

Unraveling Universal Correlations

Gaussian Characterization of Systems Near the Unitary Limit

Mario Gattobigio



Institut de Physique de Nice

Magonza, 31 July 2023



Outline

(Enlarged) Unitary Window

S-matrix

Effective Range Expansion

Zero-shape Universality

Gaussian (Eckart) characterization

Three-body sector

Efimov Effect

Level Functions - Gaussian Characterization

Moving along the universal curve

Note on DSI

More particles

LO Gaussian Potential - Two- and Three-Body Force

References

Outline

(Enlarged) Unitary Window

S-matrix

Effective Range Expansion

Zero-shape Universality

Gaussian (Eckart) characterization

Three-body sector

Efimov Effect

Level Functions - Gaussian Characterization

Moving along the universal curve

Note on DSI

More particles

LO Gaussian Potential - Two- and Three-Body Force

References

S-Matrix

- Two-body Schrödinger equation $\iff S$ -matrix

S-Matrix

- Two-body Schrödinger equation $\iff S$ -matrix
- Simplest S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

S-Matrix

- Two-body Schrödinger equation $\iff S\text{-matrix}$
- Simplest S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Physical pole $k = i/a_B$

$$B_2 = \frac{\hbar^2}{ma_B^2} \begin{cases} \text{Bound state} & \text{if } a_B > 0 \\ \text{Virtual state} & \text{if } a_B < 0 \end{cases}$$

S-Matrix

- Two-body Schrödinger equation $\iff S\text{-matrix}$
- Simplest S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- ▶ Physical pole $k = i/a_B$

$$B_2 = \frac{\hbar^2}{ma_B^2} \begin{cases} \text{Bound state} & \text{if } a_B > 0 \\ \text{Virtual state} & \text{if } a_B < 0 \end{cases}$$

- ▶ Spurious pole $k = i/r_B$

r_B = Dimensional Constant (Scale)

S-Matrix

- Two-body Schrödinger equation $\iff S\text{-matrix}$
- Simplest S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- ▶ Physical pole $k = i/a_B$

$$B_2 = \frac{\hbar^2}{ma_B^2} \begin{cases} \text{Bound state} & \text{if } a_B > 0 \\ \text{Virtual state} & \text{if } a_B < 0 \end{cases}$$

- ▶ Spurious pole $k = i/r_B$

r_B = Dimensional Constant (Scale)

- ▶ Scaling (zero-range) limit

$r_B/a_B \rightarrow 0$ with a_B fixed

$$S(k) \sim -\frac{k + i/a_B}{k - i/a_B}$$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

$$T(k) = -\frac{4\pi}{m} \left(-\frac{1}{a} + \frac{1}{2} r_e k^2 + v_2 k^4 + v_3 k^6 + \dots - ik \right)^{-1}$$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

$$T(k) = -\frac{4\pi}{m} \left(-\frac{1}{a} + \frac{1}{2} r_e k^2 + v_2 k^4 + v_3 k^6 + \dots - ik \right)^{-1}$$

- Parameter Identification

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

$$T(k) = -\frac{4\pi}{m} \left(-\frac{1}{a} + \frac{1}{2} r_e k^2 + v_2 k^4 + v_3 k^6 + \dots - ik \right)^{-1}$$

- Parameter Identification

- ▶ Shape parameters vanish - $v_n = 0$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

$$T(k) = -\frac{4\pi}{m} \left(-\frac{1}{a} + \frac{1}{2} r_e k^2 + v_2 k^4 + v_3 k^6 + \dots - ik \right)^{-1}$$

- Parameter Identification

- ▶ Shape parameters vanish - $v_n = 0$
- ▶ r_B relates the scattering lenght and the physical pole

$$r_B = a - a_B$$

Effective Range Expansion

- S -matrix

$$S(k) = \frac{k + i/a_B}{k - i/a_B} \frac{k + i/r_B}{k - i/r_B}$$

- Relation with the scattering

$$S(k) = 1 - i \frac{k m}{2\pi} T(k)$$

$$T(k) = -\frac{4\pi}{m} \left(-\frac{1}{a} + \frac{1}{2} r_e k^2 + v_2 k^4 + v_3 k^6 + \dots - ik \right)^{-1}$$

- Parameter Identification

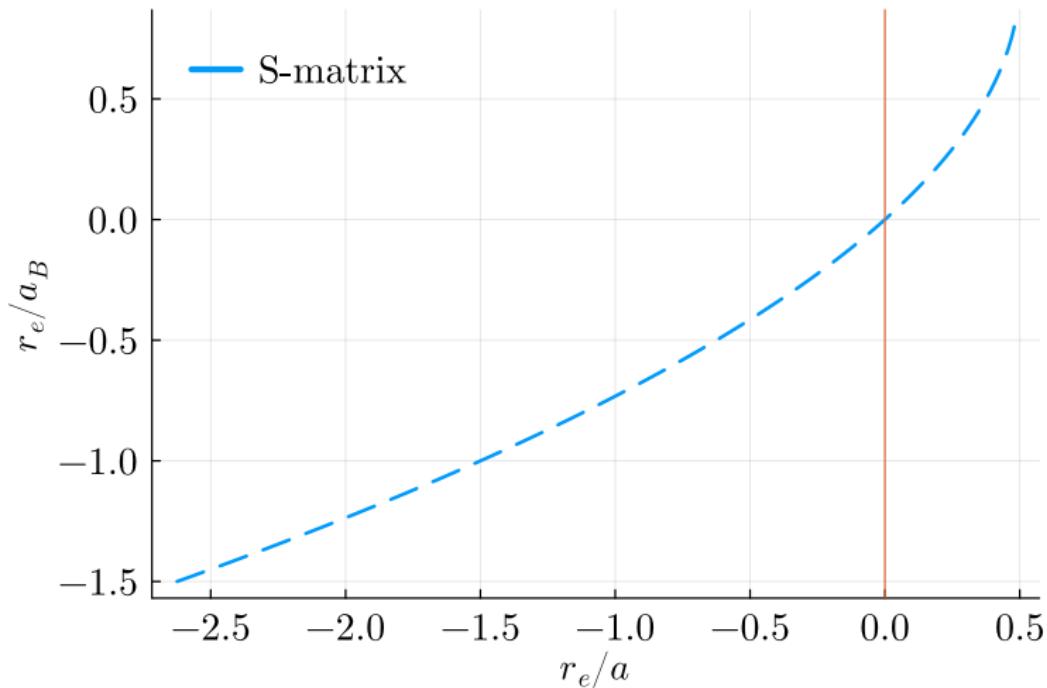
- ▶ Shape parameters vanish - $v_n = 0$
- ▶ r_B relates the scattering lenght and the physical pole

