

Hadronic vacuum polarization contribution to muon g-2 from four flavors of HISQ quarks

Ruth Van de Water
(for the Fermilab Lattice, HPQCD,
& MILC Collaborations)

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Johannes Gutenberg Universität Mainz

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Participants



Fermilab Lattice Collaboration

- Aida El Khadra (Illinois)
- Andreas Kronfeld (Fermilab)
- Ethan Neil (Colorado)
- Ruth Van de Water (Fermilab)



- Daniel Hatton (Glasgow)
- Christine Davies (Glasgow)
- Peter Lepage (Cornell)
- Craig McNeile (Plymouth)

MILC Collaboration

- | | |
|----------------------------|----------------------------|
| ● Carleton DeTar (Utah) | ● Yuzhi Liu (Indiana) |
| ● Steve Gottlieb (Indiana) | ● Doug Toussaint (Arizona) |
| ● Jack Laiho (Syracuse) | ● Alejandro Vaquero (Utah) |
| | ● Shuhei Yamamoto (Utah) |

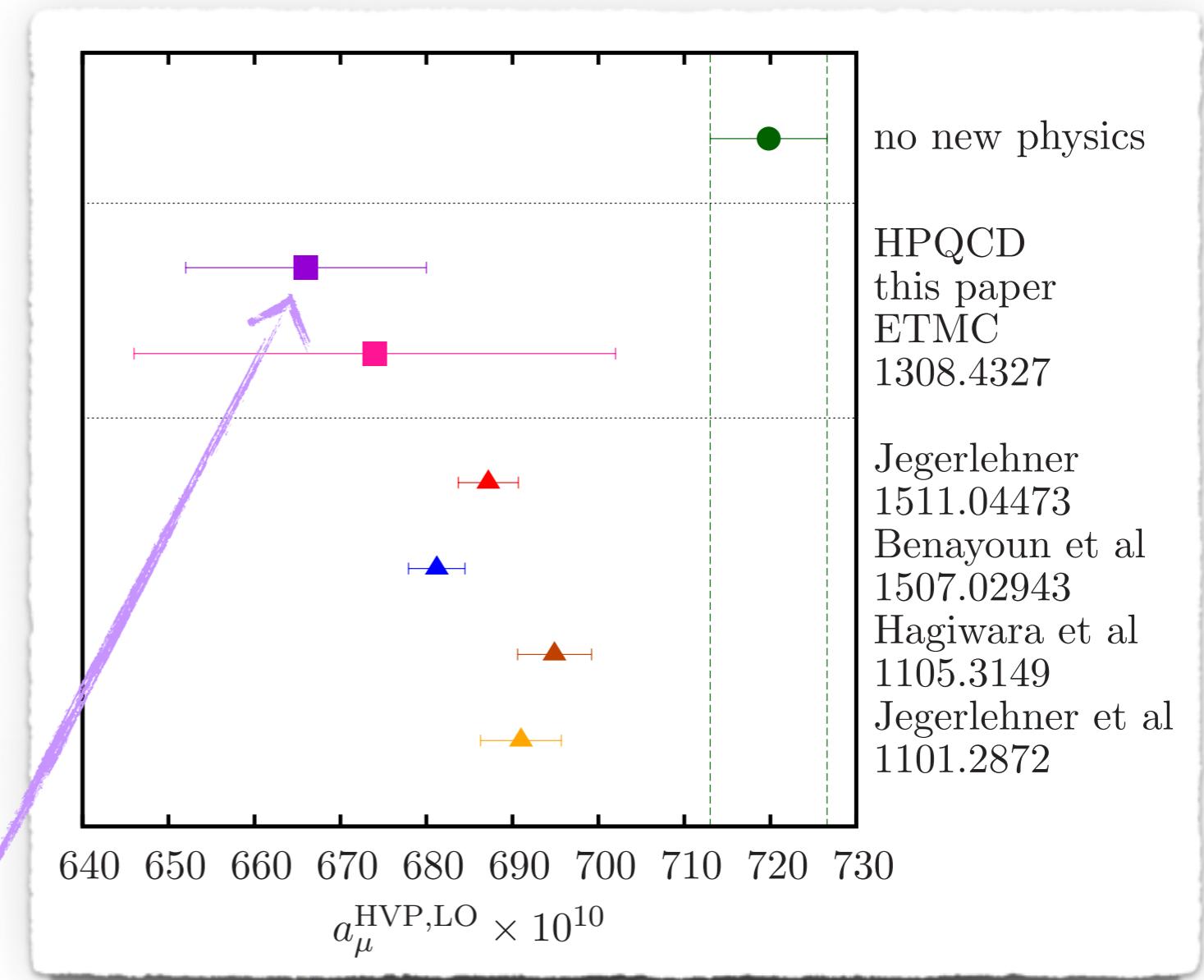
Motivation

- ♦ Muon anomalous magnetic moment ($g-2$) provides sensitive probe of physics beyond the Standard Model:
 - ❖ Mediated by quantum-mechanical loops
 - ❖ Known to very high precision of 0.54ppm
- ♦ Measurement from BNL E821 disagrees with Standard-Model theory expectations by more than 3σ
- ♦ Muon $g-2$ Experiment at Fermilab aims to reduce the experimental error by a factor of four, and *expect first results in 2019!*
- ❖ **Theory error must be reduced to a commensurate level to identify definitively whether any deviation observed between theory and experiment is due to new particles or forces**

► We are using *ab-initio* lattice-QCD to target the hadronic vacuum-polarization (HVP) contribution, which is the largest source of theory error

Background

- ♦ Build upon HPQCD's published work
 - ♦ Introduced "time-moments" method obtain a_μ^{HVP} from zero-momentum vector-current correlation functions [PRD89, no. 11, 114501 (2014)]
 - ♦ Obtained precise s, c-, and b-quark connected contributions [PRD 91, no. 7, 074514 (2015)] and bound on quark-disconnected contribution [HPQCD+HadSpec, PRD 93, no. 7, 074509 (2016)]
 - ♦ Calculated light-quark connected contribution & total $a_\mu^{\text{HVP},\text{LO}}$ with ~2% uncertainty [HPQCD+RV, PRD96 (2017) no. 3, 034516]



Dominant errors from omission of EM & isospin breaking, and from quark-disconnected contributions

Ensembles & parameters

- ♦ Employ large set of **MILC ensembles with four flavors of dynamical HISQ sea quarks with four lattice spacings, multiple spatial volumes, & physical light-quark masses**
- ♦ Vector-current correlator data from [PRD96 \(2017\) no.3, 034516](#) plus:
 - ❖ a~0.15 fm physical-mass ensemble with better-tuned quark masses & **~15×statistics**
 - ❖ **a~0.09 & 0.06 fm physical-mass data**

$\approx a$ (fm)	$am_l^{\text{sea}}/am_s^{\text{sea}}/am_c^{\text{sea}}$	w_0/a	$Z_{V,\bar{s}s}$	M_{π_5} (GeV)	$(\frac{L}{a})^3 \times (\frac{T}{a})$	$N_{\text{conf.}}$
0.15	0.01300/0.065/0.838	1.1119(10)	0.9881(10)	0.3025(16)	$16^3 \times 48$	9947
0.15	0.0064/0.064/0.828	1.12720(70)	0.9881(10)	0.2155(11)	$24^3 \times 48$	1000
0.15	0.00235/0.0647/0.831	1.13670(50)	0.9881(10)	0.13304(70)	$32^3 \times 48$	997
0.15	0.002426/0.0673/0.8447	1.13215(35)	0.9881(10)	0.13473(71)	$32^3 \times 48$	7746
0.12	0.01020/0.0509/0.635	1.3826(11)	0.99220(40)	0.3012(16)	$24^3 \times 64$	1053
0.12	0.0050724c/0.0507/0.628	1.40290(90)	0.99220(40)	0.2178(12)	$24^3 \times 64$	1020
0.12	0.0050732c/0.0507/0.628	1.40290(90)	0.99220(40)	0.2165(11)	$32^3 \times 64$	1000
0.12	0.0050740c/0.0507/0.628	1.40290(90)	0.99220(40)	0.2163(11)	$40^3 \times 64$	331
0.12	0.00184/0.0507/0.628	1.41490(60)	0.99220(40)	0.13273(70)	$48^3 \times 64$	998
0.09	0.00740/0.037/0.440	1.9006(20)	0.99400(50)	0.3076(17)	$32^3 \times 96$	1000
0.09	0.00363/0.0363/0.430	1.9330(20)	0.99400(50)	0.2191(12)	$48^3 \times 96$	298
0.09	0.00120/0.0363/0.432	1.95180(70)	0.99400(50)	0.12834(68)	$64^3 \times 96$	1557
0.06	0.0008/0.022/0.260	3.0170(23)	0.9941(11)	0.13495(72)	$96^3 \times 192$	945

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Analysis updates & improvements

