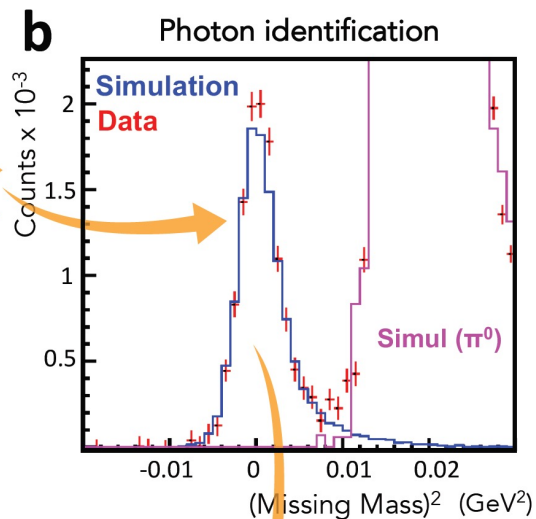
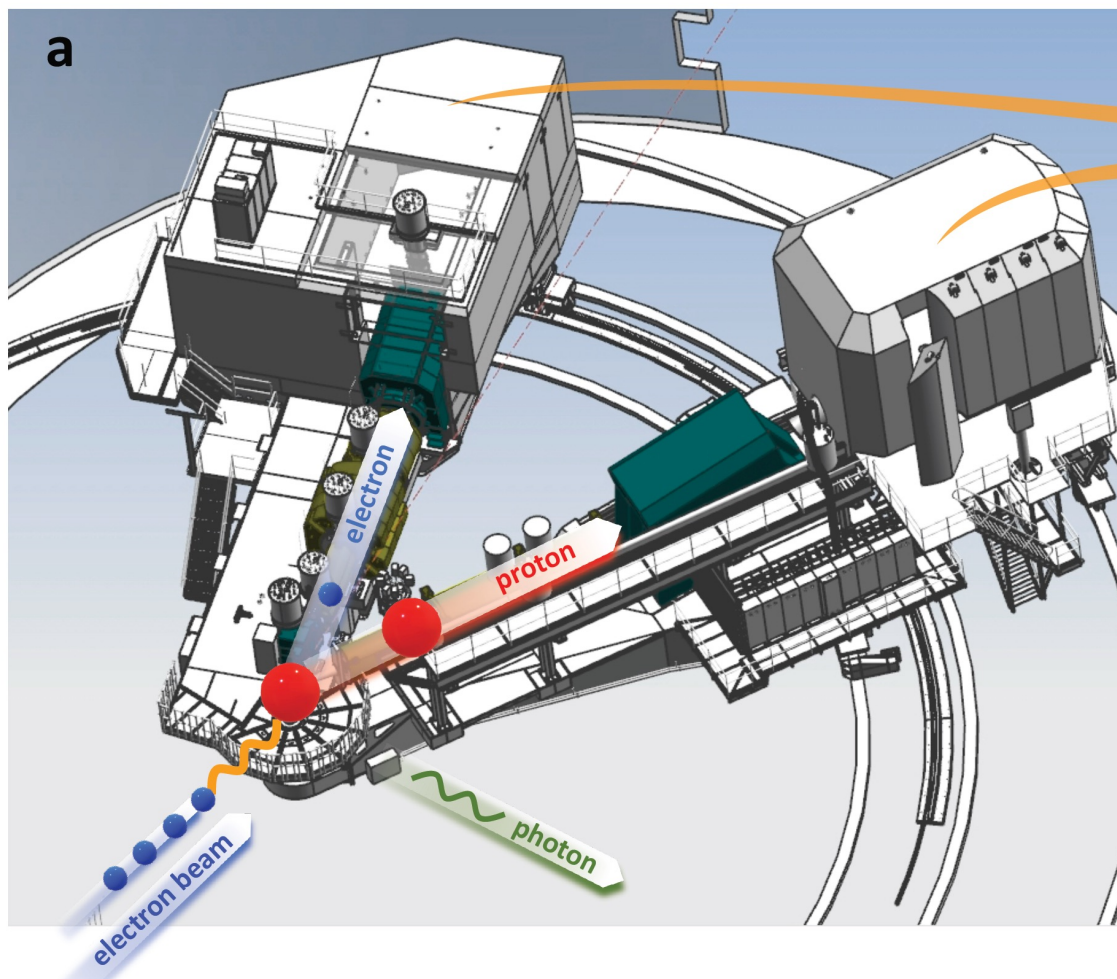


Jlab : VCS-I Experiment (E12-15-001) in Hall C

High precision measurements targeting
explicitly the kinematics of interest for a_E



The experiment



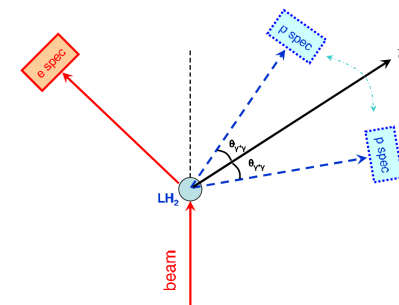
Hall C: SHMS, HMS
4.56 GeV
20 μ A
Liquid hydrogen 10 cm

cross sections & azimuthal asymmetries

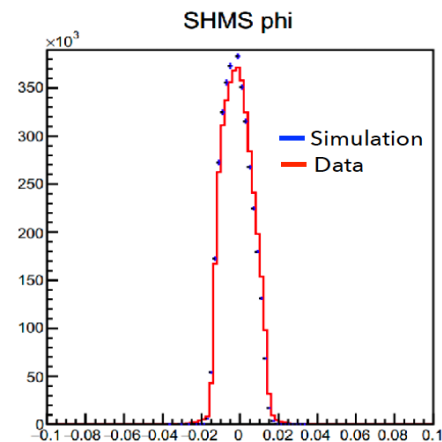
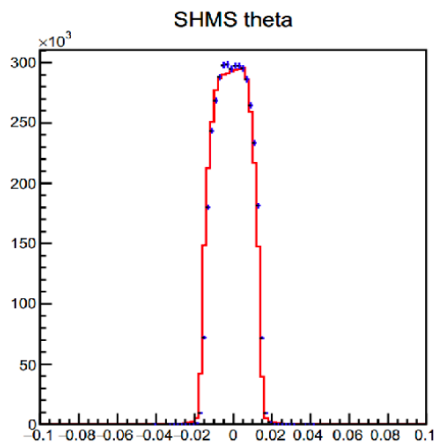
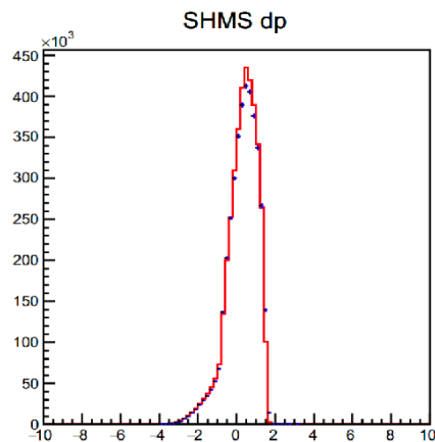
$$A_{(\phi_{\gamma^*\gamma}=0,\pi)} = \frac{\sigma_{\phi_{\gamma^*\gamma}=0} - \sigma_{\phi_{\gamma^*\gamma}=180}}{\sigma_{\phi_{\gamma^*\gamma}=0} + \sigma_{\phi_{\gamma^*\gamma}=180}}$$