$$r_B = a - a_B$$

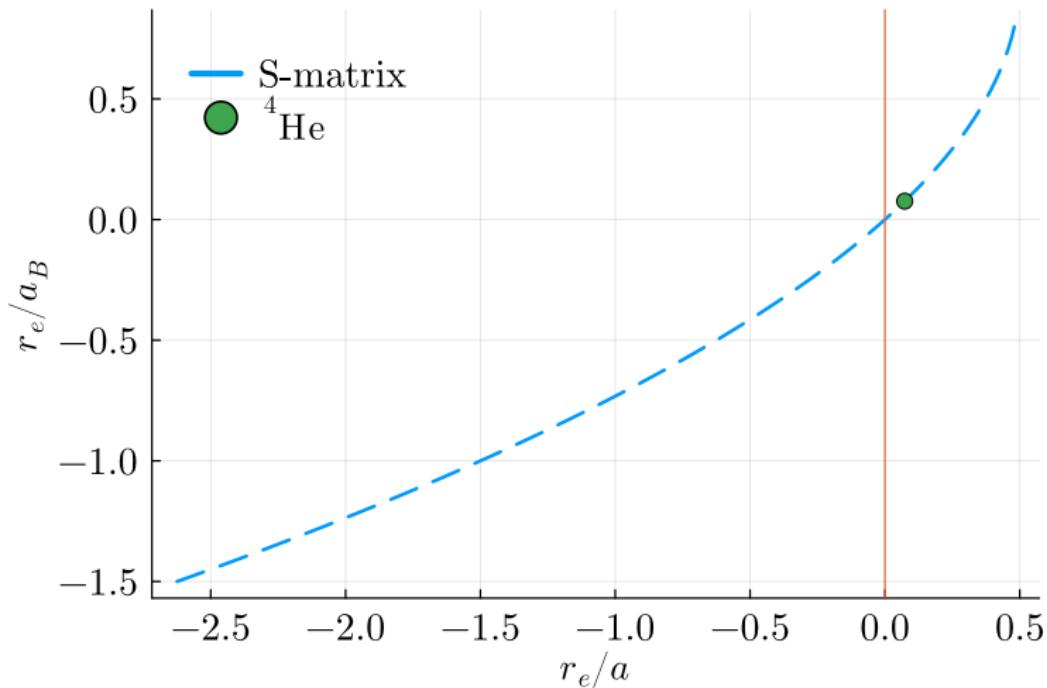
- ▶ Effective range relation

$$r_e a = 2r_B a_B$$

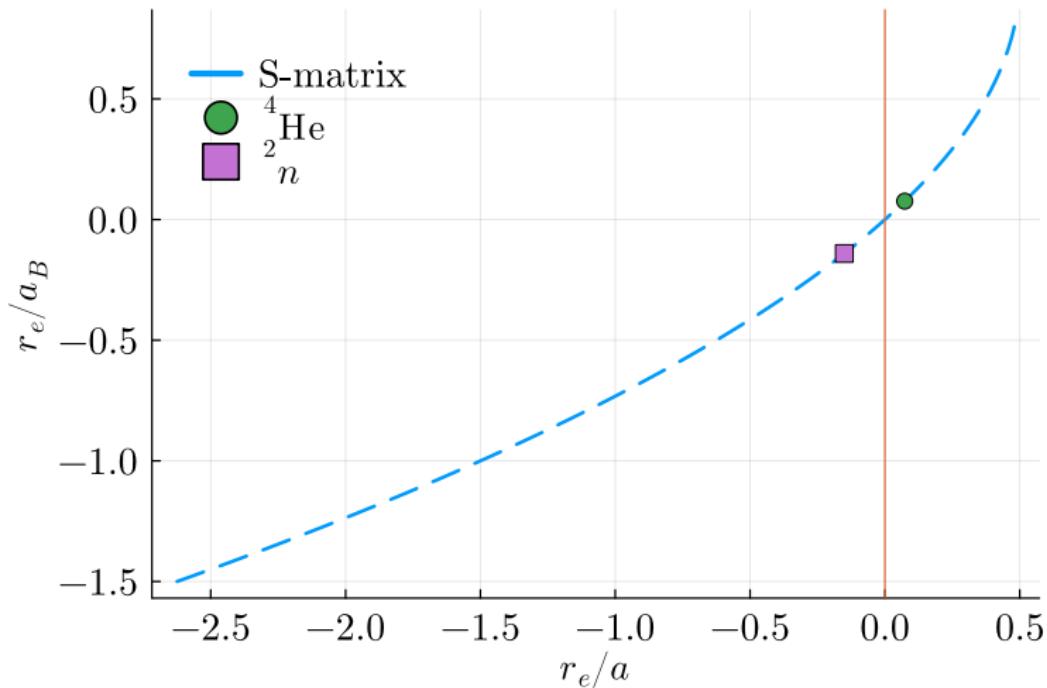
Zero-shape universality



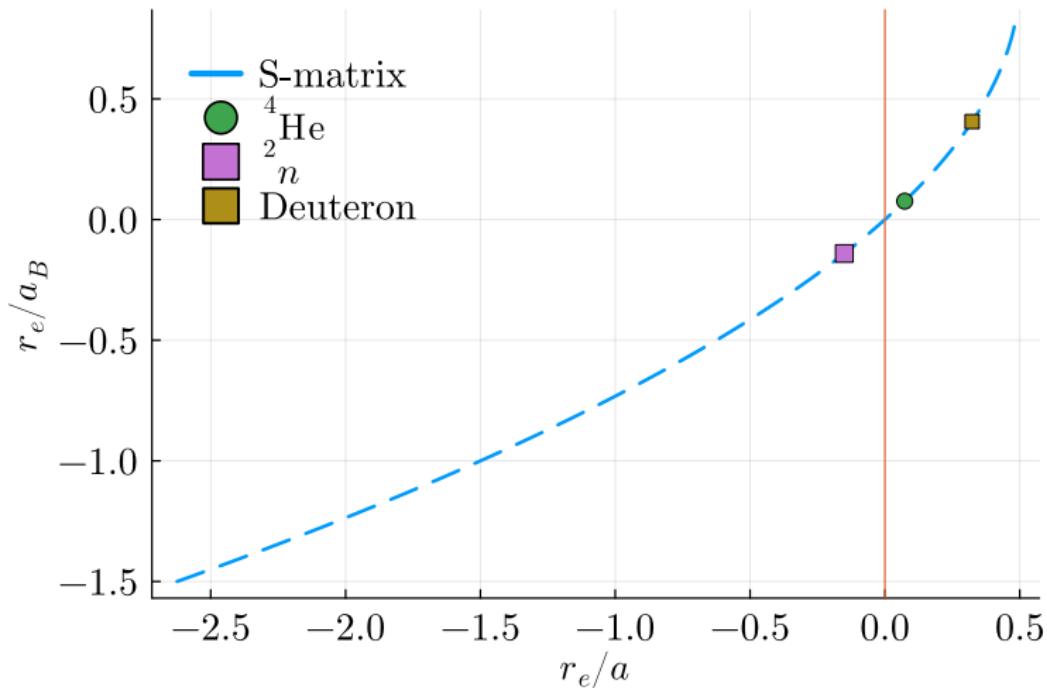
Zero-shape universality



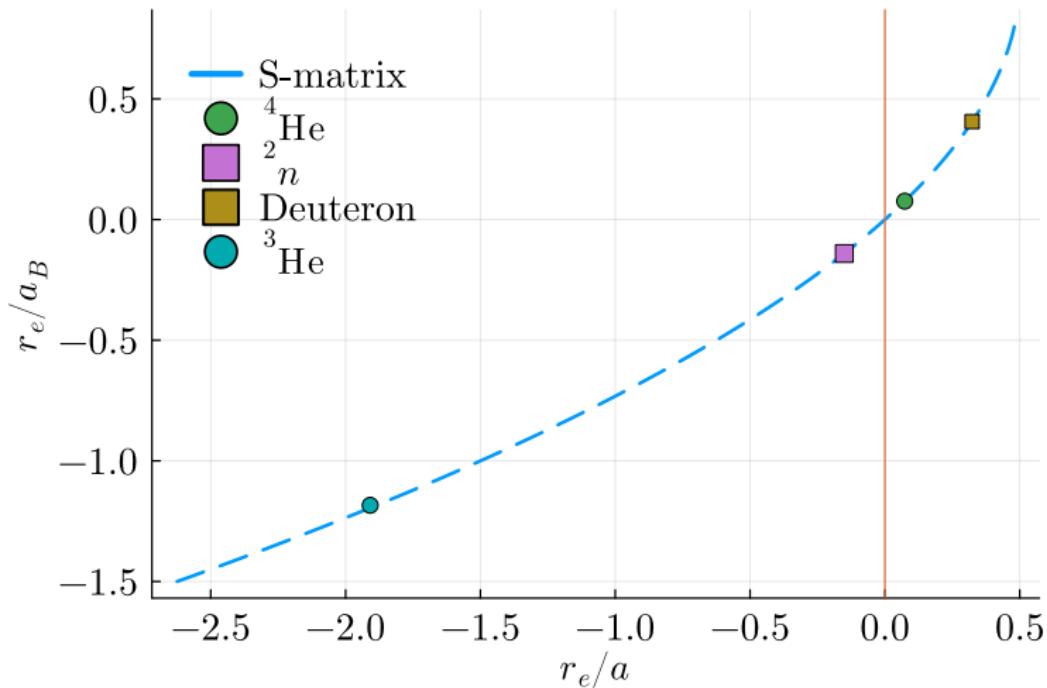
Zero-shape universality



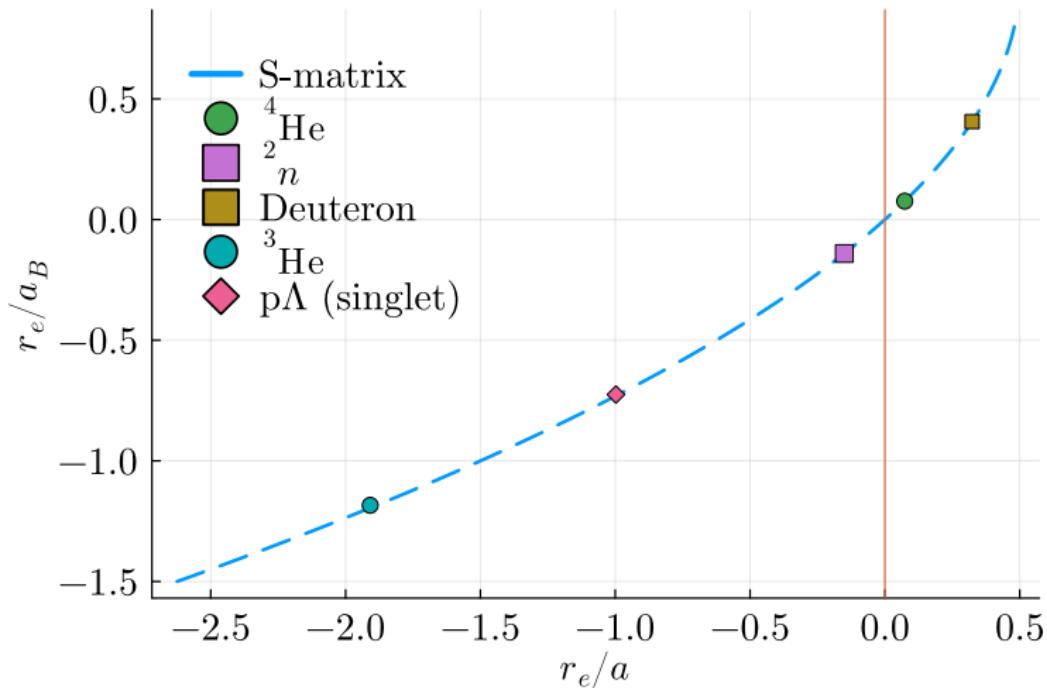
Zero-shape universality



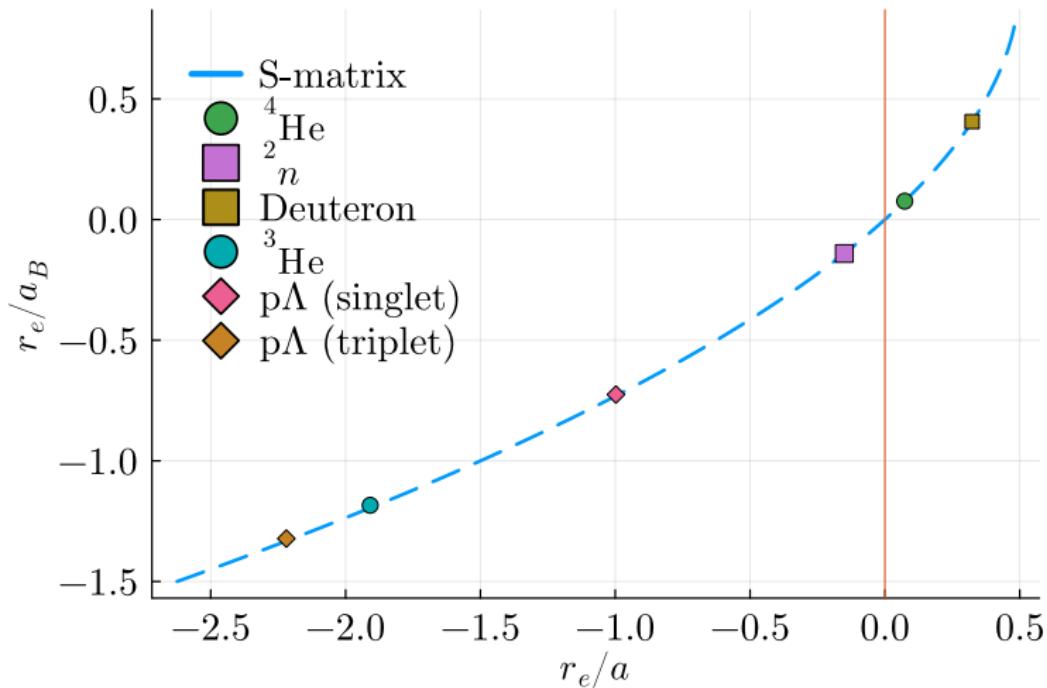
Zero-shape universality



Zero-shape universality



Zero-shape universality



Eckart Potential

$$V(r) = -2\beta\lambda^2 \frac{e^{-\lambda r}}{(1 + \beta e^{-\lambda r})^2}$$

Eckart Potential

$$V(r) = -2\beta\lambda^2 \frac{e^{-\lambda r}}{(1 + \beta e^{-\lambda r})^2}$$

- Binding length

$$a_B = \frac{1}{\lambda} \frac{2(\beta + 1)}{\beta - 1}$$

- Scattering length

$$a = \frac{4\beta}{\lambda(\beta - 1)}$$

- Effective range

$$r_e = \frac{2(\beta + 1)}{\lambda\beta}$$

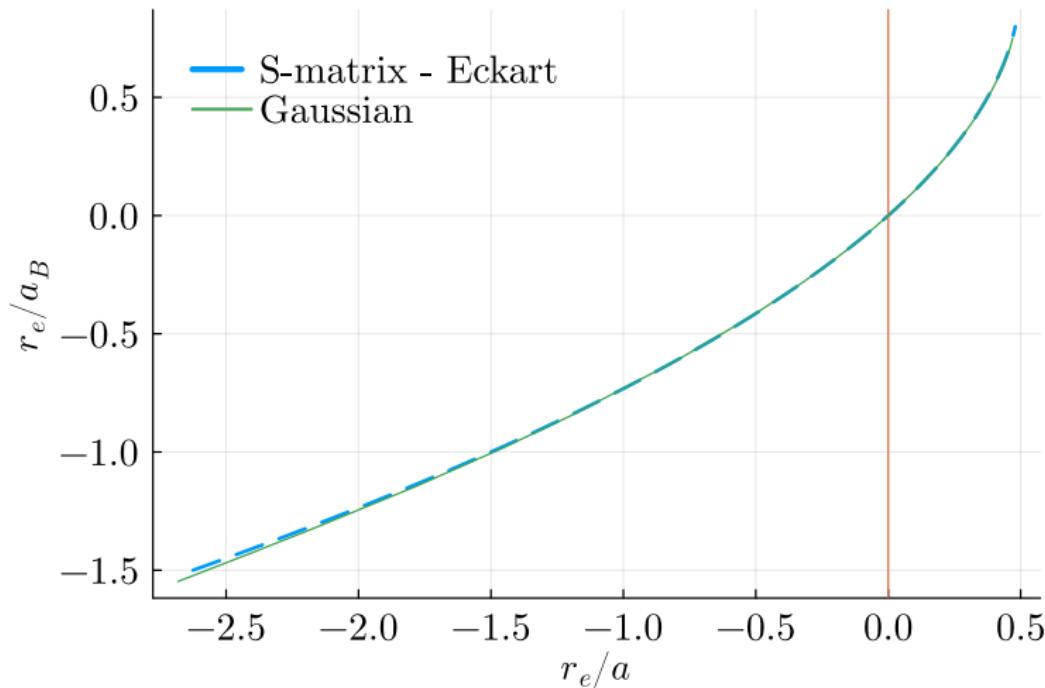
- “Interaction pole”

