

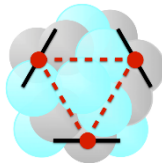
# Efficient emulator for solving 3N Faddeev equation with contact terms of chiral 3NF at N4LO

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Mainz

# Outline

Emulator for the Nd scattering:

Based on:

H.Wiśniewski et al., Few-Body Syst. 62 (2021) 23 „*Perturbative Treatment of Three Nucleon Force Contact Terms in Three-Nucleon Faddeev Equations.*”

H.Wiśniewski et al., Eur. Phys. J. A 57 (2021) 241 „*Efficient emulator for solving 3N continuum Faddeev equations with chiral 3NF comprising any number of contact terms.*”

H.Wiśniewski et al., Phys. Rev. C105 (2022) 054004 „*Significance of chiral 3NF contact terms for understanding of elastic nucleon-deuteron scattering*”

1. Formalism – new set of equations
2. Proof of concept: tests and the first results on fixing short-range 3NF parameters

# Our standard approach to Nd scattering

Prepare and solve the Faddeev equation

$$T\varphi = tP\varphi + (1+tG_0)V_{123}^{(1)}(1+P)\varphi + tPG_0T\varphi + (1+tG_0)V_{123}^{(1)}(1+P)G_0T\varphi$$

Compute amplitudes

$$U = PG_0^{-1} + V_{123}^{(1)}(1+P)\varphi + PT + V_{123}^{(1)}(1+P)G_0T$$

$$U_0 = (1+P)T$$

We work in the PWD scheme

$$3N \text{ state: } |pq\alpha\rangle \equiv |pq(ls)j(\lambda \frac{1}{2})I(jI)JM_J(t \frac{1}{2})TM_T\rangle$$

what means, that PWD of appearing in this equation operators has to be performed.

$$\langle p'q'(l's')j'(\lambda' \frac{1}{2})I'(j'I')J'M_{J'}(t' \frac{1}{2})T'M_{T'} | \hat{O} | pq(ls)j(\lambda \frac{1}{2})I(jI)JM_J(t \frac{1}{2})TM_T \rangle$$

After decomposing 3NF:

CPU time required for one run (i.e. one reaction energy, one NN+3NF potential) amounts from approx. 1-8 hrs., depending on number of partial waves, computer parameters, disk space available). Some hardware (GPU, fast memory) or software (parallelization) improvements are still possible but the cake's not worth the candle.

# Fixing parameters of 3NF

- Up to know, i.e. when working at N2LO there are only two free parameters  $c_D$  and  $c_E$ .
- Typically  $^3\text{H}$  and the  $^2a_{\text{nd}}$  or the differential Nd elastic scattering cross section at one or few energies are used.  
The latter requires solving the triton many times and the Faddeev equation 10-20 times.
- However, now we expect:
- No new 3NF free parameters at N3LO, but three new offshell LECs in the chiral NN force.
- 13 contact terms at N4LO (more precisely, due to some identities between operators, one expects in total 13 free parameters of 3NF at N4LO).
- Thus finding an efficient emulator for solving the 3N Faddeev equation seems to be essential and of high priority.

# Emulator for Nd scattering - algorithm

- The contact terms are restricted to small 3N total angular momenta and to only few partial-wave states for a given total 3N angular momentum J and parity  $\pi$

- Let us split 3NF

$$V_{123}^{(1)} = V(\theta_0) + \Delta V(\theta) \equiv V(\theta_0) + \sum_{i=1}^n c_i \Delta V_i$$

$$\theta = \{c_1, c_2, \dots\}$$

$$\theta_0 = \{0, 0, \dots\}$$

- We divide the 3N partial-wave states into two sets:
  1. The  $\beta$  set is defined by non-vanishing matrix elements of parameters dependent short-range 3NF:  $\Delta V(\theta)$ .
  2. The  $\alpha$  set comprises remaining states.