sensitivity to GPs

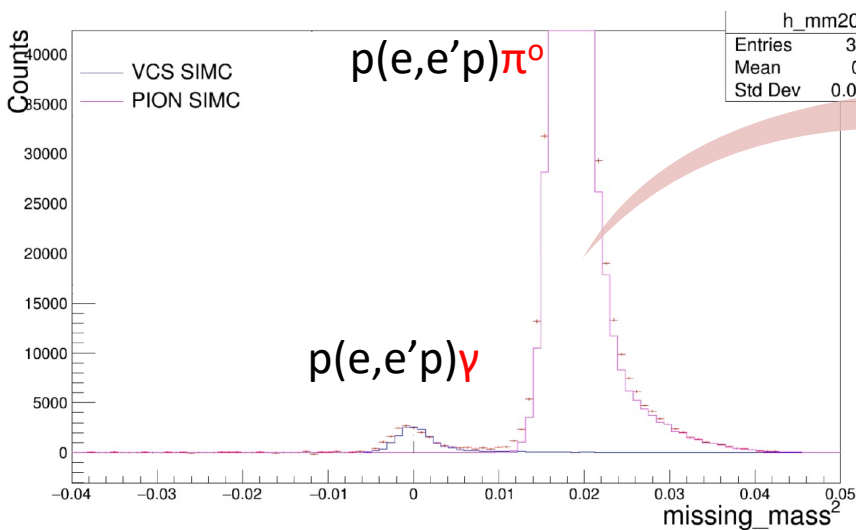
suppression of systematic asymmetries



Elastic data

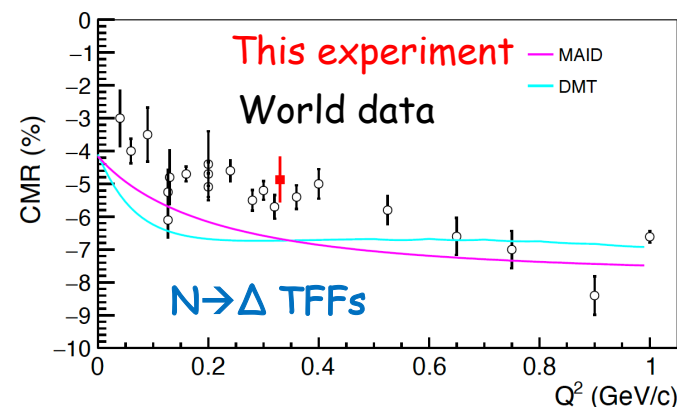
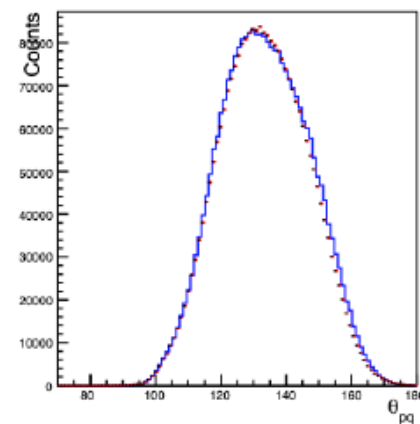
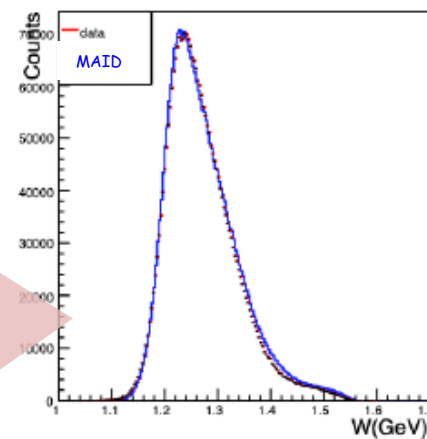


$p(e,e'p)\pi^0$

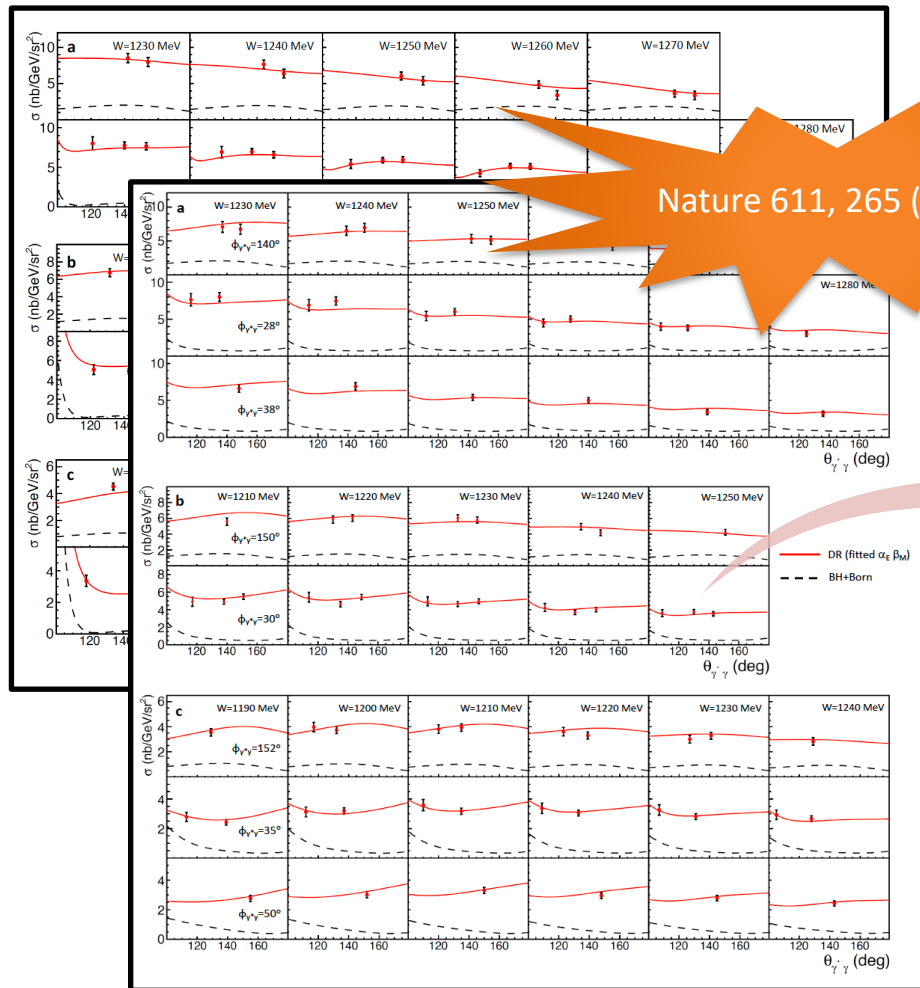


h_mm20

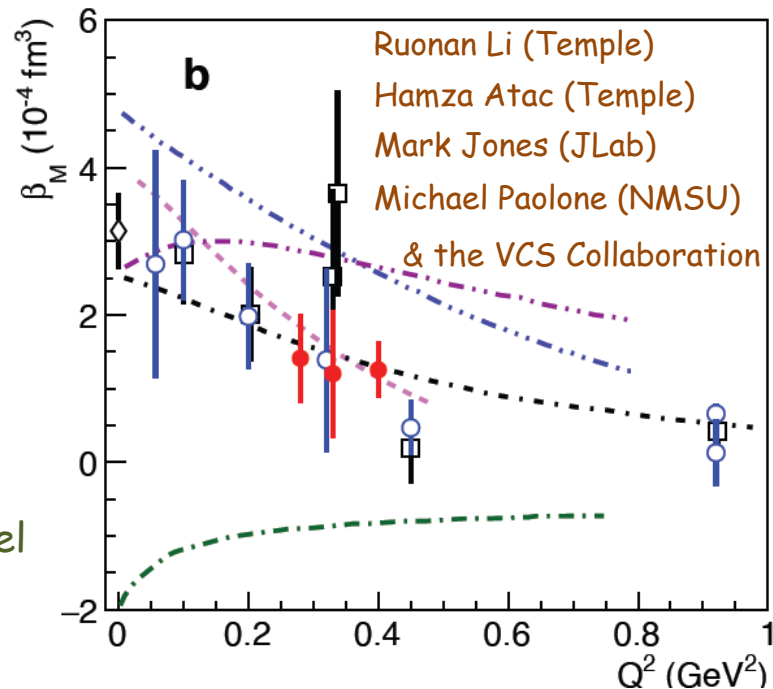
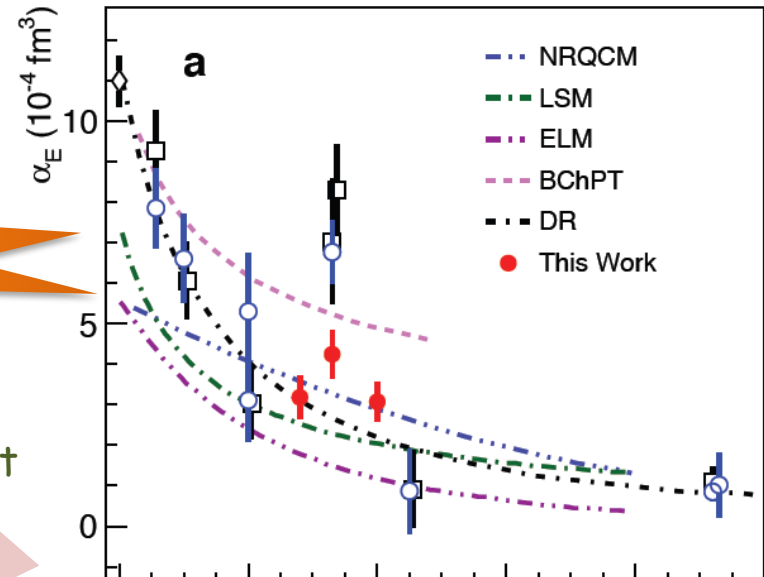
Entries	359772
Mean	0.0186
Std Dev	0.004849