$$r_B = \frac{2}{\lambda}$$

Gaussian characterization

Effective Description using Gaussian Potential

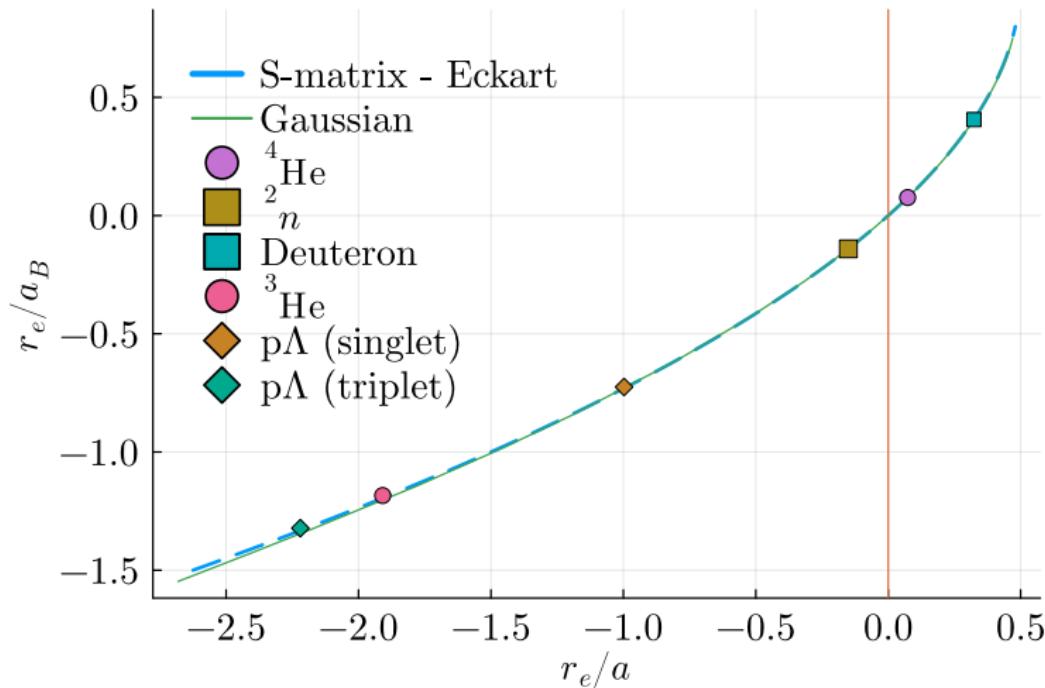
$$V(r) = V_0 e^{-(r/r_0)^2}$$



Gaussian characterization

Effective Description using Gaussian Potential

$$V(r) = V_0 e^{-(r/r_0)^2}$$



Outline

(Enlarged) Unitary Window

S-matrix

Effective Range Expansion

Zero-shape Universality

Gaussian (Eckart) characterization

Three-body sector

Efimov Effect

Level Functions - Gaussian Characterization

Moving along the universal curve

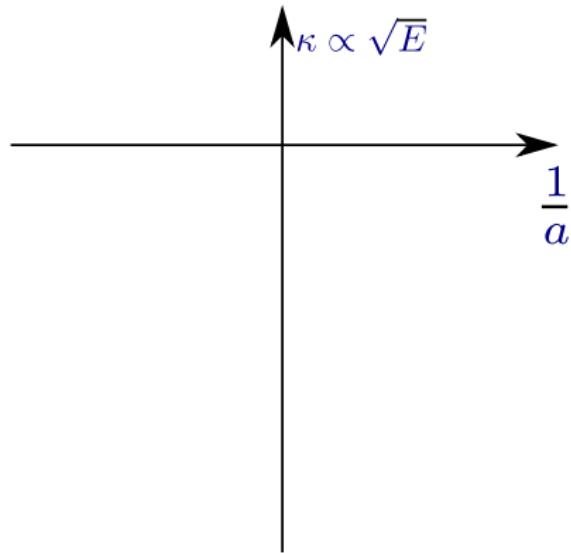
Note on DSI

More particles

LO Gaussian Potential - Two- and Three-Body Force

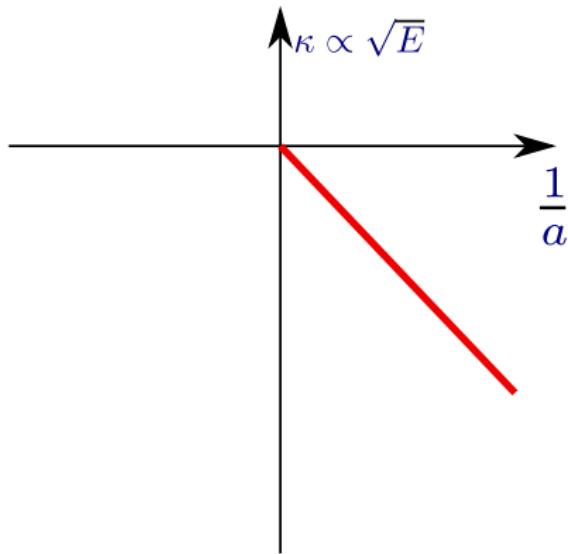
References

Efimov Effect



Efimov Effect

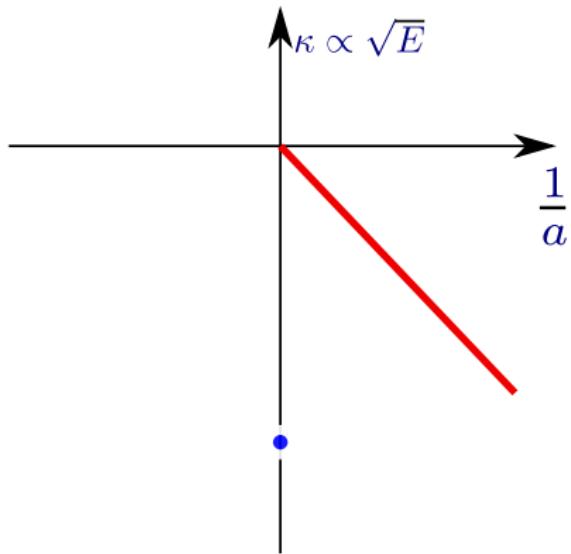
$$E_2 \propto \frac{1}{a^2}$$



Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$

$$E_3^0 \propto \frac{1}{\ell^2}$$

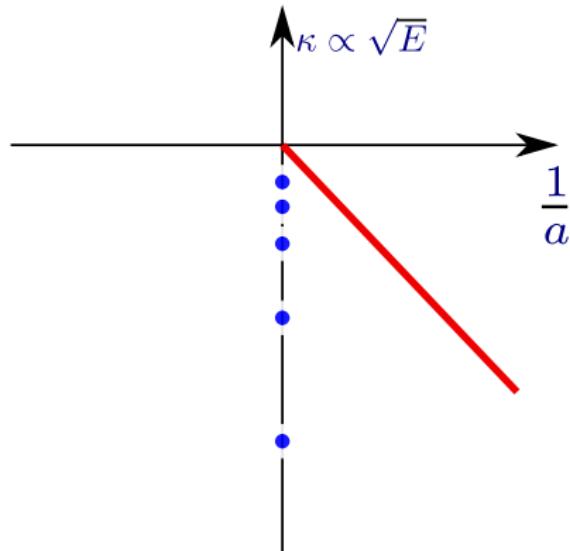


@ $1/a = 0$



Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$

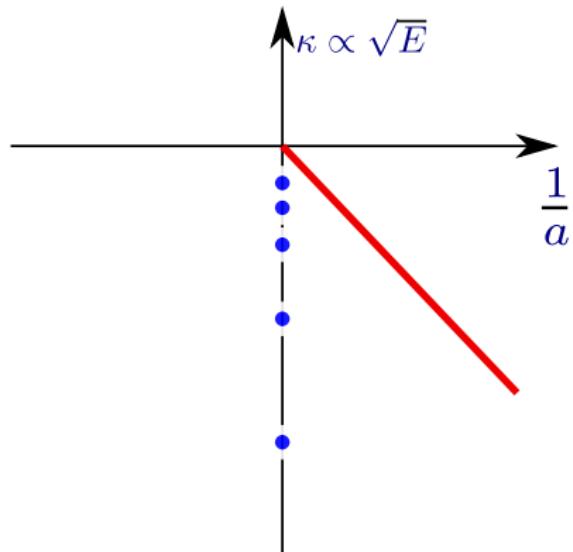


@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \end{array} \right.$$

Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$

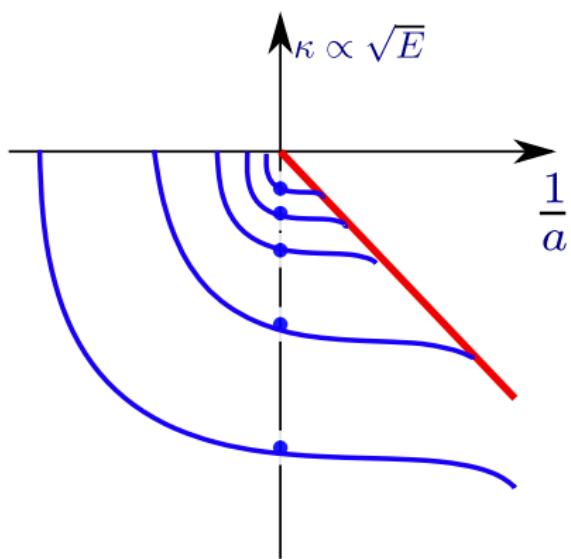


@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$

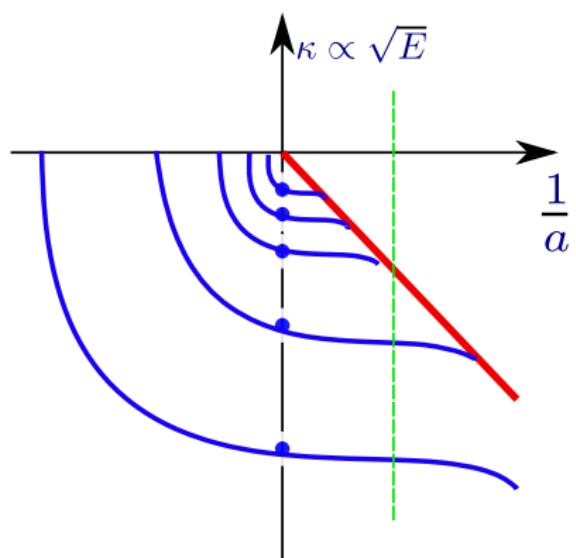


@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$



@ $1/a = 0$

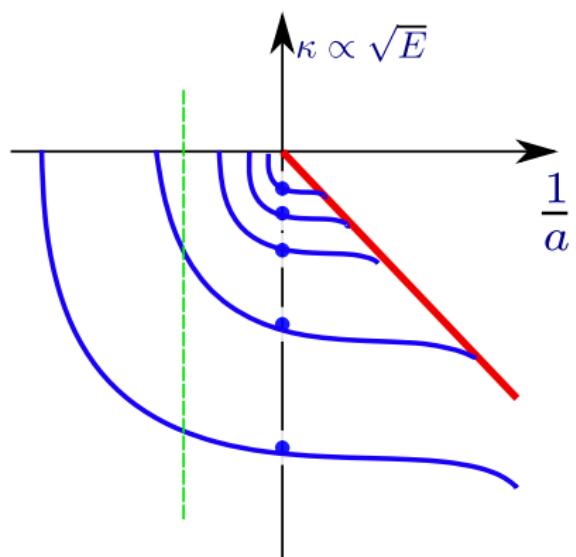
$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

@ $a > 0$

Finite # E_3 's

Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$



@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

@ $a > 0$

Finite # E_3 's

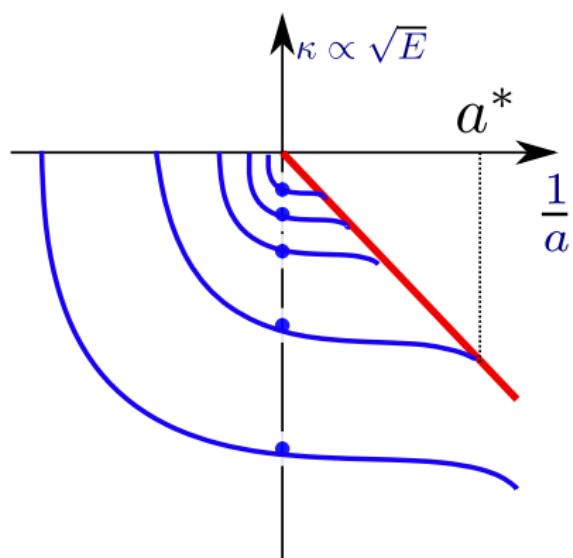
@ $a < 0$

Borromean states



Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$



@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

@ $a > 0$

Finite # E_3 's

@ $a < 0$

Borromean states

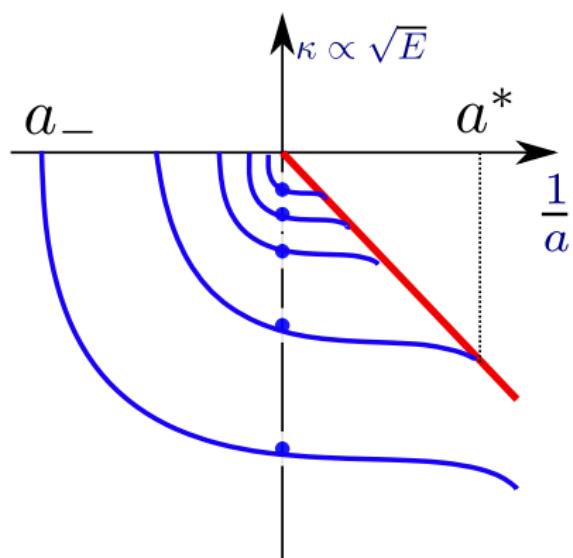


@ $a = a^*$

$$P_3 \rightarrow P_2 + P$$

Efimov Effect

$$E_2 \propto \frac{1}{a^2}$$



@ $1/a = 0$

$$\left\{ \begin{array}{l} E_3^0 \propto \frac{1}{\ell^2} \\ E_3^n \rightarrow 0 \quad n \rightarrow \infty \\ E_3^{n+1}/E_3^n \rightarrow 1/515 \\ E_3^n \sim (1/515)^n \kappa_*^2 \end{array} \right.$$

