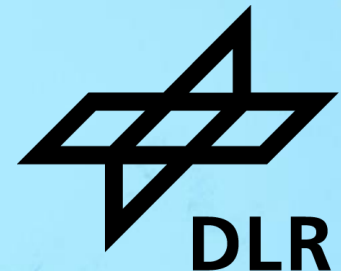


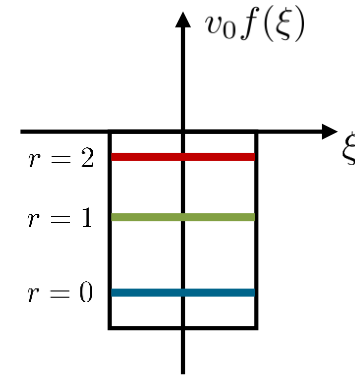
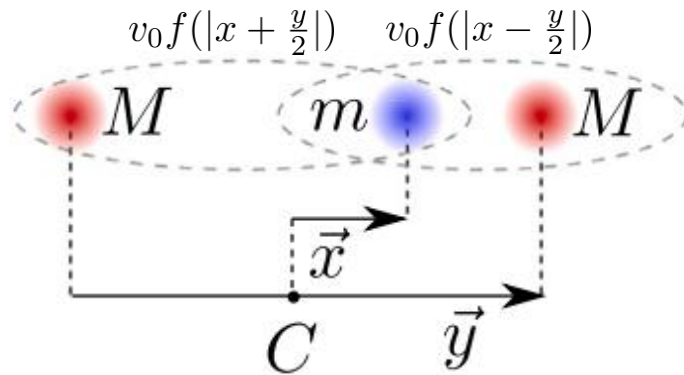
UNIVERSALITY IN LOW-DIMENSIONAL THREE-BODY SYSTEMS

Maxim Efremov

German Aerospace Center (DLR e.V.), Institute of Quantum Technologies



1D three-body problem



$$\left[-\frac{\alpha_x}{2} \frac{\partial^2}{\partial x^2} - \frac{\alpha_y}{2} \frac{\partial^2}{\partial y^2} + v_0 f\left(x + \frac{y}{2}\right) + v_0 f\left(x - \frac{y}{2}\right) \right] \psi_n = \mathcal{E}_n^{(3)} \psi_n$$

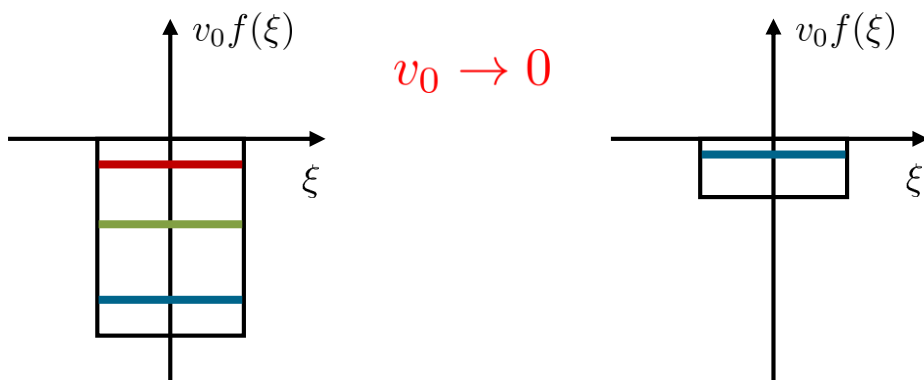
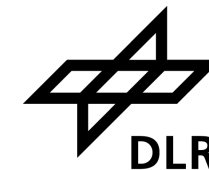
$$\alpha_x = \frac{m + 2M}{2(m + M)}$$

$$\alpha_y = \frac{2m}{m + M}$$

$$\psi_n(x, y) \Big|_{\sqrt{x^2 + y^2} \rightarrow \infty} = 0$$

Question: *is there universality, $f(\xi)$ -independence, in the system?*

The limit $v_0 \rightarrow 0$: $f(\xi) = \delta(\xi)$



Common wisdom:

$$\mathcal{E}_0^{(2)} \rightarrow 0$$

$$f(\xi) \simeq \delta(\xi)$$

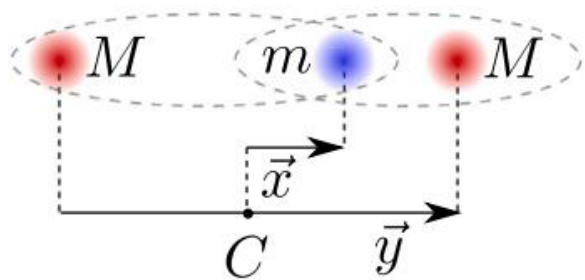
Universal constants:

$$\epsilon_n^* = \frac{\mathcal{E}_n^{(3)}}{|\mathcal{E}_0^{(2)}|} = \epsilon_n^* \left(\frac{M}{m} \right)$$

	Atomic mixture (M/m)		
	$^{87}\text{Rb}-^{40}\text{K}$ (2.2)	$^{87}\text{Rb}-^7\text{Li}$ (12.4)	$^{133}\text{Cs}-^6\text{Li}$ (22.2)
ϵ_0^*	-2.1966	-2.5963	-2.7515
ϵ_1^*	-1.0520	-1.4818	-1.6904
ϵ_2^*		-1.1970	-1.3604
ϵ_3^*		-1.0377	-1.1479
ϵ_4^*		-1.0002	-1.0525
ϵ_5^*			-1.0040

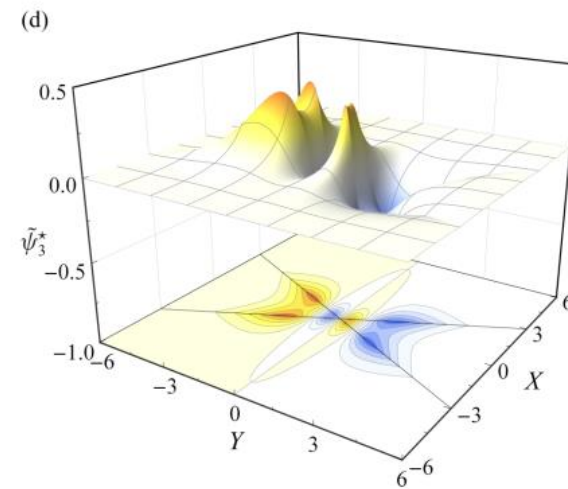
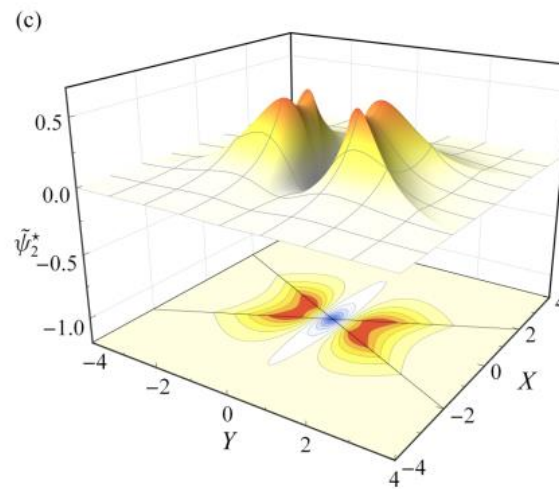
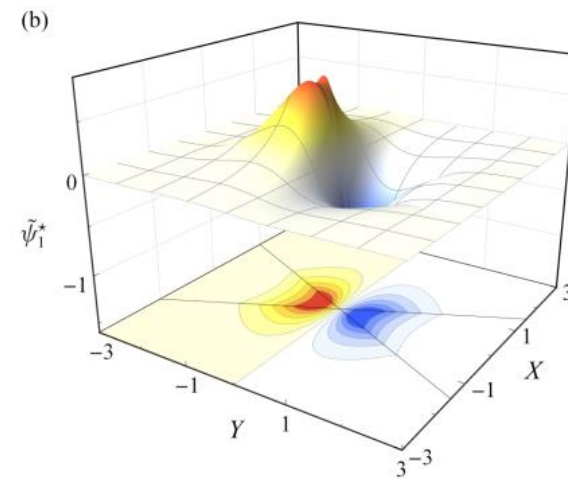
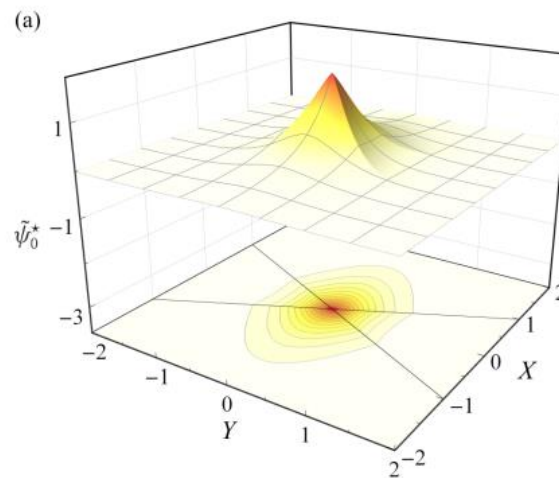
L Happ et al., Phys. Rev. A 100, 012709 (2019)

