

# The Three-Proton Correlation function

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# Introduction

- ▶ Recently the **nnn** and **ppp** systems have attracted a particular interest
- ▶ In the two-body case, the **nn** system is much less known than the **pp** (or **np**) system
- ▶ In *s*-wave they interact in the  $S = 0, T = 1$  channel.
- ▶ The associate scattering lengths are negative and large indicating the presence of a virtual state:  
 $a_{nn} \approx -18.5 \text{ fm}, a_{pp}^{sr} \approx -17.5 \text{ fm}, a_{np} \approx -23.7 \text{ fm}$
- ▶ The effective ranges are all similar:  
 $r_{nn} \approx r_{pp} \approx r_{np} \approx 2.8 \text{ fm}$
- ▶ The **nnn** and **nnnn** have been subject of intense investigations to determine if they have a low energy resonance state
- ▶ Here I discuss the **3**  $\rightarrow$  **3** scattering: nnn and ppp

# The pp correlation function

- ▶ Experimentally, the pp correlation function is defined as

$$C_{pp}(k) = \xi(k) \otimes \frac{N_{\text{same}}(k)}{N_{\text{mixed}}(k)}$$

- ▶  $\xi(k) \rightarrow$  corrections for experimental effects
- ▶  $N_{\text{same}}(k) \rightarrow$  number of detected particle pairs from the same event
- ▶  $N_{\text{mixed}}(k) \rightarrow$  number of uncorrelated pairs, the so-called mixed-event technique
  
- ▶ Theoretically the definition of the correlation function is

$$C_{pp}(k) = \int d^3y S_R(y) |\Psi|^2$$

# The pp correlation function

- ▶ The theoretical ingredients

$$C_{pp}(k) = \int d^3y S_R(y) |\Psi|^2$$

with  $S_R$  the source function defined as

$$S_R(y) = \frac{1}{(4\pi R^2)^{3/2}} e^{-(y/2R)^2}$$

and  $\Psi$  the pp scattering wave function

$$\Psi = \sum_{[LSJ]} u_{LSJ}(y) [Y_L(\hat{y}) \chi_S]_J = \Psi^0 + \sum_{[LSJ]}^{\bar{J}} \Psi_{LSJ}$$

with  $\Psi^0$  the free scattering wave function. In  $\Psi_{LSJ}$  the interaction has been considered up to  $\bar{J}$ .

## The pp free wave function

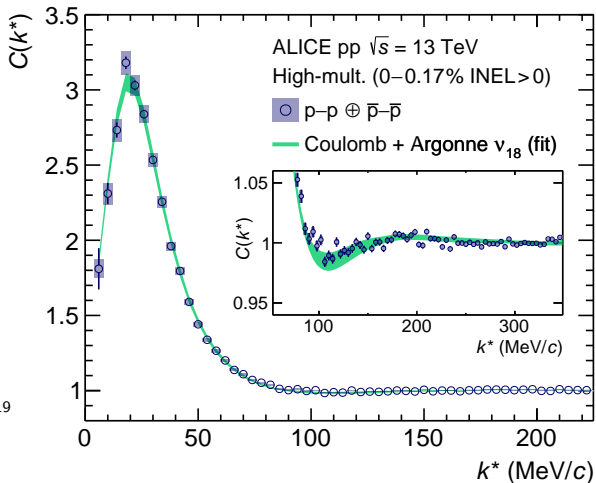
In the case of two protons, the scattering wave function has to be expanded in terms of Coulomb functions

$$\Psi_s^0 = 4\pi \sum_{[\ell S]} i^\ell (kr)^{-1} F_\ell(\eta, kr) \mathcal{Y}_{[\ell S]}(\hat{r}) \mathcal{Y}_{[\ell S]}^*(\hat{k})$$

with  $F_\ell(\eta, kr)$  a regular Coulomb wave function and  $\eta = e^2/(\hbar^2 k/m)$ . The norm results

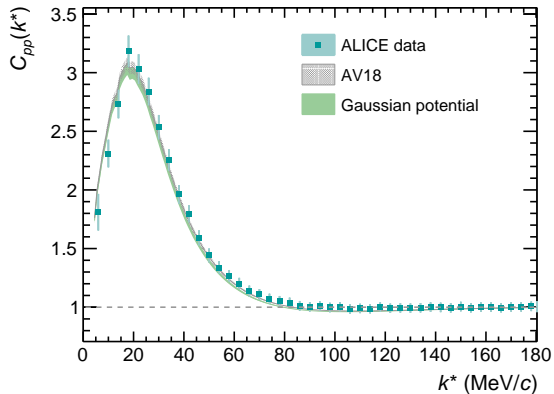
$$|\Psi_s^0|^2 = \frac{1}{2} \sum_{\ell \text{ even}} (kr)^{-2} F_\ell^2(\eta, kr)(2\ell + 1) + \frac{3}{2} \sum_{\ell \text{ odd}} (kr)^{-2} F_\ell^2(\eta, kr)(2\ell + 1)$$