@ $a > 0$

Finite # E_3 's

@ $a < 0$



Borromean states

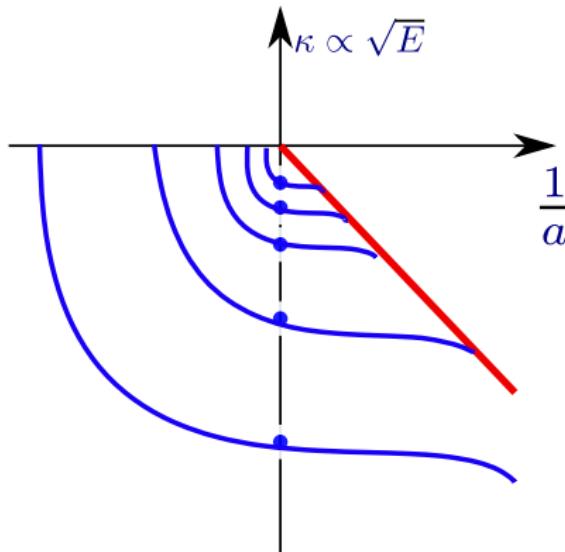
@ $a = a^*$

$$P_3 \rightarrow P_2 + P$$

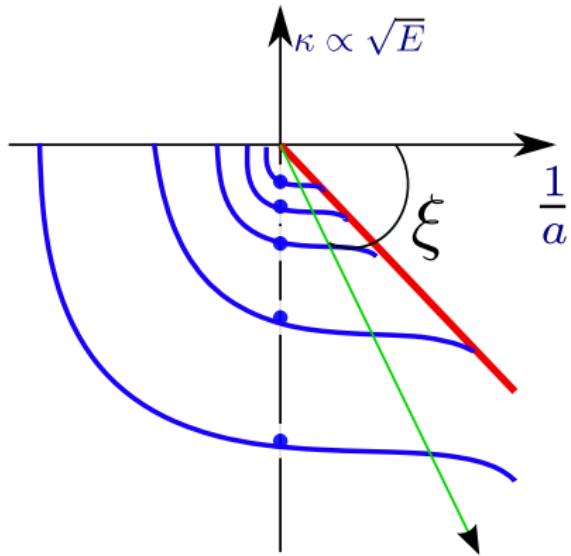
@ $a = a_-$

$$P_3 \rightarrow P + P + P$$

Efimov Levels - Discrete Scale Invariance



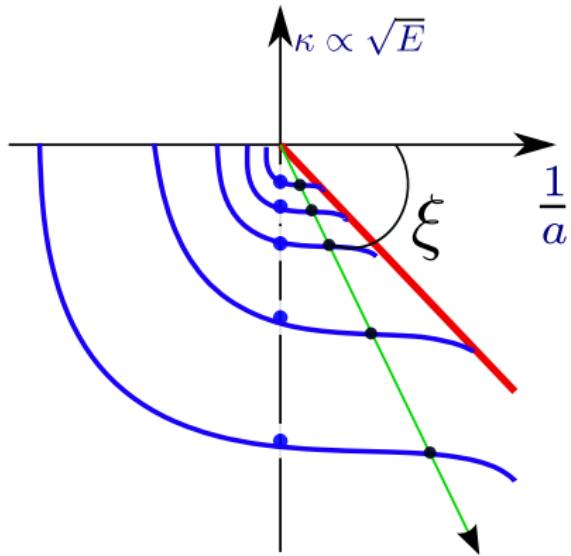
Efimov Levels - Discrete Scale Invariance



Polar coordinates

$$(H)^2 = (E_3 + E_2)/(\hbar^2/m)$$
$$\tan^2 \xi = E_3/E_2$$

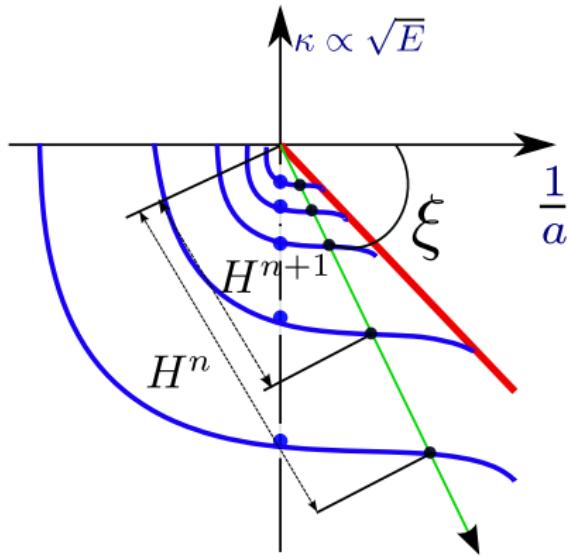
Efimov Levels - Discrete Scale Invariance



Polar coordinates

$$(H)^2 = (\textcolor{blue}{E}_3 + \textcolor{red}{E}_2)/(\hbar^2/m)$$
$$\tan^2 \xi = \textcolor{blue}{E}_3/\textcolor{red}{E}_2$$

Efimov Levels - Discrete Scale Invariance



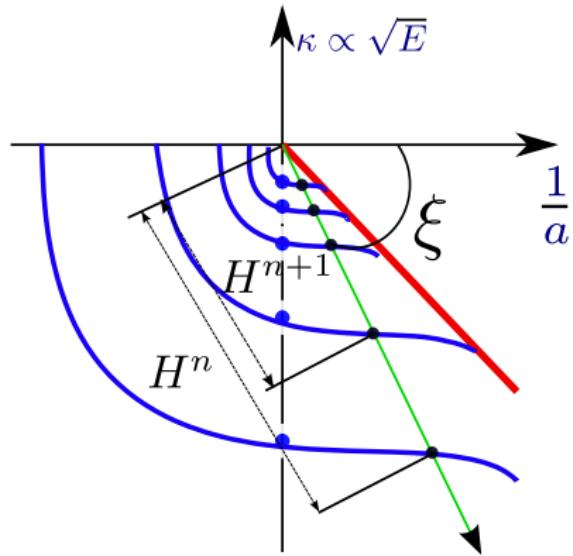
Polar coordinates

$$(H)^2 = (E_3 + E_2)/(\hbar^2/m)$$
$$\tan^2 \xi = E_3/E_2$$

For each ξ

$$H^{n+1}/H^n \rightarrow 1/22.7$$

Efimov Levels - Discrete Scale Invariance



Polar coordinates

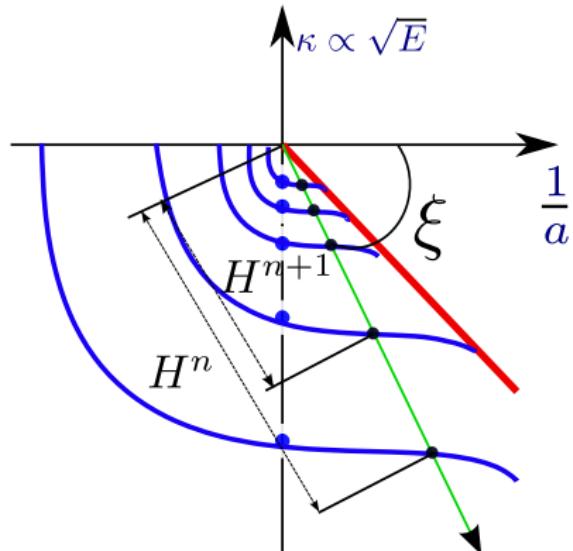
$$(H)^2 = (E_3 + E_2)/(\hbar^2/m)$$
$$\tan^2 \xi = E_3/E_2$$

For each ξ

$$H^{n+1}/H^n \rightarrow 1/22.7$$

$$E_3^n + \frac{\hbar^2}{ma^2} = \frac{\hbar^2 \kappa_*^2}{m} e^{-2(n-n^*)\pi/s_0} e^{\Delta(\xi)/s_0}$$

Efimov Levels - Discrete Scale Invariance



Polar coordinates

$$\frac{1}{a}$$

$$(H)^2 = (E_3 + E_2)/(\hbar^2/m)$$

$$\tan^2 \xi = E_3/E_2$$

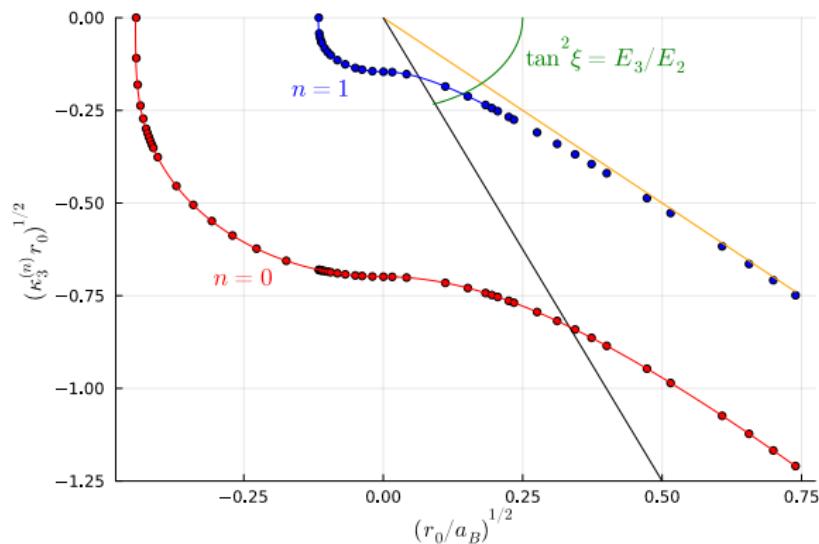
For each ξ

$$H^{n+1}/H^n \rightarrow 1/22.7$$

$$E_3^n + \frac{\hbar^2}{ma^2} = \frac{\hbar^2 \kappa_*^2}{m} e^{-2(n-n^*)\pi/s_0} e^{\Delta(\xi)/s_0}$$

$$\begin{cases} E_3^n / (\hbar^2 / ma^2) = \tan^2 \xi \\ \kappa_* a = e^{(n-n^*)\pi/s_0} \frac{e^{-\Delta(\xi)/2s_0}}{\cos \xi} \end{cases}$$

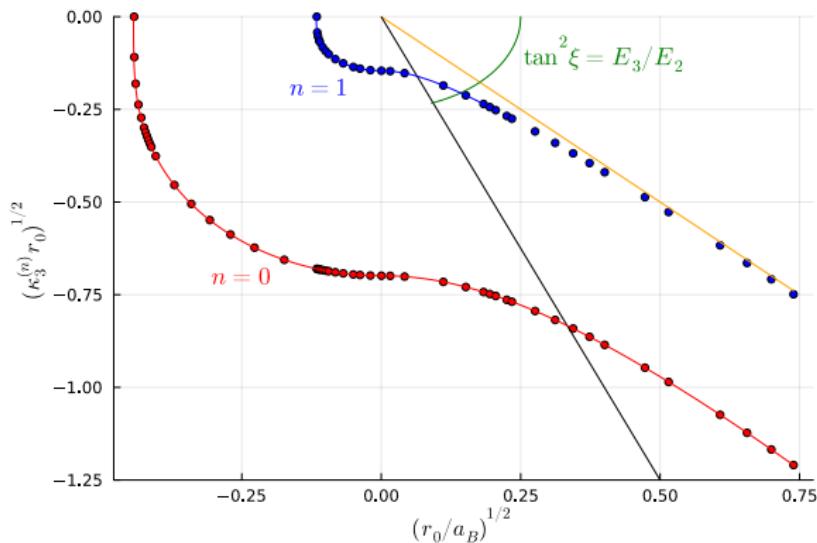
Gaussian Level Function - Universality



Gaussian Level Function - Universality

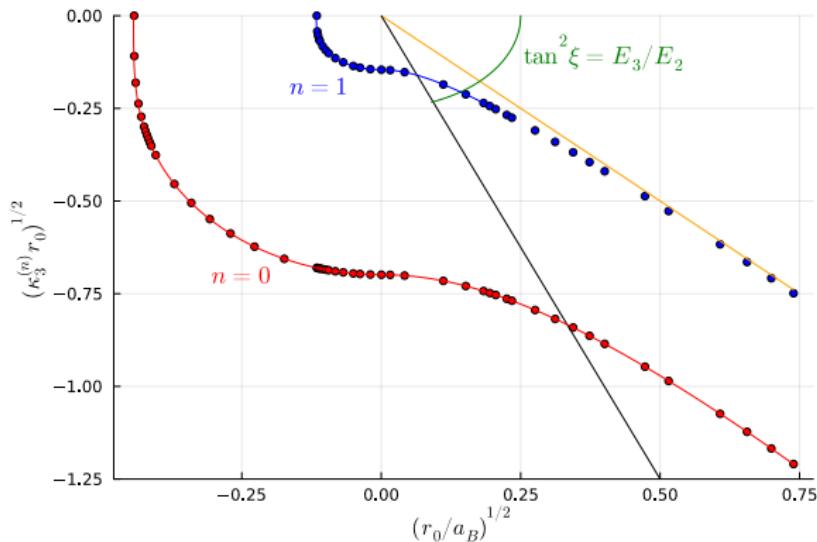
- Zero Range

$$\kappa_* a = e^{-\Delta(\xi)/2s_0} / \cos \xi$$



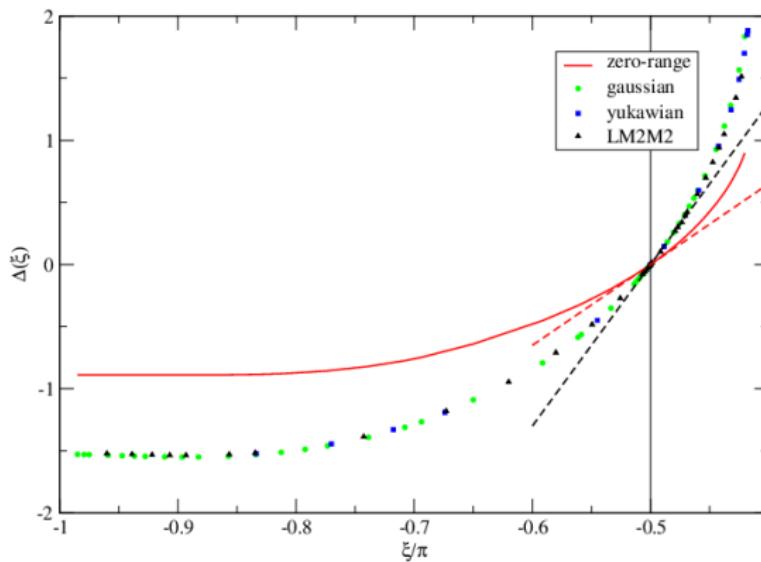
Gaussian Level Function - Universality

- Zero Range $\kappa_* a = e^{-\Delta(\xi)/2s_0} / \cos \xi$
- Potential $\kappa_* a_B = e^{-\tilde{\Delta}(\xi)/2s_0} / \cos \xi$