The limit $v_0 \rightarrow 0$: $f(\xi) = \delta(\xi)$



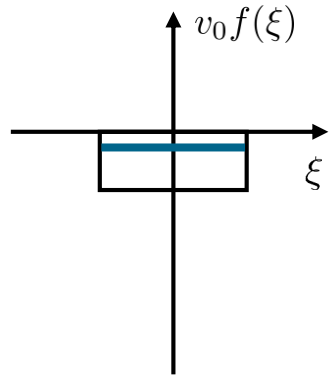
Universal wave functions:

$$\psi_n(x, y) = \psi_n^* \left(\sqrt{2|\mathcal{E}_0^{(2)}|} x, \sqrt{2|\mathcal{E}_0^{(2)}|} y \right)$$



L Happ et al., Phys. Rev. A 100, 012709 (2019)

General potential $f(\xi)$

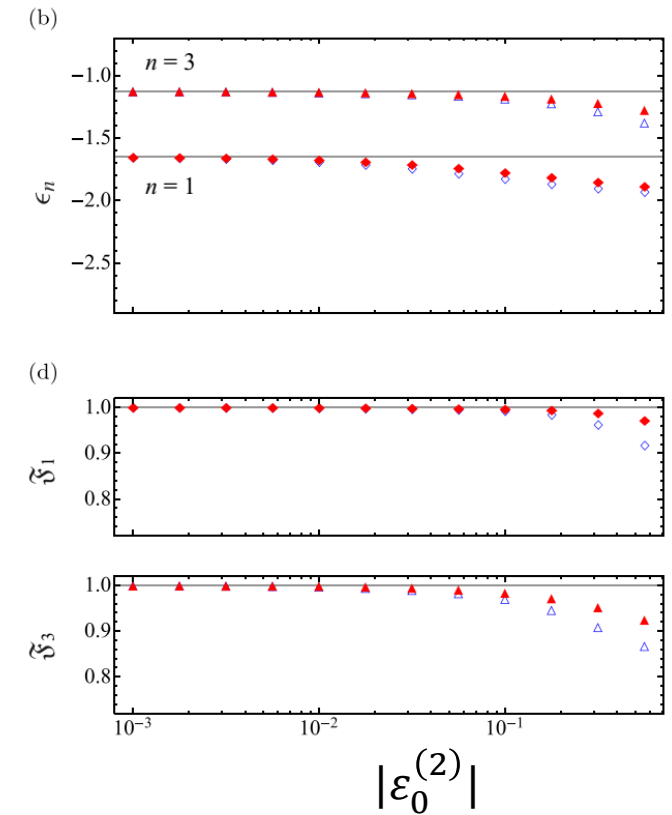
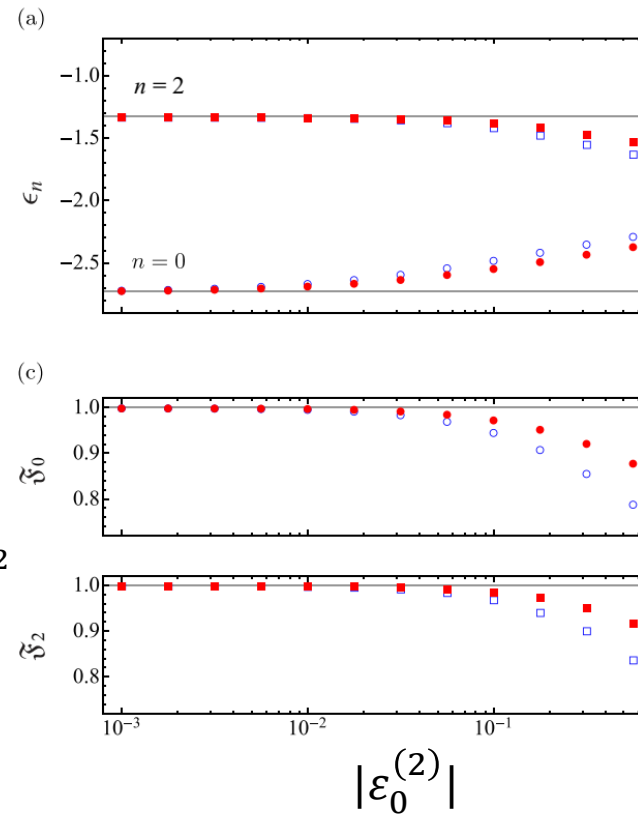


$$\epsilon_n = \frac{\epsilon_n^{(3)}}{|\epsilon_0^{(2)}|}$$

$$f_G(\xi) = e^{-\xi^2}$$

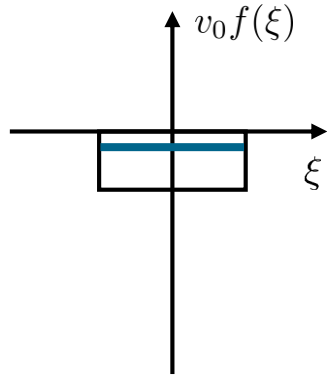
$$f_L(\xi) = \frac{1}{(1 + \xi^2)^3}$$