(1) Employ only 5 physical-mass ensembles in main analysis

- ◆ Comparison with fit to all available data provides check

(2) Drop ρ rescaling from analysis

- ◆ Number of $\pi\pi$ levels near or below the ρ resonance increases with decreasing lattice spacing → ground-state energy of correlator not the ρ mass
- ◆ Unnecessary now that analysis does not use ensembles with heavy pions

(3) Include estimate of low-energy $\pi\pi$ contributions from times $t>1.5$ fm

(4) Calculate slope and curvature of vacuum polarization function

- ◆ Enables comparison of more intermediate quantity between lattice-QCD calculations and with experiment

Noise-reduction strategy

- ♦ Reduce statistical errors in light-quark connected contribution to a_μ as in 1601.03071 by:

(1) Simultaneous fit of 4 combinations of (local, smeared) correlators G_{ij}

- ❖ Employ cosh parameterization that accounts for periodic temporal boundary conditions
- ❖ **Fit between $(t_{\min}, T-t_{\min})$** to ensure that fit describes correlator over entire lattice time extent
- ♦ Choose number of states and t_{\min} based on stability of leading energies and amplitudes & goodness-of-fit
- ❖ Employ same $t_{\min} \sim 0.45$ fm for all lattice spacings

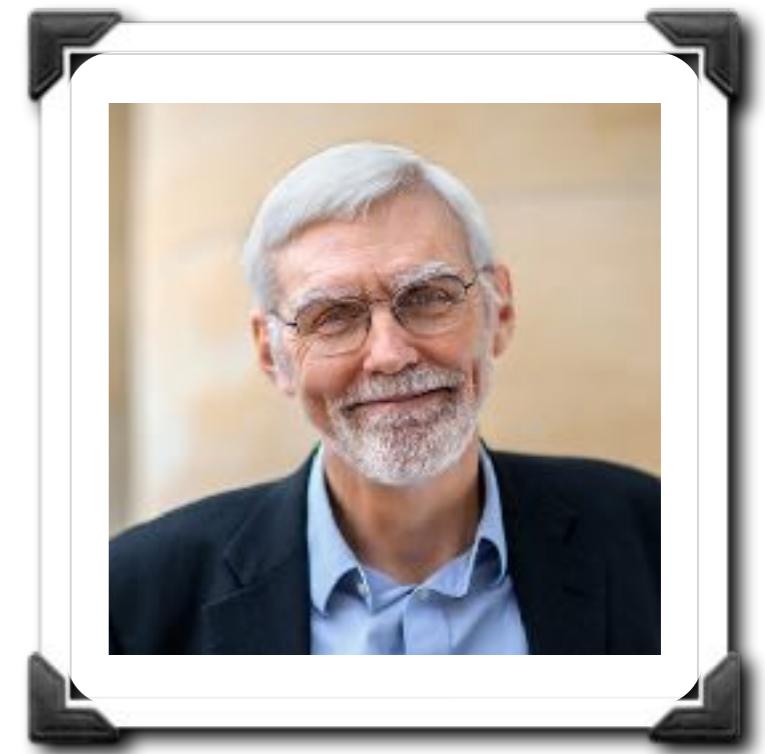
(2) Replacing $G_{\text{data}}(t)$ with $G_{\text{fit}}(t)$ for $t > t^*$

- ❖ Correct for finite temporal extent by calculating $G_{\text{fit}}(t)$ using 2-point fit parameters in exponential parameterization and extending times in $G_{\text{fit}}(t)$ to $2 \times T$

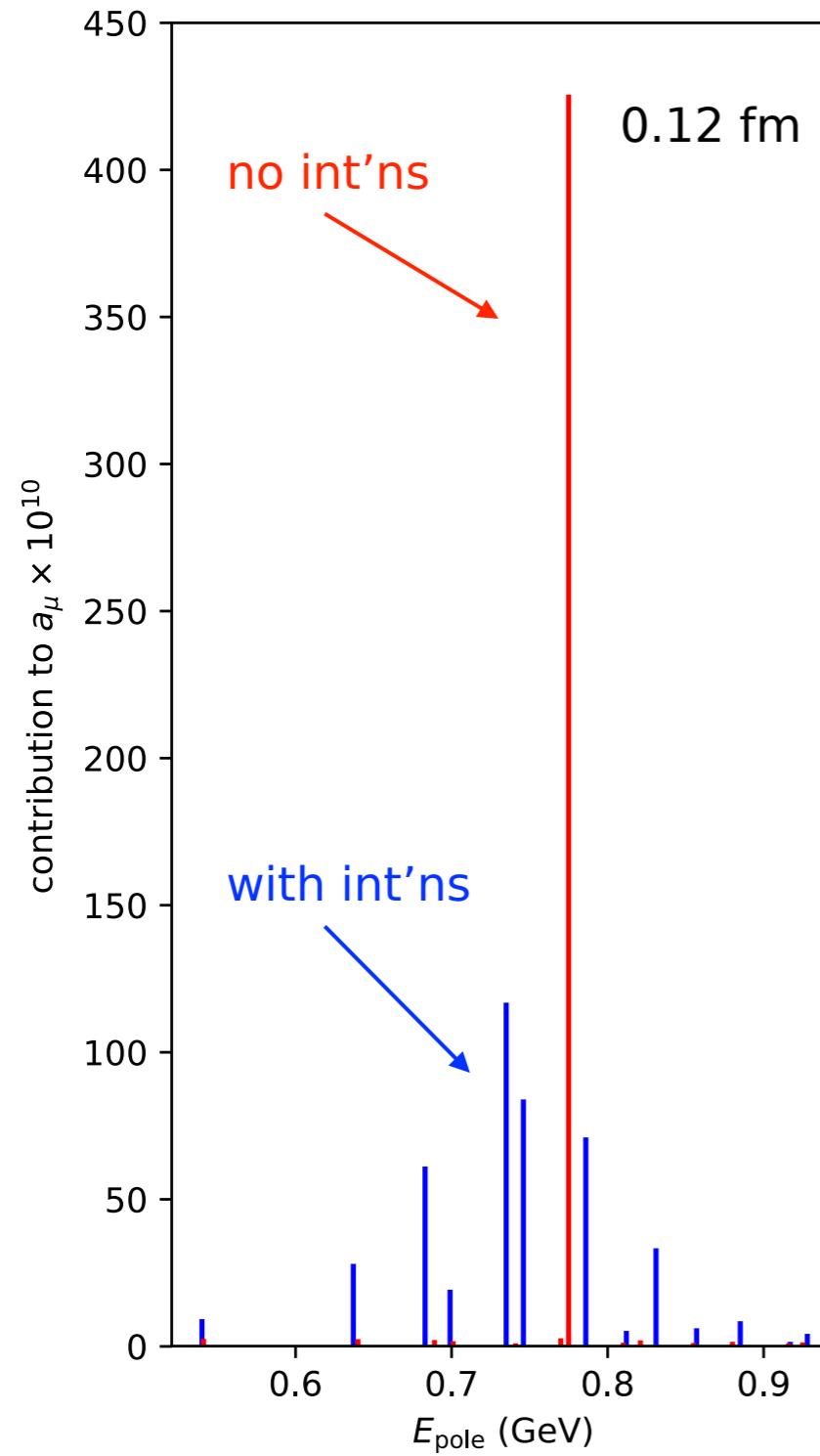
$$G(t) = \begin{cases} G_{\text{data}}(t) & t \leq t^* \\ G_{\text{fit}}(t) & t > t^* \end{cases}$$

Impact of low-energy $\pi\pi$ states

- ◆ True spectrum more complicated than simple fit parameterization employed in our analysis
 - ❖ (multiple ρ and $\pi\pi$ levels and mixing between the two)
- ◆ **Study effects of low-energy $\pi\pi$ states on a_μ^{HVP} within extended chiral perturbation theory that includes π 's, ρ 's, and γ 's [Jegerlehner & Szafron, EPJC71 1632 (2011)]**
 - ❖ (same EFT used to compute finite-volume + staggered discretization corrections)
 - ❖ Compute $\pi\pi$ levels numerically from poles and residues of $\Pi(q^2)$ in Appendix B of **Chakraborty et al., PRDD96 (2017) no.3, 034516**
 - ❖ **Include leading interactions that couple ρ and $\pi\pi$ channels**