Gaussian Level Function - Universality

- Zero Range $\kappa_* a = e^{-\Delta(\xi)/2s_0} / \cos \xi$
- Potential $\kappa_* a_B = e^{-\tilde{\Delta}(\xi)/2s_0} / \cos \xi$



Moving along the curve

- Scale Invariance

$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi)$$

Moving along the curve

- Scale Invariance

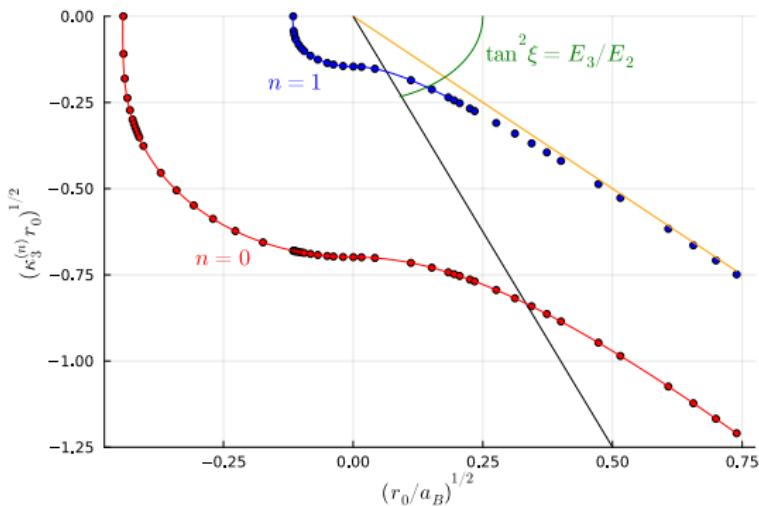
$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi) = \kappa_* a_B \Big|_{\text{Other finite range potentials}}$$

Moving along the curve

- Scale Invariance

$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi) = \kappa_* a_B \Big|_{\text{Other finite range potentials}}$$

- Unique r_0



Moving along the curve

- Scale Invariance

$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi) = \kappa_* a_B \Big|_{\text{Other finite range potentials}}$$

- Unique r_0

Potential	E_2 (mK)	E_3 (mK)	E_4 (mK)	$r_0^{(3)}(a_0)$	$r_0^{(4)}(a_0)$
HFD-HE2	0.8301	117.2	535.6	11.146	11.840
LM2M2	1.3094	126.5	559.2	11.150	11.853
HFD-B3-FCH	1.4475	129.0	566.1	11.148	11.853
CCSAPT	1.5643	131.0	571.7	11.149	11.851
PCKLJS	1.6154	131.8	573.9	11.148	11.852
HFD-B	1.6921	133.1	577.3	11.149	11.854
SAPT96	1.7443	134.0	580.0	11.147	11.850

Moving along the curve

- Scale Invariance

$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi) = \kappa_* a_B \Big|_{\text{Other finite range potentials}}$$

- Unique r_0

Potential	E_2 (mK)	E_3 (mK)	E_4 (mK)	$r_0^{(3)}(a_0)$	$r_0^{(4)}(a_0)$
HFD-HE2	0.8301	117.2	535.6	11.146	11.840
LM2M2	1.3094	126.5	559.2	11.150	11.853
HFD-B3-FCH	1.4475	129.0	566.1	11.148	11.853
CCSAPT	1.5643	131.0	571.7	11.149	11.851
PCKLJS	1.6154	131.8	573.9	11.148	11.852
HFD-B	1.6921	133.1	577.3	11.149	11.854
SAPT96	1.7443	134.0	580.0	11.147	11.850

- Energy at the unitary limit given by r_0

$$E_3^* \approx 83 \text{ mK}$$

$$E_4^* \approx 433 \text{ mK}$$

Moving along the curve

- Scale Invariance

$$\kappa_* a_B \Big|_{\text{Gaussian}} = \mathcal{F}(\xi) = \kappa_* a_B \Big|_{\text{Other finite range potentials}}$$

- Unique r_0

Potential	E_2 (mK)	E_3 (mK)	E_4 (mK)	$r_0^{(3)}(a_0)$	$r_0^{(4)}(a_0)$
HFD-HE2	0.8301	117.2	535.6	11.146	11.840
LM2M2	1.3094	126.5	559.2	11.150	11.853
HFD-B3-FCH	1.4475	129.0	566.1	11.148	11.853
CCSAPT	1.5643	131.0	571.7	11.149	11.851
PCKLJS	1.6154	131.8	573.9	11.148	11.852
HFD-B	1.6921	133.1	577.3	11.149	11.854
SAPT96	1.7443	134.0	580.0	11.147	11.850

- Energy at the unitary limit given by r_0

$$E_3^* \approx 83 \text{ mK}$$

$$E_4^* \approx 433 \text{ mK}$$

- Universal numbers

$$\kappa_3^* a_-^3 = -2.13 \quad \kappa_4^* a_-^4 = -2.32$$

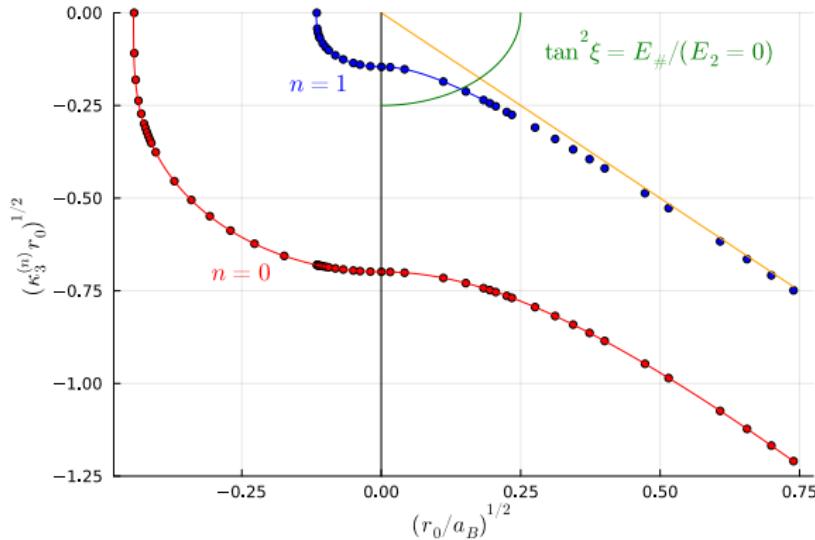
Note on DSI

- DSI \Rightarrow Log-periodic functions

Note on DSI

• DSI \Rightarrow

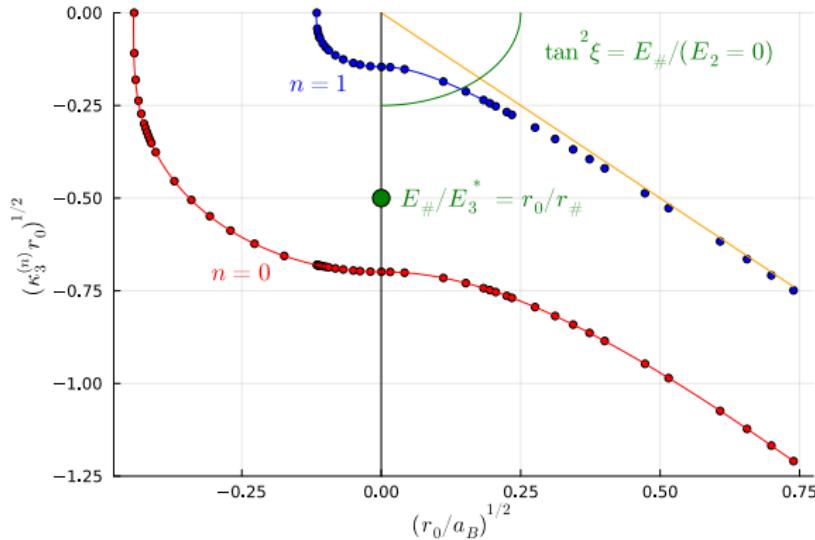
Log-periodic functions



Note on DSI

• DSI \Rightarrow

Log-periodic functions



Note on DSI

- DSI \Rightarrow Log-periodic functions

- Three-body force

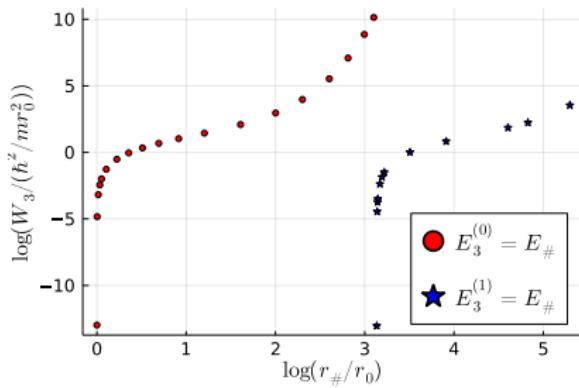
$$W = W_3 e^{-(r_{12}^2 + r_{13}^2)/r_0^2}$$

Note on DSI

- DSI \Rightarrow Log-periodic functions

- Three-body force

$$W = W_3 e^{-(r_{12}^2 + r_{13}^2)/r_0^2}$$

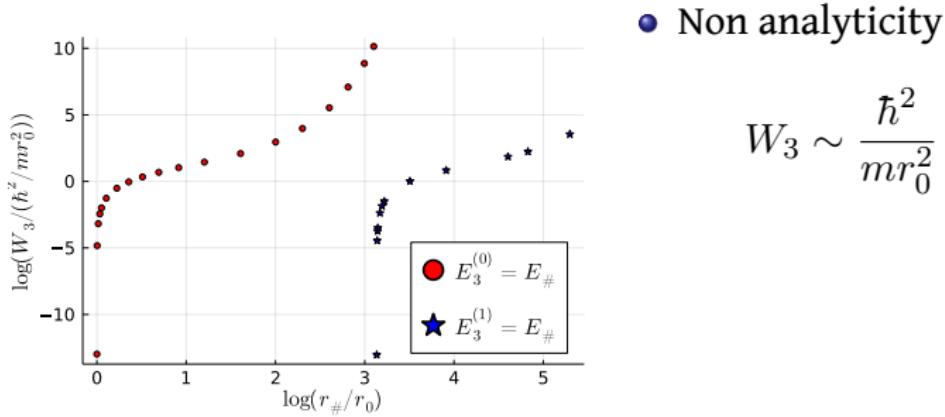


Note on DSI

- DSI \Rightarrow Log-periodic functions

- Three-body force

$$W = W_3 e^{-(r_{12}^2 + r_{13}^2)/r_0^2}$$



$$W_3 \sim \frac{\hbar^2}{mr_0^2} e^{(r_#/r_0)^{1.13}}$$

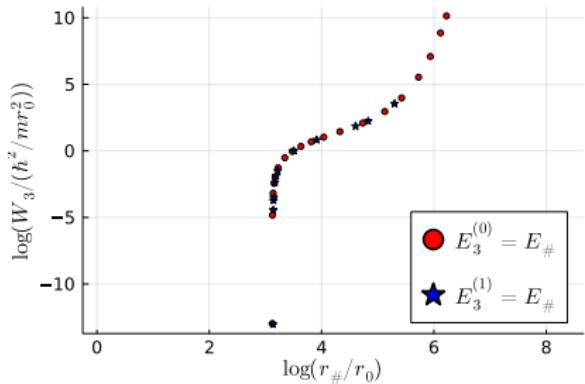
Note on DSI

- DSI \Rightarrow Log-periodic functions

- Three-body force

$$W = W_3 e^{-(r_{12}^2 + r_{13}^2)/r_0^2}$$

- Non analyticity



$$W_3 \sim \frac{\hbar^2}{mr_0^2} e^{(r_#/r_0)^{1.13}}$$

- Log-periodicity

$$\log(r_#/r_0) \rightarrow \log(r_#/r_0) - \pi/s_0$$

Outline

(Enlarged) Unitary Window

S-matrix

Effective Range Expansion

Zero-shape Universality

Gaussian (Eckart) characterization

Three-body sector

Efimov Effect

Level Functions - Gaussian Characterization

Moving along the universal curve

Note on DSI

More particles

LO Gaussian Potential - Two- and Three-Body Force

References

Effective Gaussian Description of ^4He

- “Reference” ^4He given by LM2M2 potential

$$\bar{a} = 189.415 \text{ } a_0, \bar{r}_e = 13.845 \text{ } a_0, \text{ and } r_B = 7.194 \text{ } a_0$$

N	$\bar{E}_N(\text{mK})$	$\bar{E}_N^*(\text{mK})$
2	-1.30348	
3	-126.40	-2.2706
4	-558.98 [Hiyama 2012]	-127.33 [Hiyama 2012]
5	-1300 [Bazak 2020]	
6	-2315 [Bazak 2020]	
7	-3571 [Bazak 2020]	

Effective Gaussian Description of ^4He

- “Reference” ^4He given by LM2M2 potential

$$\bar{a} = 189.415 \text{ } a_0, \bar{r}_e = 13.845 \text{ } a_0, \text{ and } r_B = 7.194 \text{ } a_0$$

