

# Recent progress in hypernuclear physics

Emiko Hiyama (Tohoku Univ./RIKEN)

# Major goals of hypernuclear physics

To understand baryon-baryon interactions

Fundamental and important for the study of nuclear physics

Total number of  
Nucleon (N) -Nucleon (N) data: 4,000



- Total number of differential cross section  
Hyperon (Y) -Nucleon (N) data: 40
- **NO** YY scattering data

YN and YY potential models so far proposed (ex. Chiral, Nijmegen, Kyoto-Niigata) have large ambiguity.

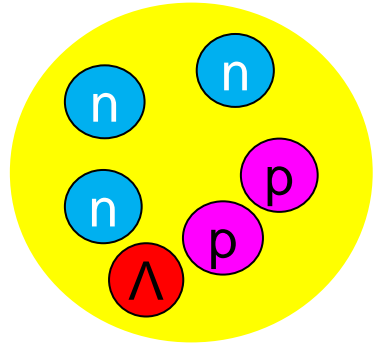
Therefore, for the study of  $YN$  and  $YY$  interactions, the systematic investigation of the structure of light hypernuclei is one of the important way.

(Some  $YN$  scattering experiments have been done and further experiment is planned at J-PARC.)

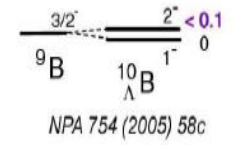
Once  $YN$  and  $YY$  interactions are determined, we can predict interesting phenomena which cannot be imagined so far. In addition, we could study inner part of neutron stars which have been observed.

Since 1998

# Hypernuclear $\gamma$ -ray data (2019)

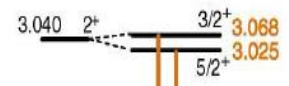


$^{10}\text{B} (K^-, \pi^+ \gamma)$  BNL E930('01)



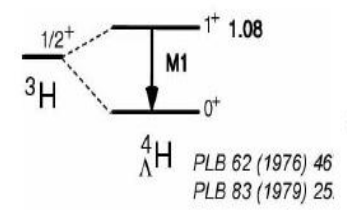
NPA 754 (2005) 58c

$^9\text{Be} (K^-, \pi^+ \gamma)$  BNL E930('98)



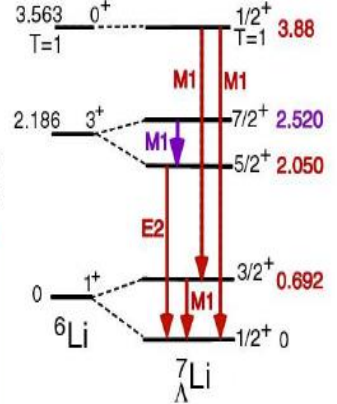
PRL 88 (2002) 082501  
NPA 754 (2005) 58c

$^7\text{Li}$  etc. ( $K^{\text{stop}}, \gamma \pi^-$ )

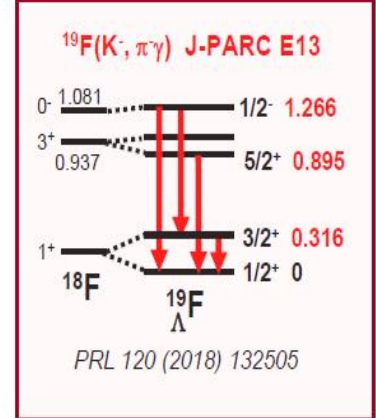
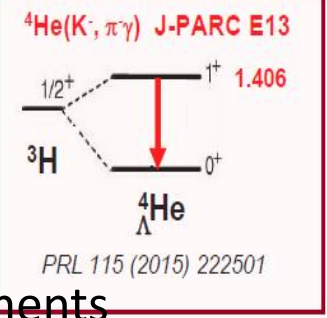


PLB 62 (1976) 46  
PLB 83 (1979) 25

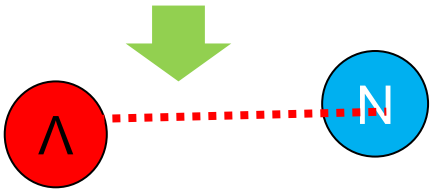
$^7\text{Li} (\pi^+, K^+ \gamma)$  KEK E419



PRL 84 (2000) 5963  
PRL 86 (2001) 1982  
PLB 579 (2004) 258  
PRC 73 (2006) 012501

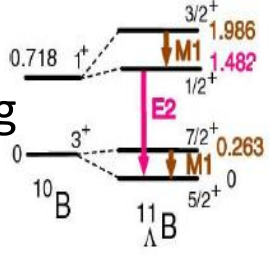


Ab-initio calculation  
Shell model calculation  
+  
High-resolution experiments



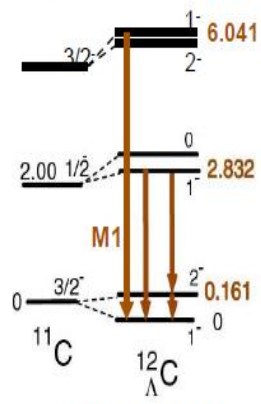
We have been obtaining information on  $\Lambda N$  two-body interaction.

$^{11}\text{B} (\pi^+, K^+ \gamma)$  KEK E518



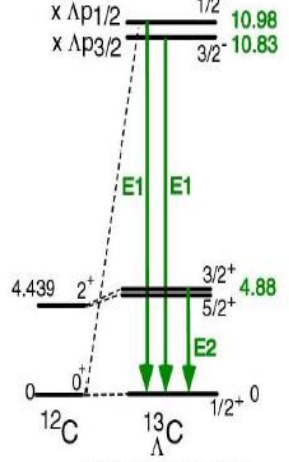
NPA835 (2010) 422

$^{12}\text{C} (\pi^+, K^+ \gamma)$  KEK E566



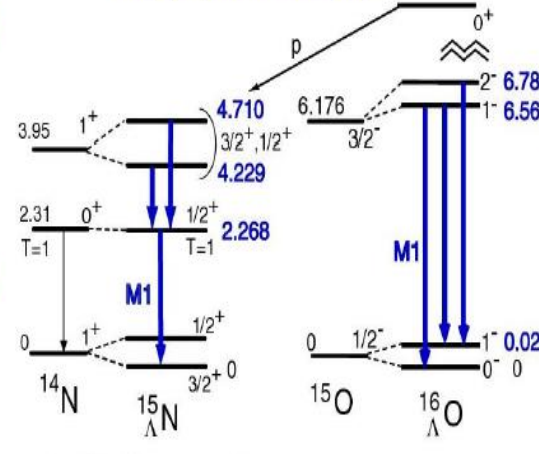
PTEP (2015) 081D01

$^{13}\text{C} (K^-, \pi^+ \gamma)$  BNL E929 (Nal)



PRL 86 (2001) 4255  
PRC 65 (2002) 034607

$^{16}\text{O} (K^-, \pi^+ \gamma)$  BNL E930('01)



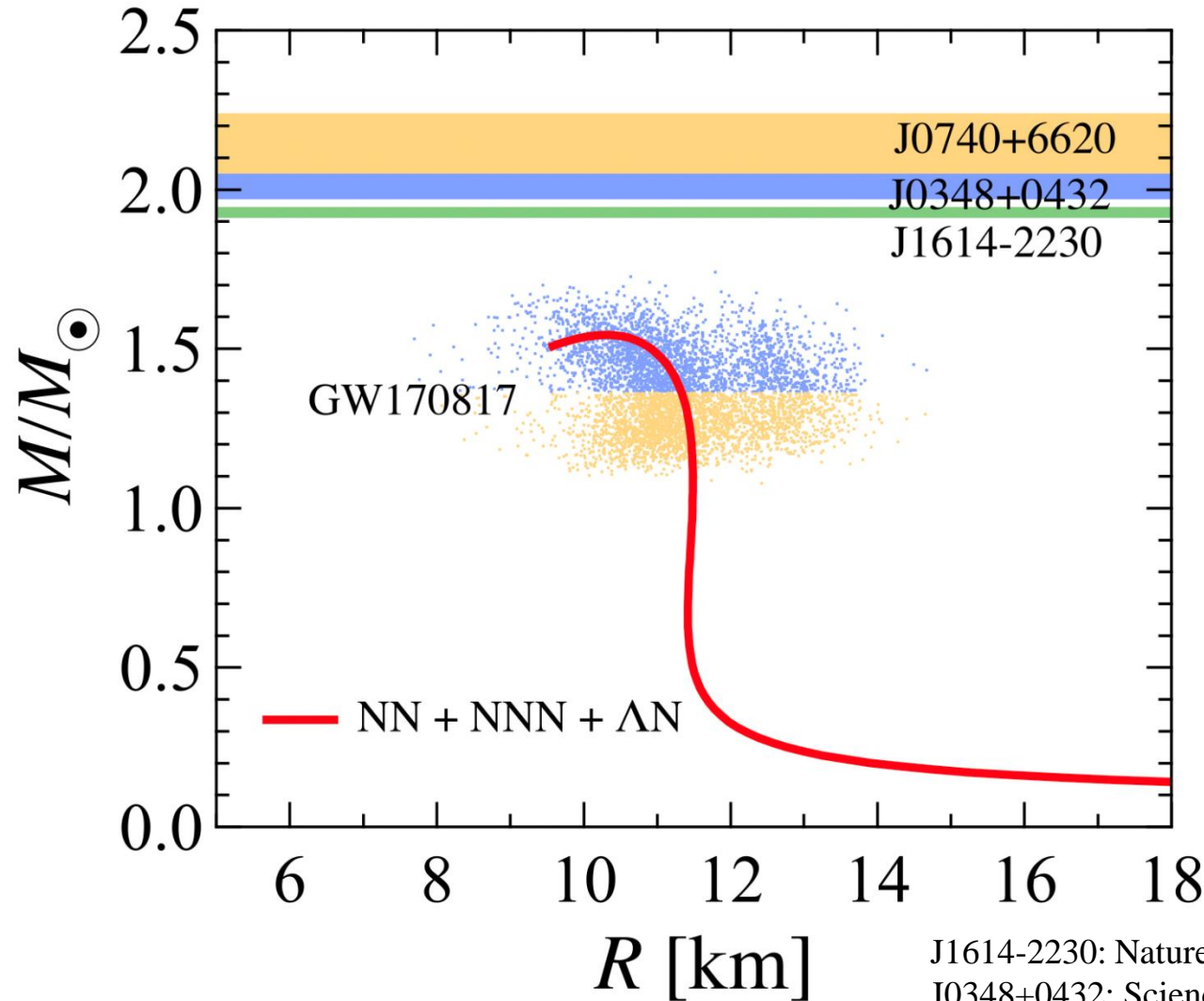
PRC 77 (2008) 054315

PRL 93 (2004) 232501  
EPJ A33 (2007) 247

$$V_{\Lambda N} = V_0 + \sigma_{\Lambda} \cdot \sigma_N V_{\sigma\sigma} + \mathbf{L} \cdot (\mathbf{s}_{\Lambda} + \mathbf{s}_N) V_{\text{SLS}} + \mathbf{L} \cdot (\mathbf{s}_{\Lambda} - \mathbf{s}_N) V_{\text{ALS}} + S_{12} V_{\text{tensor}} + \dots$$

# Mass-Radius Relation of Neutron Stars

2021



$\Lambda$ N interaction updated by structure study of  $\Lambda$  hypernuclei + EoS by cluster variational calculation done by Togashi

Still, the maximum mass of neutron star is less than twice of solar mass.

Hyperon puzzle

J1614-2230: Nature 467 (2010) 1081, APJ 832 (2016) 167

J0348+0432: Science 340 (2013) 1233232

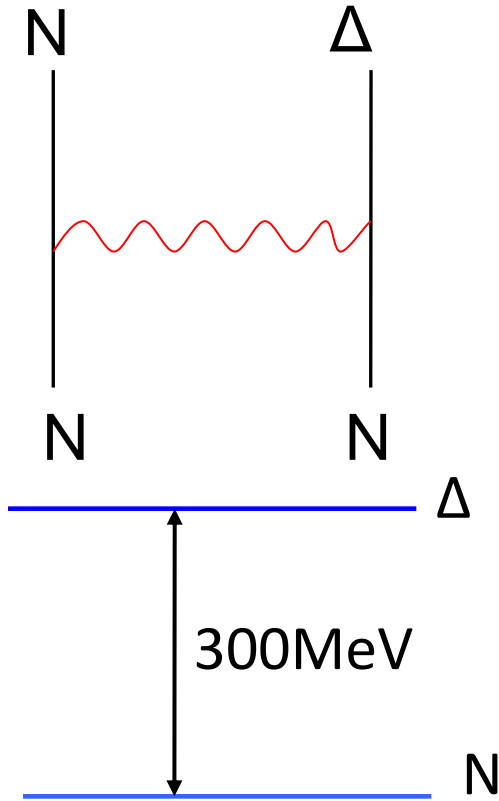
J0740+6620: Nat. Astron. (2019)

GW170817: PRL 121 (2018) 161101

missing part of YN interaction:  $\Lambda$ N- $\Sigma$ N coupling

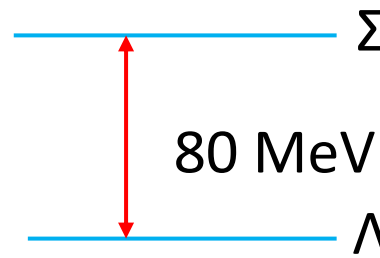
# $\Lambda N - \Sigma N$ coupling

Non-strangeness sector



Probability of  $\Delta$  in nuclei is not large.

$S = -1$

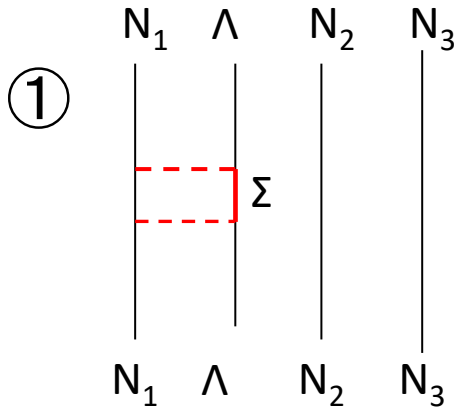


Mass is smaller.  
It is expected that  $\Lambda$ - $\Sigma$  conversion might affect in structure of  $\Lambda$  hypernuclei.

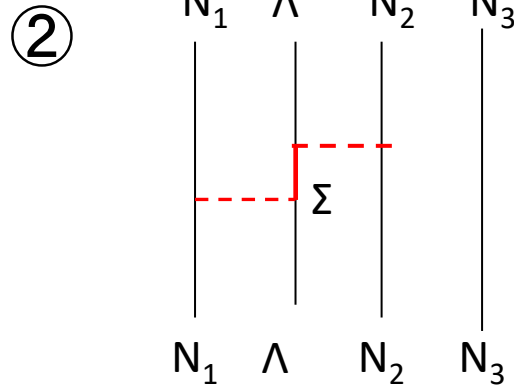
$\Lambda N$ - $\Sigma N$  coupling is key issue to construct  $YN$  two-body interaction completely.

# Role of the $\Lambda$ N- $\Sigma$ N interaction

- Three-body effect



Effective two-body force



Three-body force

In the neutron matter or neutron star, three-body force might play important role.

- Charge symmetry breaking effect

# Charge Symmetry breaking

In  $S=0$  sector

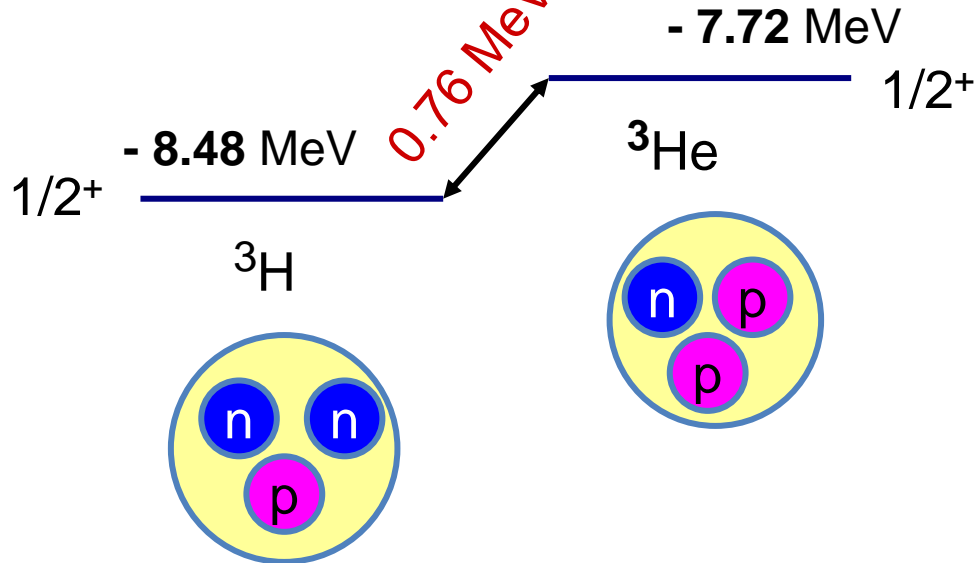
**Exp.**

$N+N+N$

Energy difference comes from dominantly Coulomb force between 2 protons.