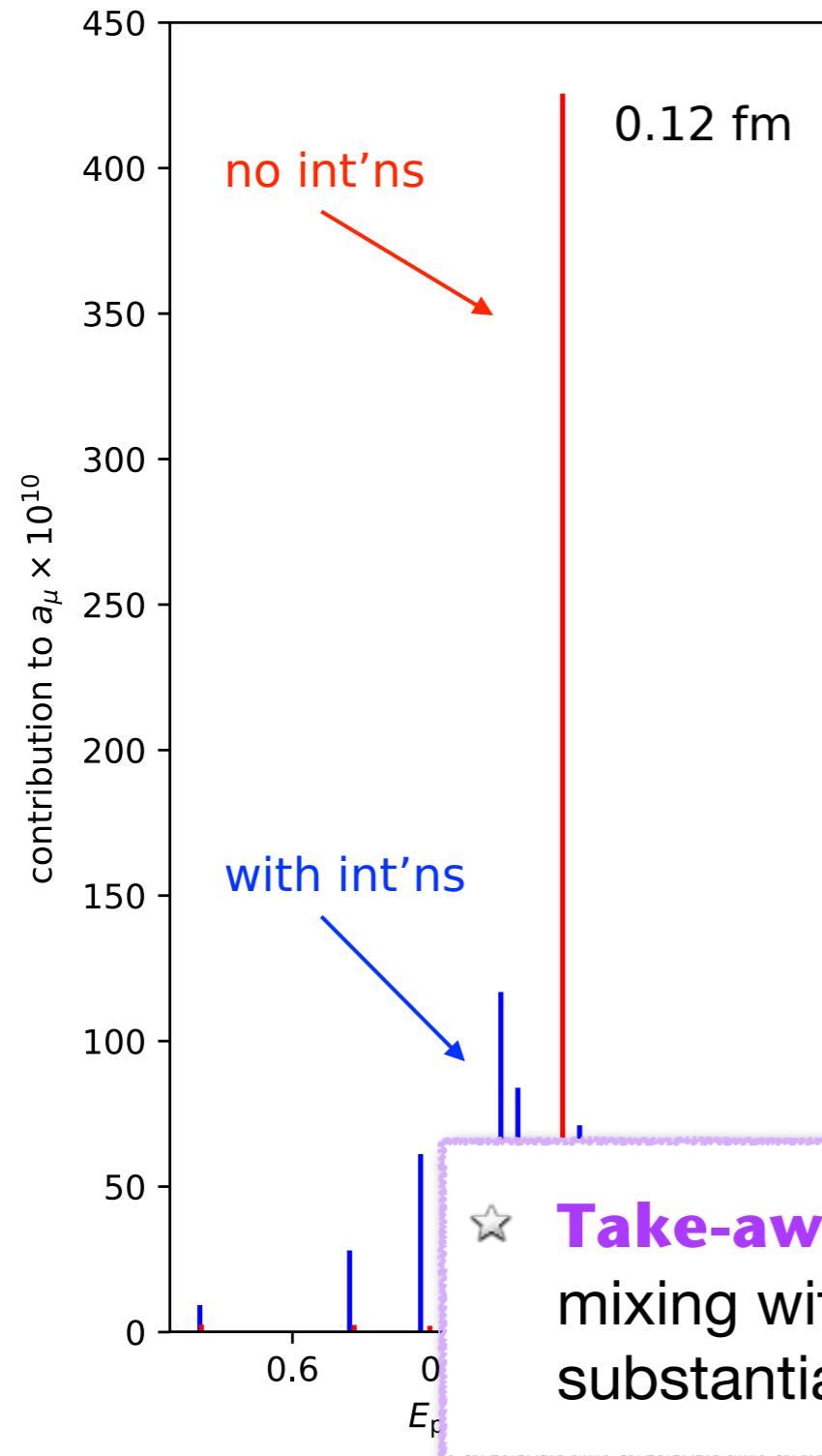
$\pi\pi$ - ρ Contributions to a_μ (Model)



- Without interactions: ρ dominates, $\pi\pi$ negligible (due to finite volume, staggered pions).
- With interactions: ρ mixes with $\pi\pi$ and its contribution spread over many states.
- Total a_μ same in both cases (<0.5%).
- Noisy data \Rightarrow fitter can't tell difference.

Model in B. Chakraborty et al 1601.03071
Phys. Rev. D96 (2017) 034516 (App. B)

$\pi\pi$ - ρ Contributions to a_μ (Model)

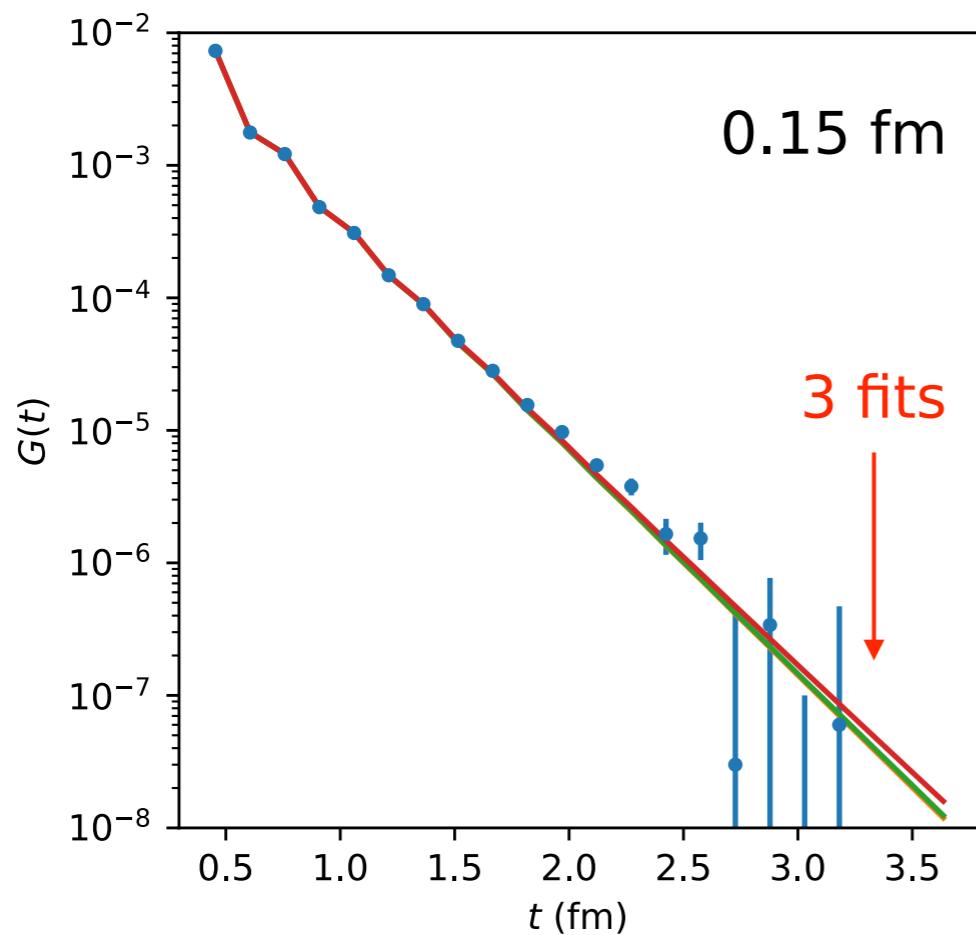


- Without interactions: ρ dominates, $\pi\pi$ negligible (due to finite volume, staggered pions).
- With interactions: ρ mixes with $\pi\pi$ and its contribution spread over many states.
- Total a_μ same in both cases ($<0.5\%$).
- Noisy data \Rightarrow fitter can't tell difference.

★ **Take-away:** $\pi\pi$ interactions and mixing with the ρ change the spectrum substantially, but not the total a_μ^{HVP}

- arXiv:1601.03071
4516 (App. B)

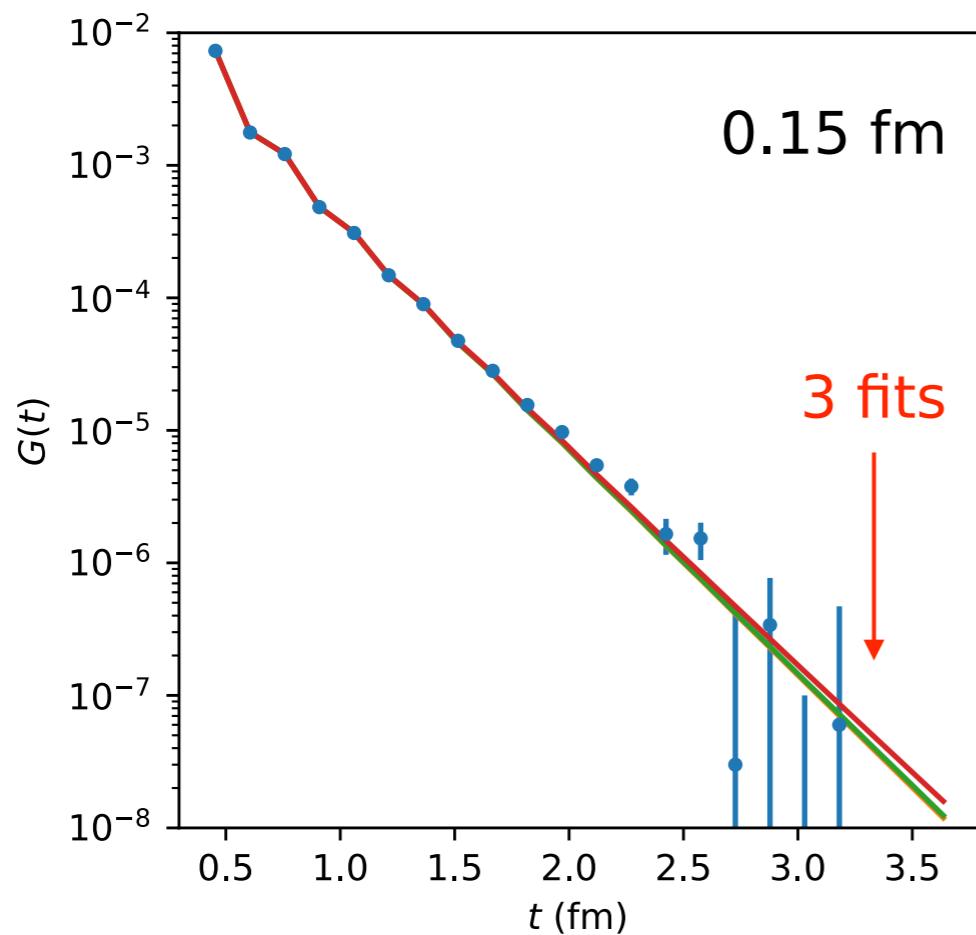
Equivalent Fits ($\chi^2/\text{dof} \leq 1$)



