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Hydrogen burning on Nitrogen isotopes in CNO and HCNO-cycles

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We investigated the radiative proton capture on nitrogen isotopes ^{12}N , ^{13}N , and ^{14}N [1 - 4] in the framework of the MPCM and obtained the reaction rates and their parametrizations. There are two stable nitrogen isotopes ^{14}N and ^{15}N , and among short-lived nitrogen isotopes, the longest-lived are ^{12}N ($t_{1/2} = 11$ ms) and ^{13}N ($t_{1/2} = 9.965$ min).

We compare the reaction rates to understand the relevance of each process at a given astrophysical temperature. The radiative proton $^{12}\text{N}(p, \gamma)^{13}\text{O}$, $^{13}\text{N}(p, \gamma)^{14}\text{O}$, $^{14}\text{N}(p, \gamma)^{15}\text{O}$, and $^{15}\text{N}(p, \gamma)^{16}\text{O}$ processes have the same Coulomb barrier, so the reaction rates will differ only due to the different values of the $S(E)$ and reduced mass μ of interacting particles in the entrance channel. The reduced masses of the pairs $p^{12}\text{N}$, $p^{13}\text{N}$, $p^{14}\text{N}$, and $p^{15}\text{N}$ are always less than the proton mass and are within the range $0.9294 \leq \mu(\text{amu}) \leq 0.9439$. Therefore, the impact of the reduced mass can be omitted, and the reaction rates depend entirely on the reaction S -factor [4].

The $p^{15}\text{N}$ reaction is the fastest, and its rate dominates up to $T_9 \sim 0.175$. $^{14}\text{N}(p, \gamma)^{15}\text{O}$ is the slowest process up to $T_9 \sim 0.1$, and it controls the rate and time of nucleosynthesis cycles. The $p^{15}\text{N}$ rate becomes dominant at temperature-explosive hydrogen burning scenarios in stars. Only in the temperature windows $0.18 \leq T_9 \leq 1.14$ and $0.66 \leq T_9 \leq 3$ the reaction $^{15}\text{N}(p, \gamma)^{16}\text{O}$ is slower than $^{13}\text{N}(p, \gamma)^{14}\text{O}$ and $^{14}\text{N}(p, \gamma)^{15}\text{O}$ reactions, respectively. Hence the $p^{15}\text{N}$ reaction controls the rate and time of cycles of nucleosynthesis in these two temperature windows [4].

The $^{15}\text{N}(p, \gamma)^{16}\text{O}$ reaction rate shows a strong interference effect but minor sensitivity to the asymptotic constant.

The main difficulty in determining reliable reaction rates of $^{12}\text{N}(p, \gamma)^{13}\text{O}$, $^{13}\text{N}(p, \gamma)^{14}\text{O}$, $^{14}\text{N}(p, \gamma)^{15}\text{O}$, $^{15}\text{N}(p, \gamma)^{16}\text{O}$ reactions for the CNO cycles is the uncertainty in the very low cross-sections at the Gamow range. Developments within the low-energy underground accelerator facility LUNA [5] and recent improvements in the detection setup [6] make taking direct measurements of nuclear reactions near the Gamow range feasible. This advantage has been demonstrated in the $^{14}\text{N}(p, \gamma)^{15}\text{O}$ reaction, which was successfully measured down to energies of 70 keV at LUNA [7].

References

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