Charge symmetry breaking (n-n,p-p) effect is small.

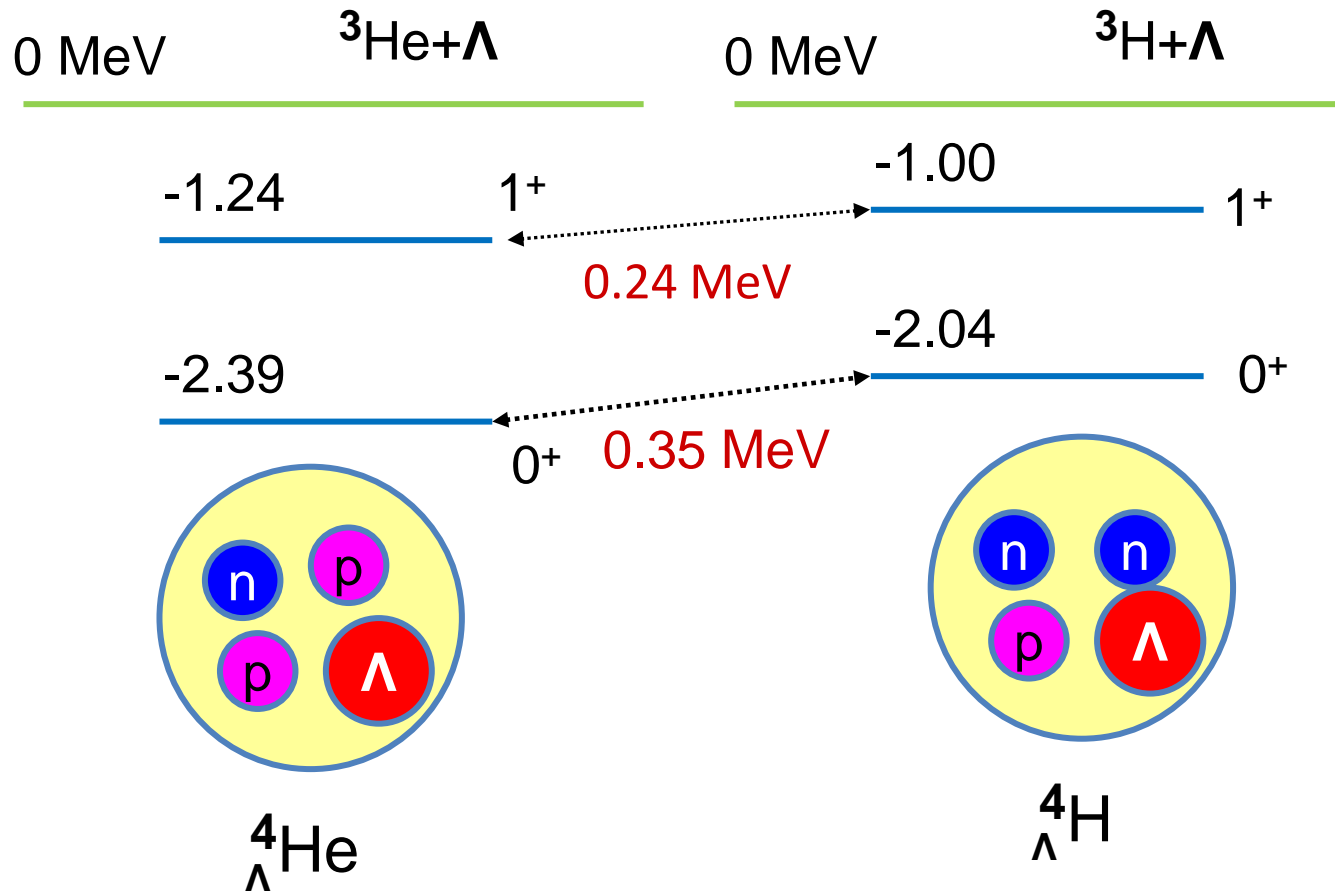
0 MeV

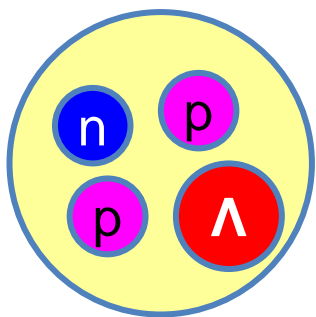




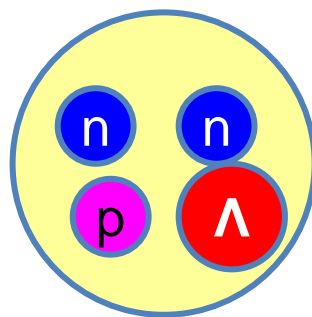
In  $S = -1$  sector

Exp.



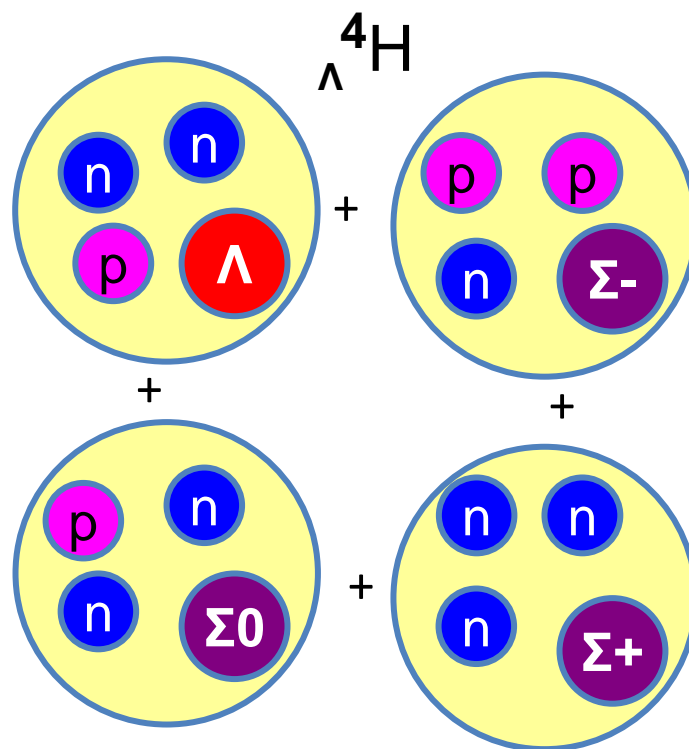
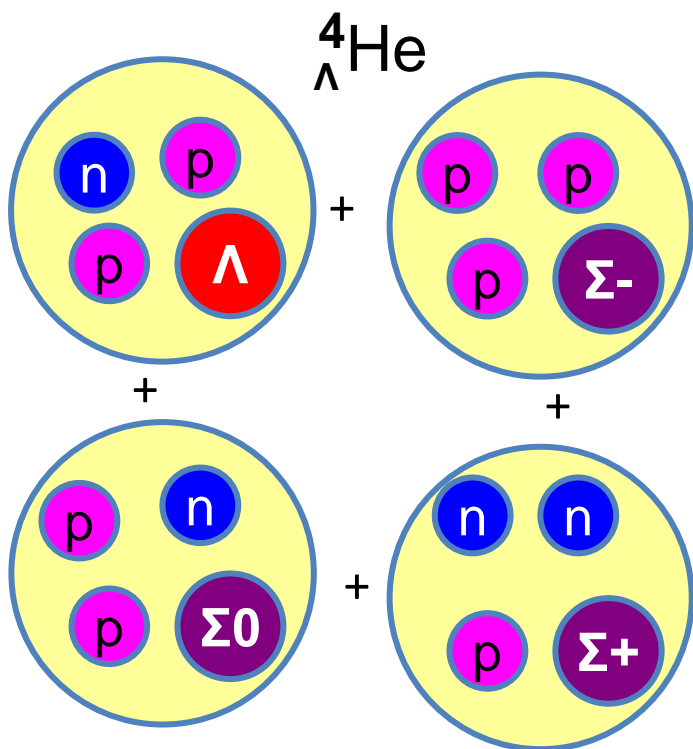


${}^4_{\Lambda}\text{He}$

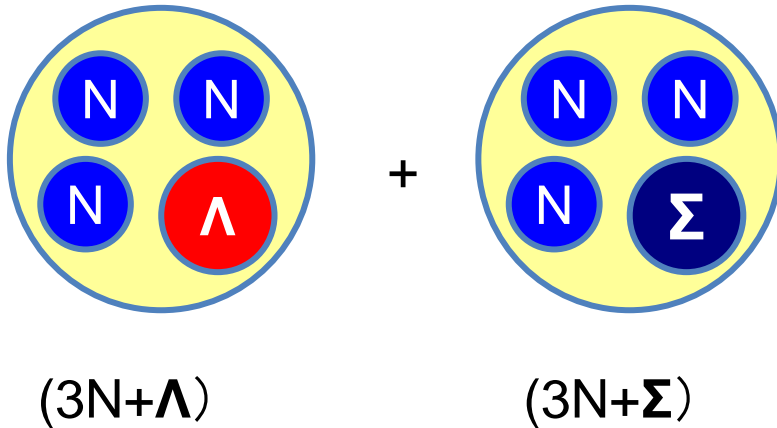


${}^4_{\Lambda}\text{H}$

However,  $\Lambda$  particle has no charge.

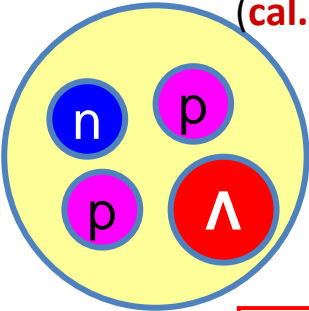
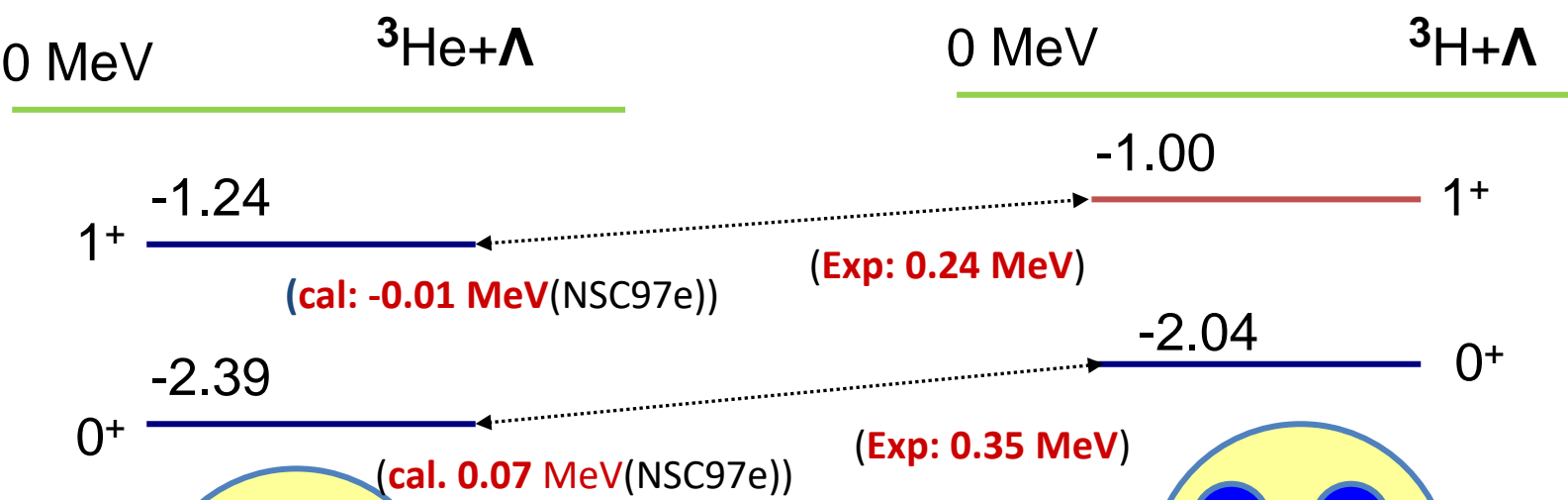


In order to explain the energy difference, **0.35 MeV**,



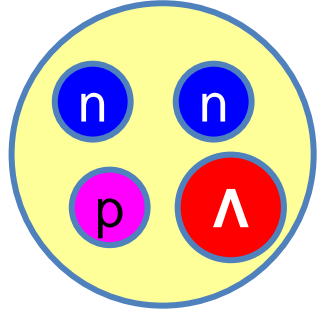
- E. Hiyama, M. Kamimura, T. Motoba, T. Yamada and Y. Yamamoto, Phys. Rev. C65, 011301(R) (2001).
- A. Nogga, H. Kamada and W. Gloeckle, Phys. Rev. Lett. 88, 172501 (2002)
- H. Nemura, Y. Akaishi and Y. Suzuki, Phys. Rev. Lett.89, 142504 (2002).

**Coulomb potentials between charged particles (p,  $\Sigma^\pm$ ) are included.**



${}^4_{\Lambda}\text{He}$

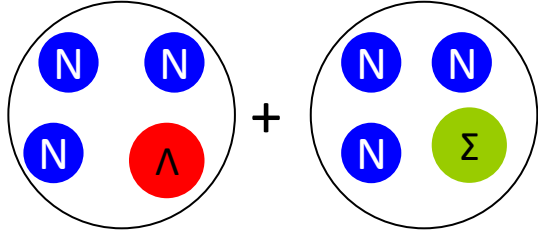
• A. Nogga, H. Kamada and W. Gloeckle, Phys. Rev. Lett. 88, 172501 (2002)