- Fit unable to resolve difference between one or multiple rho mesons; spreads contribution over multiple terms.
- Negligible difference for $t < 4\text{fm}$, and irrelevant for a_μ .

	$E_0 \setminus a_0$	$E_1 \setminus a_1$	$E_2 \setminus a_2$	$E_3 \setminus a_3$	$a_\mu \times 10^{10}$
Fit 1	0.770 \ 0.132	1.72 \ 0.35	2.5 \ 0.1		602(7)
Fit 2	0.75 \ 0.09	0.80 \ 0.09	1.75 \ 0.36		602(7)
Fit 3	0.71 \ 0.08	0.78 \ 0.04	0.84 \ 0.10	1.76 \ 0.36	608(8)

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	$E_0 \setminus a_0$	$E_1 \setminus a_1$	$E_2 \setminus a_2$	$E_3 \setminus a_3$	$a_\mu \times 10^{10}$
Fit 1	$0.770 \setminus 0$	★ Take-away: can parameterize data over a finite t range (<3-4fm) by a fit with a single ρ			602(7)
Fit 2	$0.75 \setminus 0$				602(7)
Fit 3	$0.71 \setminus 0$				608(8)

Missing $\pi\pi$ contributions

- ◆ correlator fits will catch large contributions from low-energy $\pi\pi$ states that mix substantially with the ρ
- ❖ Calculate smaller missing contributions from $\pi\pi$ states below M_ρ within (non-interacting) chiral perturbation theory, and add $G_{\pi\pi}(t)$ to $G_{\text{fit}}(t)$ for $t > t^*$ fm

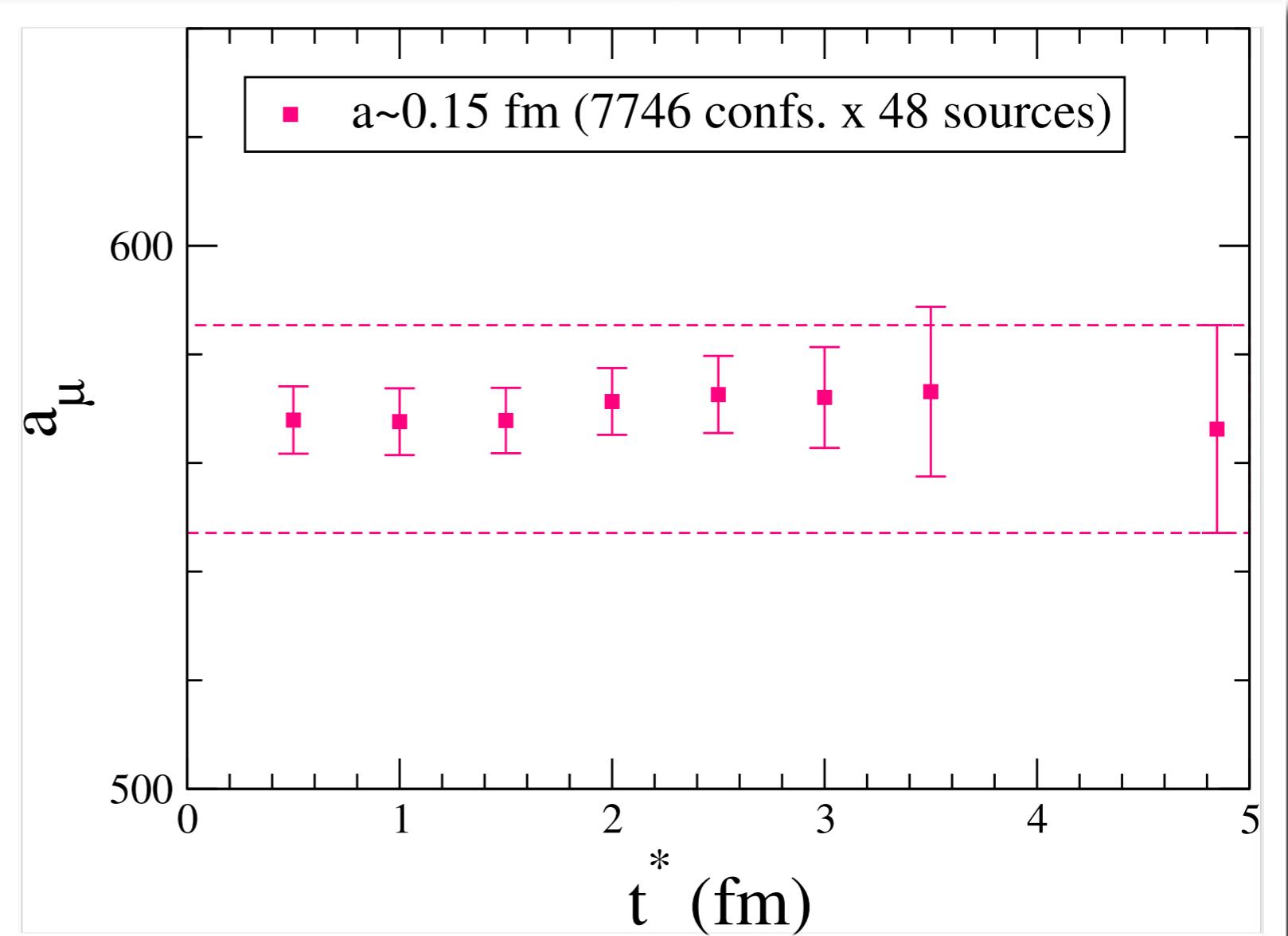
$$G_{\pi\pi}(t) = \sum_{\substack{\{\pi_1, \pi_2\} \in \\ \text{taste } V}} \frac{1}{3L^3} \sum_k \frac{k^2 \exp[-(E_1 + E_2)t]}{E_1 E_2}$$
$$E_i = \sqrt{k^2 + m_{\pi_i}^2}$$

- ❖ Estimated low-energy $\pi\pi$ contribution to a_μ^{HVP} from $t > 1.5$ fm *that is not included in the fit* ranges from about 0.2% at $a \sim 0.15$ fm to 2.5% at $a \sim 0.06$ fm

High-statistics *test*

- ◆ Test method to reduce noise and include low-energy $\pi\pi$ contributions on new high-statistics ensemble where a_μ^{HVP} is known precisely from data