# The pp Correlation Function



ALICE collaboration  
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# The pp Correlation Function



Using a gaussian potential:  $V_{pp}(r) = V_0 e^{-(r/r_0)^2} \mathcal{P}_0 + \frac{e^2}{r}$

# The pd Correlation Function

- ▶ We now consider the pd correlation function:

$$A_d C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int d^3 r_1 d^3 r_2 d^3 r_3 S_1(r_1) S_1(r_2) S_1(r_3) |\Psi_{m_2, m_1}|^2$$

- ▶ the probability of deuteron formation

$$A_d = \frac{1}{3} \sum_{m_2} \int d^3 r_1 d^3 r_2 S_1(r_1) S_1(r_2) |\phi_{m_2}|^2$$

- ▶ the single particle source function

$$S_1(r) = \frac{1}{(2\pi R_M^2)^{\frac{3}{2}}} e^{-r^2/2R_M^2}$$

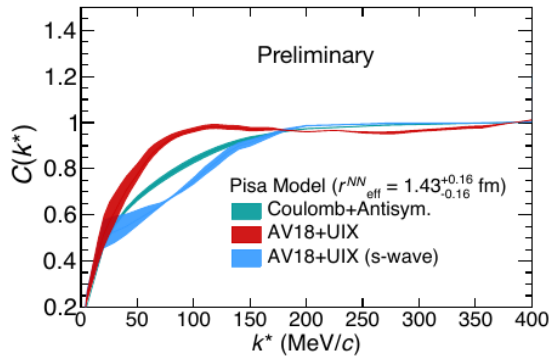


# The pd Correlation Function

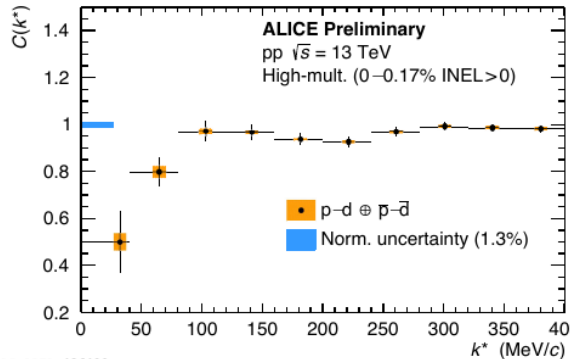
- ▶ the pd correlation function results

$$A_d C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int \rho^5 d\rho d\Omega \frac{e^{-\rho^2/4R_M^2}}{(4\pi R_M^2)^3} |\Psi_{m_2, m_1, \vec{k}}^{pd}|^2$$

$$\Psi_{m_2, m_1, \vec{k}}^{pd} = \sum_{LSJ} \sqrt{4\pi} i^L \sqrt{2L+1} e^{i\sigma_L} (1m_2 \frac{1}{2} m_1 | SJ_z)(L0SJ_z | JJ_z) \Psi_{LSJJ_z} ,$$



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M. Viviani, S. König, A. Kievsky, L.E. Marcucci, B. Singh, O. Vázquez Doce, arXiv:2306.02478 [nucl-th]  
ALICE collaboration+theoreticians, in preparation

## The ppp correlation function

- ▶ Now we consider the ppp correlation function:

$$C_{ppp}(Q) = \int \rho^5 d\rho d\Omega S_{\rho_0}(\rho) |\Psi_{ppp}|^2$$

with  $Q$  the hyper-momentum,  $S_{\rho_0}$  the source function defined as

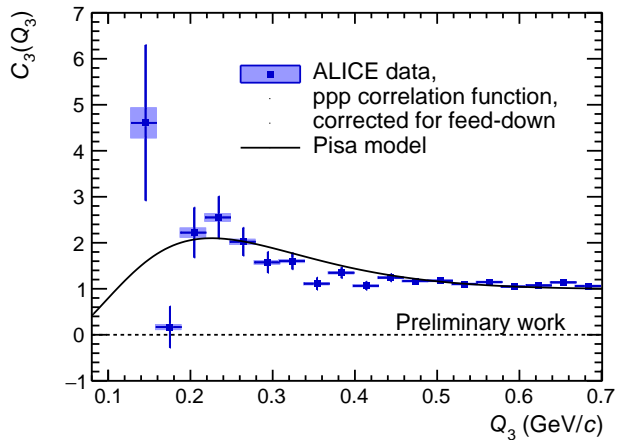
$$S_{\rho_0}(\rho) = \frac{1}{\pi^3 \rho_0^6} e^{-(\rho/\rho_0)^2}$$