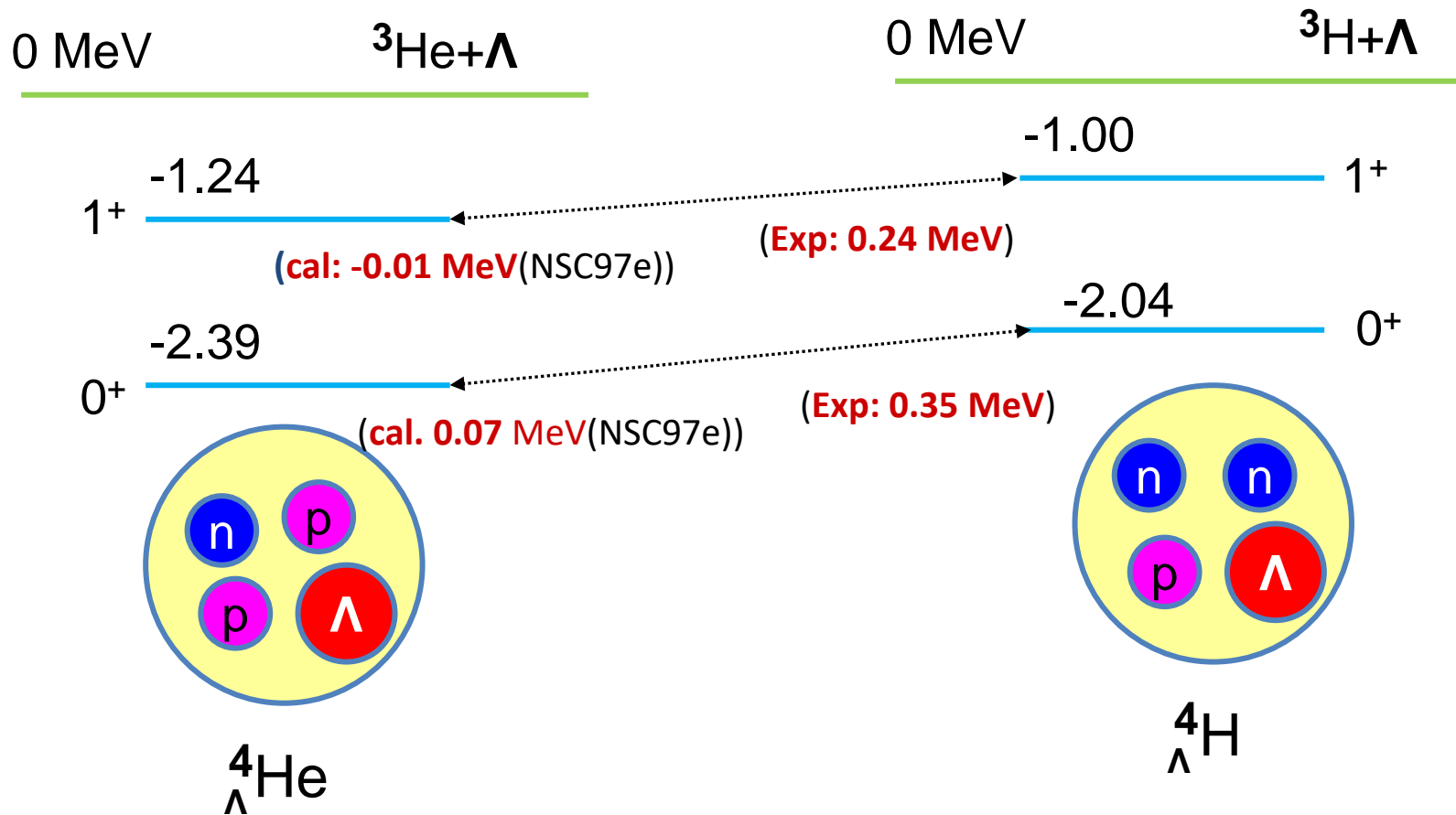


${}^4_{\Lambda}\text{H}$

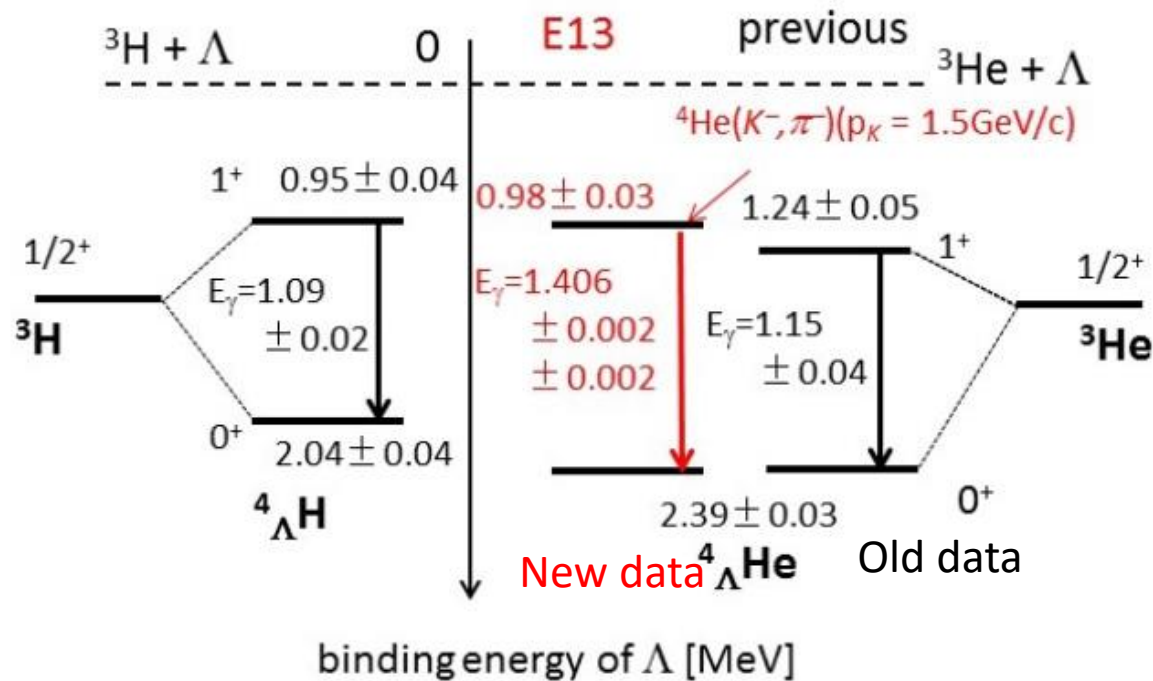
• E. Hiyama, M. Kamimura, T. Motoba, T. Yamada and Y. Yamamoto, Phys. Rev. C65, 011301(R) (2001).

• H. Nemura, Y. Akaishi and Y. Suzuki, Phys. Rev. Lett.89, 142504 (2002).





There has been exist NO YN interaction to reproduce the data.



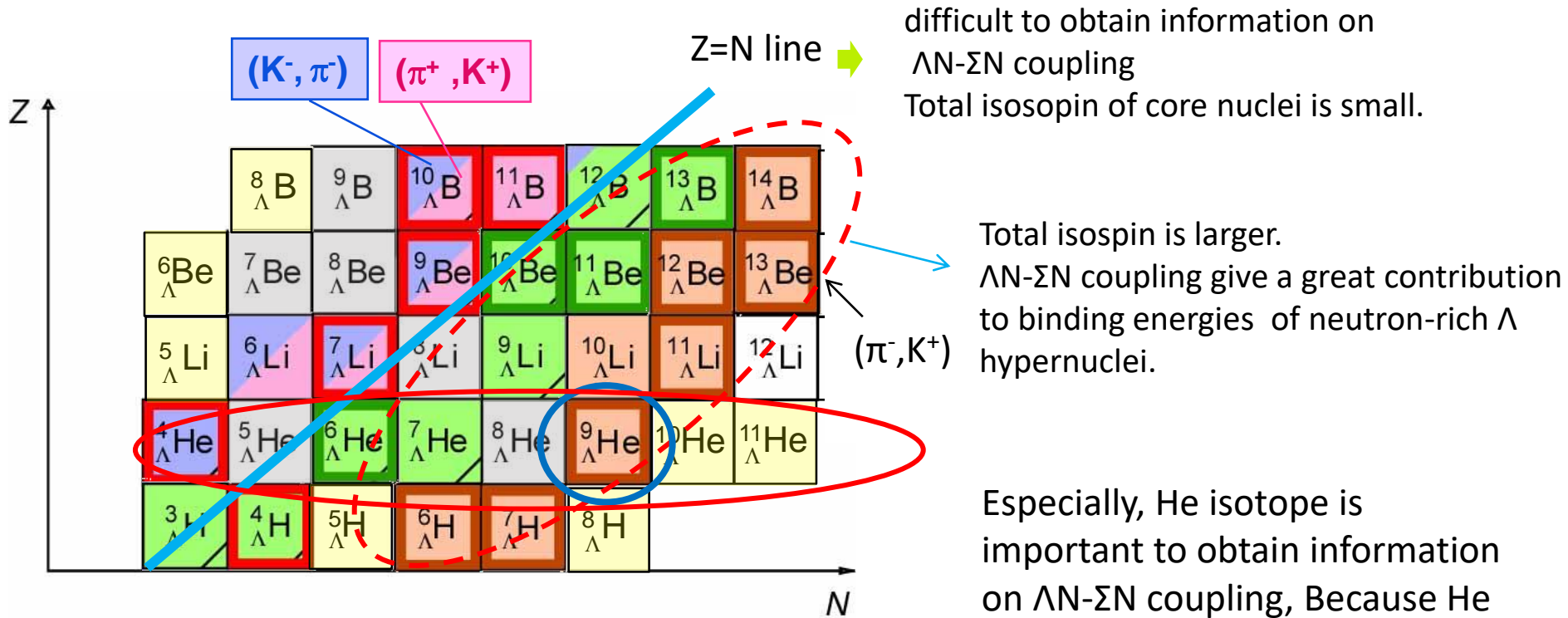
T. O. Yamamoto, Phys. Rev. Lett.115, 2225 (2015).

M. Schafer et al., PRC106, L031001(2022)

Still it is difficult to reproduce the data for the study of CSB which is related to  $\Lambda N$ - $\Sigma N$  coupling. We need more data related to  $\Lambda N$ - $\Sigma N$  coupling.

# How do we obtain information on $\Lambda N$ - $\Sigma N$ coupling?

- (1)  $\gamma N$  scattering experiment at J-PARC
- (2) To study neutron-rich  $\Lambda$  hypernuclei at J-PARC



These neutron-rich  $\Lambda$  hypernuclei are important.

# Structure of neutron-rich He $\Lambda$ Hypernuclei using the cluster orbital shell model

T. Myo<sup>1,2</sup>

<sup>1</sup>General Education, Faculty of Engineering, Osaka Institute of Technology, Osaka 535-8585, Japan and

<sup>2</sup>Research Center for Nuclear Physics (RCNP), Osaka University, Ibaraki 567-0047, Japan

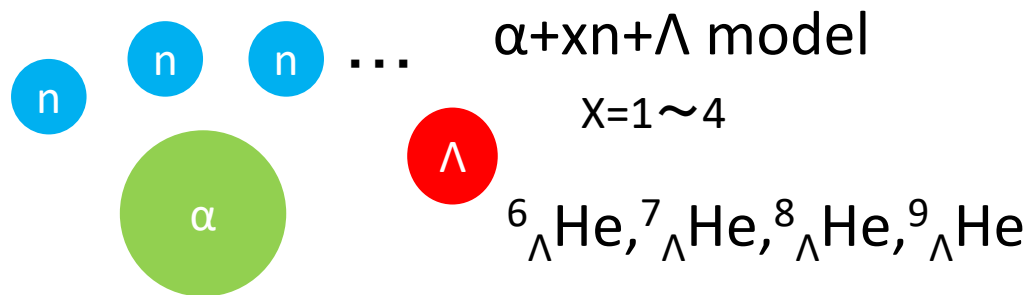
E. Hiyama<sup>3,4</sup>

<sup>3</sup>Department of Physics, Tohoku University, Sendai, 980-8578, Japan and

<sup>4</sup>RIKEN, Nishina Center, Wako, Saitama, 351-0198, Japan

We calculated the energy spectra of the neutron-rich He  $\Lambda$  hypernuclei with  $A = 6$  to  $9$  within the framework of an  $\alpha + \Lambda + Xn$  ( $X = 1 \sim 4$ ) cluster model using the cluster orbital shell model. The employed constituent particles reproduce their observed properties. For resonant states of core nuclei such as  ${}^5\text{He}$ ,  ${}^6\text{He}$  and  ${}^7\text{He}$ , the complex scaling method is employed to obtain energies and decay widths. The calculated ground states of  ${}^6_{\Lambda}\text{He}$  and  ${}^7_{\Lambda}\text{He}$  are in good agreement with published data. The energy levels of  ${}^8_{\Lambda}\text{He}$  and  ${}^9_{\Lambda}\text{He}$  are predicted. In  ${}^9_{\Lambda}\text{He}$ , we find one deeply bound state and two excited resonant states, which are proposed to be produced at J-PARC by the double-charge-exchange reaction ( $\pi^-$ ,  $K^+$ ) using a  ${}^9\text{Be}$  target.

T. Myo and E. Hiyama,  
Phys. Rev. C107, 054302(2023)

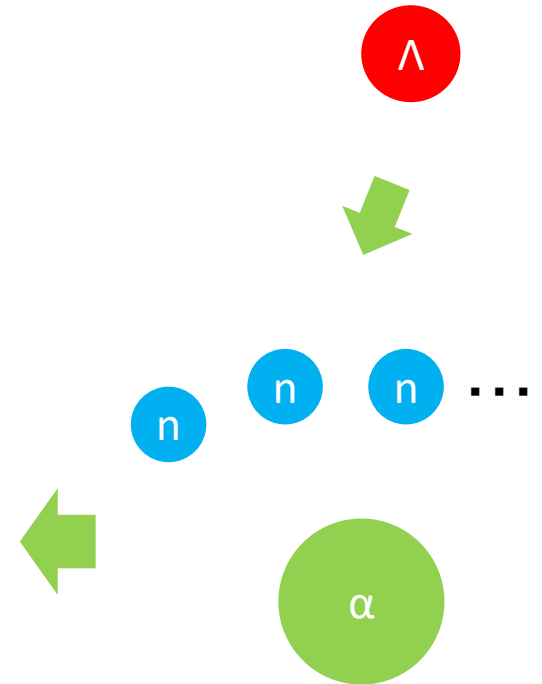
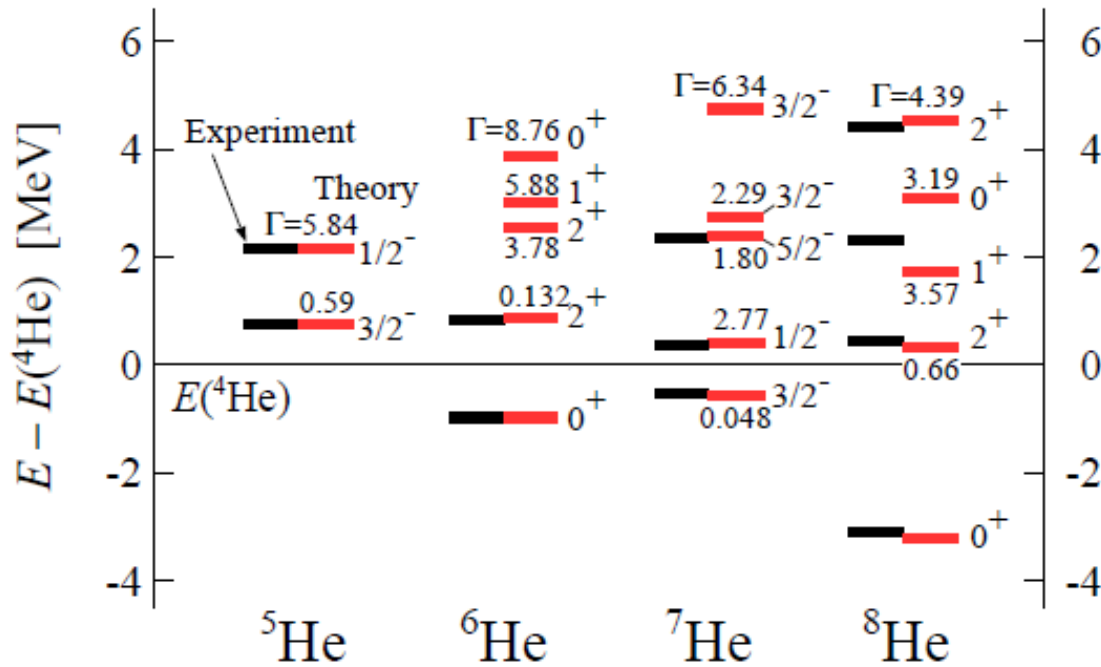


For  ${}^9_{\Lambda}\text{He}$ , there have been no observed data. To predict binding energy of this hypernucleus, it is necessary to reproduce the energy spectra of core nucleus,  ${}^8\text{He}$ .

Cf. R. Wirth and R. Roth, Phys. Lett. B779, 336 (2019).: Non-core shell model+NN+NNN+YN  
Since they did not focus on reproducing observed binding energies of hypernuclei and core nuclei, it was difficult to predict the energy spectra of  ${}^9_{\Lambda}\text{He}$ .