- Similarly to 3NF

$$T = T(\theta_0) + \Delta T(\theta)$$

# Emulator for Nd scattering - algorithm

- Inserting this to the Faddeev equation leads to sets of equations with one equation for  $T(\theta_0)$  which is a standard Faddeev equation but with  $V(\theta_0)$  and

$$\langle \alpha | \Delta T(\theta) | \phi \rangle = \langle \alpha | t P G_0 \Delta T(\theta) | \phi \rangle + \langle \alpha | (1 + t G_0) V(\theta_0) (1 + P) G_0 \Delta T(\theta) | \phi \rangle$$

$$\begin{aligned} \langle \beta | \Delta T(\theta) | \phi \rangle = & \langle \beta | (1 + t G_0) \Delta V(\theta) (1 + P) | \phi \rangle + \langle \beta | (1 + t G_0) \Delta V(\theta) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ & + \langle \beta | (1 + t G_0) [V(\theta_0) + \Delta V(\theta)] (1 + P) G_0 \Delta T(\theta) | \phi \rangle \\ & + \langle \beta | t P G_0 \Delta T(\theta) | \phi \rangle . \end{aligned}$$

- We neglect term  $\sim \Delta V \Delta T$ , which allows to separate contributions from different  $\Delta V_i$  and leads to set of equations for a corresponding  $\Delta T_i$ .  
Next, for single parameter dependent component of  $V$ :  $V_i = c_i V$  we may solve that equation separately (at  $c_i = 1$ ) obtaining corresponding  $\langle \beta | \Delta T_i | \phi \rangle$  and build

$$\langle \beta | \Delta T(\theta) | \phi \rangle = \sum_{i=1}^N c_i \langle \beta | \Delta T_i | \phi \rangle$$

and find also  $\langle \alpha | \Delta T_i | \phi \rangle$ .

# Emulator for Nd scattering - algorithm

- In this way we have matrix elements of T

$$\begin{aligned}\langle\alpha|T(\theta)|\phi\rangle &= \langle\alpha|T(\theta_0)|\phi\rangle + \sum_i c_i \langle\alpha|\Delta T_i|\phi\rangle, \\ \langle\beta|T(\theta)|\phi\rangle &= \langle\beta|T(\theta_0)|\phi\rangle + \sum_i c_i \langle\beta|\Delta T_i|\phi\rangle.\end{aligned}\quad (10)$$

- Let us now come back to the scattering amplitudes

$$U = P G_0^{-1} + V_{123}^{(1)}(1+P)\phi + P T + V_{123}^{(1)}(1+P)G_0 T$$

$$U_0 = (1+P)T$$

- They are linear in T: the dependence on the  $c_i$  parameters carries over to them