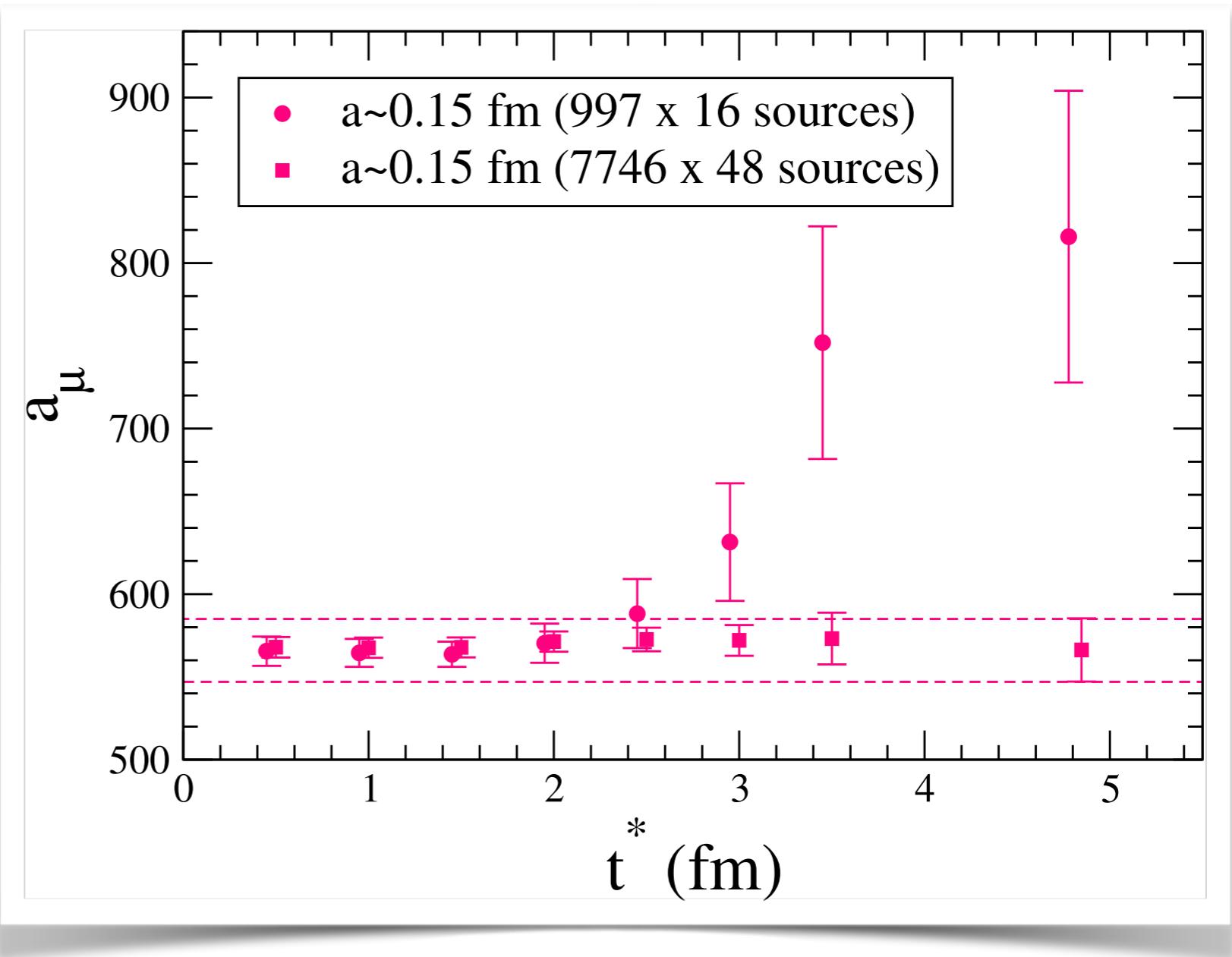
★ a_μ^{HVP} obtained from data + fit + estimated $\pi\pi$ contributions consistent with data-only value for all values of t^*



Comparison with noisy data

- ◆ Compare with results on independent ensemble with same lattice spacing and slightly different quark masses, but substantially lower statistics
- ◆ Answers consistent until $t^* \sim 2$ fm, but diverge substantially beyond $t^* \sim 3$ fm

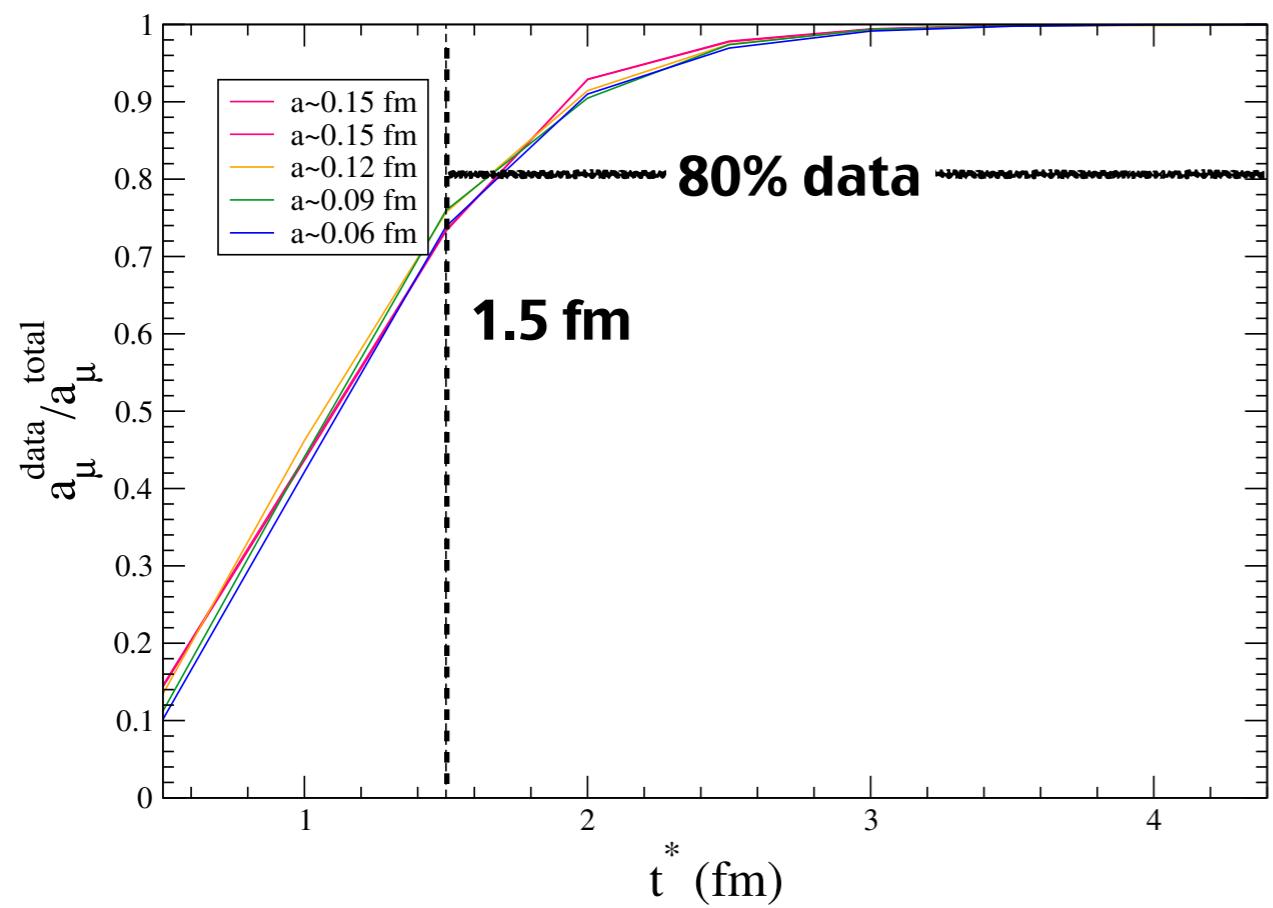
★ Can obtain a_μ^{HVP} from data + fit for t^* below $\sim 2\text{--}2.5$ fm