$$\mathfrak{F}_n = (\langle \psi_n | \psi_n^* \rangle)^2$$

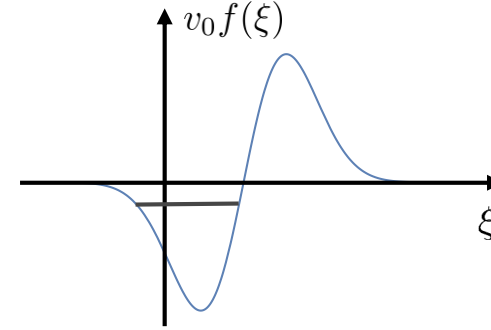


L Happ et al., Phys. Rev. A 100, 012709 (2019)

The limit $v_0 \rightarrow 0$: general $f(\xi)$



$$F(p) \equiv \int e^{-ip\xi} f(\xi) d\xi$$



$$\int v_0 f(\xi) d\xi < 0$$

$$\int v_0 f(\xi) d\xi = 0$$

$$\mathcal{E}_0^{(2)}(v_0 \rightarrow 0) = -v_0^2 [F(0)]^2$$

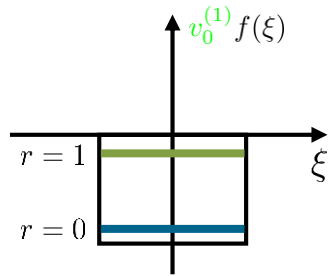
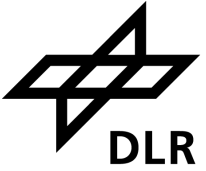
$$\mathcal{E}_0^{(2)}(v_0 \rightarrow 0) = -v_0^4 \left[\int \frac{dp}{\pi} \frac{|F(p)|^2}{p^2} \right]^2$$