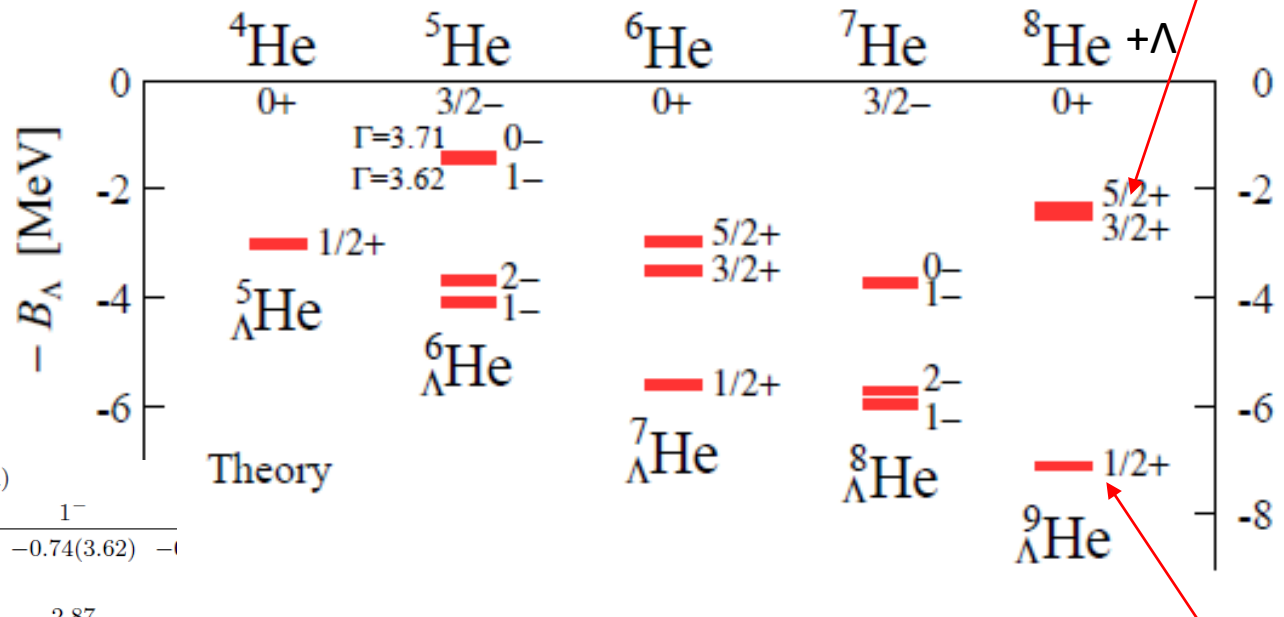


# Energy spectra of He Isotope core nuclei:



The theoretical results are in good agreement with data.  
Let's add a  $\Lambda$  particle into He isotope nuclei.

$$B_{\Lambda} = 2.49 \text{ MeV}$$



$$B_{\Lambda} = 7.09 \text{ MeV}$$

	${}^6\text{He}(\alpha n \Lambda)$		
$J^{\pi}$	$1^{-}$	$2^{-}$	$1^{-}$
$E$ (MeV)	-3.35	-2.94(0.06)	-0.74(3.62)
$E^{\text{exp}}$ (MeV)			
$B_{\Lambda}$ (MeV)	4.10	3.34	2.87
$B_{\Lambda}^{\text{exp}}$ (MeV)	4.18		
	${}^7_{\Lambda}\text{He}(\alpha n n \Lambda)$		
$J^{\pi}$	$1/2^{+}$	$3/2^{+}$	$5/2^{+}$
$E$ (MeV)	-6.48	-4.37	-3.84
$E^{\text{exp}}$ (MeV)			
$B_{\Lambda}$ (MeV)	5.51	3.69	1.42
$B_{\Lambda}^{\text{exp}}$ (MeV)	5.55	3.65	

In  ${}^9_{\Lambda}\text{He}$ , three bound states are predicted.

By  $(\pi^{-}, K^{+})$  reaction at J-PARC using  ${}^9\text{Be}$  target, it is possible to produce these hypernucleus.

This would be observation of the most heavy He isotope  $\Lambda$  hypernucleus.

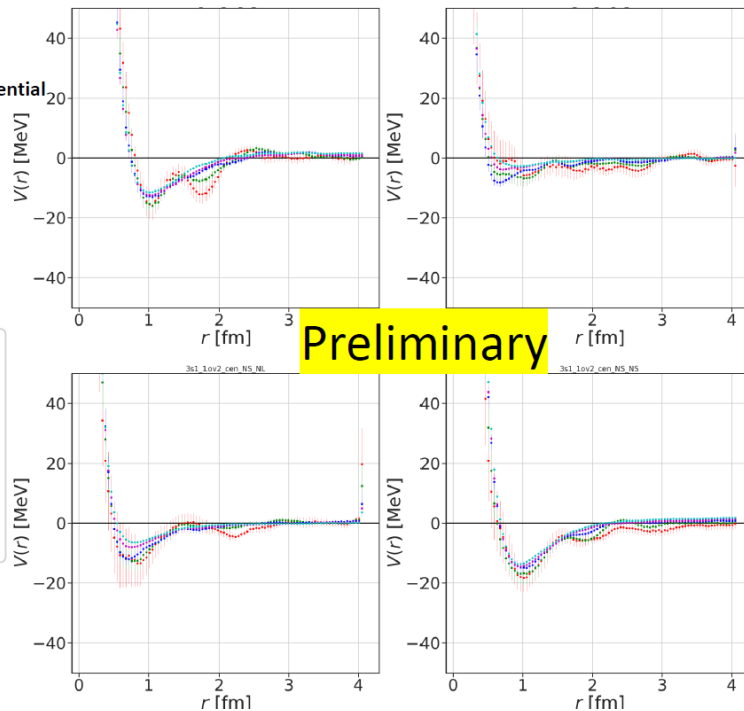
- YN scattering experiment  
proposal: J-PARC-E90



Improvement of potentials  
Chiral potential etc.

- YN interaction from view point of Ab-initio calculation such as Lattice QCD

$N\Lambda - N\Sigma$   
coupled channel potential  
3S1, l=1/2  
central  
binsize=80  
Nconf=800  
w/ Misner

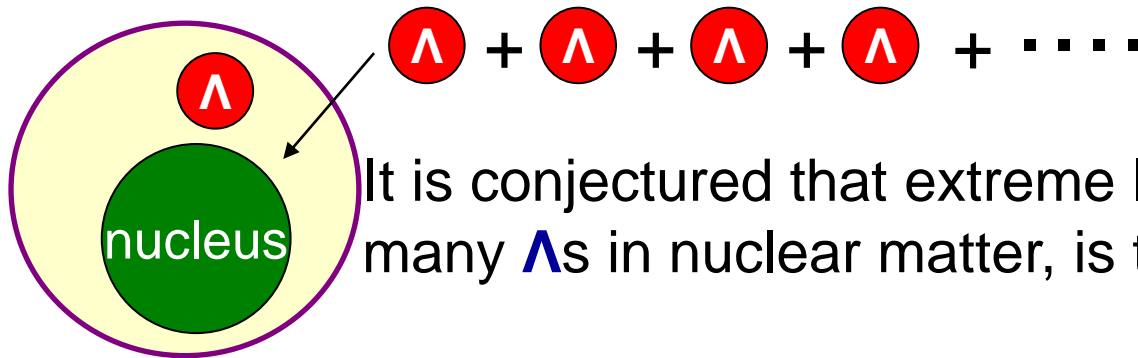


YN interaction by HAL QCD  
It will be possible to employ  
the interaction.

slide by T. Doi

S=-2 hypernuclei  
and  
YY interaction

What is the structure when one or more  $\Lambda$ s are added to a nucleus?



It is conjectured that extreme limit, which includes many  $\Lambda$ s in nuclear matter, is the **core of a neutron star**.

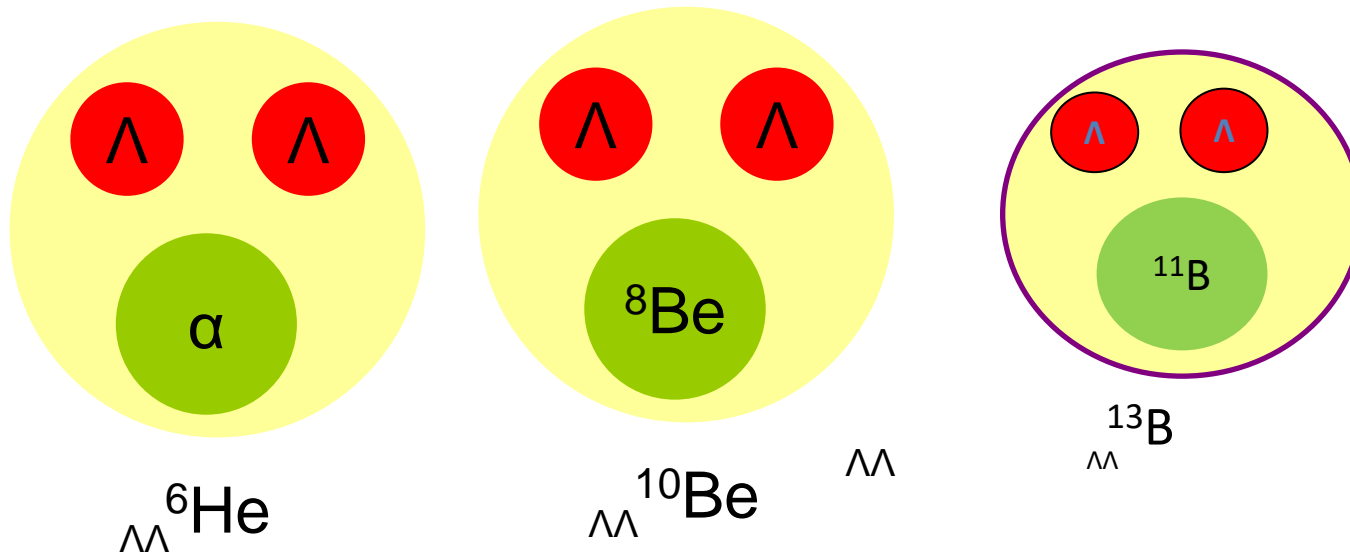
In this meaning, the sector of  $S=-2$  nuclei , double  $\Lambda$  hypernuclei and  $\Xi$  hypernuclei is just the entrance to the **multi-strangeness** world.

However, we have hardly any knowledge of the  $YY$  interaction because there exist no  $YY$  scattering data.

Then, in order to understand the  $YY$  interaction, it is crucial to study the structure of double  $\Lambda$  hypernuclei and  $\Xi$  hypernuclei.

Before 2000

Only three double  $\Lambda$  hypernuclei



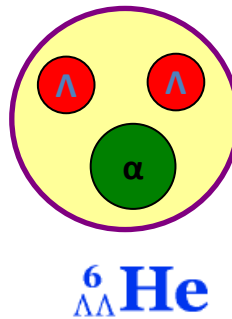
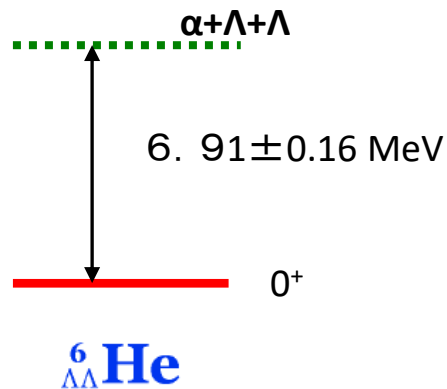
Ambiguity for identifying these double  $\Lambda$  hypernuclei

There was NO observed double  $\Lambda$  hypernuclei  
**without ambiguity.**

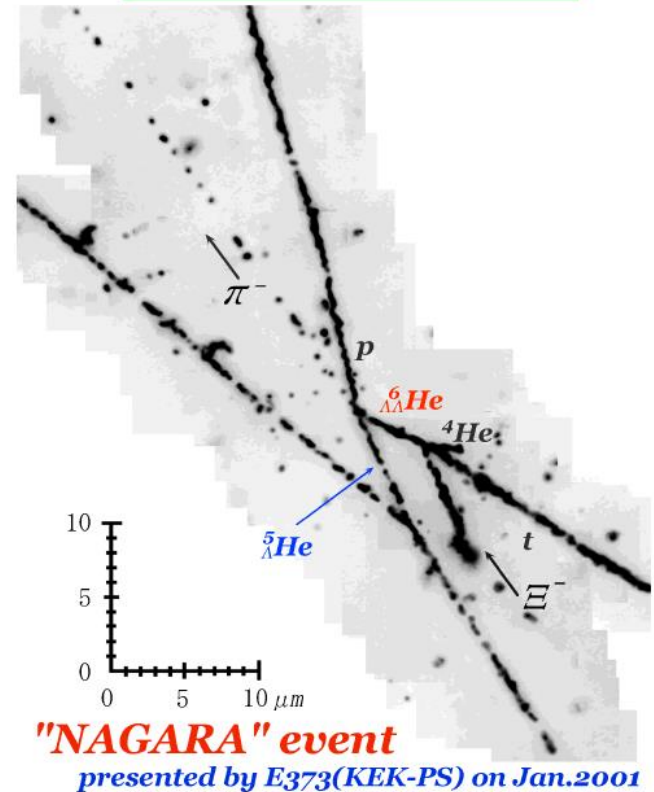
In 2001, the epoch-making data has been reported by the KEK-E373 experiment.

Observation of  ${}^6_{\Lambda\Lambda}\text{He}$

Uniquely identified without ambiguity for the first time



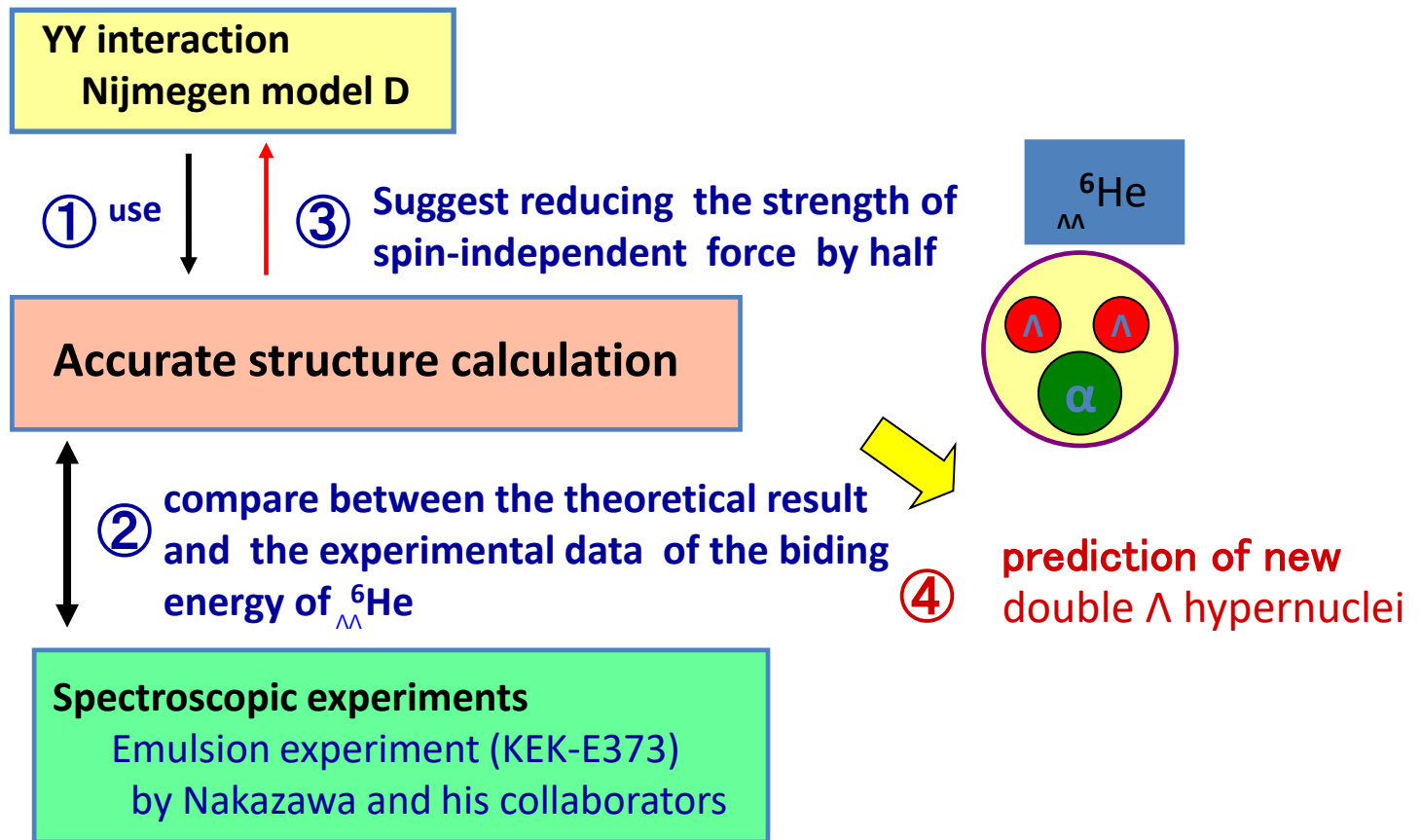
${}^6_{\Lambda\Lambda}\text{He}$  double-hypernucleus  
Unique interpretation!!