$$U = U(\theta_0) + \sum_i c_i U_i + \sum_{i,k} c_i c_k U_{ik}$$

$$U_0 = U_0(\theta_0) + \sum_i c_i U_{0i}$$

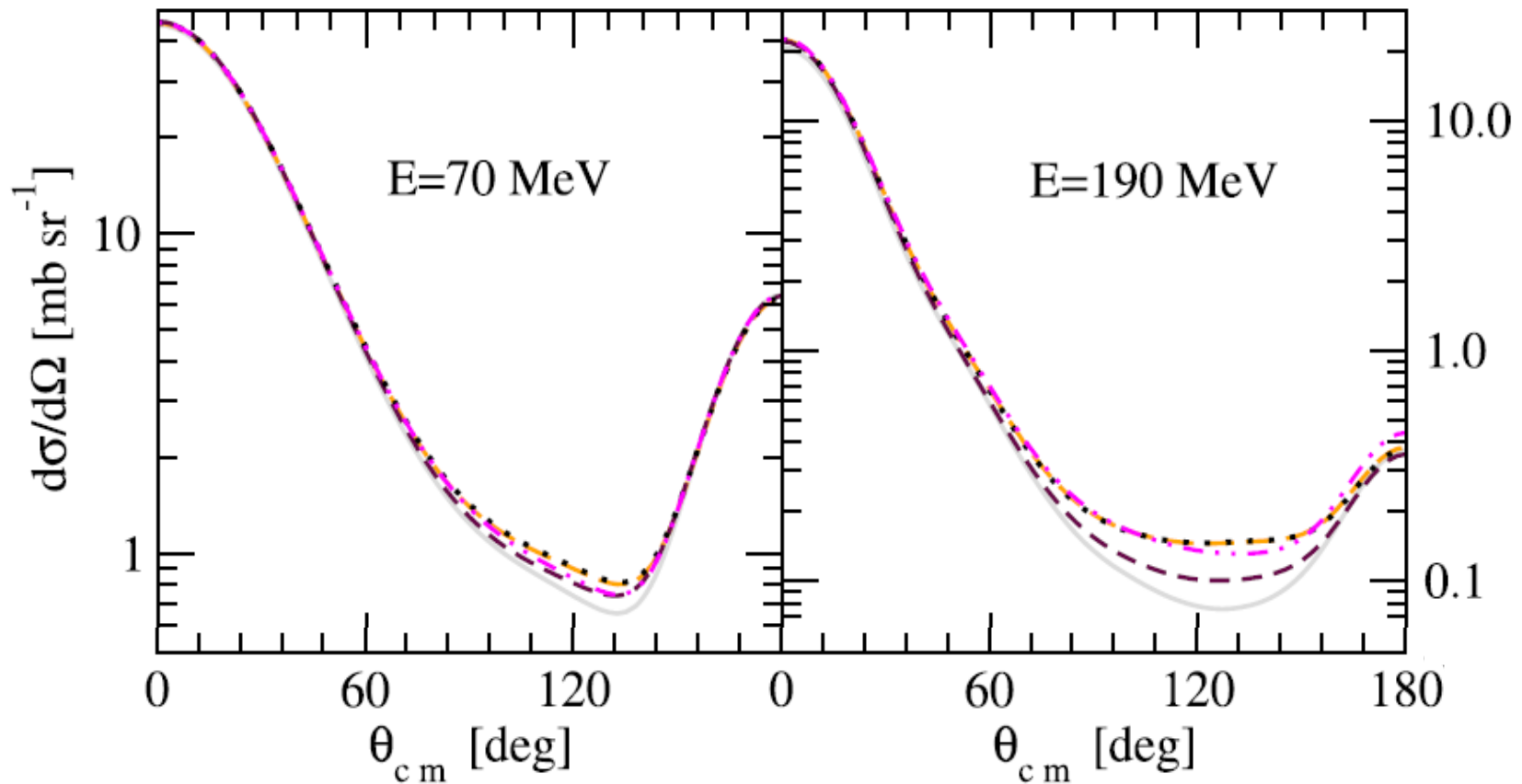
- **Summarizing: one needs to solve N+1 Faddeev equations (one for  $T(\theta_0)$  and N for  $\langle\beta|\Delta T_i|\phi\rangle$ ) and next build transition amplitudes for any set of  $c_i$ .**

# Emulator for Nd scattering – algorithm - application

- We used SMS N4LO+ NN potential at  $\Lambda=450$  MeV, combined with the N2LO chiral 3NF and supplemented by all subleading N4LO 3NF contact terms from:
  1. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 84, 014001 (2011).,
  2. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 102, 019903(E) (2020).
- All terms are regulated with the non-local regulator.
- Such a Hamiltonian comprises altogether 15 short-range contributions to 3NF, two from N2LO with the strengths  $c_D$  and  $c_E$ , and thirteen from N4LO with the strengths  $E_i$ ,  $i = 1, \dots, 13$ . However, for two pairs of the  $E_i$  terms matrix elements are identical, thus finally there are 13 unknown parameters.