N	$\bar{E}_N(\text{mK})$	$\bar{E}_N^*(\text{mK})$
2	-1.30348	
3	-126.40	-2.2706
4	-558.98 [Hiyama 2012]	-127.33 [Hiyama 2012]
5	-1300 [Bazak 2020]	
6	-2315 [Bazak 2020]	
7	-3571 [Bazak 2020]	

- Effective Gaussian Potential

$$V_{\text{LO}}(r) = V_0 e^{-(r/r_0)^2}$$

Effective Gaussian Description of ^4He

- “Reference” ^4He given by LM2M2 potential

$$\bar{a} = 189.415 \text{ } a_0, \bar{r}_e = 13.845 \text{ } a_0, \text{ and } r_B = 7.194 \text{ } a_0$$

N	$\bar{E}_N(\text{mK})$	$\bar{E}_N^*(\text{mK})$
2	-1.30348	
3	-126.40	-2.2706
4	-558.98 [Hiyama 2012]	-127.33 [Hiyama 2012]
5	-1300 [Bazak 2020]	
6	-2315 [Bazak 2020]	
7	-3571 [Bazak 2020]	

- Effective Gaussian Potential

$$V_{\text{LO}}(r) = V_0 e^{-(r/r_0)^2}$$

- Small parameter

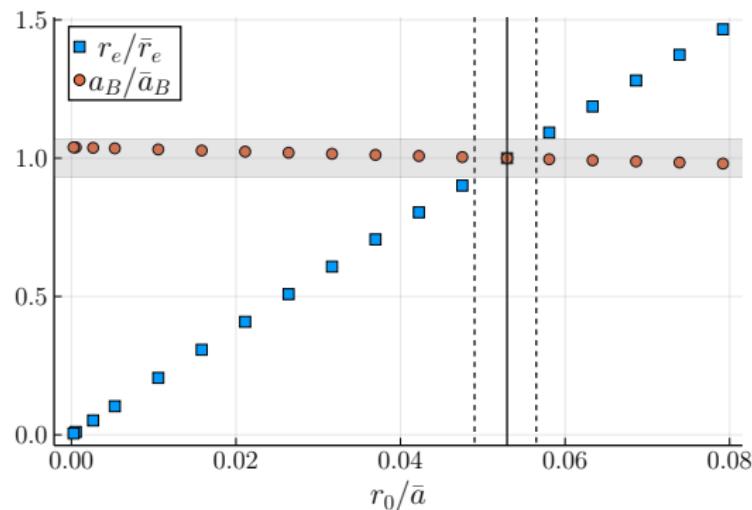
$$\varepsilon = \bar{r}_e/\bar{a} \approx 7\%$$

Two Body

- Effective Gaussian Potential

$$V_{\text{LO}}(r) = V_0 e^{-(r/r_0)^2}$$

- Fix only \bar{a}

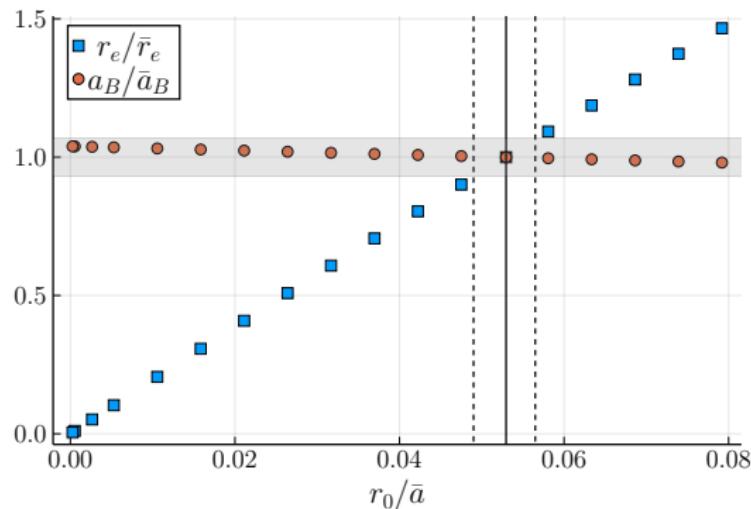


Two Body

- Effective Gaussian Potential

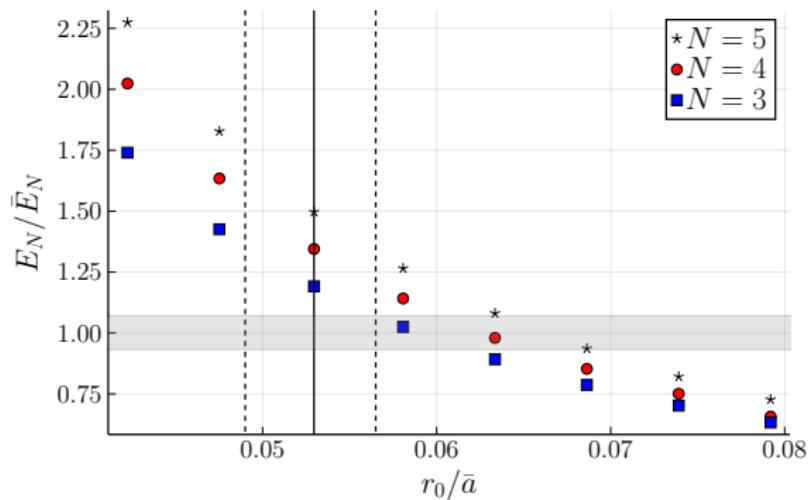
$$V_{\text{LO}}(r) = V_0 e^{-(r/r_0)^2}$$

- Fix only \bar{a}

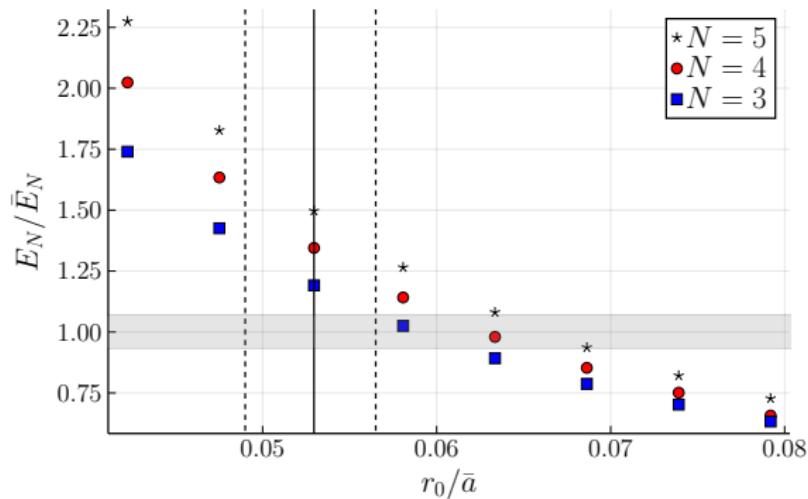


- Look for $\varepsilon = \bar{r}_e/\bar{a} \approx 7\%$ description also for r_e

Few Body

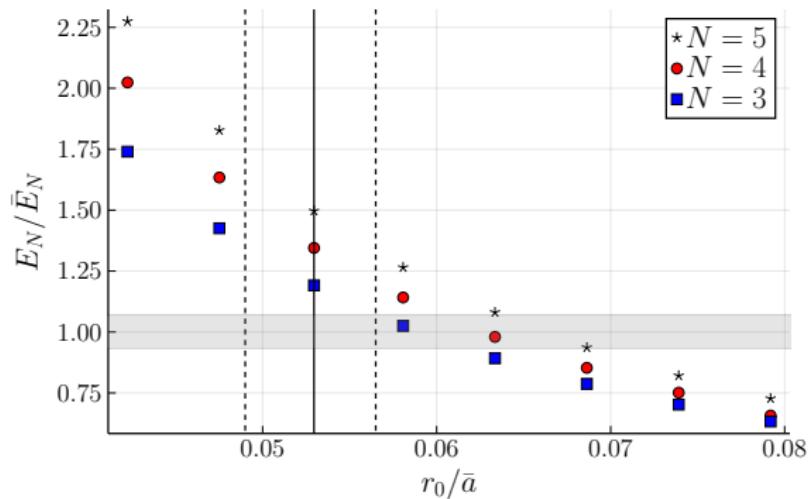


Few Body



- Not inside the $\varepsilon = 7\%$ band

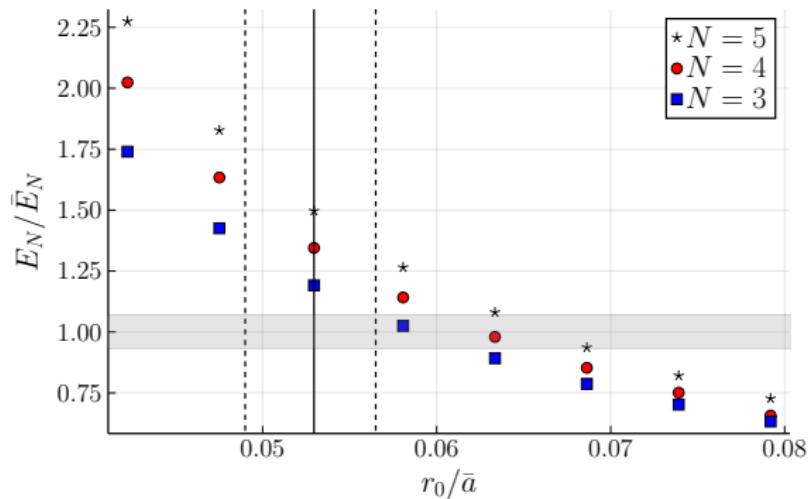
Few Body



- Not inside the $\varepsilon = 7\%$ band
- Collapse as $N \rightarrow \infty$

$$\frac{E_N}{N} = \frac{V_0}{2} N$$

Few Body



- Not inside the $\varepsilon = 7\%$ band
- Collapse as $N \rightarrow \infty$

$$\frac{E_N}{N} = \frac{V_0}{2} N$$

- Need for a three-body force

Few Body

- Three-body force

$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$

Few Body

- Three-body force

$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$

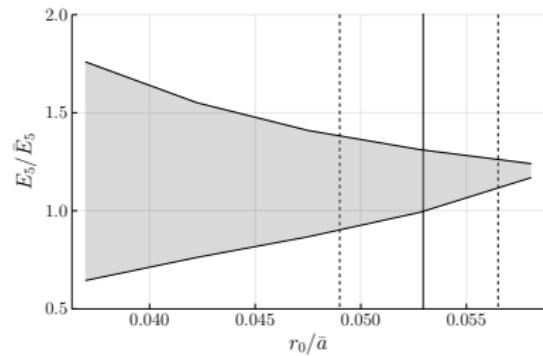
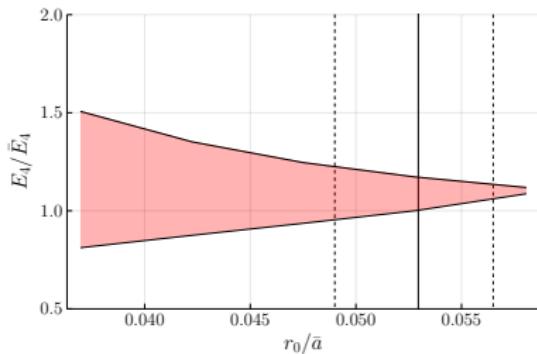
- A family of values (W_0, ρ_0) which fix \bar{E}_3

Few Body

- Three-body force

$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$

- A family of values (W_0, ρ_0) which fix \bar{E}_3
- Variation in \bar{E}_N

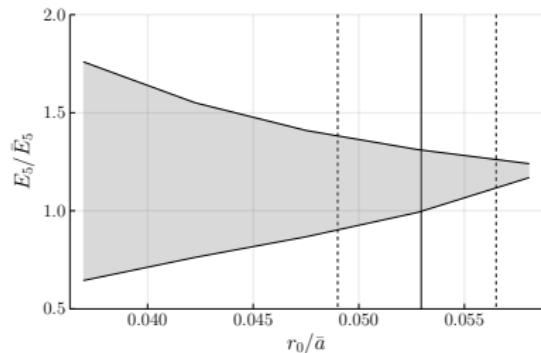
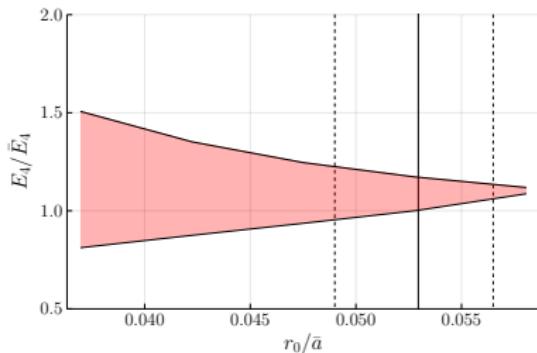