VCS-I results: GPs



DR fit



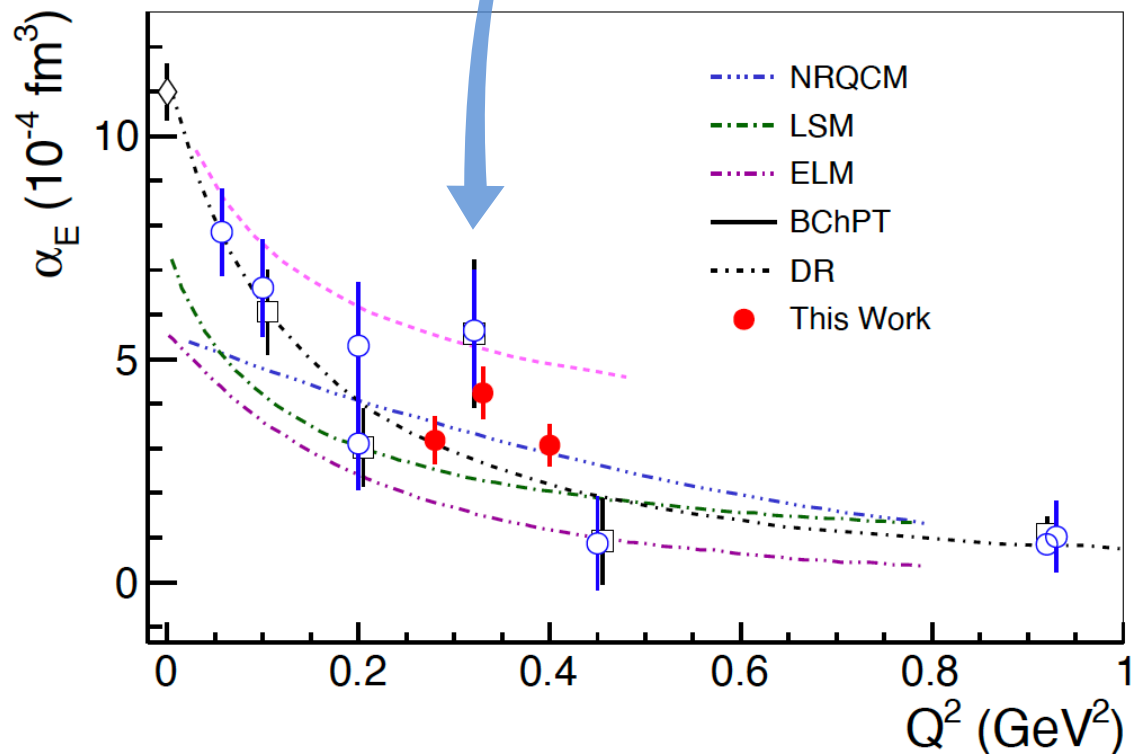
Experimental cross sections are compared to the DR model predictions for all possible values for the GPs

→ $\alpha_E(Q^2)$ and $\beta_M(Q^2)$ are fitted by a χ^2 minimization

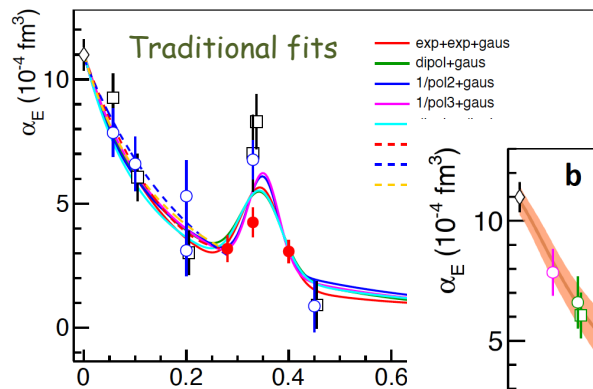
Electric GP (Q^2)

Is there a non-trivial structure?

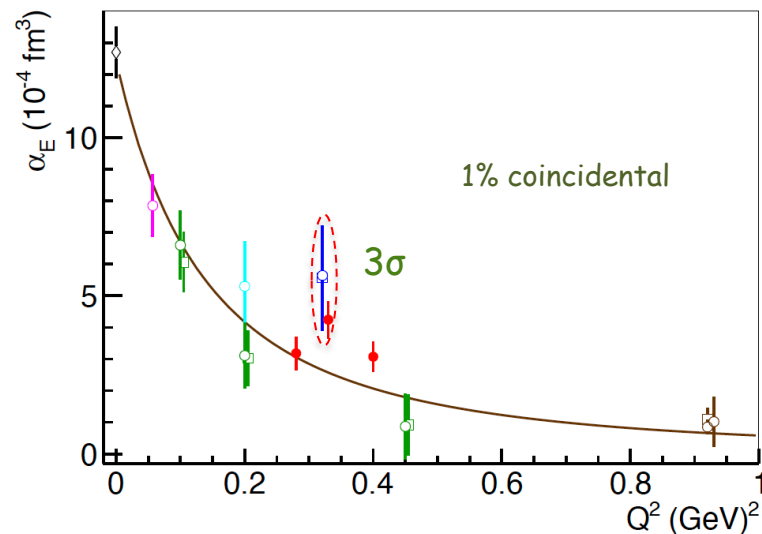
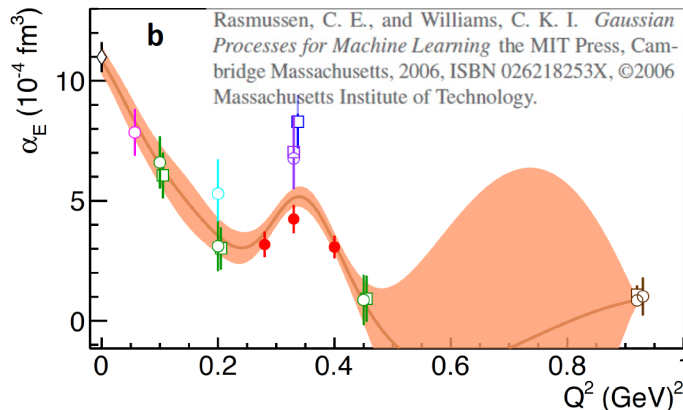
MAMI-I re-analysis (unpublished)



Electric GP



Data-driven techniques:
no underlying functional
form is assumed



Is the observed α_E structure coincidental or not?

If true: Measure the shape precisely \rightarrow input to theory

If not: We are able to show it with more measurements

Strong tension between world data (?)

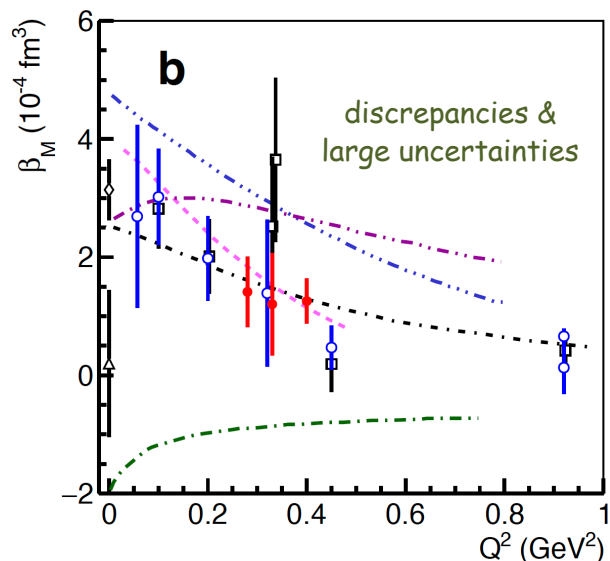
Things we do not yet understand well?

Underestimated uncertainties? ...

Magnetic GP: Large uncertainties & discrepancies
Needed to disentangle diamagnetism vs
paramagnetism in the proton

Ability to measure α_E and β_M with superb precision
and with consistent systematics across Q^2

Magnetic GP



Theory: BxPT

Generalized polarizabilities of the nucleon in baryon chiral perturbation theory

Vadim Lensky^{1,2,3,a}, Vladimir Pascalutsa¹, Marc Vanderhaeghen¹

¹ Institut für Kernphysik, Cluster of Excellence PRISMA, Johannes Gutenberg Universität Mainz, 55128 Mainz, Germany

² Institute for Theoretical and Experimental Physics, Moscow 117218, Russia

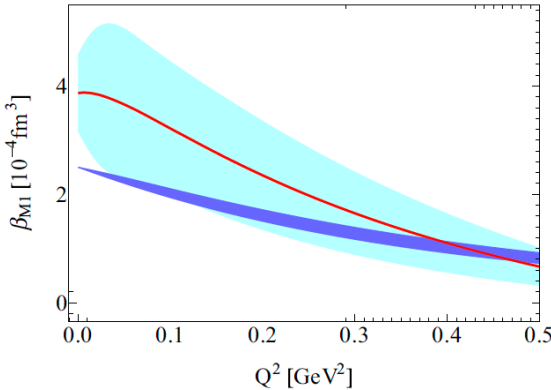
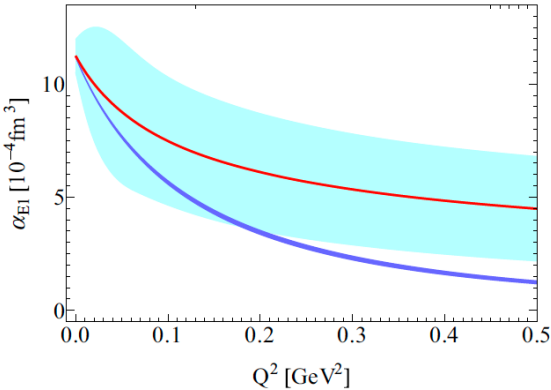
³ National Research Nuclear University MEPhI (Moscow Engineering Physics Institute), Moscow 115409, Russia