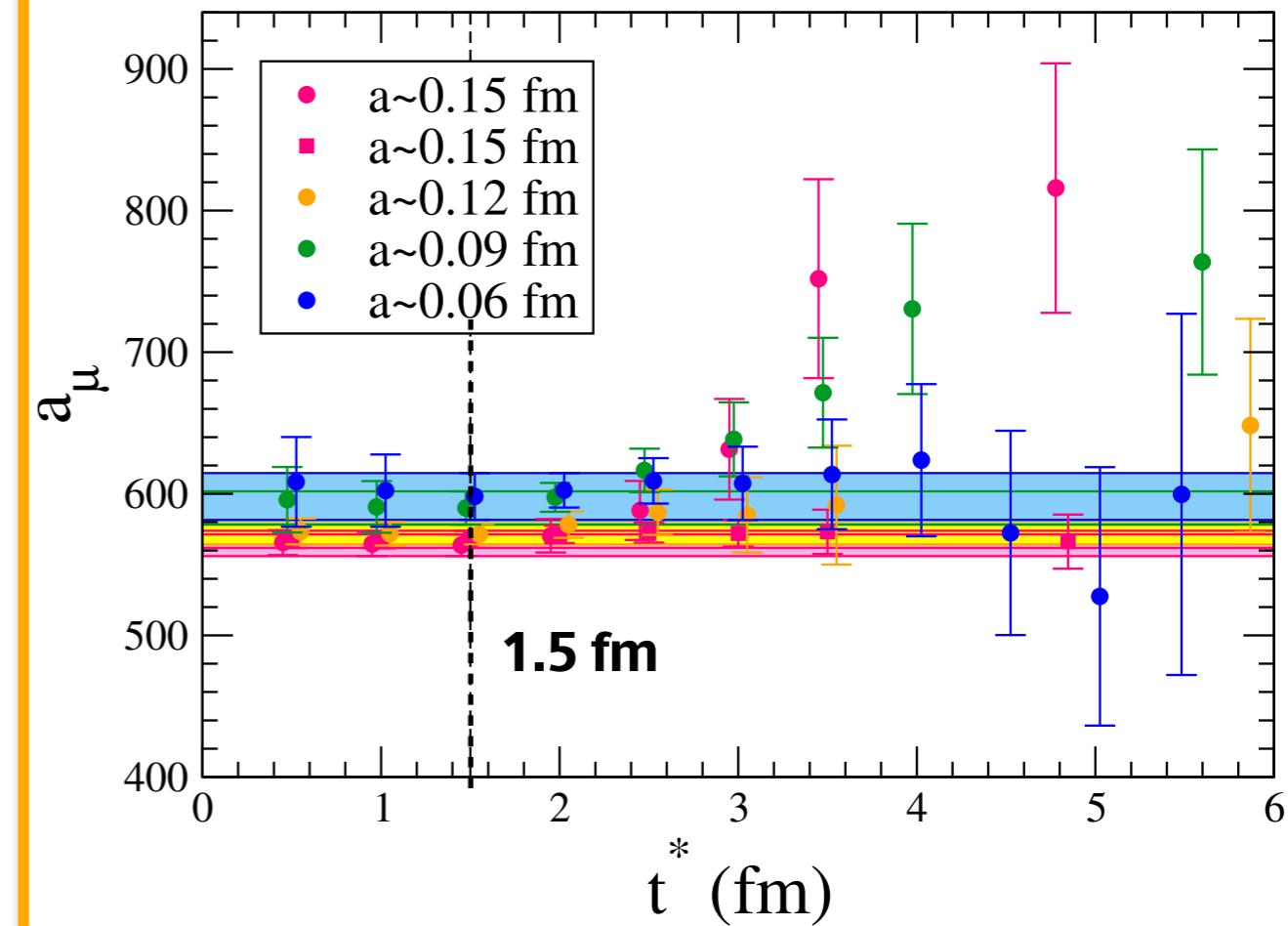
Selection of t^*

- ♦ For $t^* = 1.5 \text{ fm}$, a_μ^{HVP} comes primarily from data region, and errors still controlled

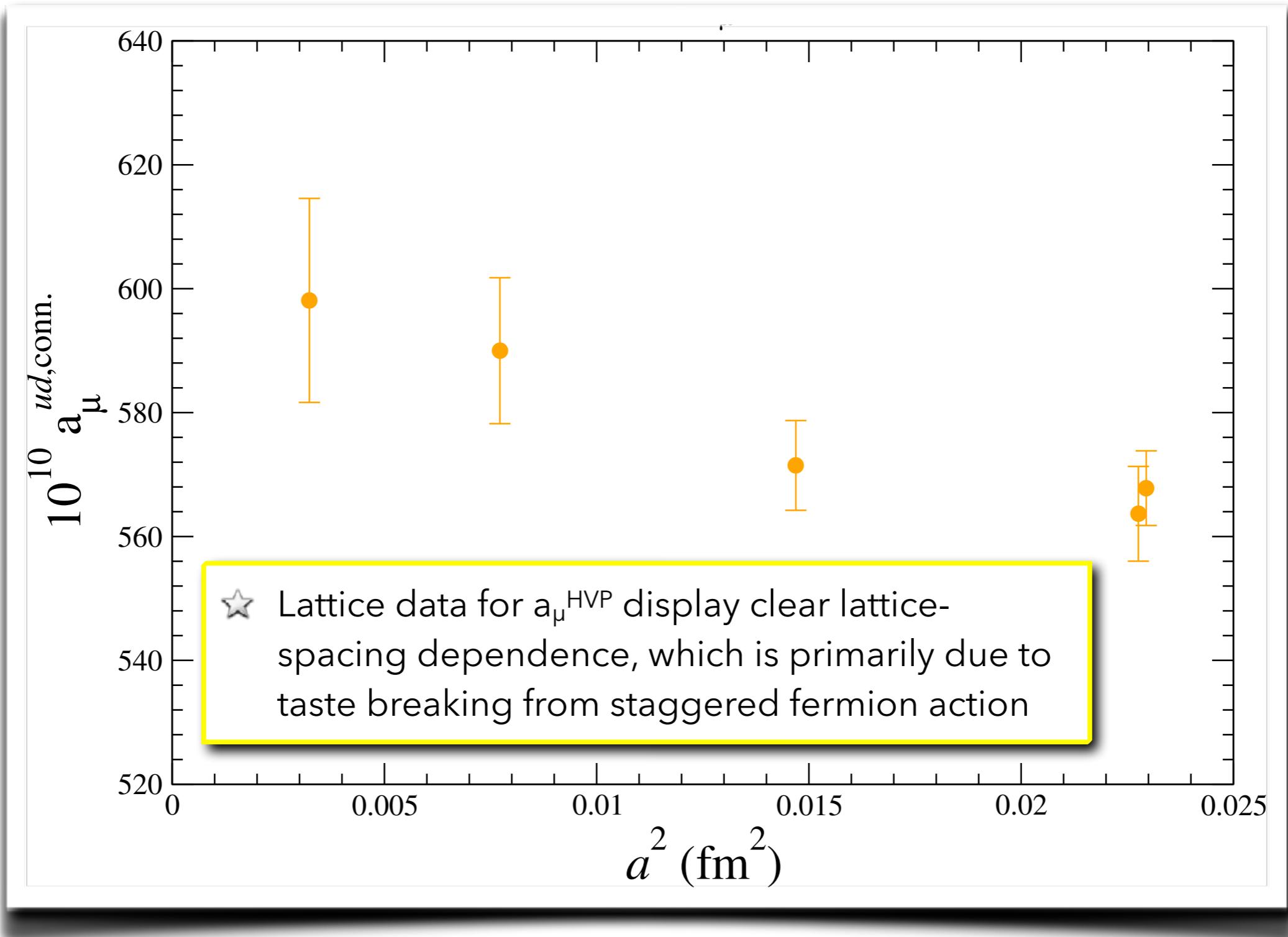
Data contribution



Stability



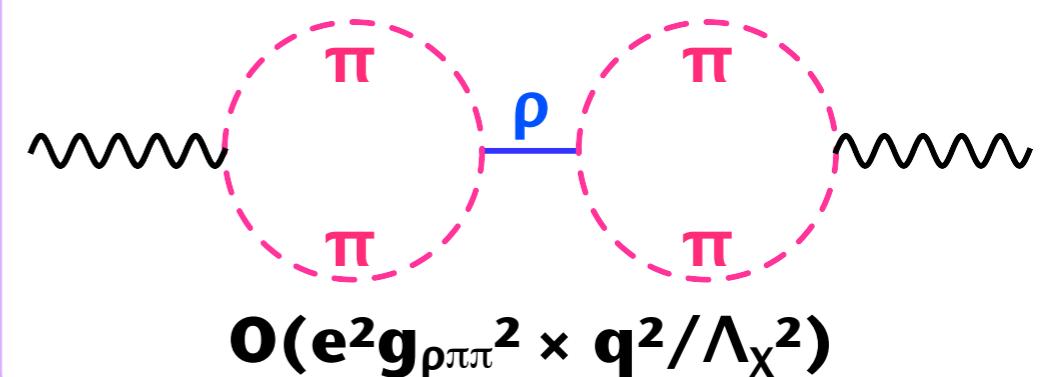
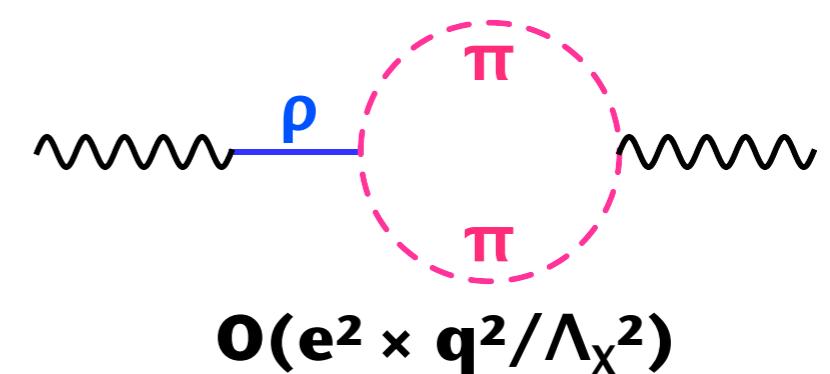
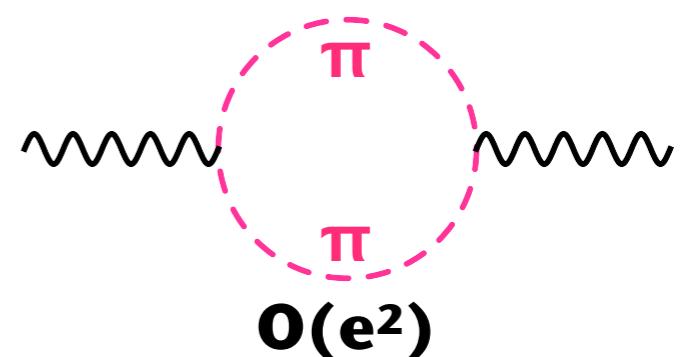
Approach to continuum