Few Body

- Three-body force

$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$

- A family of values (W_0, ρ_0) which fix \bar{E}_3
- Variation in \bar{E}_N

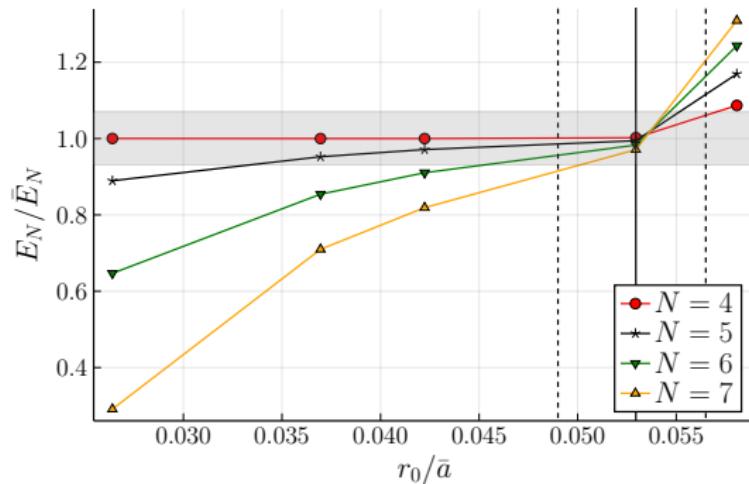


- We can use (W_0, ρ_0) to best fix \bar{E}_4

LO Gaussian Description

- LO Potential

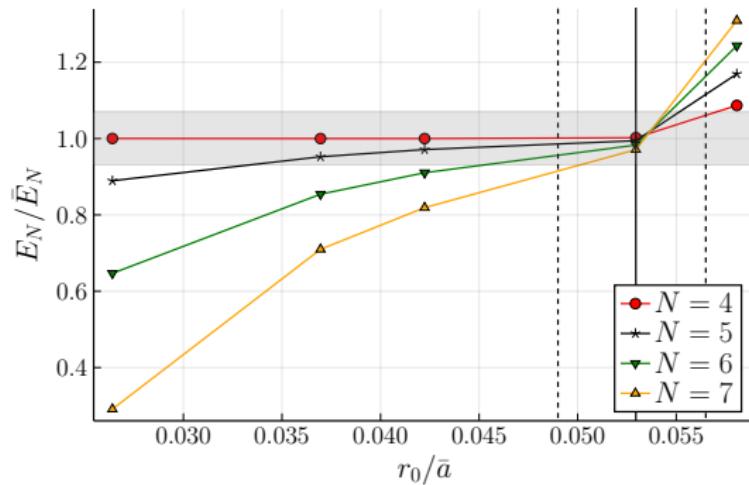
$$V_0 e^{-(r/r_0)^2} + W_0 e^{-(r_{12}^2+r_{13}^2+r_{23}^2)/\rho_0^2}$$



LO Gaussian Description

- LO Potential

$$V_0 e^{-(r/r_0)^2} + W_0 e^{-(r_{12}^2+r_{13}^2+r_{23}^2)/\rho_0^2}$$



- Best point is where we reproduce \bar{a} , \bar{a}_B , and r_e !!

LO Gaussian Description

- Description within the ε -LO band up to liquid

	Physical point	
	SGP	HFD-HE2
$r_0[a_0]$	10.0485	
$V_0[\text{K}]$	1.208018	
$\rho_0[a_0]$	8.4853	
$W_0[\text{K}]$	3.011702	
$E_4[\text{K}]$	0.536	0.536
$E_5[\text{K}]$	1.251	1.266
$E_6[\text{K}]$	2.216	2.232
$E_{10}/10[\text{K}]$	0.792(2)	0.831(2)
$E_{20}/20[\text{K}]$	1.525(2)	1.627(2)
$E_{40}/40[\text{K}]$	2.374(2)	2.482(2)
$E_{70}/70[\text{K}]$	3.07(1)	3.14(1)
$E_{112}/112[\text{K}]$	3.58(2)	3.63(2)
$E_N/N(\infty)[\text{K}]$	7.2(3)*	7.14(2)
HFD-B [K]		7.33(2)

Outline

(Enlarged) Unitary Window

S-matrix

Effective Range Expansion

Zero-shape Universality

Gaussian (Eckart) characterization

Three-body sector

Efimov Effect

Level Functions - Gaussian Characterization

Moving along the universal curve

Note on DSI

More particles

LO Gaussian Potential - Two- and Three-Body Force

References

References

- For the finite-range universality

- [1] *Efimov Physics and Connections to Nuclear Physics*
A. Kievsky, M. Gattobigio, L. Girlanda, M. Viviani
A. Kievsky, M. Gattobigio, L. Girlanda, M. Viviani
Annual Review of Nuclear and Particle Science **7**, 465-490 (2021) [▶ Link](#)
- [2] *Gaussian characterization of the unitary window for $N = 3$: Bound, scattering, and virtual states*
A. Deltuva, M. Gattobigio, A. Kievsky, and M. Viviani
Phys. Rev. C **102**, 064001 (2020) [▶ Link](#)

- For the universal Level functions

- [3] *Universality and scaling in the N -body sector of Efimov physics*
M. Gattobigio and A. Kievsky
Phys. Rev. A **90**, 012502 (2014) [▶ Link](#)
- [4] *Matching universal behavior with potential models*
R. Álvarez-Rodríguez, A. Deltuva, M. Gattobigio, and A. Kievsky
Phys. Rev. A **93**, 062701 (2016) [▶ Link](#)
- [5] *Gaussian Parametrization of Efimov Levels: Remnants of Discrete Scale Invariance*
Recchia, P., Kievsky, A., Girlanda, and M. Gattobigio
Few-Body Syst **63**, 8 (2022) [▶ Link](#)

- For the LO- and NLO-Gaussian description of ${}^4\text{He}$

- [6] *Few bosons to many bosons inside the unitary window: A transition between universal and nonuniversal behavior*
A. Kievsky, A. Polls, B. Juliá-Díaz, N. K. Timofeyuk, and M. Gattobigio
Phys. Rev. A **102**, 063320 (2020) [▶ Link](#)
- [7] *Subleading contributions to N -boson systems inside the universal window*
P. Recchia, A. Kievsky, L. Girlanda, and M. Gattobigio
Phys. Rev. A **106**, 022812 (2022) [▶ Link](#)

Collaborators



Paolo Recchia



Natalia Timofeyuk



Alejandro Kievsky



Artur Polls



Michele Viviani



Bruno Julia Diaz



Luca Girlanda

Collaborators



Paolo Recchia



Natalia Timofeyuk



Alejandro Kievsky



Artur Polls



Michele Viviani



Bruno Julia Diaz



Luca Girlanda

Thanks!

NLO Gaussian Description - Two body

- NLO two-body force

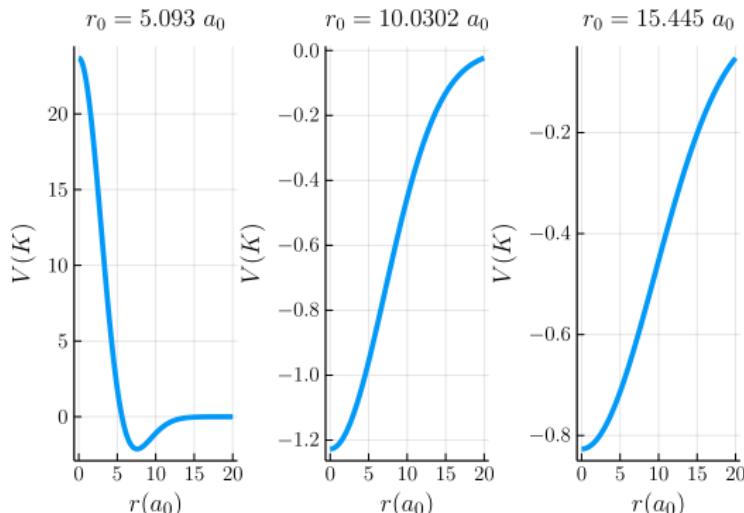
$$V_{\text{NLO}}(r) = V_0 e^{-(r/r_0)^2} + V_1 \frac{r^2}{r_0^2} e^{-(r/r_0)^2}$$

NLO Gaussian Description - Two body

- NLO two-body force

$$V_{\text{NLO}}(r) = V_0 e^{-(r/r_0)^2} + V_1 \frac{r^2}{r_0^2} e^{-(r/r_0)^2}$$

- We fix both \bar{a} and \bar{r}_e

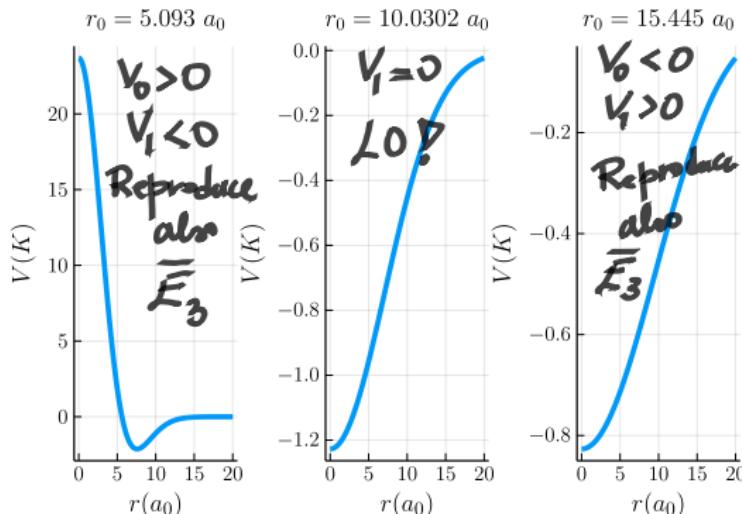


NLO Gaussian Description - Two body

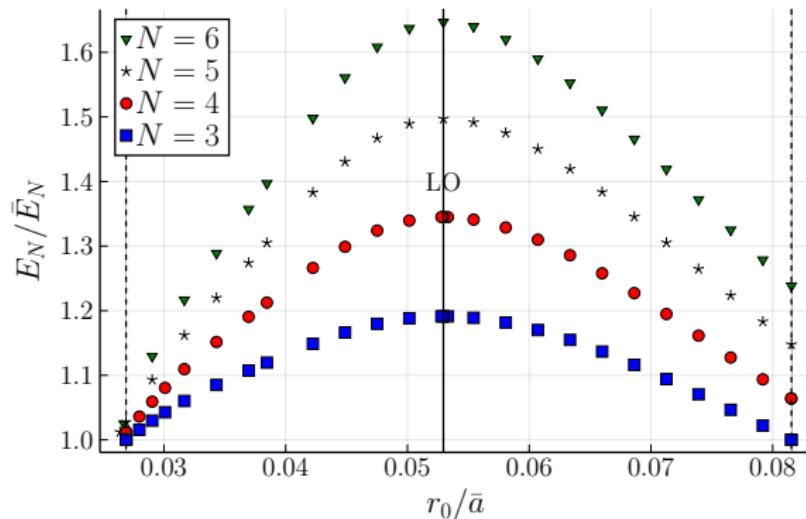
- NLO two-body force

$$V_{\text{NLO}}(r) = V_0 e^{-(r/r_0)^2} + V_1 \frac{r^2}{r_0^2} e^{-(r/r_0)^2}$$

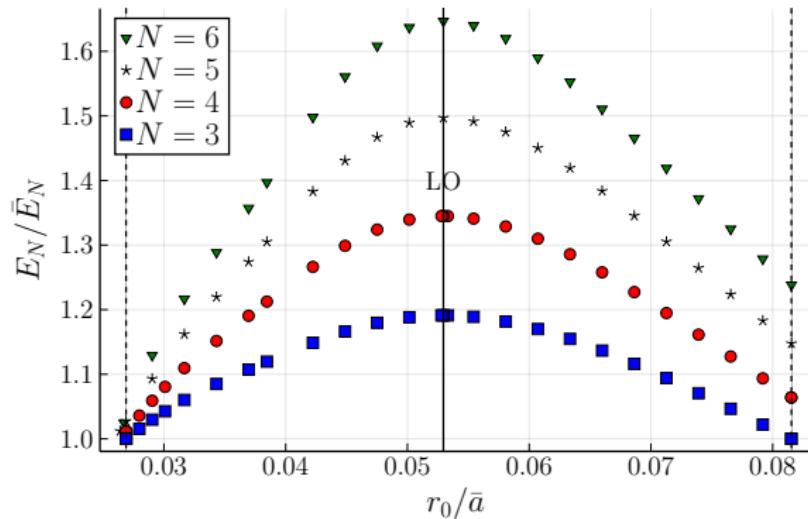
- We fix both \bar{a} and \bar{r}_e



NLO Two body - Few-body energies



NLO Two body - Few-body energies



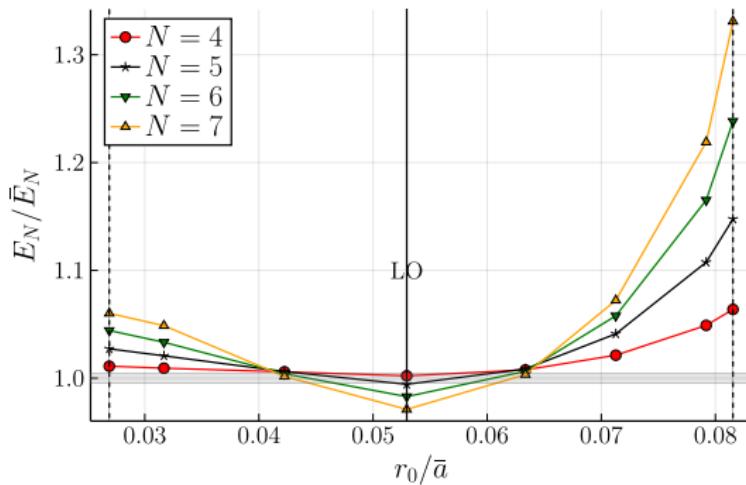
- Without 3-body force the system is unstable

$$\frac{E_N}{N} \propto N$$

NLO Two body + LO Three body

- With the LO 3-body force

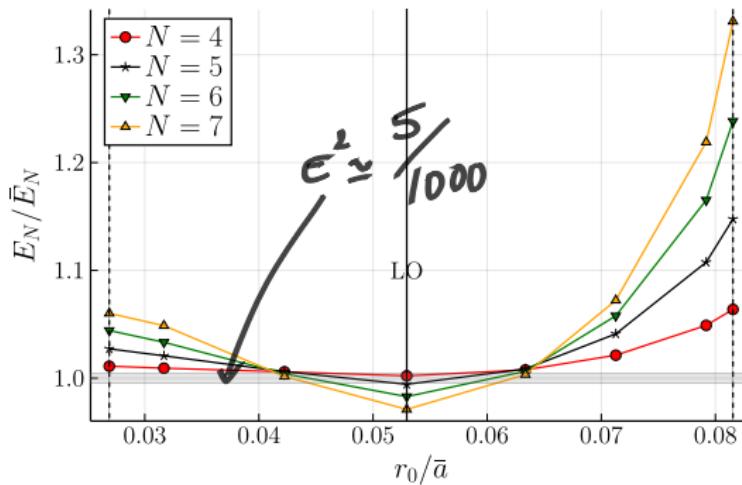
$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$