B. Simon, *Ann. Phys.* **97**, 279 (1976)

Universality: $\mathcal{E}_n^{(3)} = -\epsilon_n^* |\mathcal{E}_0^{(2)}|$
 $(v_0 \rightarrow 0)$ $\psi_n(x, y) = \psi_n^* \left(\sqrt{2|\mathcal{E}_0^{(2)}|} x, \sqrt{2|\mathcal{E}_0^{(2)}|} y \right)$

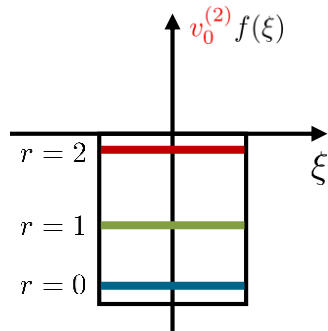
L Happ and M Efremov, *J. Phys. B* **54**, 21LT01 (2021)

1D three-body problem: excited states



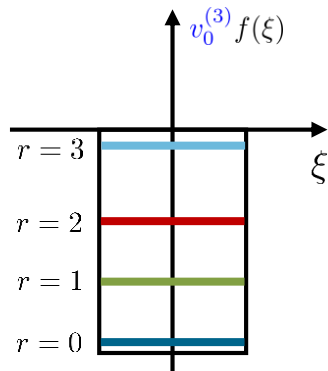
$$\mathcal{E}_{r=1}^{(2)}(v_0^{(1)} \rightarrow v_1) \rightarrow 0$$