# Emulator for Nd scattering – test



Exact:

NN N4LO+

NN N4LO+ + 3NF N2LO ( $c_D=c_E=c_{E_7}=0.0$ )

NN N4LO+ + 3NF N2LO+E<sub>7</sub>

( $c_D = -8.2053$ ,  $c_E = -1.0019$ ,  $c_{E_7} = 2.0$ )

Emulator:

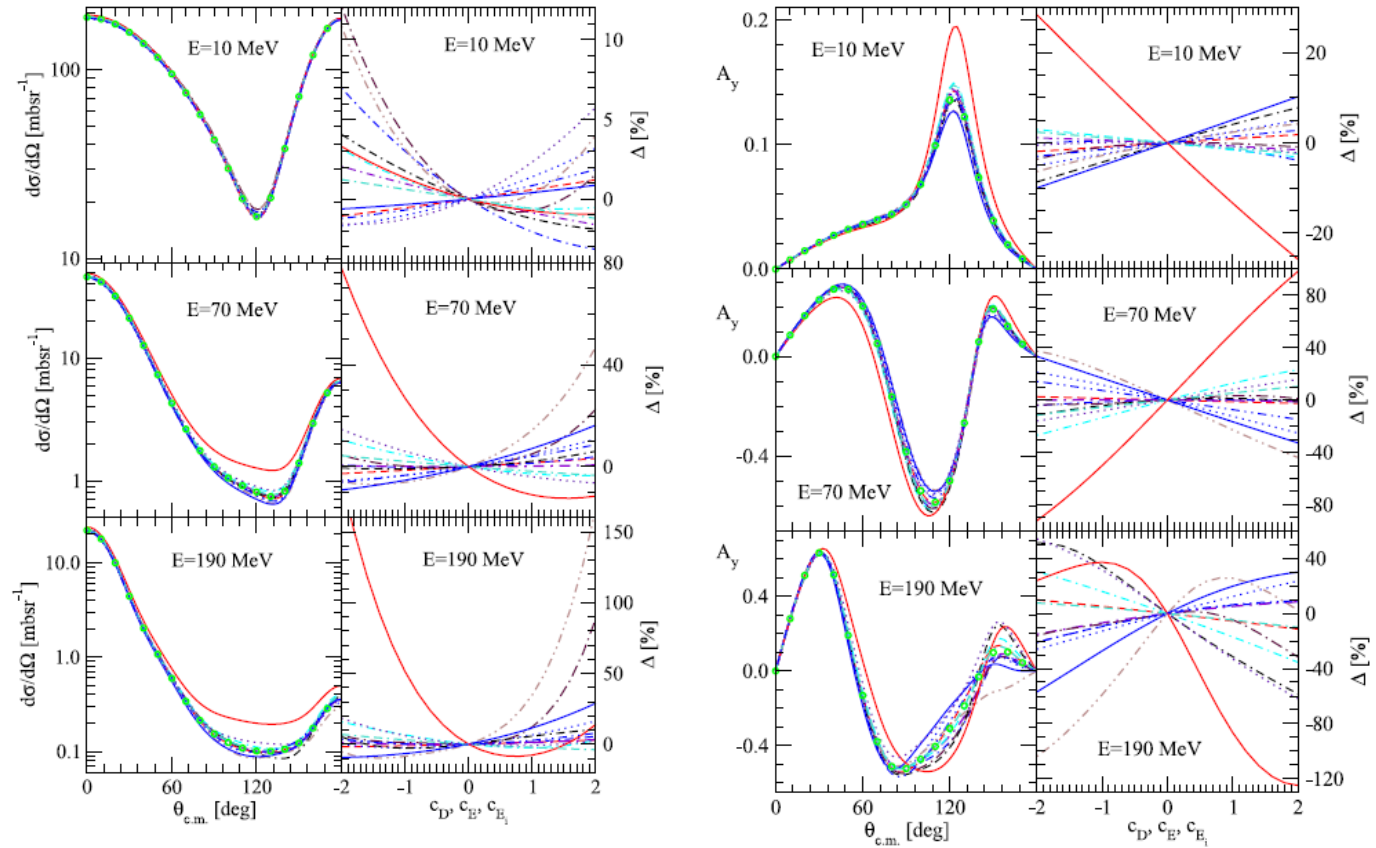
NN N4LO+ + 3NF N2LO+E<sub>7</sub> ( $\beta = {}^1S_0, {}^3S_1-{}^3D_1$ )