BxPT calculation to NLO
in the δ -counting scheme



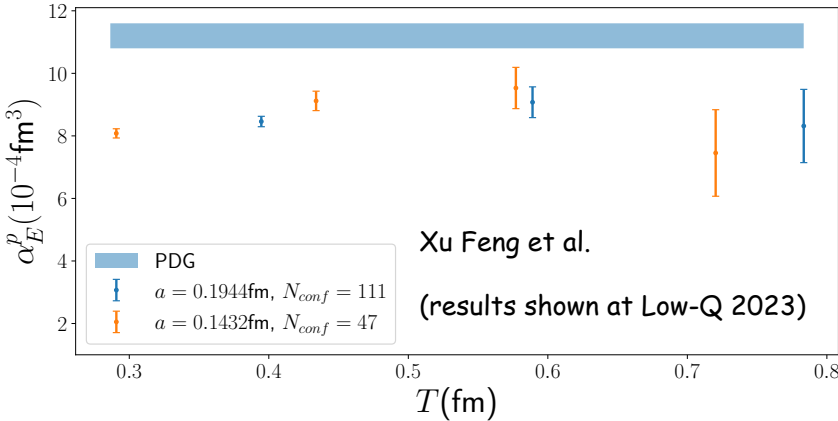
DR calculation
D. Drechsel, B. Pasquini, M. Vanderhaeghen,
Phys. Rep. 378,99 (2003)



Theory: Lattice QCD

Lattice QCD results for
the static polarizabilities

Next step: Lattice QCD
calculations for the GPs



Spatial dependence of induced polarizations

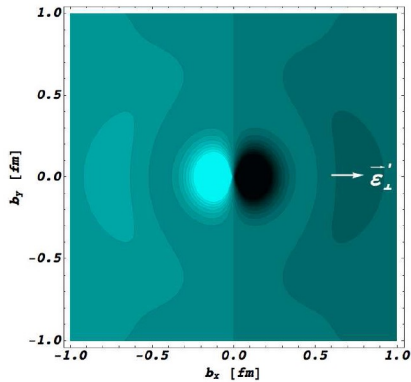
Nucleon form factor data → light-front quark charge densities

Formalism extended to the deformation of these quark densities when applying an external e.m. field:

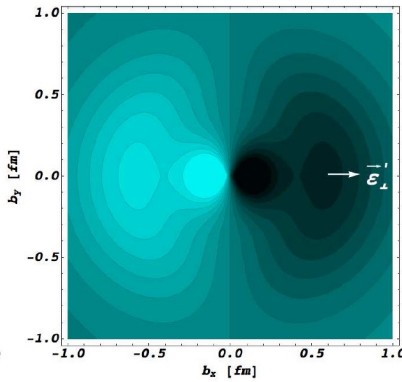
GPs → spatial deformation of charge & magnetization densities under an applied e.m. field

Induced polarization in a proton when submitted to an e.m. field

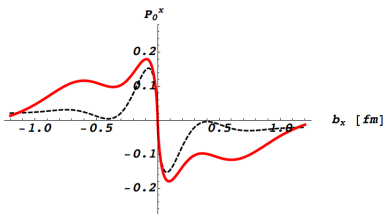
GP I



GP II



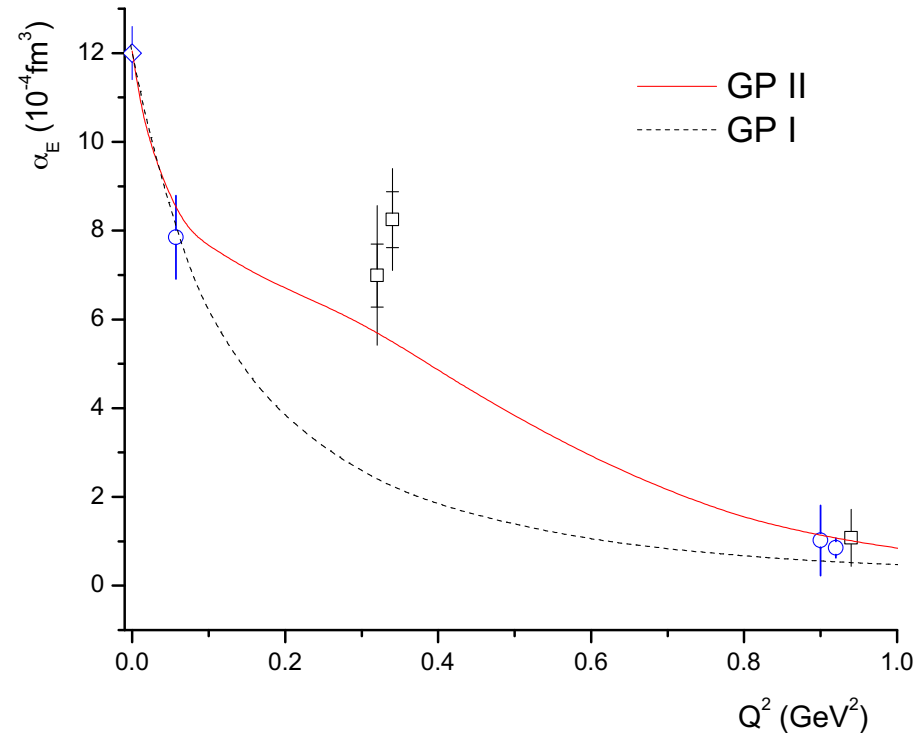
Light (dark) regions → largest (smaller) values
(photon polarization along x-axis, as indicated)



Induced polarization along $b_y=0$

Phys. Rev. Lett. 104, 112001 (2010)

M. Gorchtein, C. Lorce, B. Pasquini, M. Vanderhaeghen

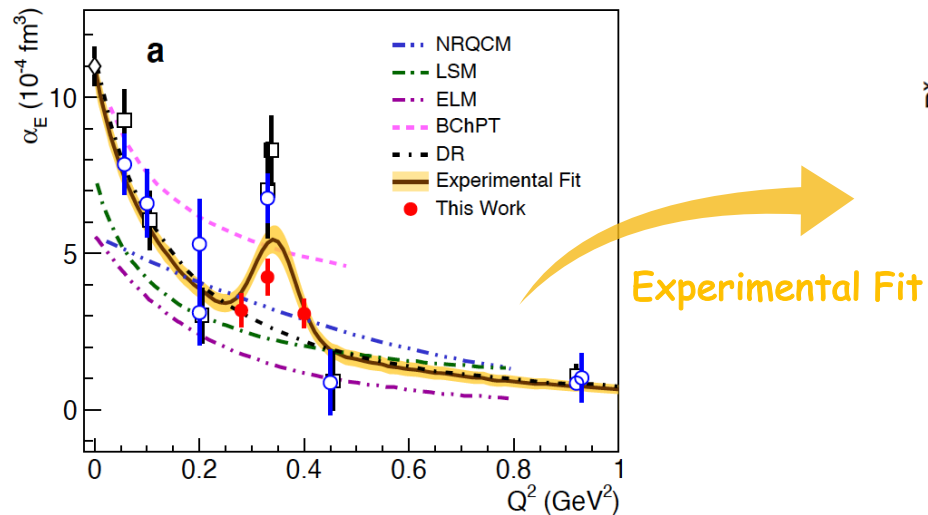


Spatial dependence of induced polarizations

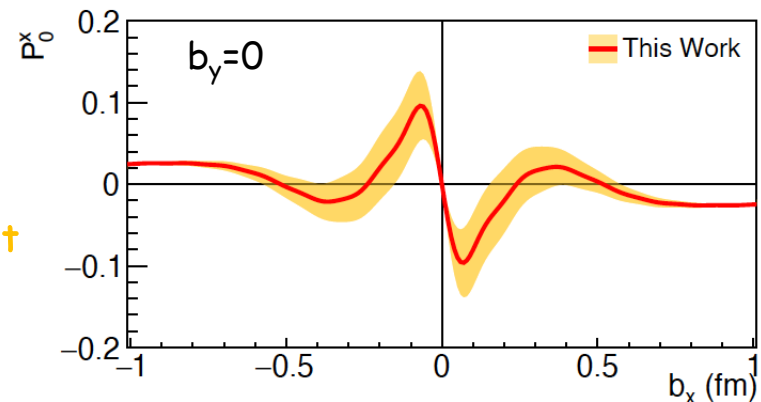
Nucleon form factor data → light-front quark charge densities