"NAGARA" event  
presented by E373(KEK-PS) on Jan.2001

H. Takahashi et al., PRL 87, 212502-1 (2001)

# Strategy of how to determine YY interaction from the study of light hypernuclear structure

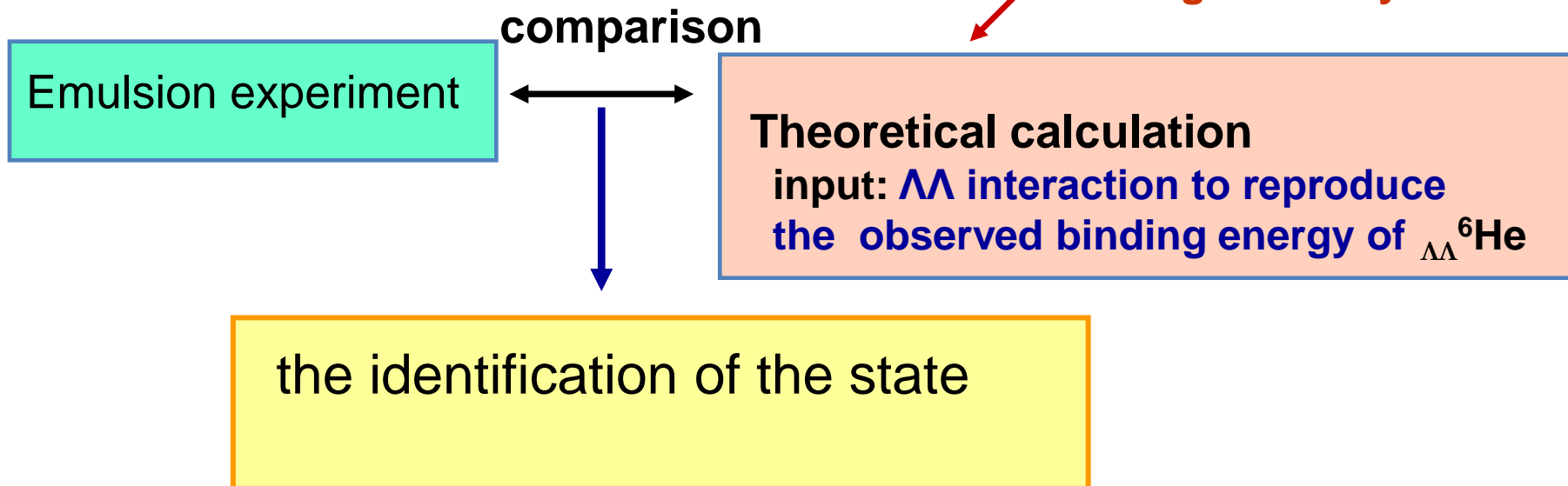




- E07 Approved proposal at J-PARC  
“Systematic Study of double strangeness systems at J-PARC”  
by Nakazawa and his collaborators(done in 2018)

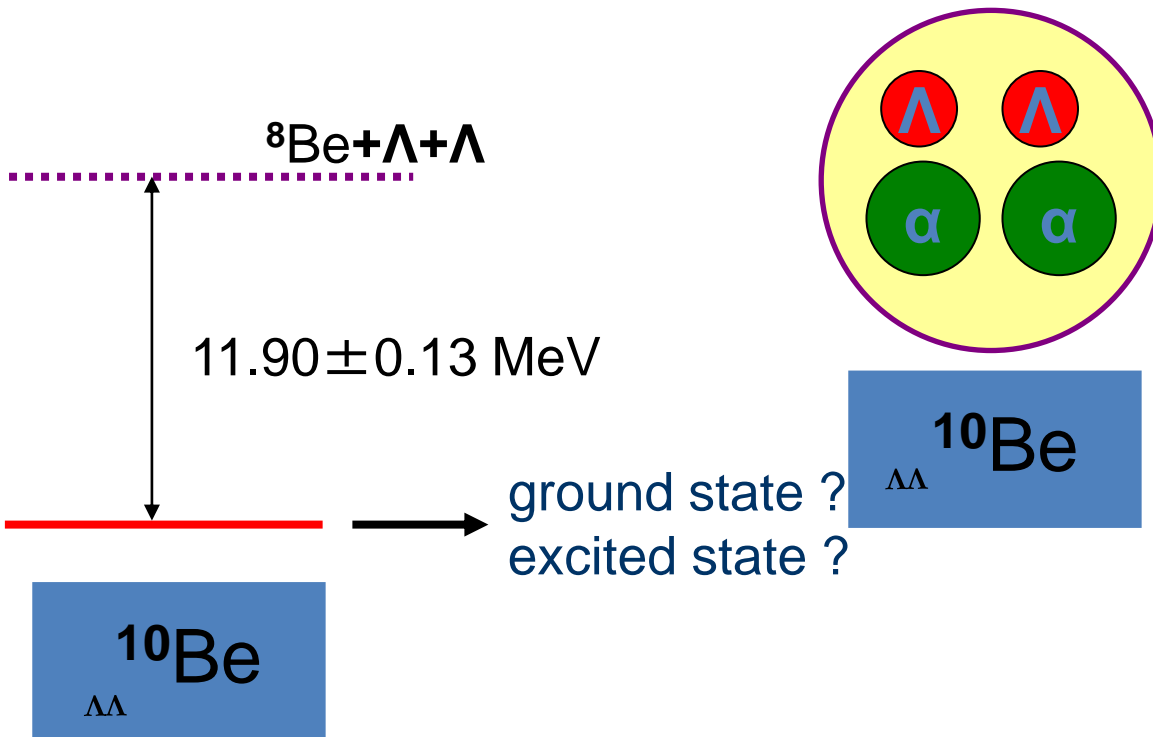
It is difficult to determine { (1) spin-parity  
(2) whether the observed state is the ground state or an excited state

My theoretical contribution using few-body calculation

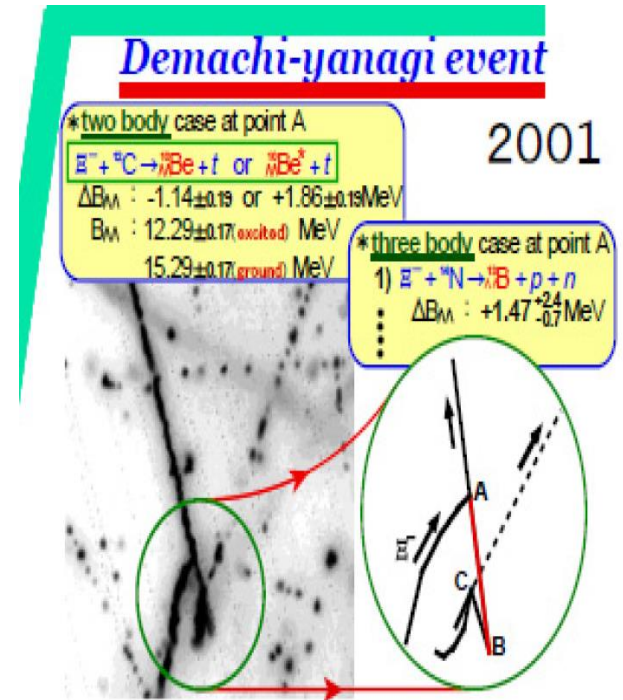


Successful example to determine spin-parity of double  $\Lambda$  hypernucleus --- Demachi-Yanagi event for  ${}^{10}_{\Lambda\Lambda}\text{Be}$

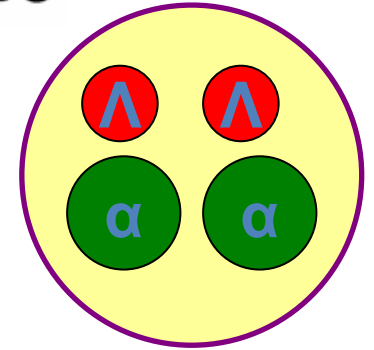
Observation of  ${}^{10}_{\Lambda\Lambda}\text{Be}$  --- KEK-E373 experiment



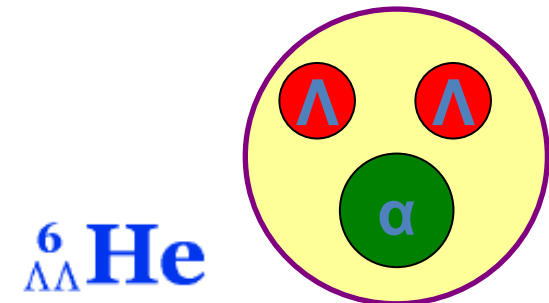
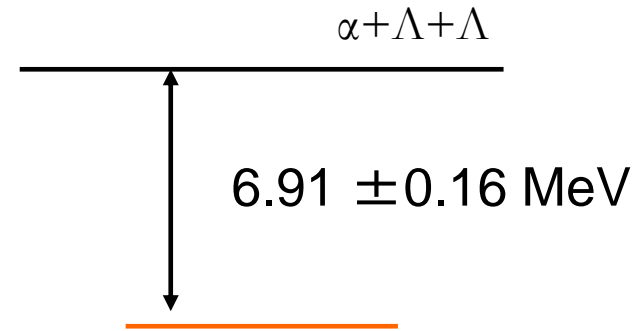
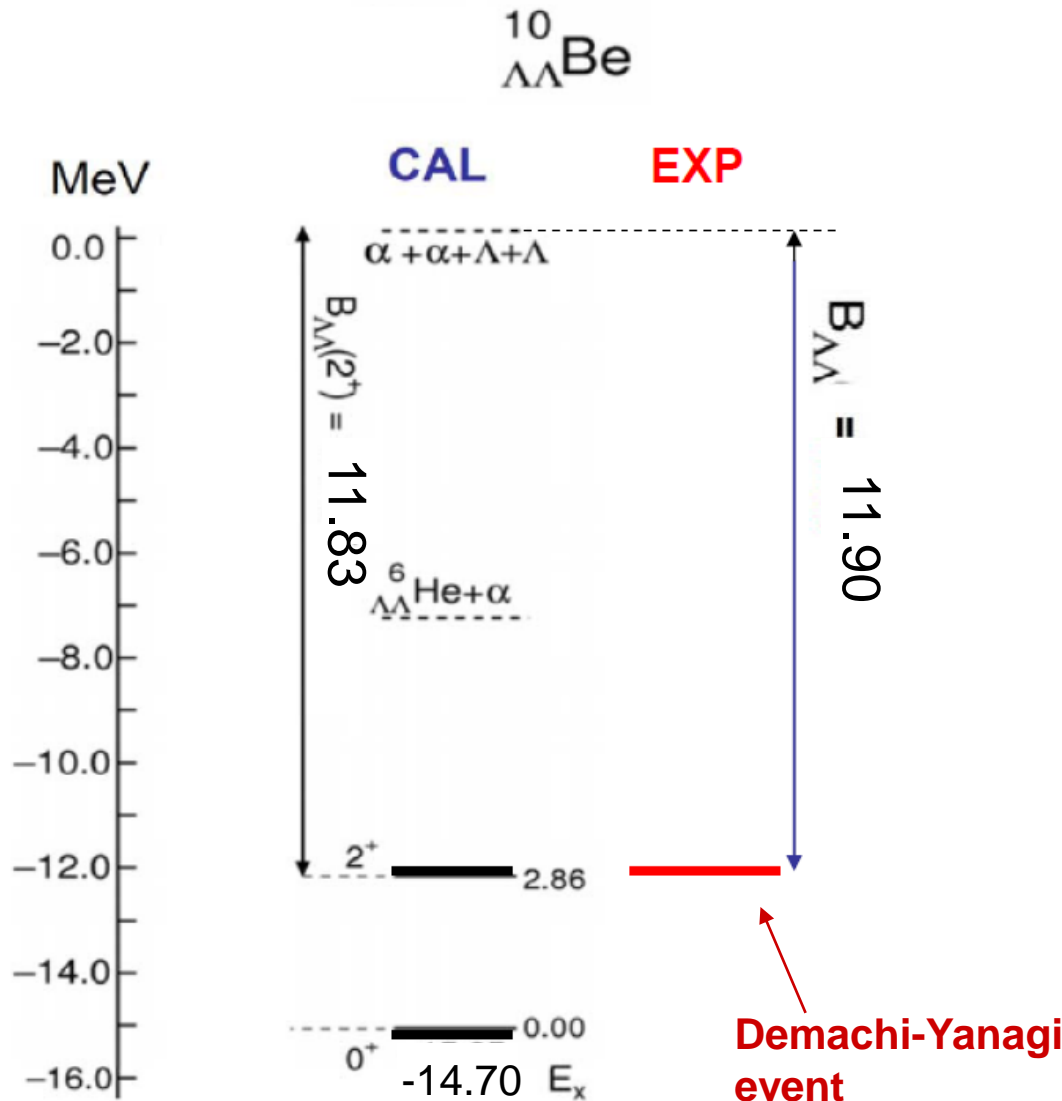
Demachi-Yanagi event



# Successful interpretation of spin-parity of $^{10}_{\Lambda\Lambda}\text{Be}$



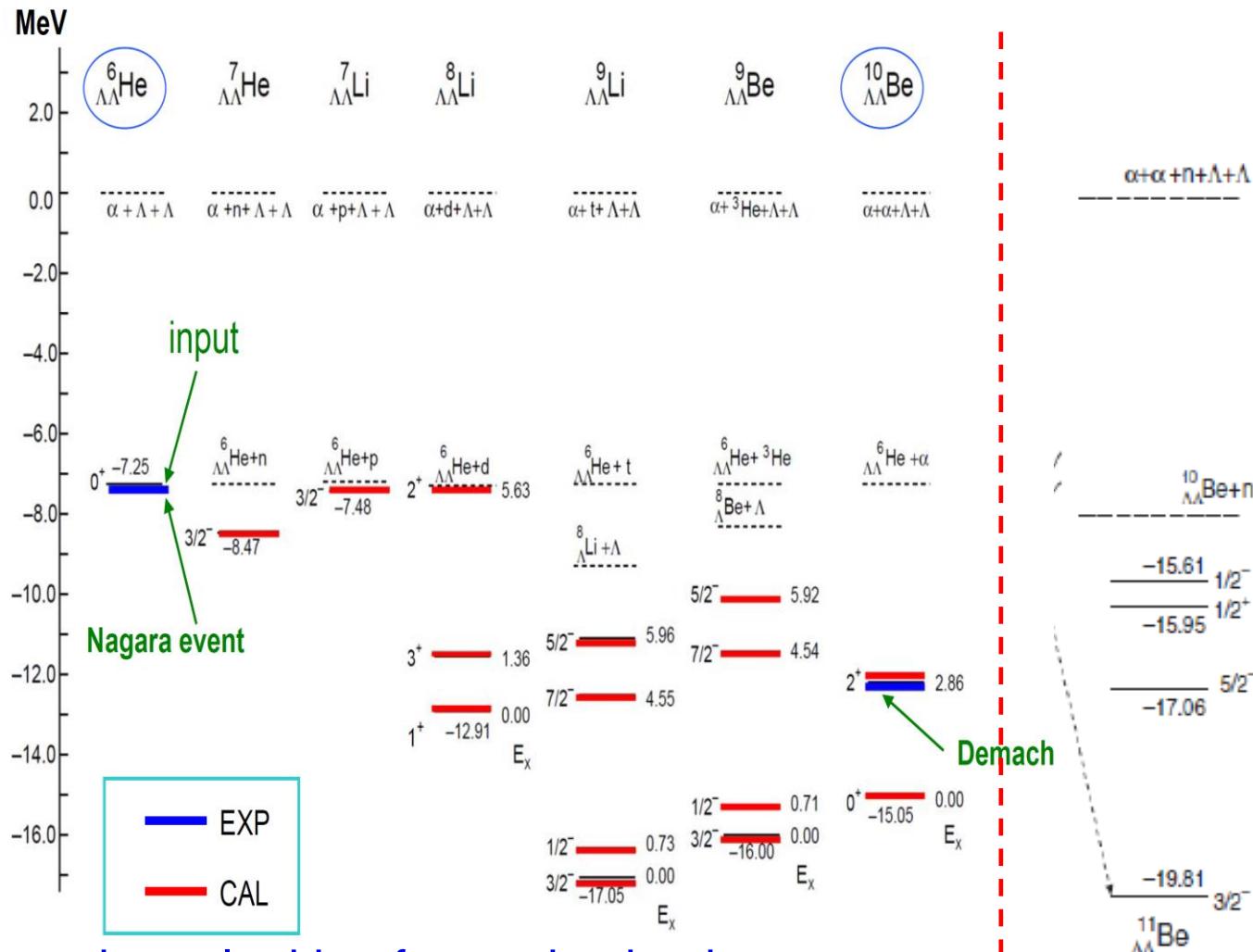
E. Hiyama, M. Kamimura, T. Motoba,  
T. Yamada and Y. Yamamoto  
Phys. Rev. 66 (2002), 024007



${}^6_{\Lambda\Lambda}\text{He}$

# Spectroscopy of $\Lambda\Lambda$ -hypernuclei

E. Hiyama, M. Kamimura, T. Motoba, T. Yamada and Y. Yamamoto  
 Phys. Rev. 66 (2002), 024007

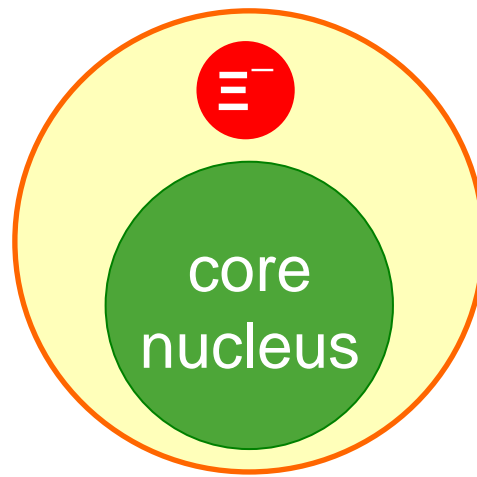


New data:  
 -20.64 MeV

H. Ekawa et al.  
 PTEP 2019, 021D  
 (2019)

I have been looking forward to having new data in this mass-number region.