NN N4LO+ + 3NF N2LO+E<sub>7</sub> ( $\beta = j \leq 2$ )

# Emulator for Nd scattering – algorithm - application

Sensitivity of 3N scattering observables to  $E_i$  terms

Green circles  $V(\theta_0)$   
 $C_i = -1$  (left)  
 red solid  $E_8$ ,  
 blue solid  $E_7$ ,  
 brown dashed-  
 double dotted  $E_5$



- N2LO D and E terms do not dominate
- Some observables are more sensitive to specific terms, e.g.  $T_{22}$  to  $E_{10}$

$$\Delta \equiv \Delta(c_i) = \frac{1}{N_\theta} \sum_{\theta_k} \frac{Obs(c_i, \theta_k) - Obs(\theta_0, \theta_k)}{Obs(\theta_0, \theta_k)}$$

# Emulator for Nd scattering – application

## $V_i$ expectation values in ${}^3\text{H}$ at $c_i=1.0$

TABLE I. Contributions of the  $\text{N}^2\text{LO}$  and  $\text{N}^4\text{LO}$  contact terms to the potential energy of the three nucleons in the triton. These expectation values were obtained for the  ${}^3\text{H}$  wave function calculated with the SMS chiral  $\text{N}^4\text{LO}^+$   $NN$  potential ( $\Lambda = 450$  MeV) and assuming strengths of contact terms  $c_i = 1.0$ .

$V_i$	$\langle \psi_{3H}   V_i   \psi_{3H} \rangle$ [MeV]
$V_D$	0.1661
$V_E$	-1.4294
$V_{E1}$	0.3463
$V_{E2}$	-0.4173
$V_{E3}$	-0.2754
$V_{E4}$	-1.0390
$V_{E5}$	-0.9559
$V_{E6}$	-1.0699
$V_{E7}$	$0.1798 \times 10^{-4}$
$V_{E8}$	$0.8817 \times 10^{-2}$
$V_{E9}$	-0.2407
$V_{E10}$	1.0571
$V_{E11}$	-0.2407
$V_{E12}$	1.0571
$V_{E13}$	0.3060

Relative strengths

Nearly all terms are important (for  $c_i=1$ ) with exception of  $E_7$  and  $E_8$  terms

# Emulator for Nd scattering – fit to the true data at 10,70, and 135 MeV (786 data points)

TABLE III. The values of strengths  $c_i$  found in the least squares fit to the data from Table II at the three energies  $E = 10, 70$ , and 135 MeV.

$c_D$	$-1.49 \pm 0.06$
$c_E$	$-1.27 \pm 0.06$
$c_{E_1}$	$6.40 \pm 0.33$
$c_{E_2}$	$7.80 \pm 0.36$
$c_{E_3}$	$6.97 \pm 0.34$
$c_{E_4}$	$-2.06 \pm 0.13$
$c_{E_5}$	$-0.36 \pm 0.05$
$c_{E_6}$	$0.52 \pm 0.03$
$c_{E_7}$	$-7.40 \pm 0.14$
$c_{E_8}$	$-2.61 \pm 0.05$
$c_{E_9}$	$-4.59 \pm 0.22$
$c_{E_{10}}$	$-0.98 \pm 0.05$
$c_{E_{13}}$	$-1.14 \pm 0.05$