Formalism extended to the deformation of these quark densities when applying an external e.m. field:

GPs → spatial deformation of charge & magnetization densities under an applied e.m. field



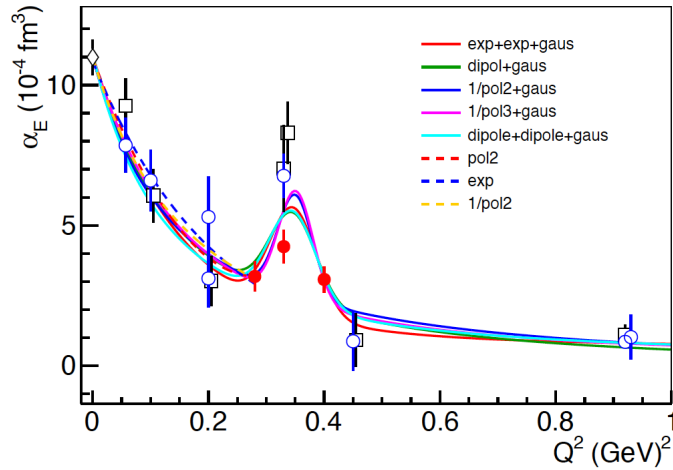
Induced polarization in a proton
when submitted to an e.m. field



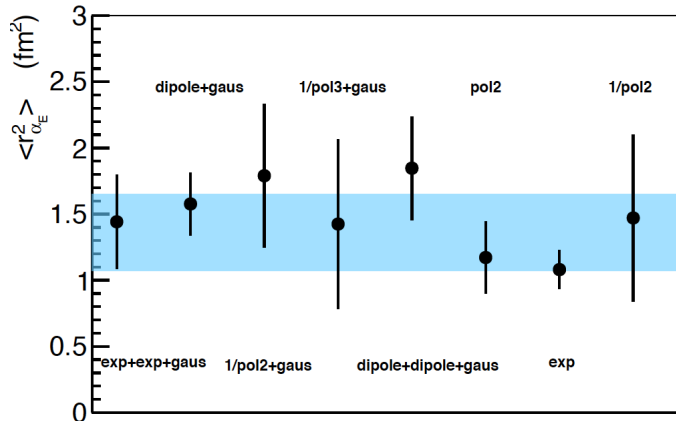
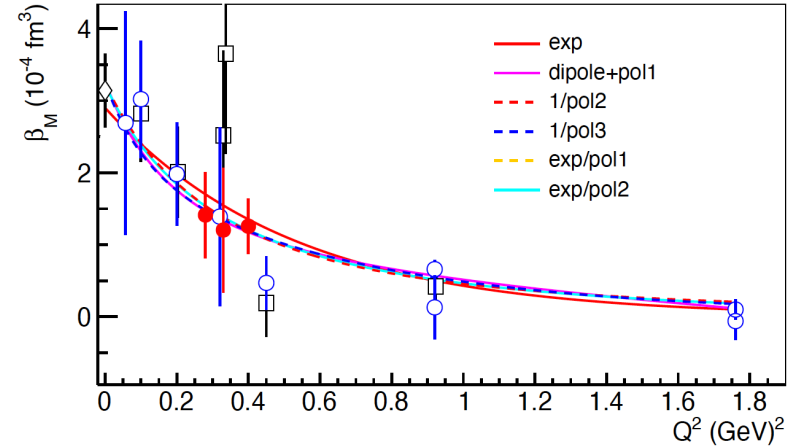
x-y defines the transverse plane with
the z-axis being the direction of the
fast-moving proton

Polarizability radii

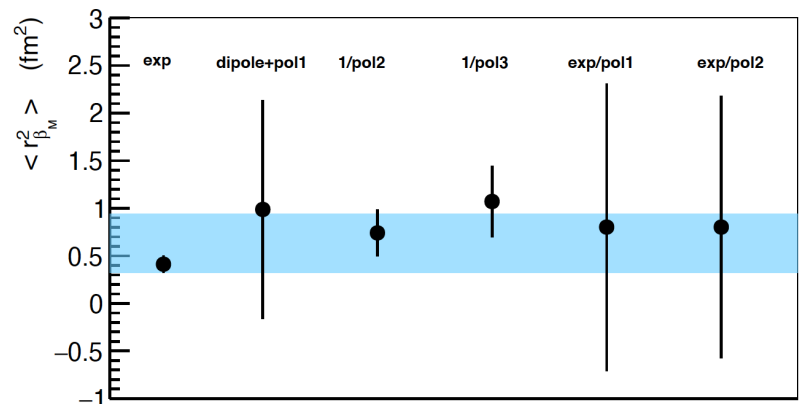
$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$



$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$



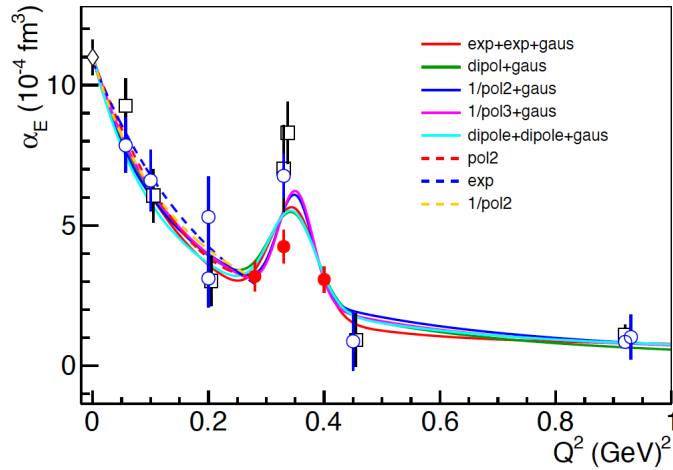
$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$



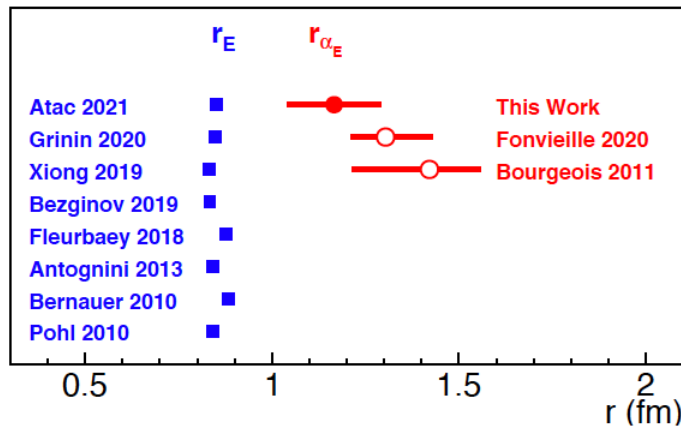
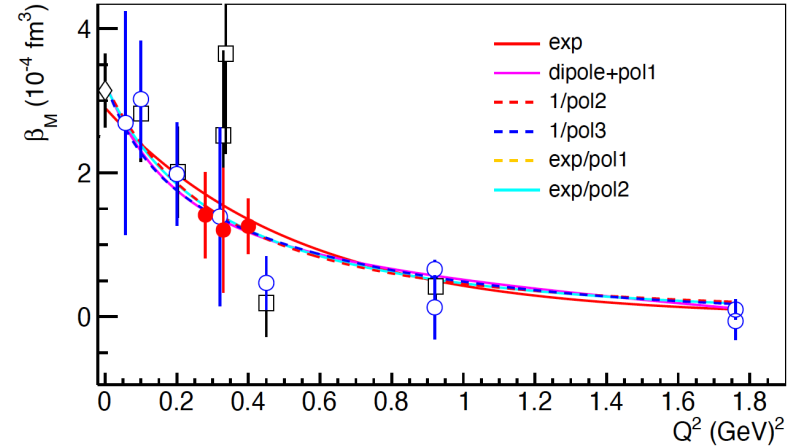
$$\langle r_{\beta_M}^2 \rangle = 0.63 \pm 0.31 \text{ fm}^2$$

Polarizability radii

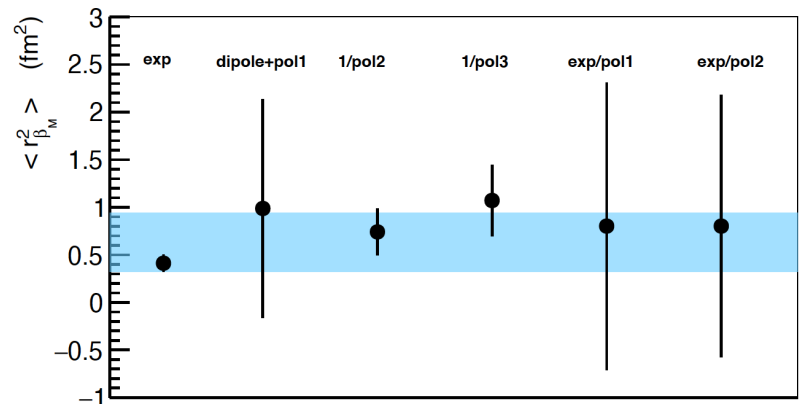
$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$