$\Xi$  hypernuclei

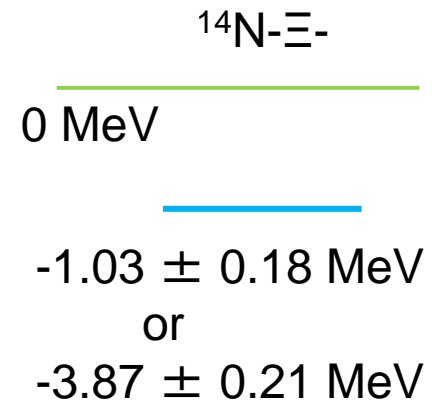
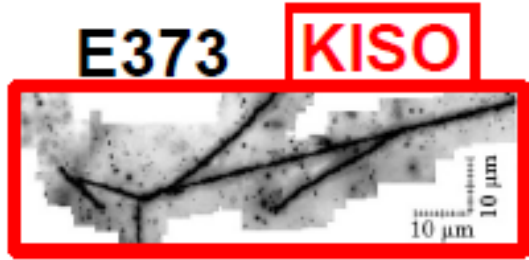


For the study of  $\Xi N$  interaction, it is important to study the structure of  $\Xi$  hypernuclei.

However, so far there was no observed  $\Xi$  hypernucleus. Therefore, we do not know that  $\Xi N$  interaction is attractive or repulsive.

If we observe  $\Xi$  hypernuclei as bound states, we understand  $\Xi N$  interaction should be attractive. Thus, we have been searching bound  $\Xi$  hypernuclei experimentally.

# The first measurement of bound $\Xi$ hypernucleus, $^{14}\text{N}-\Xi$ .



**PTEP**

Prog. Theor. Exp. Phys. **2015**, 033D02 (11 pages)  
DOI: 10.1093/ptep/ptv008

## The first evidence of a deeply bound state of $\Xi^- - ^{14}\text{N}$ system

K. Nakazawa<sup>1,\*</sup>, Y. Endo<sup>1</sup>, S. Fukunaga<sup>2</sup>, K. Hoshino<sup>1</sup>, S. H. Hwang<sup>3</sup>, K. Imai<sup>3</sup>, H. Ito<sup>1</sup>, K. Itonaga<sup>1</sup>, T. Kanda<sup>1</sup>, M. Kawasaki<sup>1</sup>, J. H. Kim<sup>4</sup>, S. Kinbara<sup>1</sup>, H. Kobayashi<sup>1</sup>, A. Mishina<sup>1</sup>, S. Ogawa<sup>2</sup>, H. Shibuya<sup>2</sup>, T. Sugimura<sup>1</sup>, M. K. Soe<sup>1</sup>, H. Takahashi<sup>5</sup>, T. Takahashi<sup>5</sup>, K. T. Tint<sup>1</sup>, K. Umehara<sup>1</sup>, C. S. Yoon<sup>4</sup>, and J. Yoshida<sup>1</sup>

<sup>1</sup>Physics Department, Gifu University, 1-1 Yanagido, Gifu 501-1193, Japan

<sup>2</sup>Department of Physics, Toho University, Funabashi 274-8510, Japan

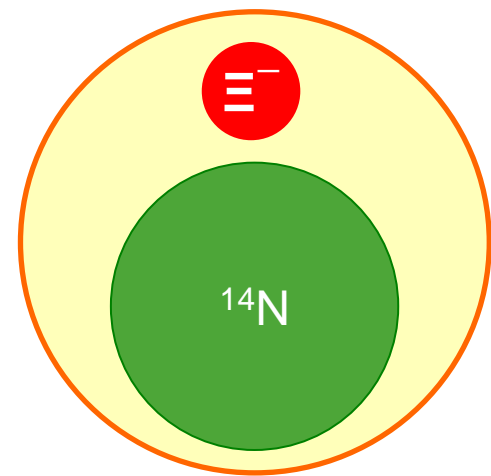
<sup>3</sup>Advanced Science Research Center, JAEA, Tokai 319-1195, Japan

<sup>4</sup>Department of Physics, Gyeongsang National University, Jinju 660-701, Korea

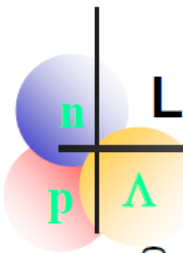
<sup>5</sup>Institute of Particle and Nuclear Studies, KEK, Tsukuba 305-0801, Japan

\*E-mail: nakazawa@gifu-u.ac.jp

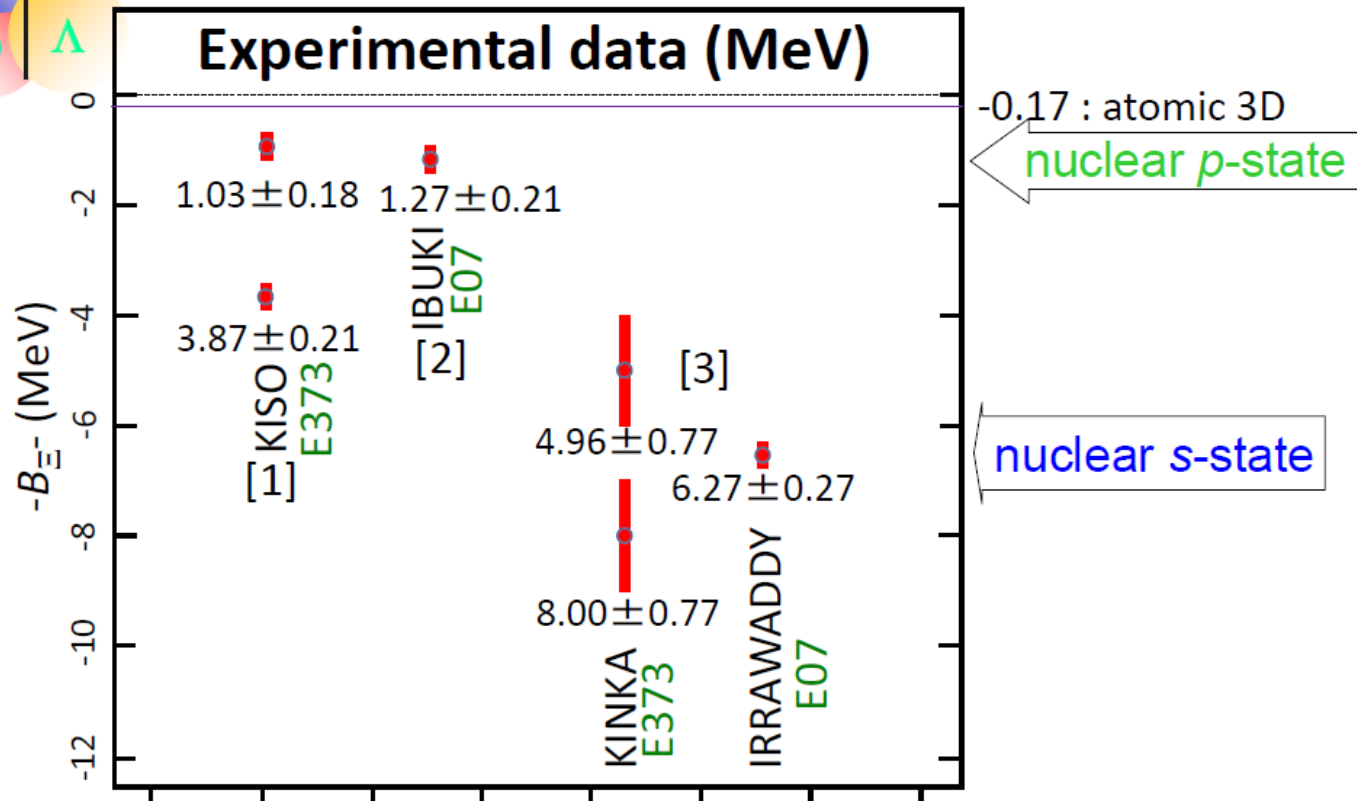
Received October 27, 2014; Revised December 25, 2014; Accepted January 9, 2015; Published March 5, 2015



We understood  $\Xi$ -nuclear potential should be attractive.



# Level scheme of $\Xi$ hypernucleus ( ${}^{15}_{\Xi}\text{C}$ [ $\Xi^{-}$ - ${}^{14}\text{N}$ ])



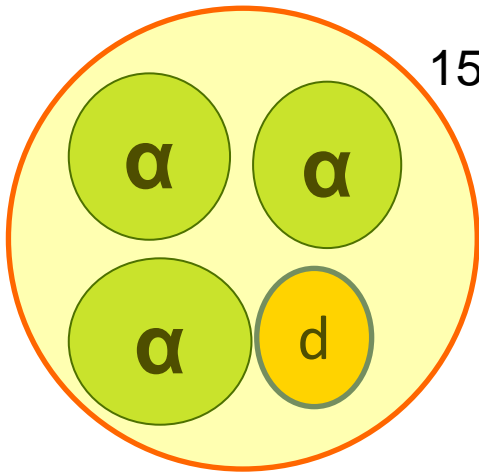
- [1] K. Nakazawa, et. al., Prog. Theor. Exp. Phys. 2015, 033D02 (2015), E. Hiyama and K. Nakazawa, Ann. Rev. Nucl. Part. Sci. 68, 131 (2018).
- [2] S. Hayakawa, et. al., Phy. Rev. Lett., 126, 062501 (2021).
- [3] M. Yoshimoto, et. al., Prog. Theor. Exp. Phys. 2021, 073D02 (2021).

Slide by Nakazawa

After observation of Kiso event, they observed several events of  ${}^{14}\text{N}-\Xi$  hypernucleus. Some are observed as excited state and some are observed as ground state.



$$V_{\Xi N} = V_0 + \sigma \cdot \sigma V_{\sigma \cdot \sigma} + \tau \cdot \tau V_{\tau \cdot \tau} + (\sigma \cdot \sigma)(\tau \cdot \tau) V_{\sigma \cdot \sigma \tau \cdot \tau}$$



$^{15}\text{C}$

By observation of  $^{15}\text{C}(^{14}\text{N}-\Xi)$ , we find that  $V_{\Xi N}$  itself is attractive.

Because,

All of the terms contribute to binding energy of  $^{15}\text{C}$  ( $^{14}\text{N}$  is not spin-, isospin- saturated).

Next,

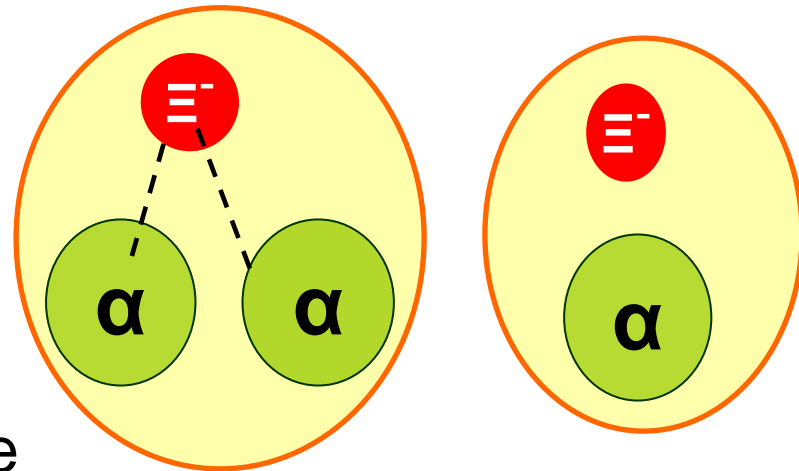
we want to know desirable strength of  $V_0$ , the spin-, isospin-independent term.

$$V_{\Xi N} = V_0 + \sigma \cdot \sigma V_{\sigma \cdot \sigma} + \tau \cdot \tau V_{\tau \cdot \tau} + (\sigma \cdot \sigma)(\tau \cdot \tau) V_{\sigma \cdot \sigma \tau \cdot \tau}$$

In order to obtain useful information about  $V_0$ , the following systems are suited, because

the  $(\sigma \cdot \sigma)$ ,  $(\tau \cdot \tau)$  and  $(\sigma \cdot \sigma)(\tau \cdot \tau)$  terms of  $V_{\Xi N}$  vanish

by folding them into the  $\alpha$ -cluster wave function that are spin-, isospin-saturated.



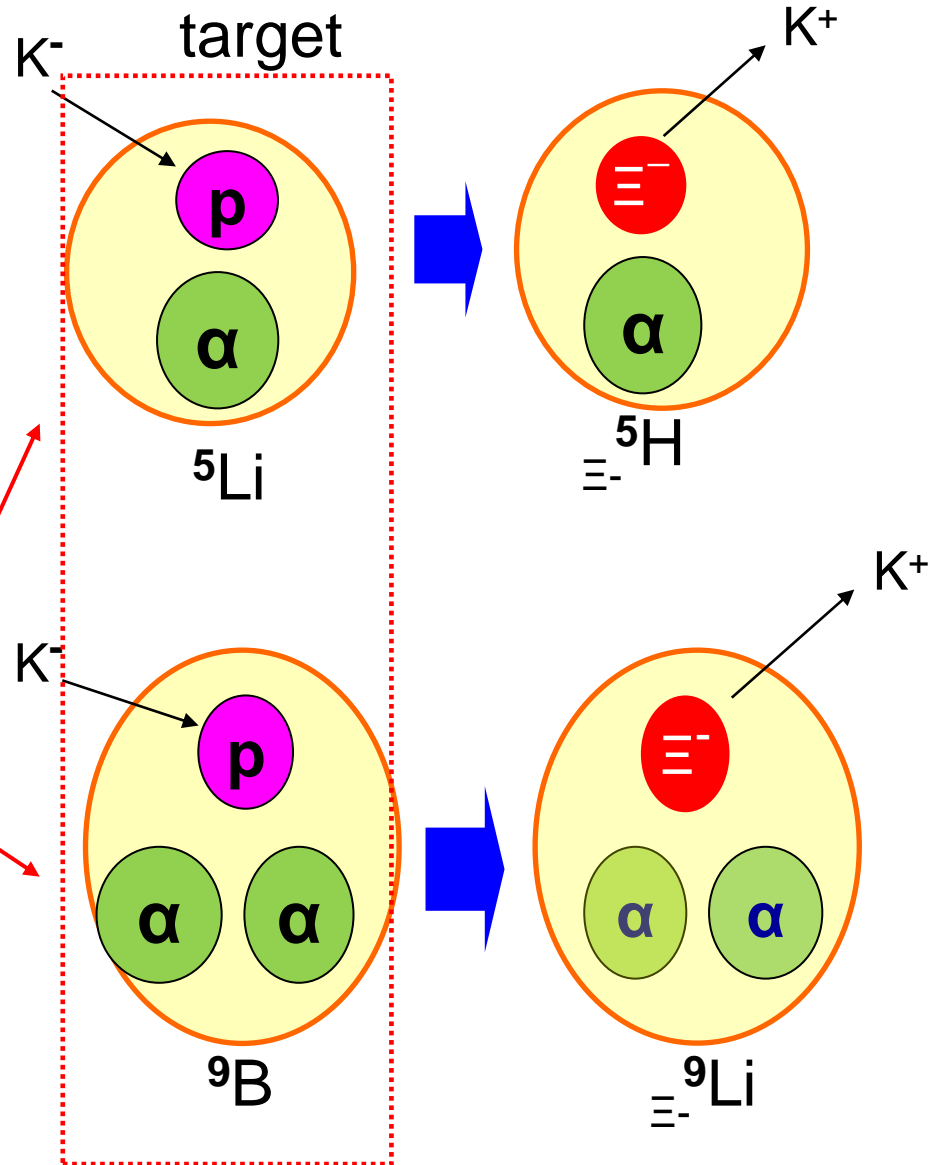
problem : there is NO target to produce them by the  $(K^-, K^+)$  experiment .

Because, ...