NLO Two body + LO Three body

- With the LO 3-body force

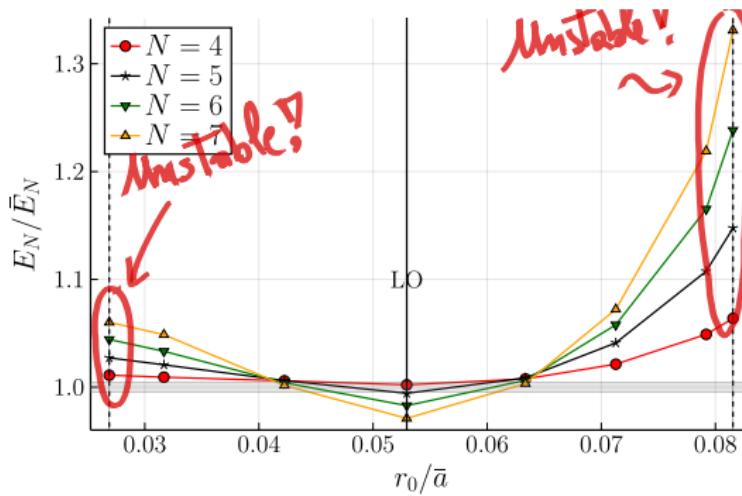
$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$



NLO Two body + LO Three body

- With the LO 3-body force

$$W_{\text{LO}} = W_0 e^{-(r_{12}^2 + r_{13}^2 + r_{23}^2)/\rho_0^2}$$

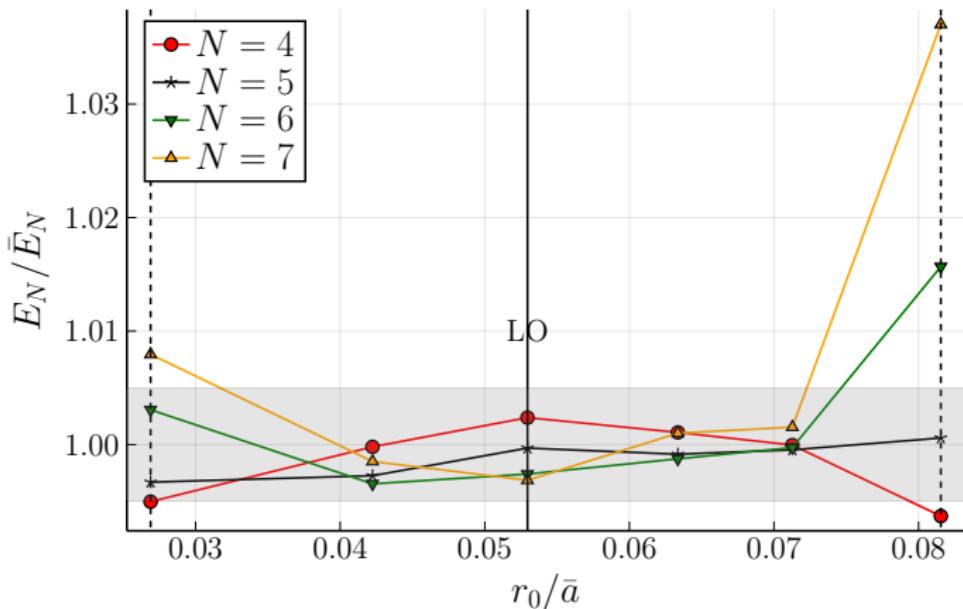


- Need another force at NLO!!!

Analysis with NLO 3-Body

- NLO Three-body force

$$W_{\text{NLO}} = W_0 e^{-r_{123}^2/\rho_0^2} + W_1 \left(\frac{r_{123}}{\rho_0} \right)^2 e^{-r_{123}^2/\rho_0^2}$$



Analysis with NLO 3-Body

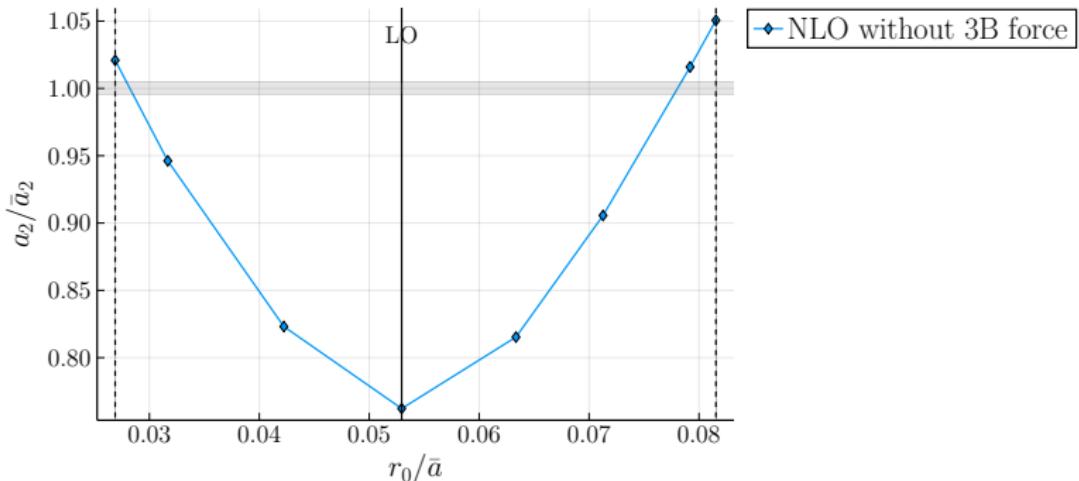
- What happens to different three-body observables?

Analysis with NLO 3-Body

- What happens to different three-body observables?
- Atom Dimer scattering length $\bar{a}_2 = 218 a_0$

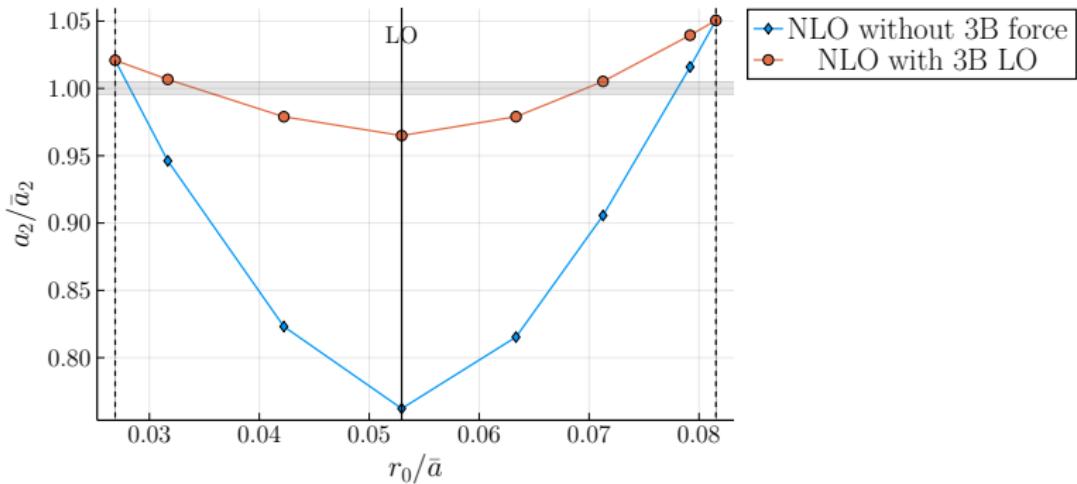
Analysis with NLO 3-Body

- What happens to different three-body observables?
- Atom Dimer scattering length $\bar{a}_2 = 218 a_0$



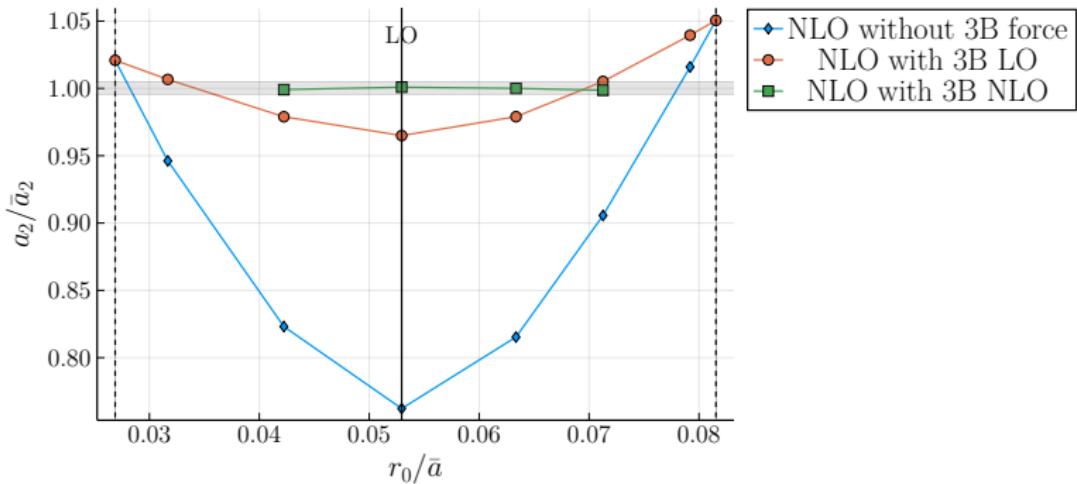
Analysis with NLO 3-Body

- What happens to different three-body observables?
- Atom Dimer scattering length $\bar{a}_2 = 218 a_0$



Analysis with NLO 3-Body

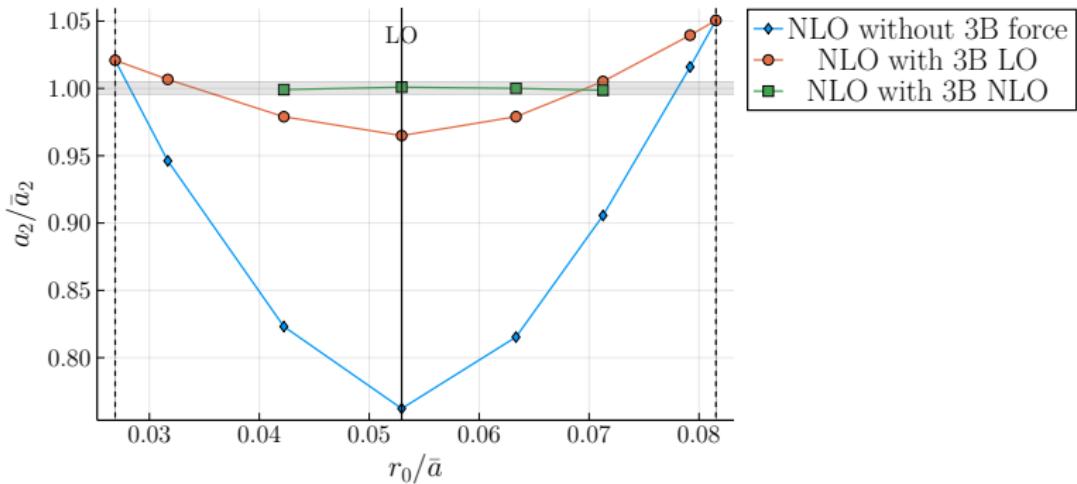
- What happens to different three-body observables?
- Atom Dimer scattering length $\bar{a}_2 = 218 a_0$



- Different 3-Body potential strengths

Analysis with NLO 3-Body

- What happens to different three-body observables?
- Atom Dimer scattering length $\bar{a}_2 = 218 a_0$



- Different 3-Body potential strengths
- Space for a NLO 4-Body potential