$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$

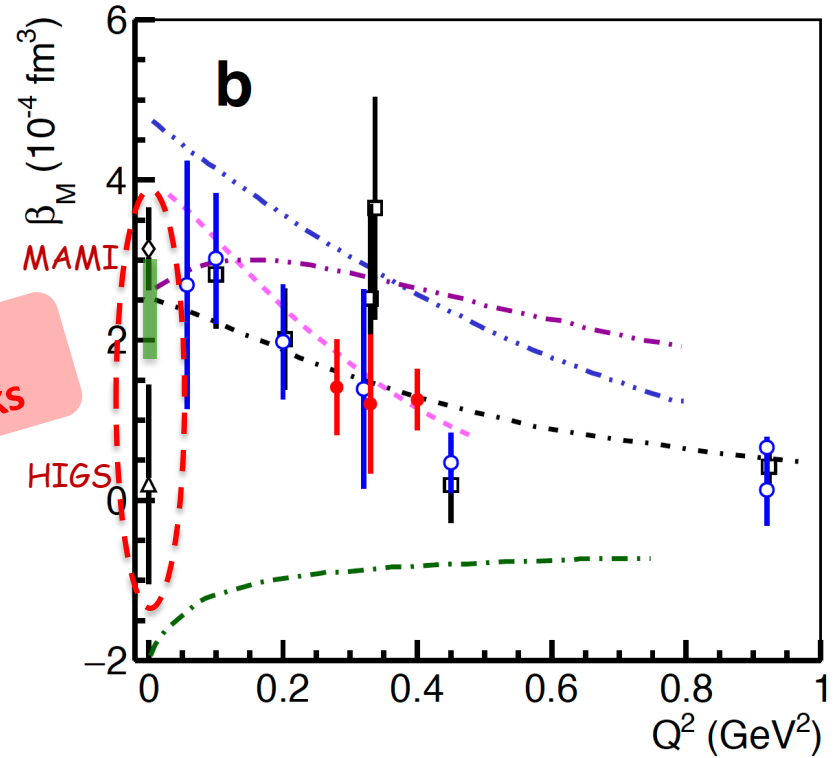
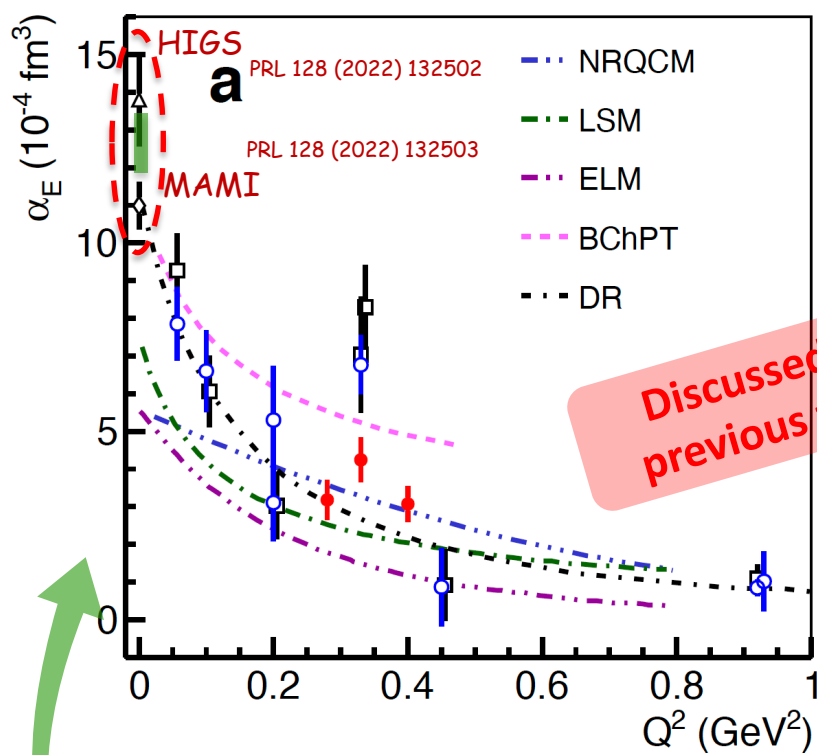


$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$



$$\langle r_{\beta_M}^2 \rangle = 0.63 \pm 0.31 \text{ fm}^2$$

Static Polarizabilities



PHYSICAL REVIEW LETTERS 129, 102501 (2022)

First Concurrent Extraction of the Leading-Order Scalar and Spin Proton Polarizabilities

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³Dipartimento di Fisica, Università degli Studi di Pavia, I-27100 Pavia, Italy

⁴INFN Sezione di Pavia, I-27100 Pavia, Italy

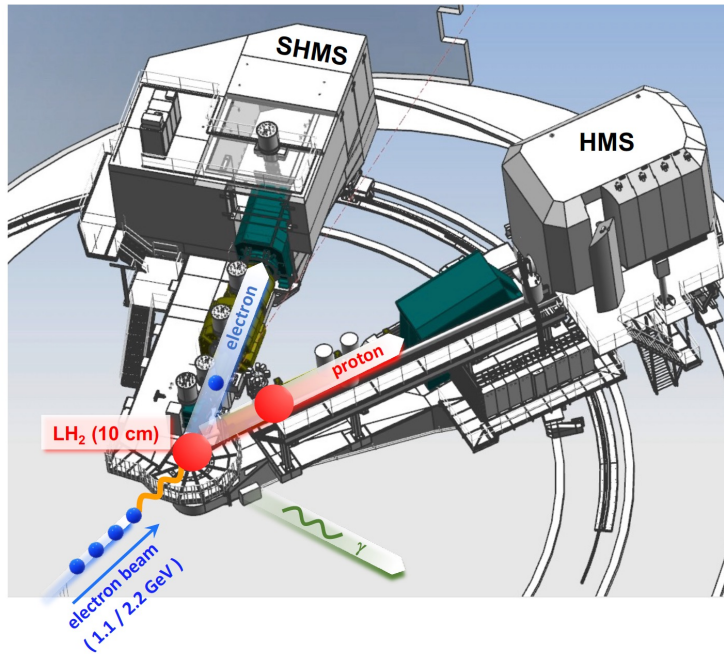
(Received 3 May 2022; revised 11 July 2022; accepted 2 August 2022; published 31 August 2022)

We performed the first simultaneous extraction of the six leading-order proton polarizabilities. We reached this milestone thanks to both new high-quality experimental data and an innovative bootstrap-based fitting method. These new results provide a self-consistent and fundamental benchmark for all future theoretical and experimental polarizability estimates.

$$\begin{aligned}\alpha_{E1} &= [12.7 \pm 0.8(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^3, \\ \beta_{M1} &= [2.4 \pm 0.6(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^3, \\ \gamma_{E1E1} &= [-3.0 \pm 0.6(\text{fit}) \pm 0.4(\text{model})] \times 10^{-4} \text{ fm}^4, \\ \gamma_{M1M1} &= [3.7 \pm 0.5(\text{fit}) \pm 0.1(\text{model})] \times 10^{-4} \text{ fm}^4, \\ \gamma_{E1M2} &= [-1.2 \pm 1.0(\text{fit}) \pm 0.3(\text{model})] \times 10^{-4} \text{ fm}^4, \\ \gamma_{M1E2} &= [2.0 \pm 0.7(\text{fit}) \pm 0.4(\text{model})] \times 10^{-4} \text{ fm}^4,\end{aligned}$$