$\Psi_{ppp}$  is the ppp scattering wave function

$$\Psi_{ppp} = \sum_{[K]} u_{[K]}(\rho) \mathcal{B}_{[K]}(\Omega) = \Psi^0 + \sum_{J, [K]}^{\bar{J}, \bar{K}} \Psi_{[K]}^J$$

To be noticed that  $\Psi^0$  is not well known. In  $\Psi_{[K]}^J$  the interaction has been considered up to  $\bar{J}$  and  $\bar{K}$

## Comparison to data (preliminary)



## Some remarks

- ▶ In the pp correlation function the main ingredient is the  $\Psi_{pp}$  scattering function. Accordingly  $C_{pp}(k)$  is sensitive to the NN interaction.
- ▶ In the pd and ppp correlation functions the main ingredients are the  $\Psi_{pd}$  and  $\Psi_{ppp}$  scattering functions. Accordingly  $C_{pd}(k)$  and  $C_{ppp}(Q)$  should be sensitive to the NN interaction (and NNN interaction).
- ▶ The ppp wave function is an open problem due to long-range coulomb interaction

# The ppp Wave Function

The total wave function is

$$\Psi(\vec{x}, \vec{y}) = \sum_i \psi(\vec{x}_i, \vec{y}_i) = \rho^{-5/2} \sum_{[K]} u_{[K]}(\rho) \mathcal{B}_{[K]}^{J\pi}(\Omega)$$

with  $\mathcal{B}_{[K]}^{J\pi}$  antisymmetric HH-spin functions

The ppp wave is completely determined from the hyperradial functions  $u_{[K]}(\rho)$ . And they are determined from the boundary conditions as  $\rho \rightarrow \infty$ .

For a given energy,  $E = \hbar^2 Q^2 / m$ , **and in the nnn case**

$$u_{[K]}(\rho \rightarrow \infty) \rightarrow \sqrt{Q\rho} [J_{K+2}(Q\rho) + \tan \delta_K Y_{K+2}(Q\rho)]$$

In the ppp case the asymptotic equations are coupled not allowing this simple picture

## ppp Correlation Analysis

Using the property of the HH functions

$$\Psi_s^0 = e^{i\vec{Q}\cdot\vec{\rho}} = \frac{(2\pi)^3}{(Q\rho)^2} \sum_{[K]} i^K J_{K+2}(Q\rho) \mathcal{Y}_{[K]}(\Omega) \mathcal{Y}_{[K]}^*(\hat{Q})$$

where  $\vec{Q} \cdot \vec{\rho} = \vec{k}_1 \cdot \vec{x} + \vec{k}_2 \cdot \vec{y}$  and  $J_{K+2}$  a Bessel function.

- ▶ For the case of three nucleons we have to include the correct symmetrization.
- ▶ For the case of three protons we have to include the correct asymptotics

The nnn case:

$$\Psi_s^0 = \frac{(2\pi)^3}{(Q\rho)^2} \sum_{[K]} i^K J_{K+2}(Q\rho) \mathcal{B}_{[K]}(\Omega) \mathcal{B}_{[K]}^*(\hat{Q})$$

with  $\mathcal{B}_{[K]}(\Omega)$  antisymmetric in the hyperangle-spin space.

## ppp Correlation Analysis

For three protons the asymptotic form changes (and it is not known in a close form)  
The Coulomb interaction coupled the asymptotic equations through the term

$$\sum_{ij} \frac{e^2}{r_{ij}}$$

In this preliminary study we perform an average of the Coulomb interaction on the hyperangles

$$V_c(\rho) = \int d\Omega \sum_{ij} \frac{e^2}{r_{ij}} |\mathcal{Y}_0(\Omega)|^2 = \frac{16}{\pi} \frac{e^2}{\rho}$$

and the plane wave takes the form

$$e^{i\vec{Q}\cdot\vec{p}} \rightarrow \Psi_s^0 = \frac{1}{C_{3/2}(0)} \frac{(\pi)^3}{(Q\rho)^{5/2}} \sum_{[K]} i^K F_{K+3/2}(\eta, Q\rho) \mathcal{B}_{[K]}(\Omega) \mathcal{B}_{[K]}^*(\hat{Q})$$