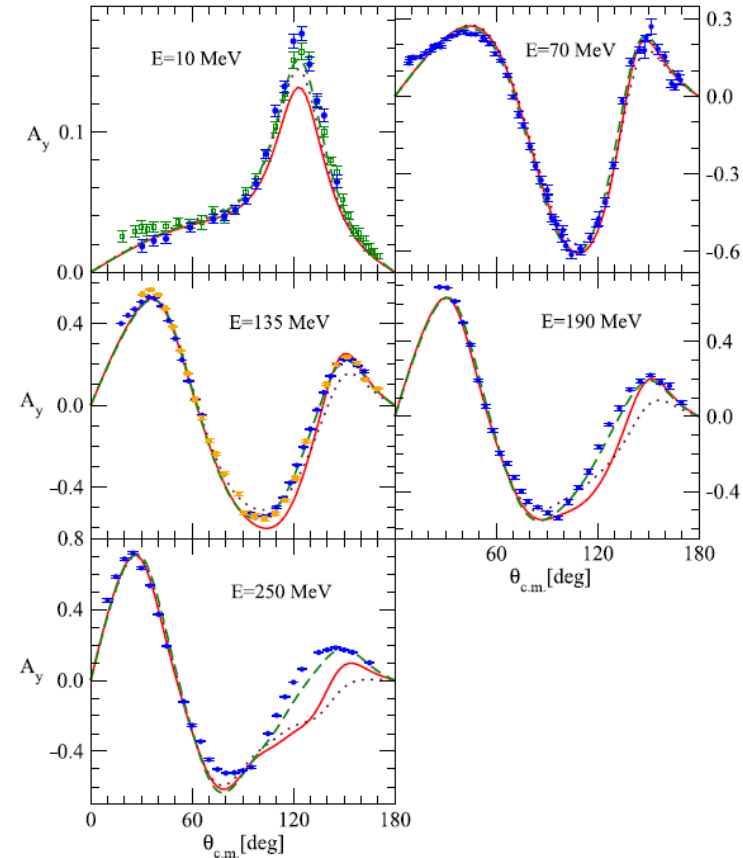
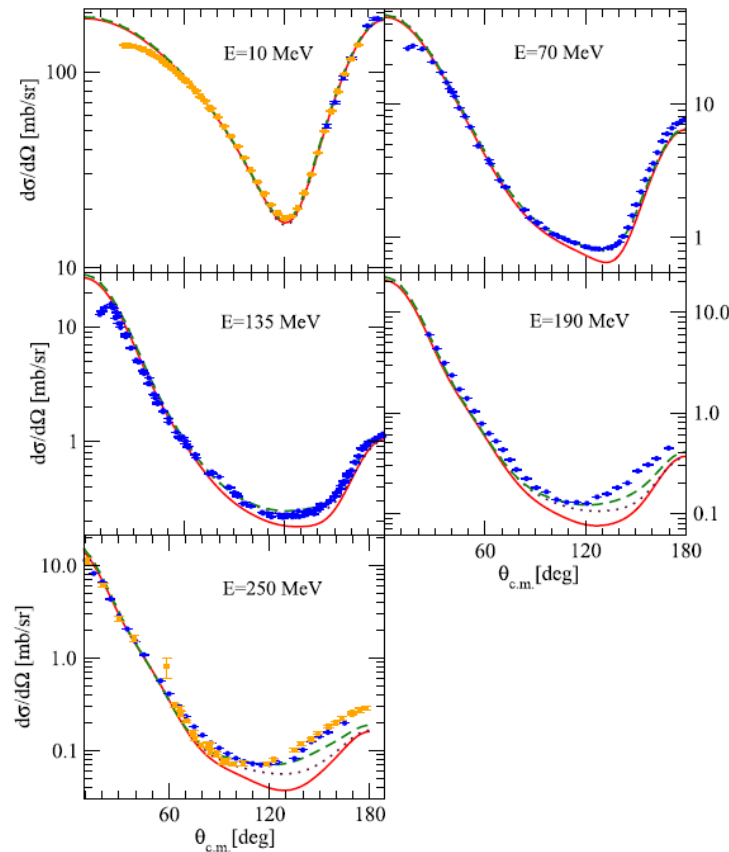
TABLE IV. The covariance matrix for the strengths  $c_i$  determined by the least squares fit of data from Table II at the three energies  $E = 10, 70$ , and 135 MeV [the values shown are  $\text{Cov}(c_i, c_j) \times 1000$ ].