Moving Forward

VCS-II (E12-23-001) @ JLab

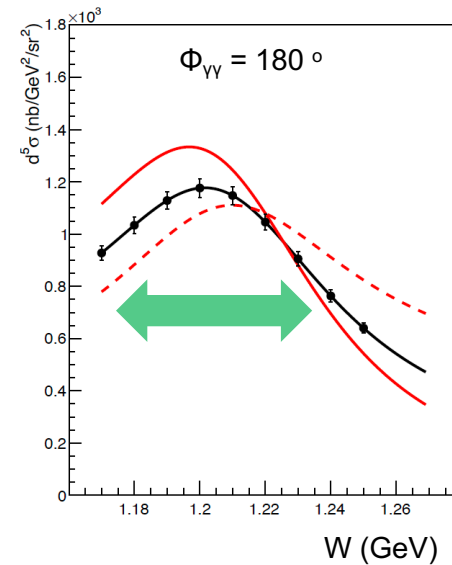
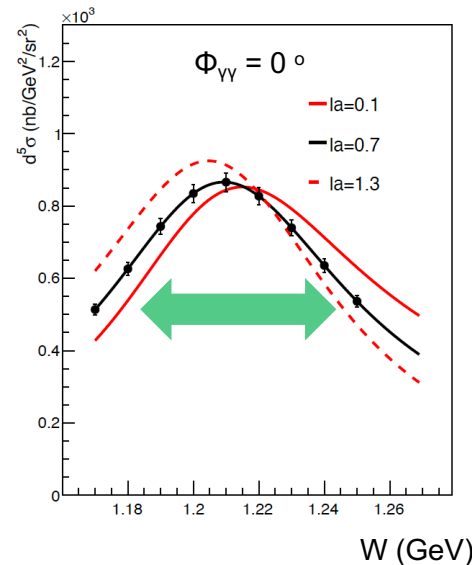
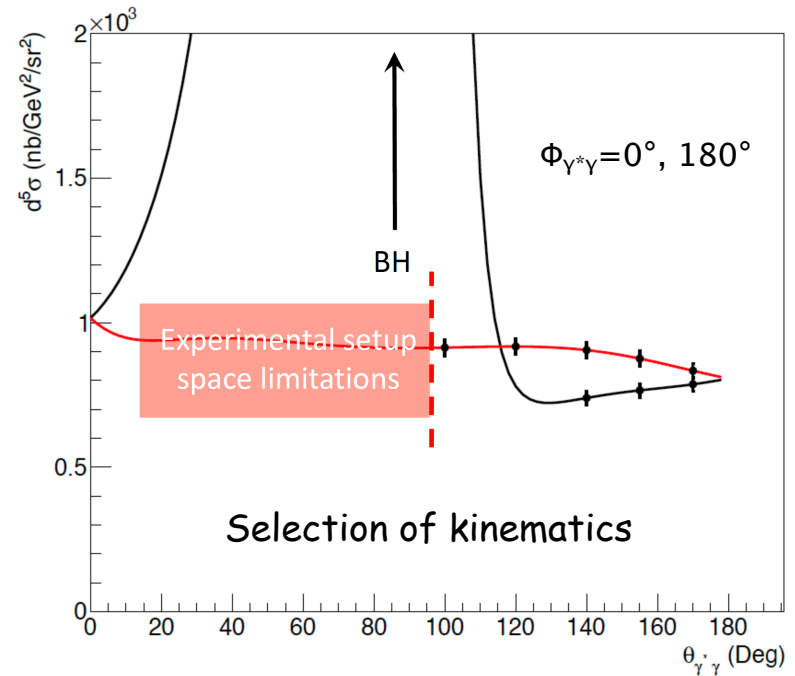


Extend Q^2 range & targeted measurements to fully exploit the sensitivity to the GPs

**APPROVED
PAC 51**

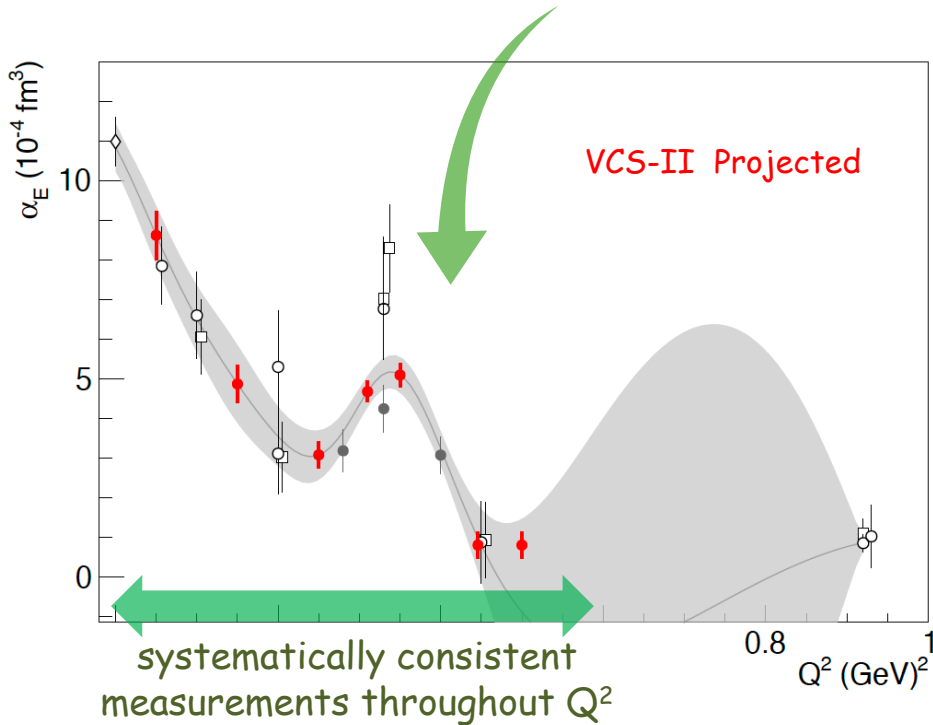
Production ($E_0 = 1.1 \text{ GeV}$) : 6 days
 Production ($E_0 = 2.2 \text{ GeV}$) : 53 days
 Studies (optics/dummy/calibrations) : 3 days

Total: 62 days

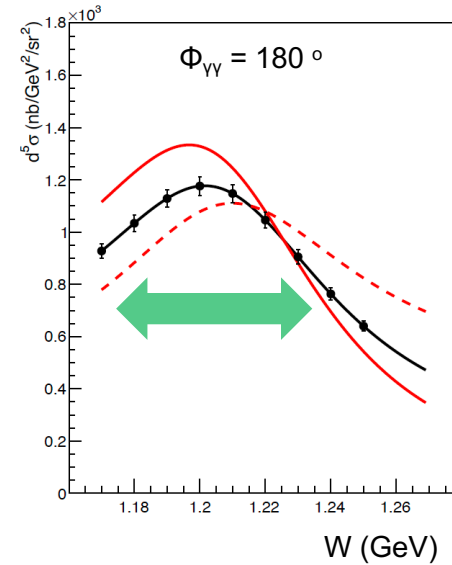
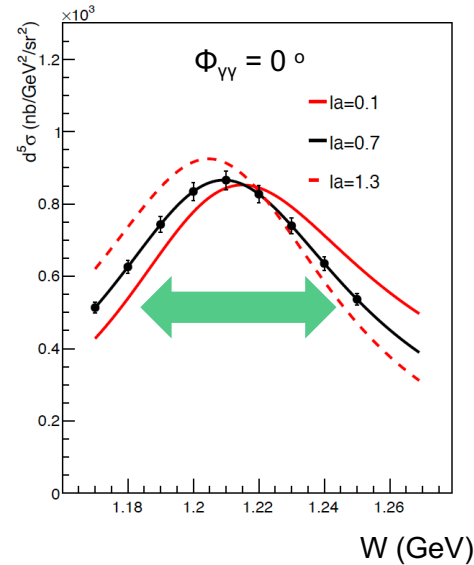


VCS-II Projected Measurements

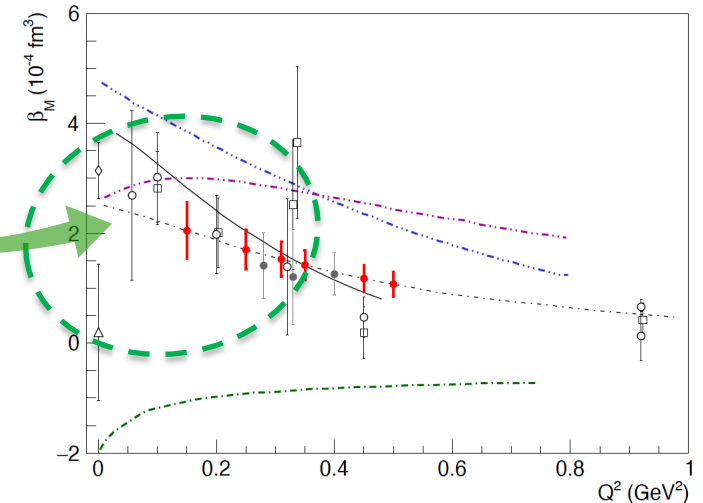
High precision measurements
combined with a fine mapping in Q^2



Targeted measurements to fully
exploit the sensitivity to the GPs



Improve upon β_M :
Pin down the competing
para/dia-magnetic contributions
in the nucleon



Can we measure with a different method ?