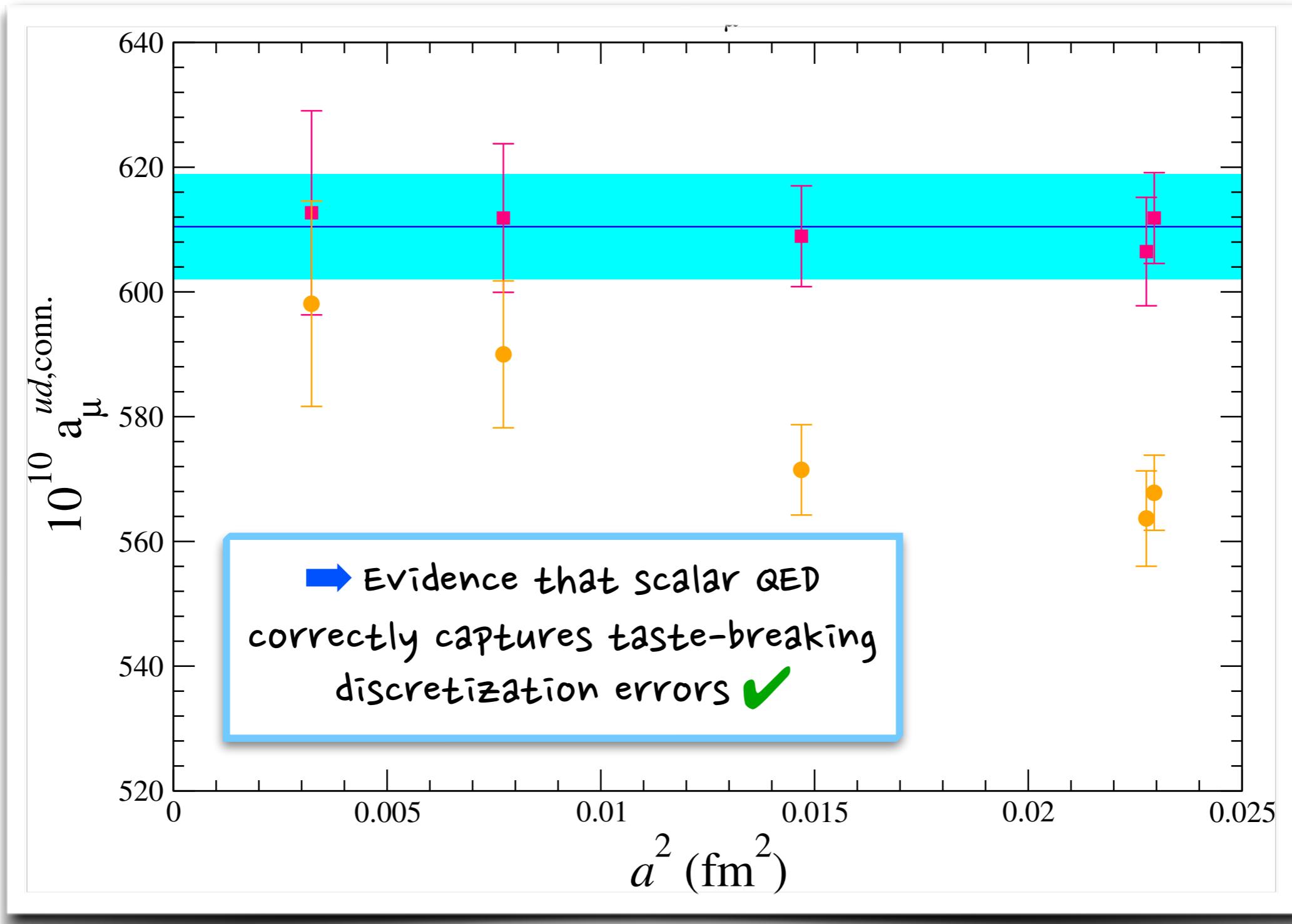
Lattice corrections

- ◆ Correct a_μ^{HVP} for finite-volume and taste-breaking discretization effects before continuum extrapolation
- ◆ Calculate 1-pion-loop contributions to a_μ^{HVP} to all orders in leading interactions that couple ρ^0 - $\gamma\pi^+\pi^-$ channels
 - ❖ Corrections given by difference between results in continuum, infinite volume, and in finite volume with lattice artifacts
 - ❖ corrections largely from taste splittings, and decrease with lattice spacing

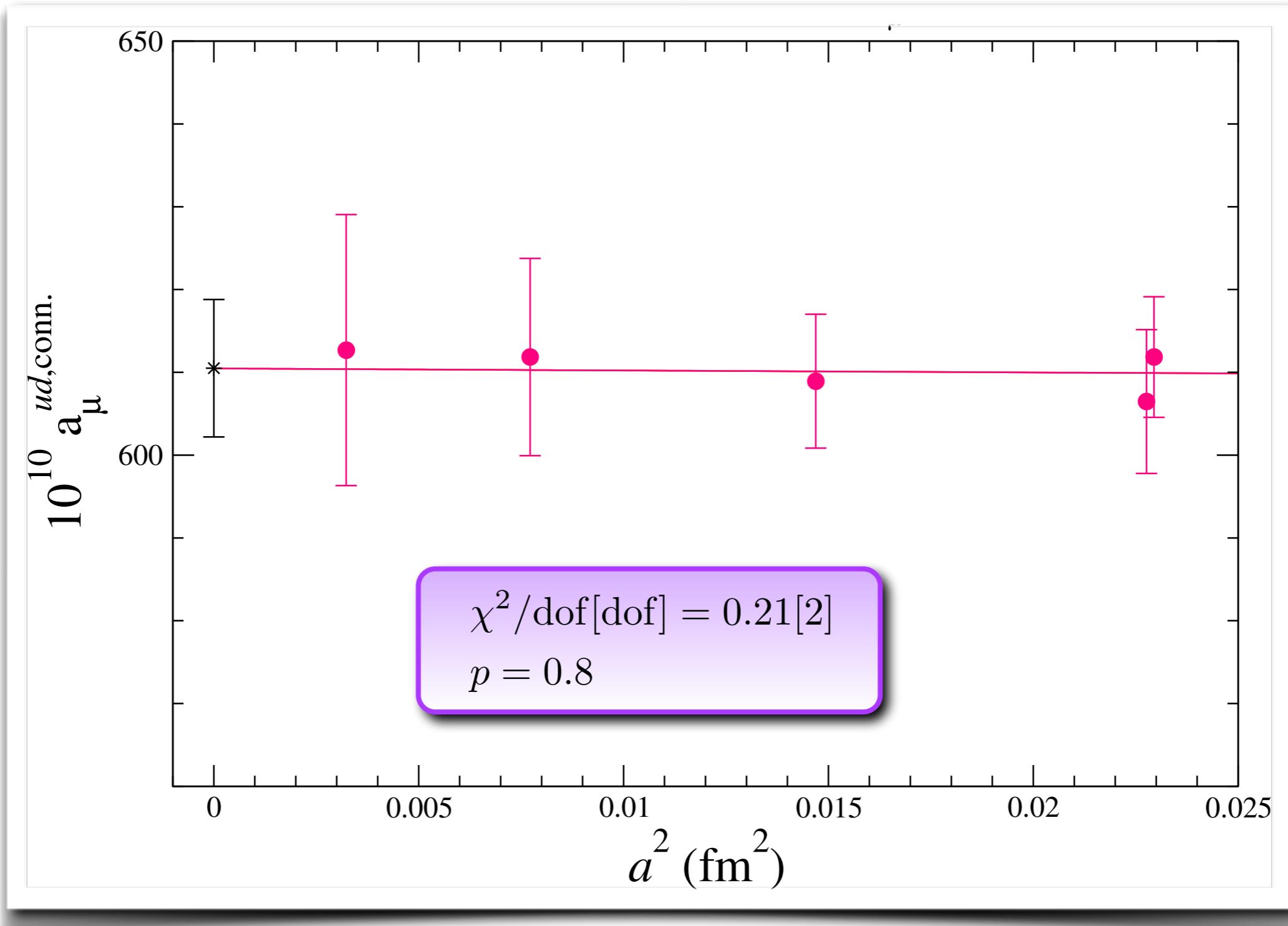
(diagrams below + all iterations of these diagrams)



Residual a^2 dependence



Continuum extrapolation



Error budget for $a_\mu^{ud,\text{conn.}}$

- ◆ Shown for isospin-limit quantity → include QED & isospin-breaking errors only in total LO HVP contribution

	$a_\mu^{ud,\text{HVP}} \text{ (%)}$	
	HPQCD+RV	Fermilab-MILC-HPQCD
	1601.03071	(preliminary)
Continuum ($a \rightarrow 0$) extrapolation	0.2	0.8
$\pi\pi$ states (t^*)	0.5	0.6
Finite-volume & discretization corrections	0.7	0.4
Statistics + 2pt fit	0.4	0.3
Sea (m_s) adjustment	0.2	0.1
Current renormalization (Z_V)	0.2	0.1
Lattice-spacing (a^{-1}) uncertainty	<0.05	0.8
Padé approximants	0.4	0.0
Fit total	1.1%	1.4%