	$c_D$	$c_E$	$c_{E_1}$	$c_{E_2}$	$c_{E_3}$	$c_{E_4}$	$c_{E_5}$	$c_{E_6}$	$c_{E_7}$	$c_{E_8}$	$c_{E_9}$	$c_{E_{10}}$	$c_{E_{13}}$
$c_D$	3.914	-0.456	1.412	4.573	0.843	0.844	-0.729	-0.892	1.109	0.267	-0.726	0.123	-0.207
$c_E$		3.560	0.947	-3.571	1.345	-0.633	-0.172	-0.217	-2.416	-0.809	-1.702	0.393	0.571
$c_{E_1}$			108.9	112.8	108.9	-35.13	1.409	-2.418	25.92	7.513	12.99	3.861	0.443
$c_{E_2}$				130.7	113.4	-35.15	-1.995	-3.241	32.43	9.561	-0.534	0.763	-3.332
$c_{E_3}$					112.9	-38.92	1.617	-1.814	27.52	8.068	8.366	1.598	-0.193
$c_{E_4}$						15.97	-1.966	-0.362	-10.50	-3.198	-4.866	0.345	-0.222
$c_{E_5}$							2.415	0.669	0.791	0.281	9.892	1.311	1.766
$c_{E_6}$								0.635	-0.874	-0.226	1.426	-0.226	0.210
$c_{E_7}$									20.33	6.455	3.464	-0.324	-1.463
$c_{E_8}$										2.071	1.041	-0.158	-0.462
$c_{E_9}$											50.23	9.133	8.813
$c_{E_{10}}$												2.625	1.910
$c_{E_{13}}$													2.499

- Big values of  $c_{E_1}, c_{E_2}, c_{E_3}, c_{E_7}, c_{E_9}$
- Correlation coefficients close to  $\pm 1$ :  $\rho(E_1, E_2)$ ,  $\rho(E_2, E_3)$ ,  $\rho(E_1, E_3)$ ,  $\rho(E_3, E_4)$ ,  $\rho(E_7, E_8)$
- Correlation coefficients close to 0:  $(c_D, c_E), (c_D, c_{E_i}), (c_E, c_{E_i})$
- $\chi^2/\text{data} \approx 35$