3-body energies
(Faddeev equations): $\mathcal{E}_{n,1}^{(3)}$



$$\mathcal{E}_{r=2}^{(2)}(v_0^{(2)} \rightarrow v_2) \rightarrow 0$$

3-body energies
(Faddeev equations): $\mathcal{E}_{n,2}^{(3)}$



$$\mathcal{E}_{r=3}^{(2)}(v_0^{(3)} \rightarrow v_3) \rightarrow 0$$

3-body energies
(Faddeev equations): $\mathcal{E}_{n,3}^{(3)}$

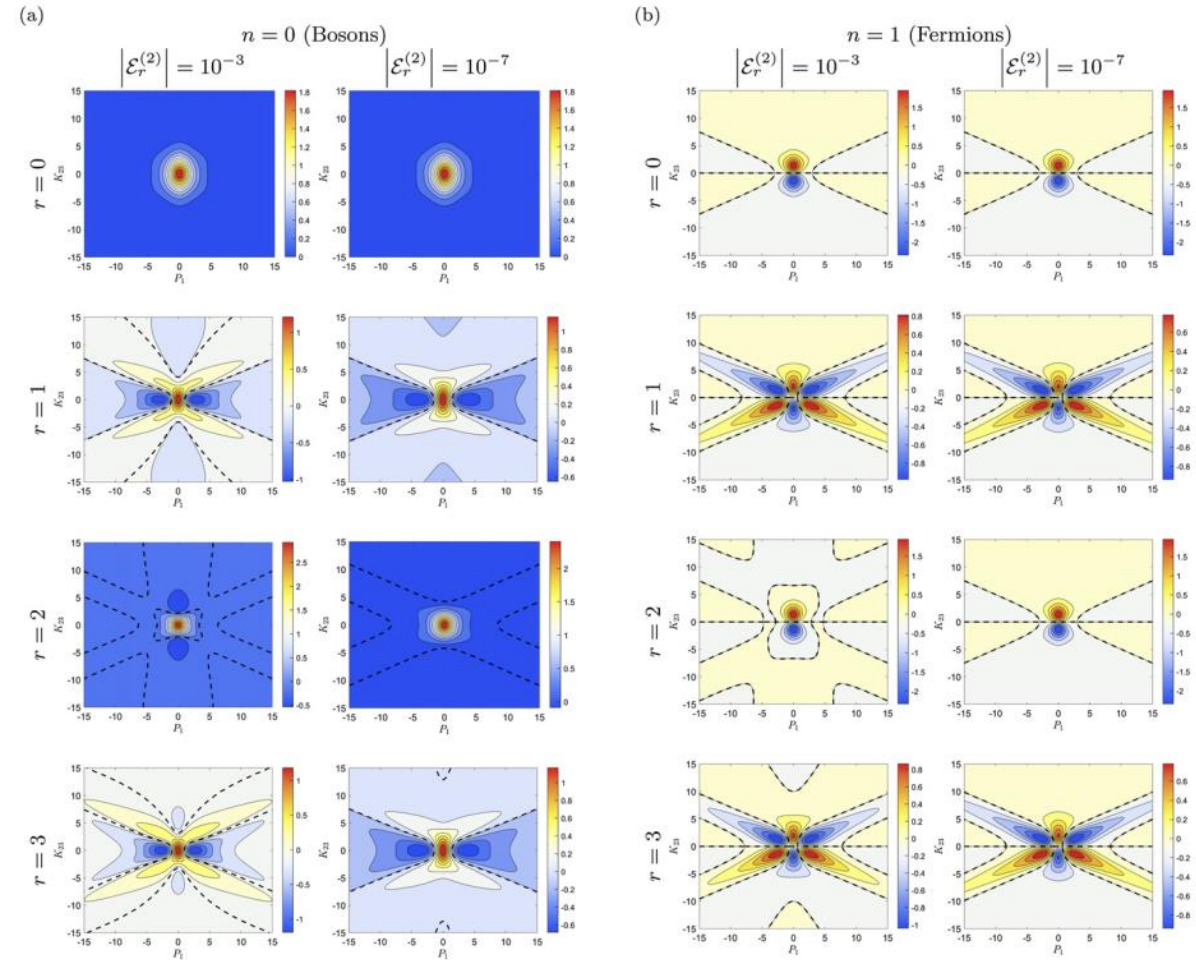
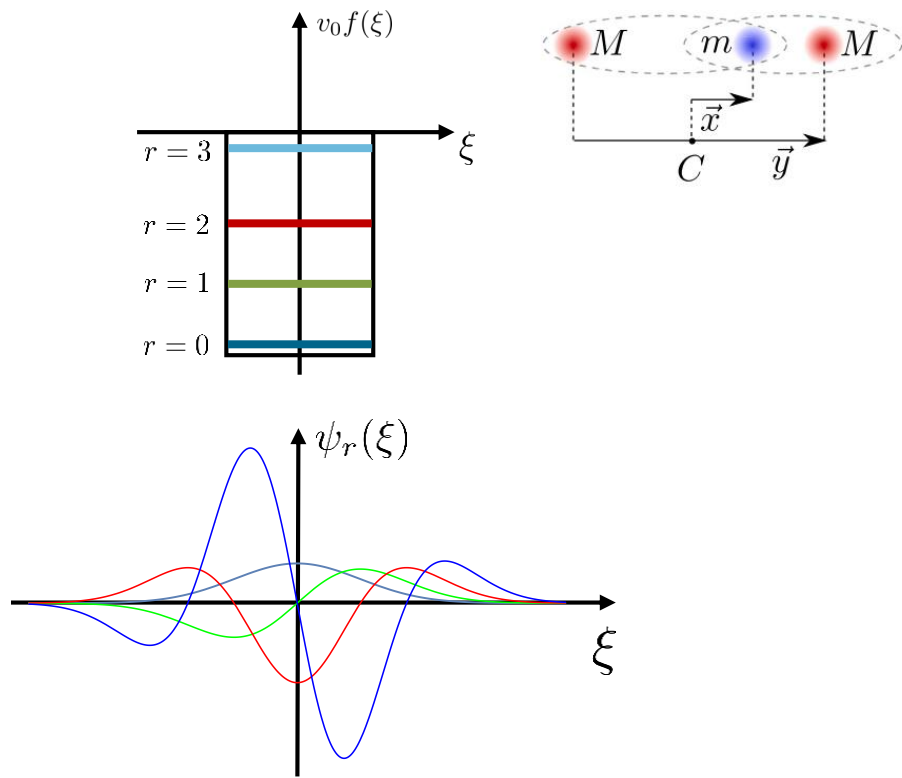
Universality:

$$\frac{\mathcal{E}_{n,r}^{(3)}}{|\mathcal{E}_r^{(2)}|} = -\epsilon_n^*$$

as for $f^*(\xi) = \delta(\xi)$

L Happ, M Zimmermann and M Efremov, J. Phys. B 55, 015301 (2022)

1D three-body problem: excited states



Universality for even or odd r : $\mathcal{E}_{n,r}^{(3)} = -\epsilon_n^* |\mathcal{E}_r^{(2)}|$, $\psi_{n,2l}(x, y) = \psi_n^* \left(\sqrt{2|\mathcal{E}_{2l}^{(2)}} |x, \sqrt{2|\mathcal{E}_{2l}^{(2)}} |y \right)$
 $\psi_{n,2l+1}(x, y) = \psi_{n,2l+3}(x, y)$

1D M - m - M system:

- ☑ If M - m subsystem has a shallow ground state:

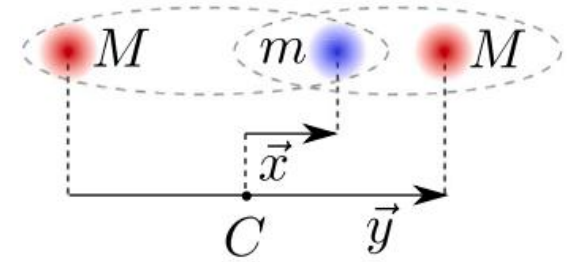
$$\text{As } \mathcal{E}_0^{(2)} \rightarrow 0: \quad \mathcal{E}_{n,0}^{(3)} = -\epsilon_n^* |\mathcal{E}_0^{(2)}|, \quad \psi_{n,0}(x, y) = \psi_n^* \left(\sqrt{2|\mathcal{E}_0^{(2)}|} x, \sqrt{2|\mathcal{E}_0^{(2)}|} y \right)$$

- ☑ If M - m subsystem has a shallow excited state (see talk by L. Happ, 01.08, Few-body systems):

$$\text{As } \mathcal{E}_{2l}^{(2)} \rightarrow 0: \quad \mathcal{E}_{n,2l}^{(3)} = -\epsilon_n^* |\mathcal{E}_{2l}^{(2)}|, \quad \psi_{n,2l}(x, y) = \psi_n^* \left(\sqrt{2|\mathcal{E}_{2l}^{(2)}|} x, \sqrt{2|\mathcal{E}_{2l}^{(2)}|} y \right)$$

$$\text{As } \mathcal{E}_{2l+1}^{(2)} \rightarrow 0: \quad \mathcal{E}_{n,2l+1}^{(3)} = -\epsilon_n^* |\mathcal{E}_{2l+1}^{(2)}|, \quad \psi_{n,2l+1}(x, y) = ?$$

$$f^*(\xi) = \frac{d}{d\xi} \delta(\xi)?$$



2D M - m - M system:

- ☑ If M - m subsystem has a shallow ground state: *universality for energies and wave-functions*

J Thies, MT Hof, M Zimmermann and M Efremov, J. Comp. Science **64**, 101859 (2022)