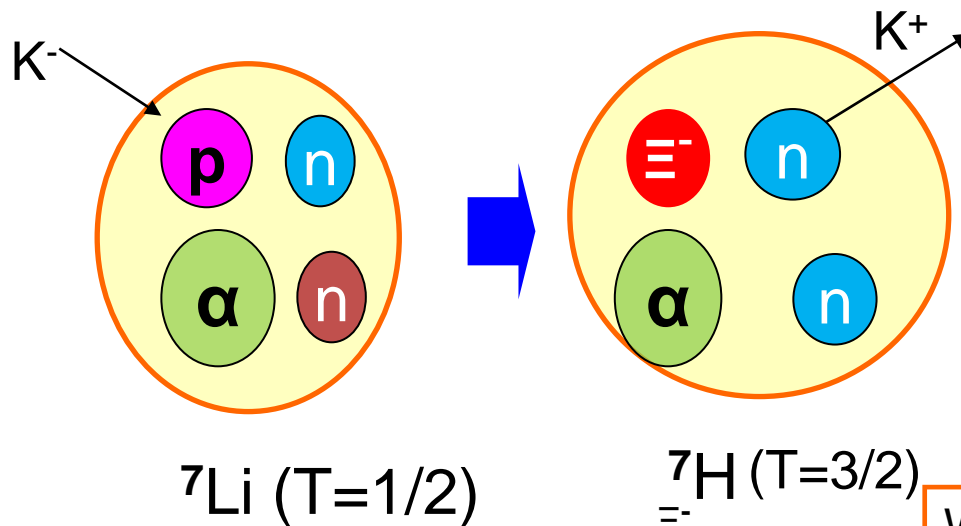
To produce  $\alpha\Xi^-$  and  $\alpha\alpha\Xi^-$  systems by  $(K^-, K^+)$  reaction,

These systems are unbound.

Then, we cannot use them as targets.

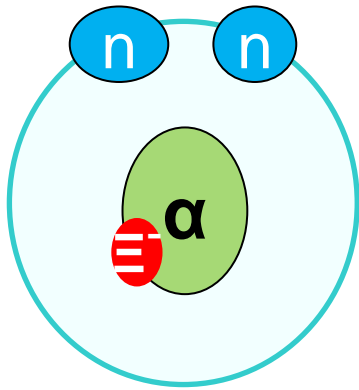


As the second best candidates to extract information about the spin-, isospin-independent term  $V_0$ , we propose to perform...



Why they are suited for investigating  $V_0$ ?

(more realistic illustration)



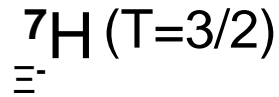
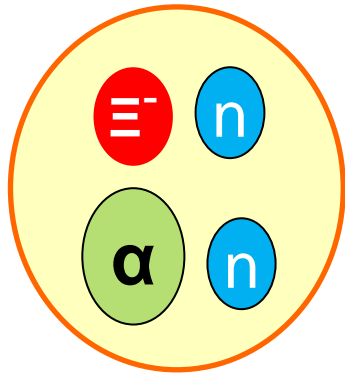
Core nucleus  ${}^6\text{He}$  is known to be halo nucleus. Then, valence neutrons are located far away from  $\alpha$  particle.

Valence neutrons  $n$  are located in p-orbit, whereas  ${}^3\text{H}$  particle  ${}^3\text{H}$  is located in 0s-orbit.

${}^7\text{H}$  ( $T=3/2$ )  
 ${}^3\text{H}$

Then, distance between  ${}^3\text{H}$  and  $n$  is much larger than the interaction range of  ${}^3\text{H}$  and  $n$ .

Then,  $\alpha$ - ${}^3\text{H}$  potential, in which only  $V_0$  term works, plays a dominant role in the binding energies of this system.



Before the experiments will be done,  
we should predict whether this  $\Xi$  hypernucleus will be observed as bound states or not.

Namely, we calculate the binding energies of this hypernucleus.

## $\Xi$ N interaction

• **ESC04** (Nijmegen soft core) and **ND** (Nijmegen Model D)

**HAL potential (based on Lattice QCD)**

$$V_{\Xi N} = V_0(r) + (\sigma_{\Xi} \cdot \sigma_N) V_s(r) + (\tau_{\Xi} \cdot \tau_N) V_t(r) + (\sigma_{\Xi} \cdot \sigma_N)(\tau_{\Xi} \cdot \tau_N) V_{ts}(r)$$

All terms are central parts only.

## Property of the spin- and isospin-components of ESC04, ND, HAL

V(T,S)	ESC04	ND	HAL
T=0, S=1	strongly attractive (a bound state)	} weakly attractive	Weakly attractive
T=0, S=0	weakly repulsive		Strongly attractive
T=1, S=1	weakly attractive		Weakly attractive
T=1, S=0	weakly repulsive		Weakly repulsive

Although the spin- and isospin-components of these models are very different (due to the different meson contributions), we find that the spin- and isospin-averaged property,

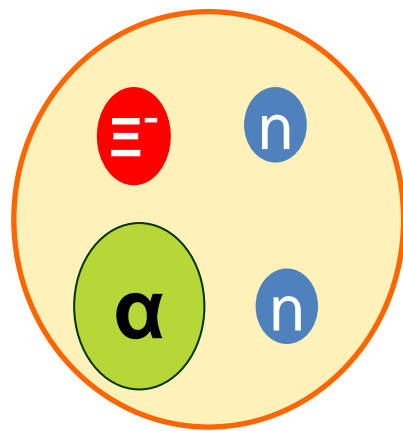
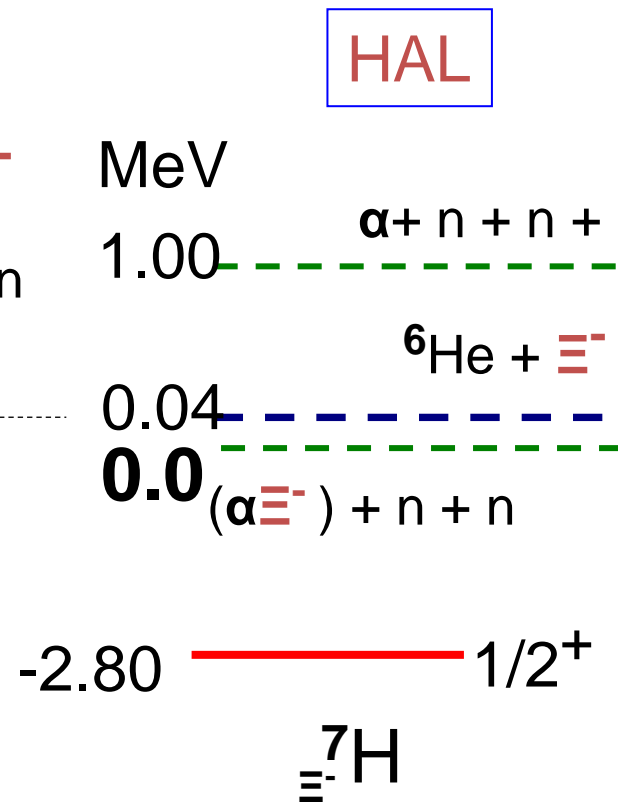
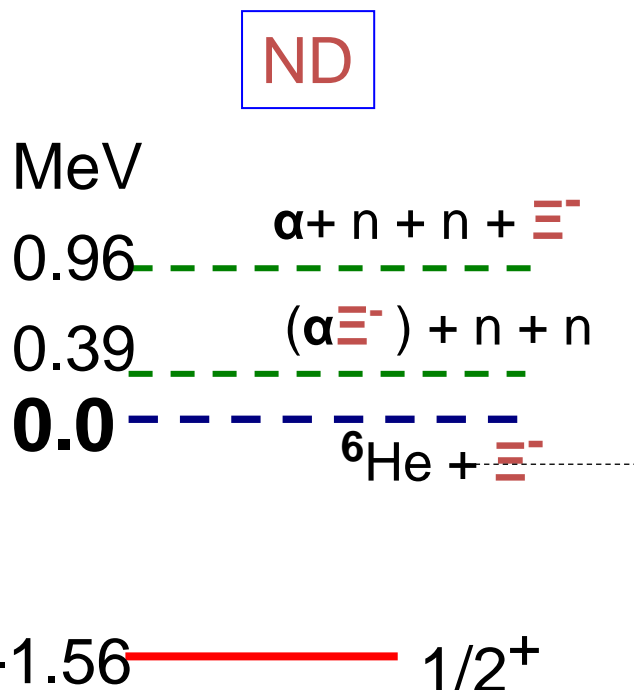
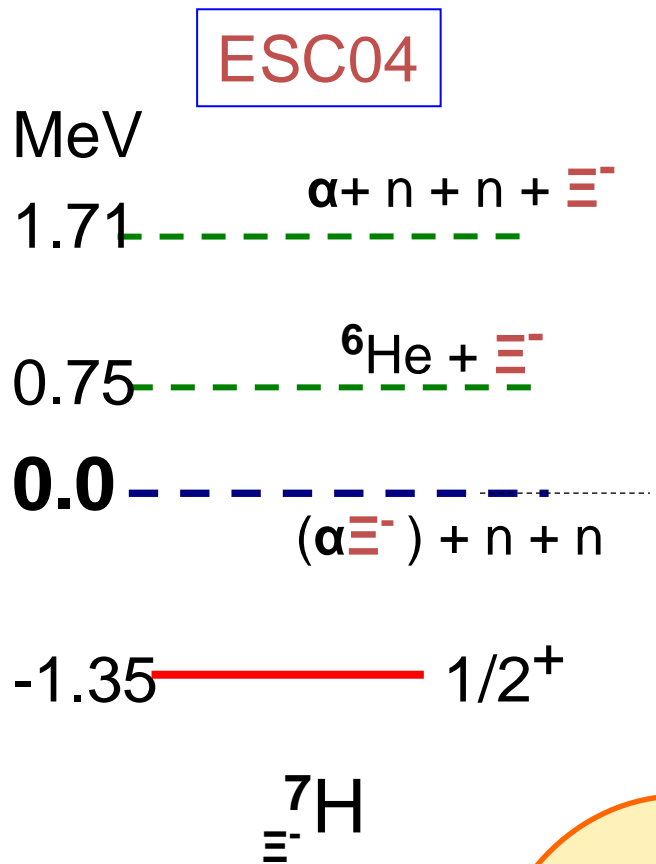
$$V_0 = [ V(0,0) + 3V(0,1) + 3V(1,0) + 9V(1,1) ] / 16,$$

namely, strength of the  $V_0$ - term is similar to each other.



# 4-body calculation of ${}^7_{\Xi^-}\text{H}$

E. Hiyama et al.,  
PRC78 (2008) 054316



In experiments,  
we can expect  
a bound state.

Similar binding  
energies using ND and  
ESC04.

$$V_0 = [ V(0,0) + 3V(0,1) + 3V(1,0) + 9V(1,1) ] / 16,$$

which partial contribution makes attractive for  $V_0$  ?

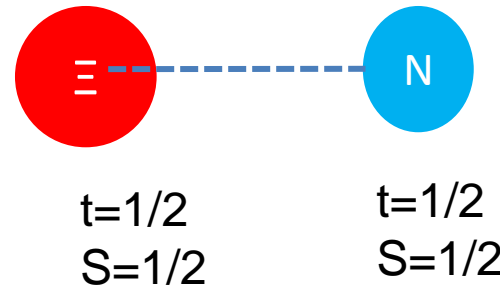
$\Xi$ N interaction:

$T=0, S=0$

$T=0, S=1$

$T=1, S=0$

$T=1, S=1$



Cf. NN interaction

we have a two-body bound state for  $\Xi$ N system?

No idea



$T=0, S=0, l=\text{odd}$

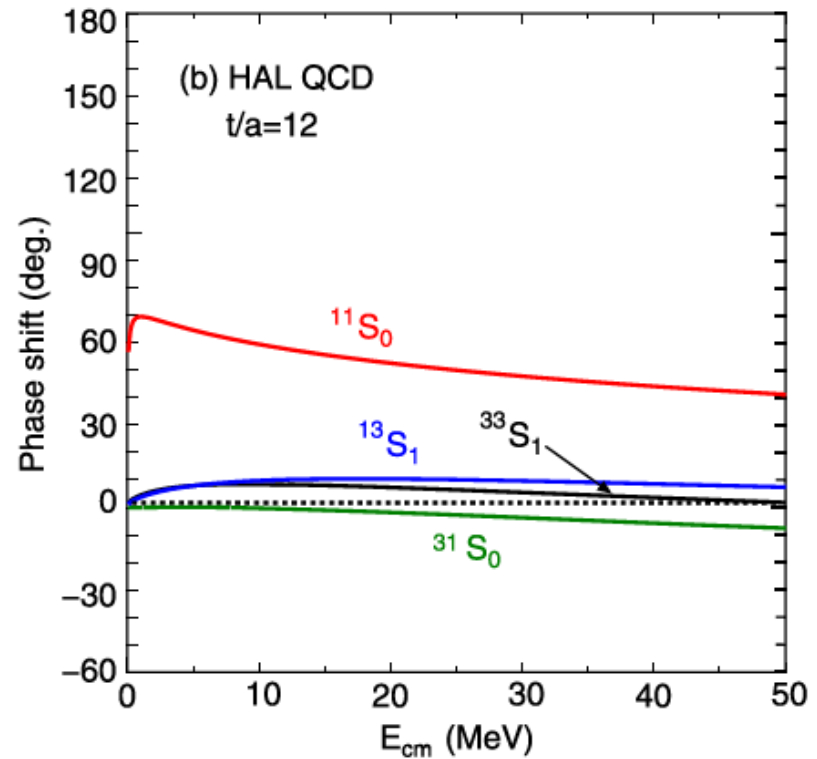
$T=0, S=1$   strong attraction to have a bound state as a deuteron

$T=1, S=0$

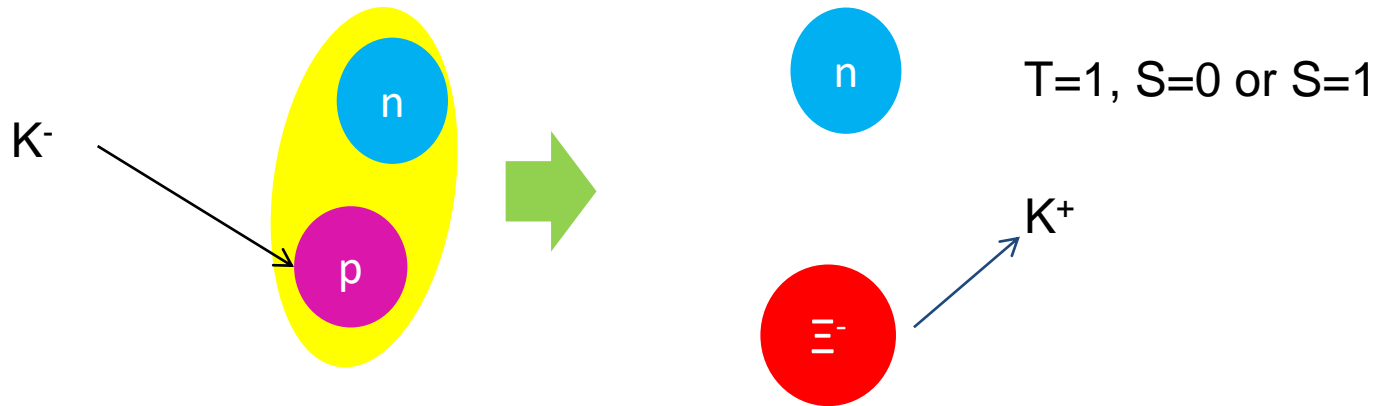
$T=1, S=1, l=\text{odd}$

# Property of the spin- and isospin-components of HAL

$V(T,S)$	HAL
$T=0, S=1$ ( $^{13}S_1$ )	Weakly attractive
$T=0, S=0$ ( $^{11}S_0$ )	Strongly attractive
$T=1, S=1$ ( $^{33}S_1$ )	Weakly attractive
$T=1, S=0$ ( $^{31}S_0$ )	Weakly repulsive



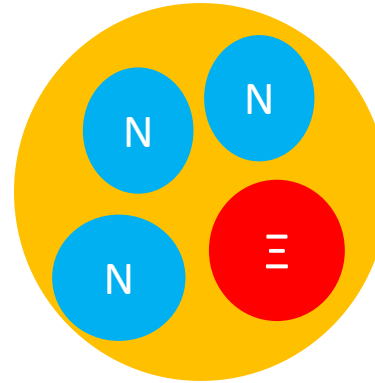
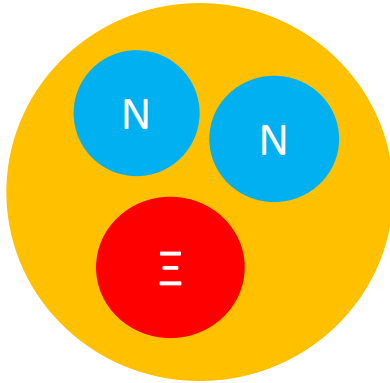
To investigate bound state of  $\Xi N$  system, it might be possible to perform the following experiment:



It would be difficult to obtain information on  $\Xi N$  interaction ( $T=1, S=0$  or  $1$ ). Because, there might be no bound state for this system.