## ppp Correlation Analysis

The ppp wave function is

$$\Psi_{ppp} = \sum_{[K]} u_{[K]}(\rho) \mathcal{B}_{[K]}(\Omega) = \Psi^0 + \sum_{J,K}^{\bar{J},\bar{K}} \Psi_K^J$$

To determine  $\Psi^J$  we use the Adiabatic Hyperspherical Harmonic basis:

$$-\frac{\hbar^2}{m} \frac{\Lambda^2(\Omega)}{\rho^2} = H_\Omega \phi_\nu(\rho, \Omega) = U_\nu(\rho) \phi_\nu(\rho, \Omega)$$

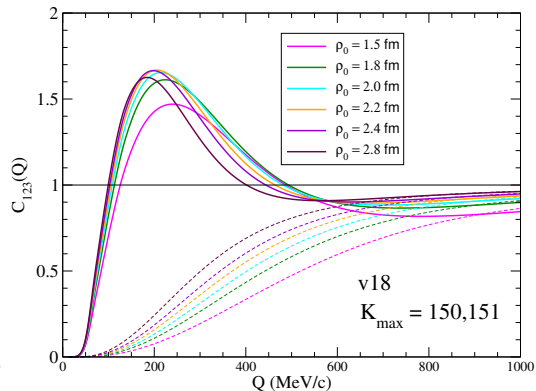
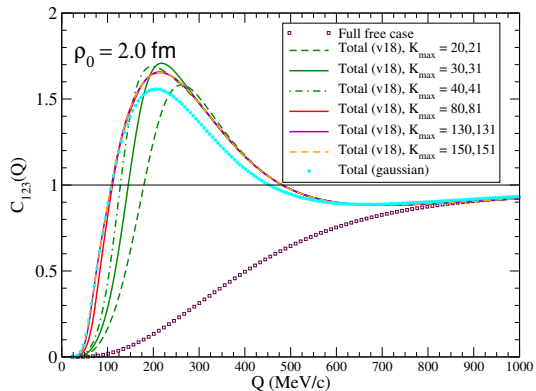
$$\Psi_K^J = \rho^{-5/2} \sum_{\nu} w_\nu^{J,K}(\rho) \phi_\nu(\rho, \Omega)$$

with the adiabatic functions  $\phi_\nu(\rho \rightarrow \infty, \Omega) \rightarrow \mathcal{B}_{[K]}(\Omega)$

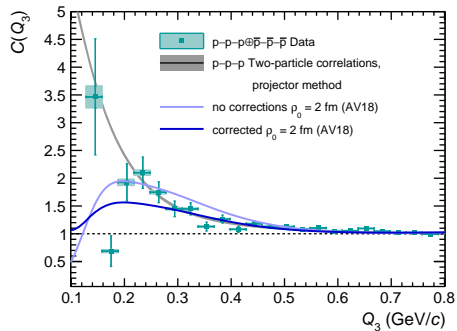
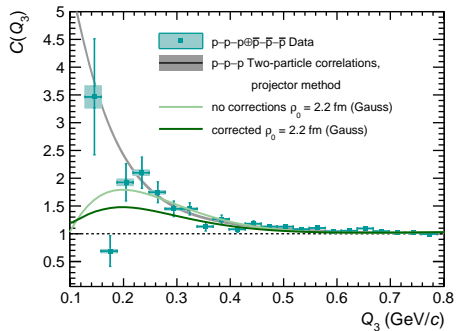
and the hyperradial functions  $w_\nu^{J,K}(\rho) \rightarrow \sqrt{Q\rho} [\delta_{KK'} F_{K+3/2}(Q\rho) + T_{KK'} \mathcal{O}_{K'+3/2}(Q\rho)]$

The ppp correlation function  $C_{123}(Q) = \int \rho^5 d\rho d\Omega S_{123} |\Psi_{ppp}|^2$

Convergence and preliminary results with different size sources



# Comparison to data (preliminary)



# Summary

- ▶ Although its apparent simplicity, the three-nucleon problem is of great complexity
- ▶ Measurements of the correlation function allow for new tests of the NN and NNN interactions
- ▶ In the ppp case the Coulomb interaction couples the asymptotic dynamics increasing the difficulties of the numerical treatment
- ▶ In this preliminary study the Coulomb interaction was averaged on the hyperangles
- ▶ The Adiabatic expansion was used in the lowest channels
  
- ▶ The next work is to include more channels and to relax the average of the Coulomb force
- ▶ Studies on the  $pp\Lambda$  correlation function has been started