# Emulator for Nd scattering – fit to the data: cross section and $A_Y(N)$

- Data at 10, 70 and 135 MeV
- Results at 190 and 250 MeV are predictions



NN N4LO+

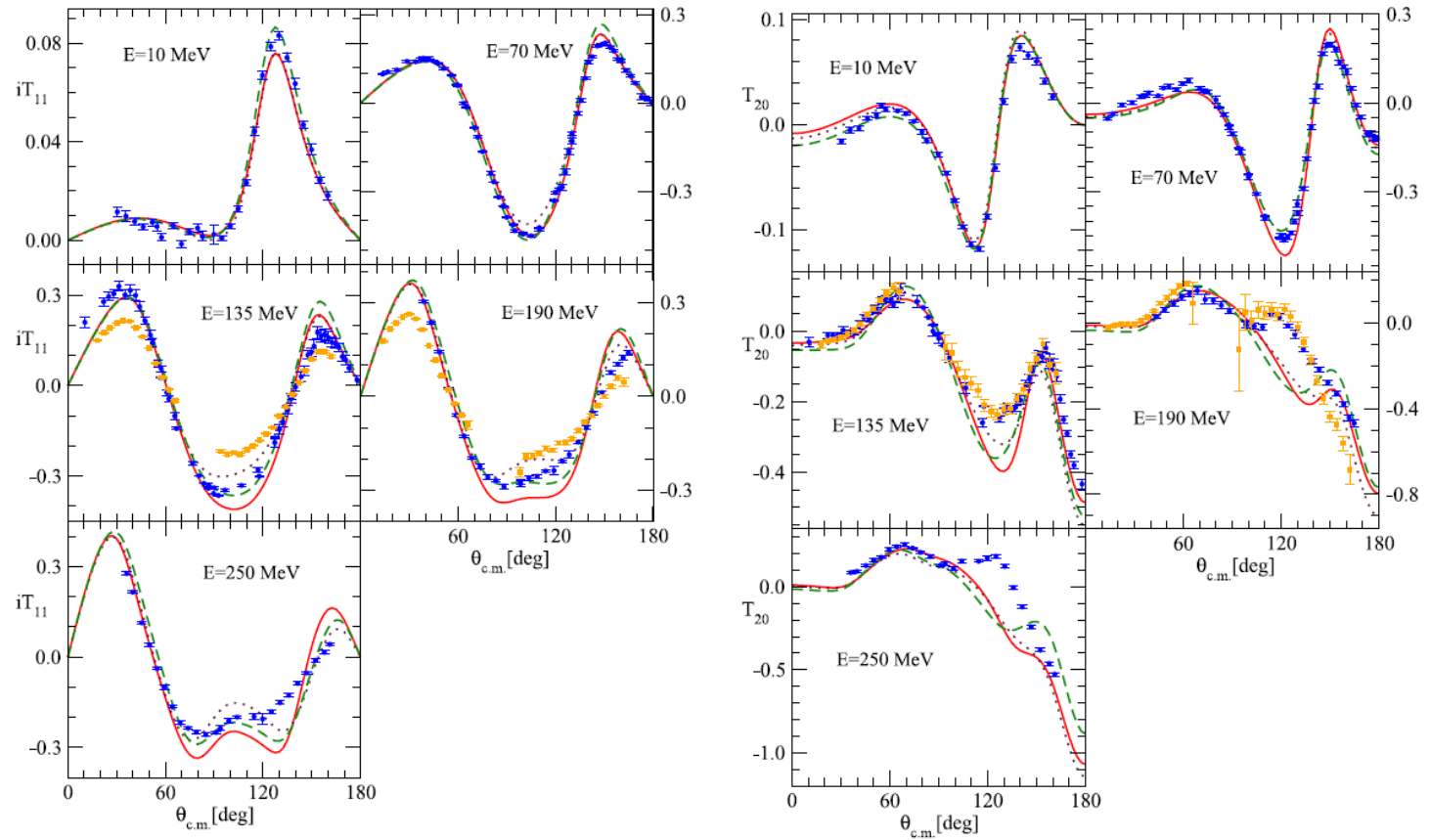
NN N4LO+ + 3NF N2LO

NN N4LO+ + 3NF N2LO +  $E_i$

# Emulator for Nd scattering – fit to the data: $iT_{11}$ and $T_{20}$

■ Data at 10, 70 and 135 MeV

■ Results at 190 and 250 MeV are predictions



NN N4LO+

NN N4LO+ + 3NF N2LO

NN N4LO+ + 3NF N2LO +  $E_i$

# Summary

- We constructed and tested an efficient and accurate emulator for solving 3N Faddeev equation.
- We applied it to the Nd scattering up to  $E=250$  MeV, using the chiral SMS NN potential at N4LO+ supplemented by 3NF at N2LO and 13 N4LO contact terms.
- Our emulator allows us to fix free parameters of all short-range terms in the 3NF. We found that even at low energies some observables are sensitive to N4LO 3NF contact terms.
- In general, sensitivity of predictions to N4LO 3NF contact terms depends on observable, energy and scattering angle.
- Usually we observe improvements in data description, but very likely above  $\approx 200$  MeV 3NF is not sufficient to explain discrepancies with the data.
- The deuteron breakup data can be used in fitting as well.
- Coulomb correction (if needed) and 3NF at N3LO has to be included for final conclusions.