To obtain  $\Xi N$  two-body interaction, the suited systems to study are s-shell  $\Xi$  hypernuclei such as  $NN\Xi$  and  $NNN\Xi$  systems.

E. Hiyama et al., PRL124, 092501 (2020)



I show my results of these light systems.

NN interaction: AV8 potential

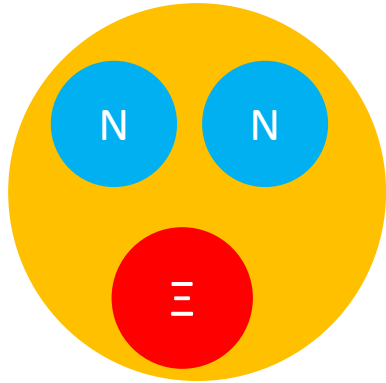
$\Xi N$  interaction :

Nijmegen extended soft core potential (ESC08c)

Realistic potential (only  $\Xi N$  channel)

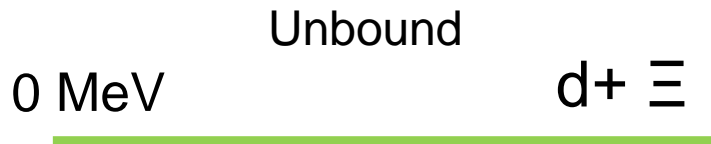
$\Xi N$  interaction by HAL collaboration (Lattice QCD calculation)

The potential was made by K. Sasaki, Miyamoto, Hatsuda and Aoki.

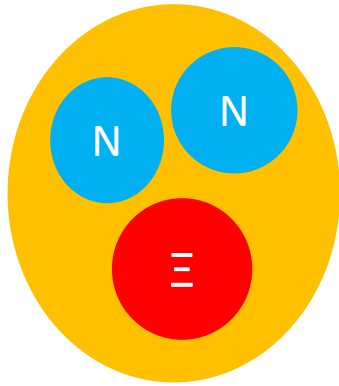


$T=1/2, J=1/2^+$  and  $J=3/2^+$

ESC08c



However, I also have two bound states in three-body system.



$T=1/2$ ,  $J=1/2^+$  and  $J=3/2^+$

HAL potential

0 MeV

$d + \Xi$

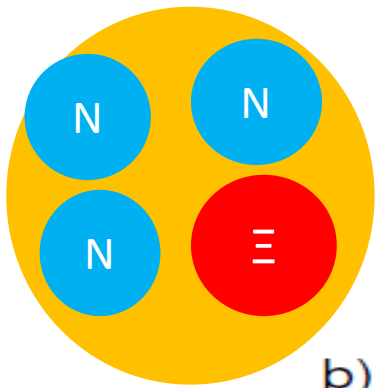
0 MeV

$d + \Xi$

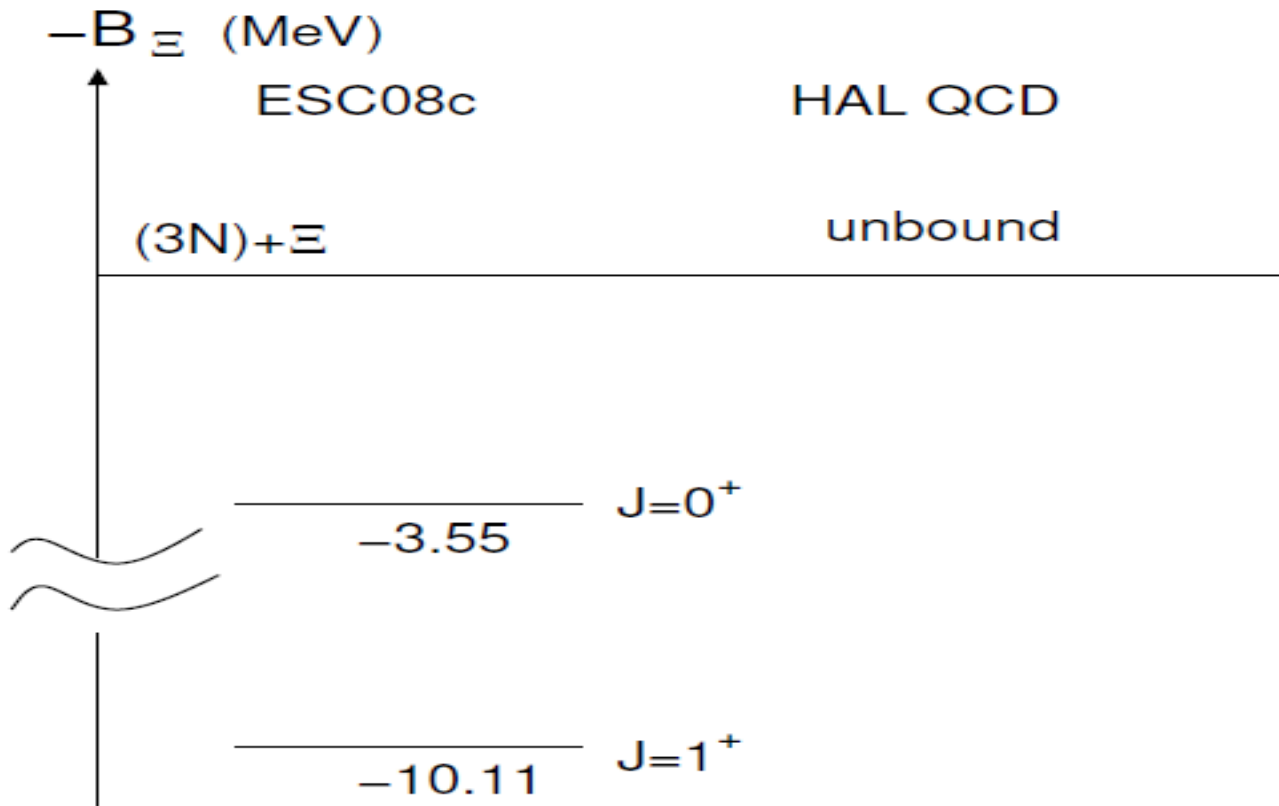
No bound state

$J=1/2^+$

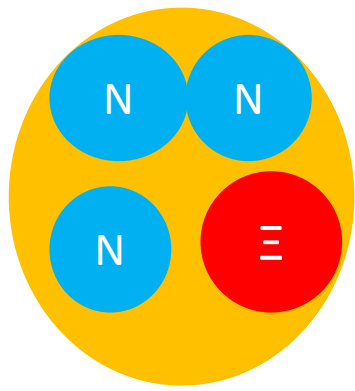
$J=3/2^+$



b)  $T=1$







a)  $T=0$

$-B_{\Xi}$  (MeV)

ESC08c

HAL QCD

$(3N)+\Xi$

0.36 (36)

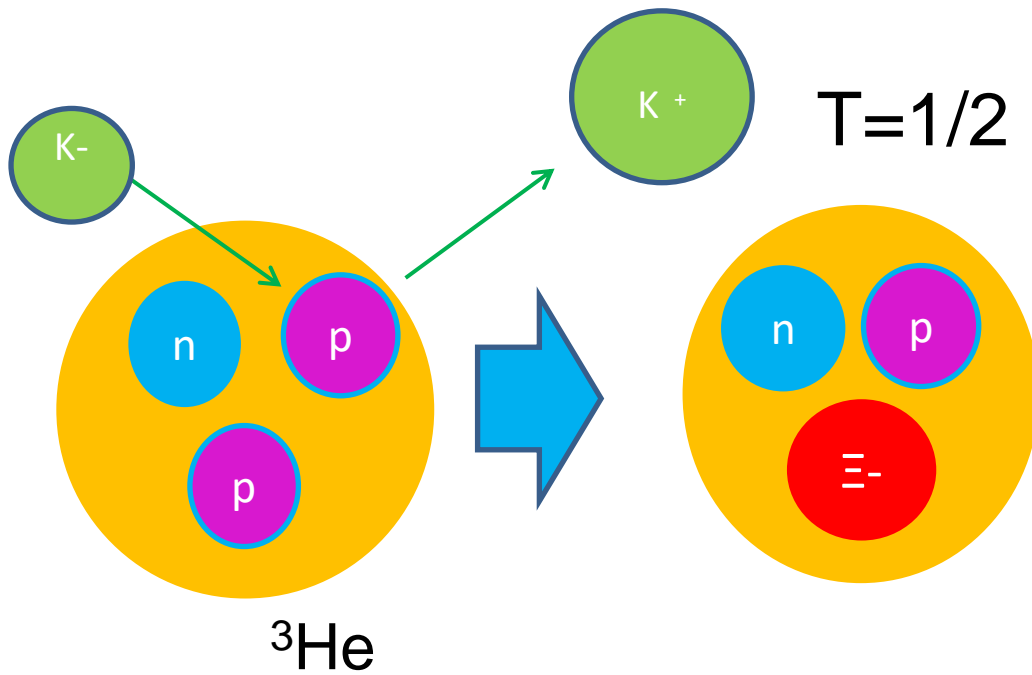
$J=1^+$



-10.20  $J=1^+$

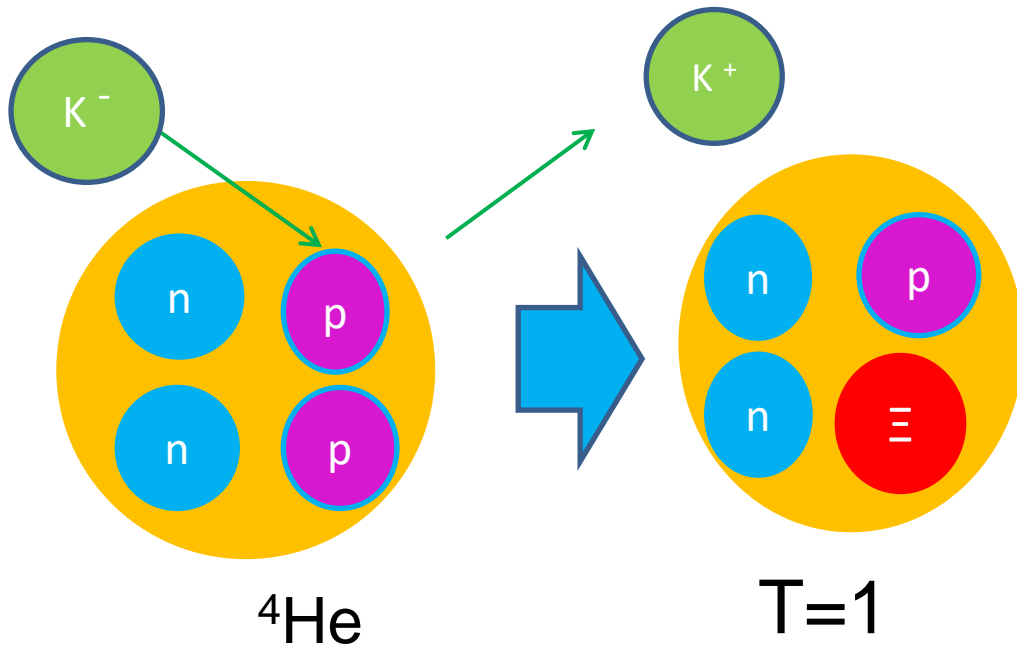
H.Le et al., EPJA57, 339 (2021).

		$B_{\Xi}$ [MeV]	$\Gamma$ [MeV]
${}^4_{\Xi}\text{H}(1^+, 0)$	NNN $\Xi$	$0.48 \pm 0.01$	0.74
${}^4_{\Xi}\text{n}(0^+, 1)$		$0.71 \pm 0.08$	0.2
${}^4_{\Xi}\text{n}(1^+, 1)$		$0.64 \pm 0.11$	0.01
${}^4_{\Xi}\text{H}(0^+, 0)$		—	—



Using  ${}^3\text{He}$  and  ${}^4\text{He}$  target,  
It might be possible  
to produce  $\text{NN}\Xi$  and  $\text{NNN}\Xi$   
systems by  $(K^-, K^+)$  reaction.

Another tool is to use  
Heavy ion collision.

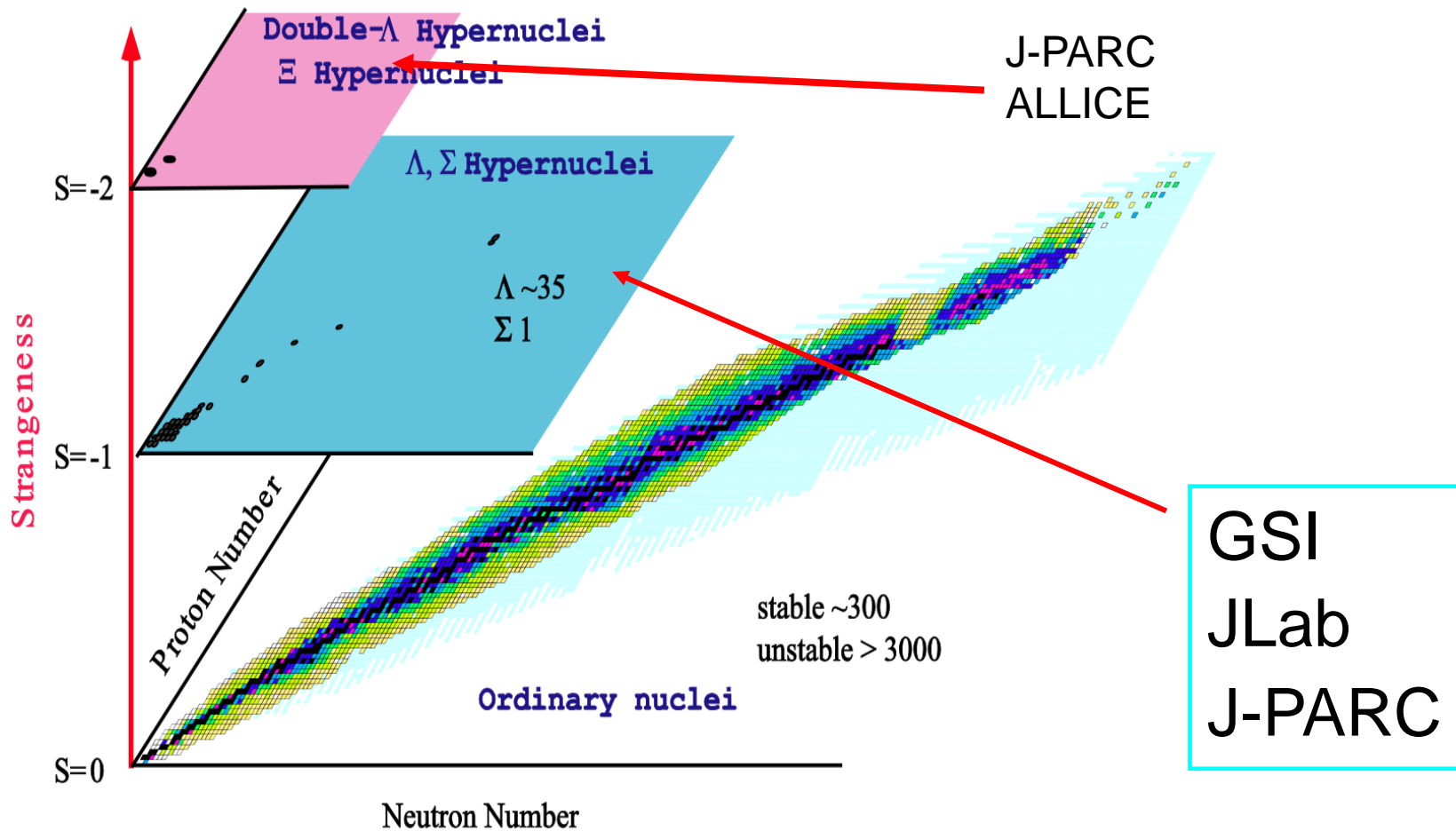


In the future,  
we hope to observe  
these light  $\Xi$  hypernuclei.

# Concluding remark

Multi-strangeness system  
such as Neutron star

Three-Dimensional Nuclear Chart



Thank you