Yes: positrons and/or beam spin asymmetries

Positrons allow for an independent path to access experimentally the GPs

Eur. Phys. J. A 57 (2021) 11, 316

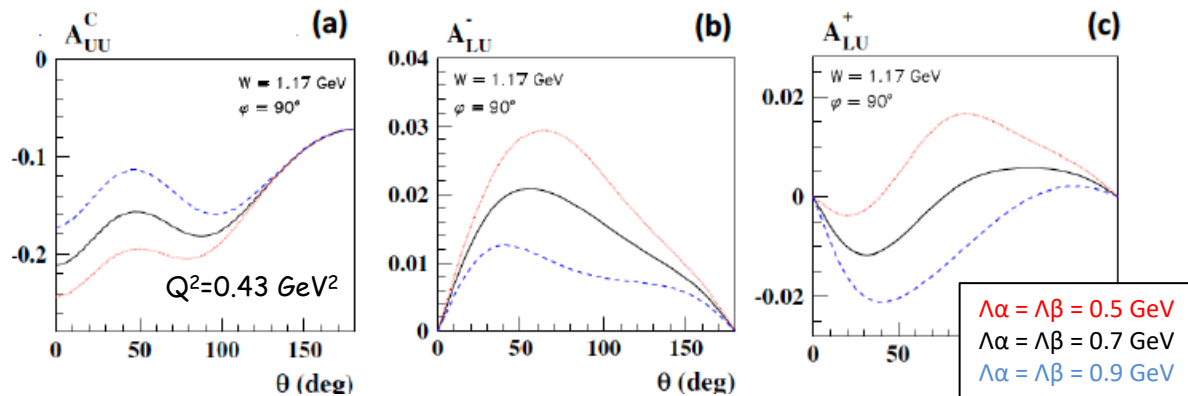
Virtual Compton scattering at low energies with a positron beam

Barbara Pasquini^{a,1,2}, Marc Vanderhaeghen^{b,3}

¹Dipartimento di Fisica, Università degli Studi di Pavia, 27100 Pavia, Italy

²Istituto Nazionale di Fisica Nucleare, Sezione di Pavia, 27100 Pavia, Italy

³Institut für Kernphysik and PRISMA⁺ Cluster of Excellence, Johannes Gutenberg Universität, D-55099 Mainz, Germany



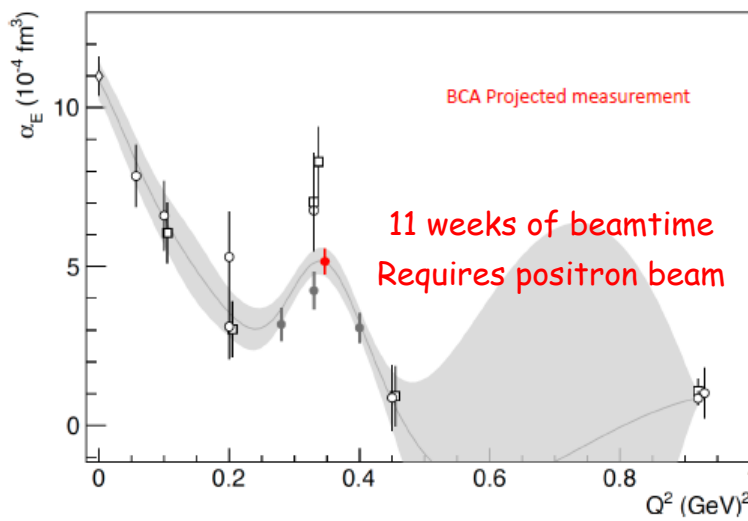
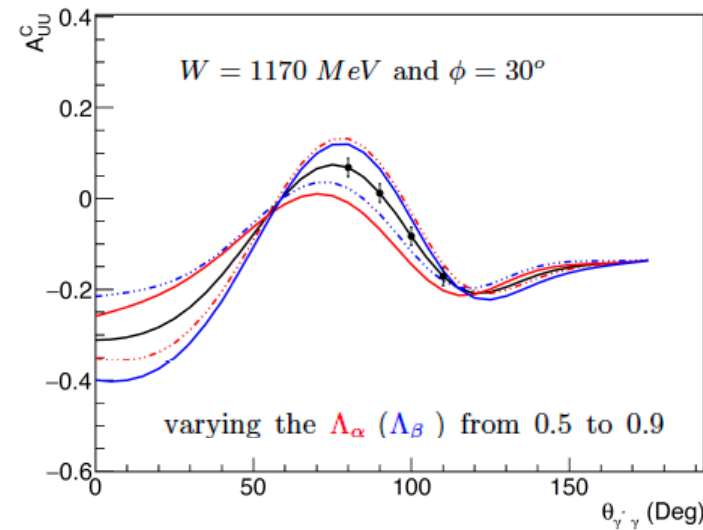
(a): The beam-charge asymmetry as a function of the photon scattering angle at $Q^2 = 0.43$ GeV².

(b) & (c): The electron and positron beam-spin asymmetry as a function of the photon scattering angle for out-of-plane kinematics.

Unpolarized beam charge asymmetry (BCA):
$$A_{UU}^C = \frac{(d\sigma_+^+ + d\sigma_-^+) - (d\sigma_+^- + d\sigma_-^-)}{d\sigma_+^+ + d\sigma_-^+ + d\sigma_+^- + d\sigma_-^-}$$

Lepton beam spin asymmetry (BSA):
$$A_{LU}^e = \frac{d\sigma_+^e - d\sigma_-^e}{d\sigma_+^e + d\sigma_-^e}$$

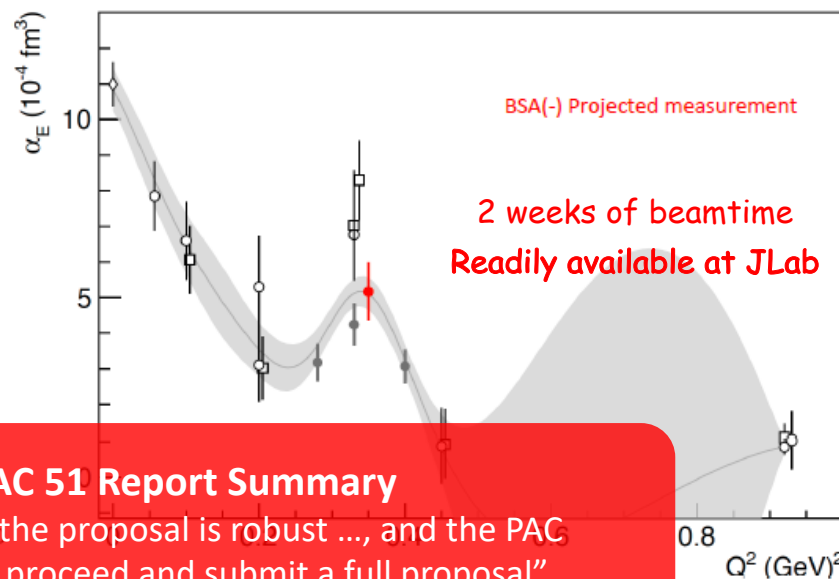
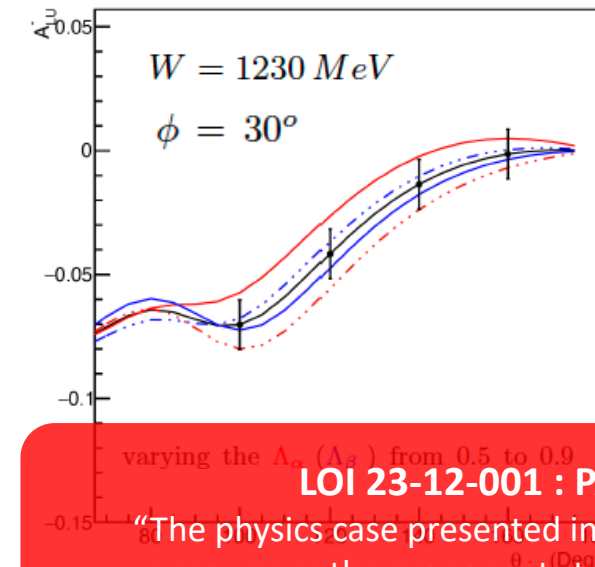
BCA (electrons & positrons)



Hall C (SHMS / HMS)

e^- : ~ 1 week @ 50 μA
and
 e^+ : ~ 10 weeks @ 5 μA

BSA (electrons or positrons)



e^- (pol. 85% @ 70 μA)
~ 2 weeks of beamtime

or

e^+ (pol. 60% @ 50 nA)
~ 3 orders of magnitude
more beamtime

LOI 23-12-001 : PAC 51 Report Summary

"The physics case presented in the proposal is robust ..., and the PAC encourages the proponents to proceed and submit a full proposal"

Summary

Progress measuring proton's fundamental properties / response to an EM field

Insight to spatial deformation of the nucleon densities under an applied EM field, interplay of para/dia-magnetic mechanisms in the proton, polarizability radii, ...

Electric GP: $\left\{ \begin{array}{l} \text{possibility for a non-trivial (non-monotonic) behavior in } a_E(Q^2) \\ \text{(albeit with a smaller magnitude than originally suggested)} \\ \text{or} \\ \text{at minimum: strong tension between world data} \end{array} \right.$

Experiment ahead of theory:

Stringent constraints to theoretical predictions / can improve further

High precision benchmark data for upcoming LQCD calculations

Future measurements:

Pin down precisely the shape of the a_E structure
(if it exists) - important input for the theory

Independent cross-check

Measure via a different channel (BS asymmetries & positrons)

Thank you!