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Two finer
lattice
spacings

Error budget for $a_\mu^{ud,\text{conn.}}$

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	$a_{ud}^{\text{HVP}} (\%)$	C-HPQCD
HPQCD	160	160
Continuum ($a \rightarrow 0$) extrapolation	0.2	0.8
$\pi\pi$ states (t^*)	0.5	0.6
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Include missing $\pi\pi$
 states directly in $G(t)$

Error budget for $a_\mu^{ud,\text{conn.}}$

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Use [3,3]
Padés

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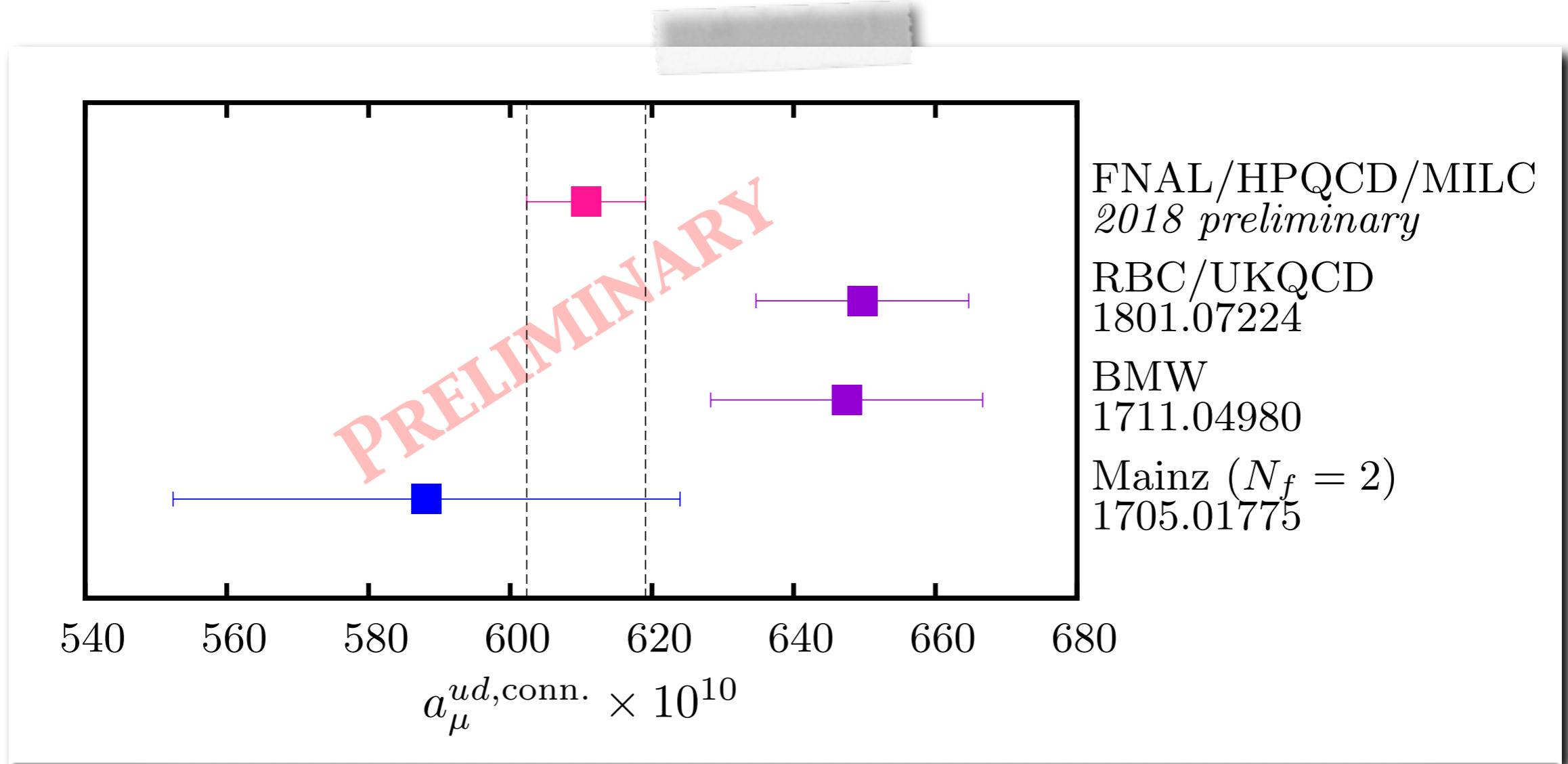
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Padé approximants	0.4		0.0
Fit total	1.1%		1.4%

Use [3,3]
Padés

- Increased (preliminary!) uncertainty primarily from dropping unphysical-mass ensembles and E_0 rescaling, but analysis still in progress

Comparison with other work

- ◆ Results shown are for isospin-limit quantity without QED & isospin-breaking corrections

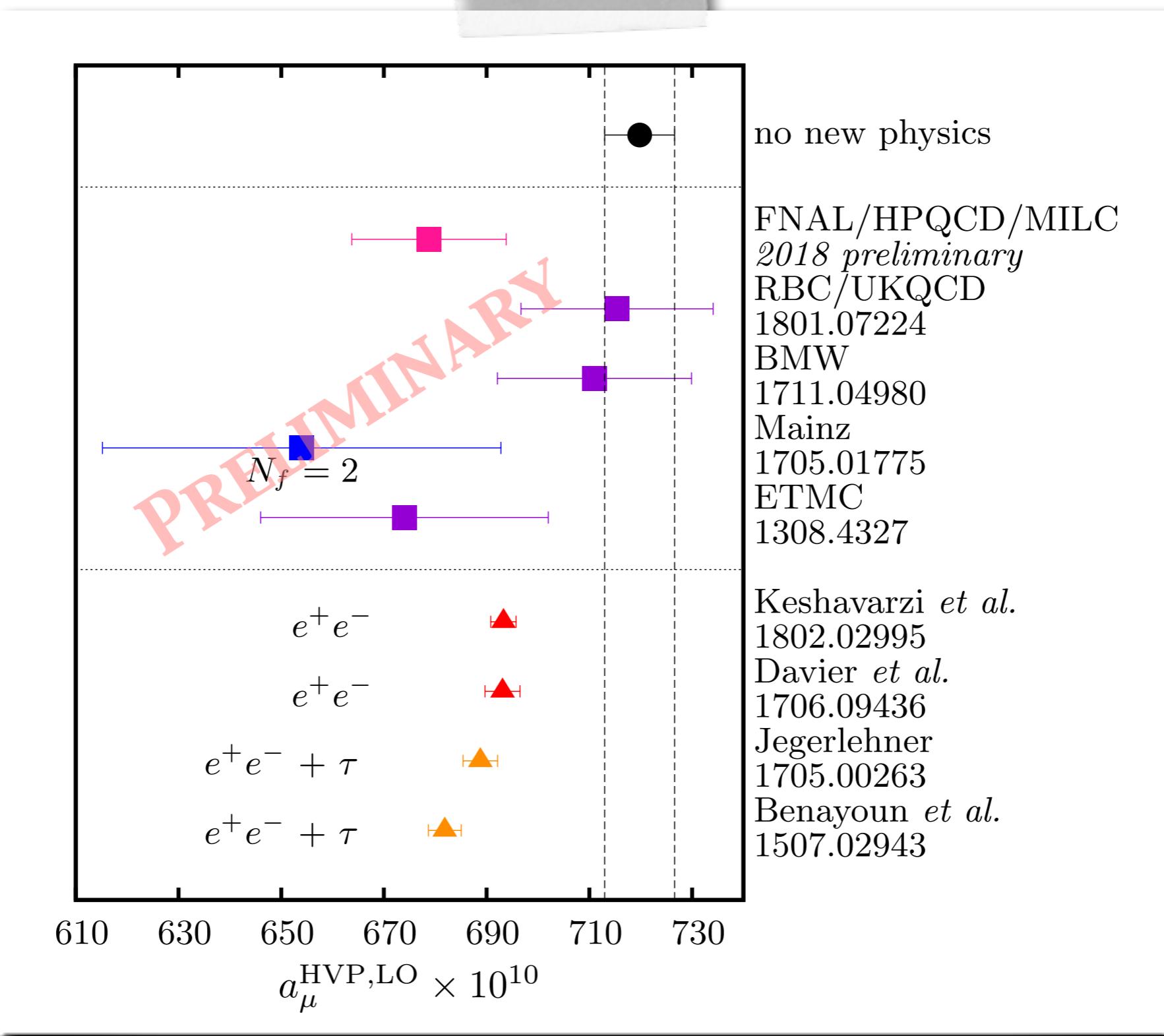


Difference between BMW & RBC/UKQCD reduced, but still ~ 20

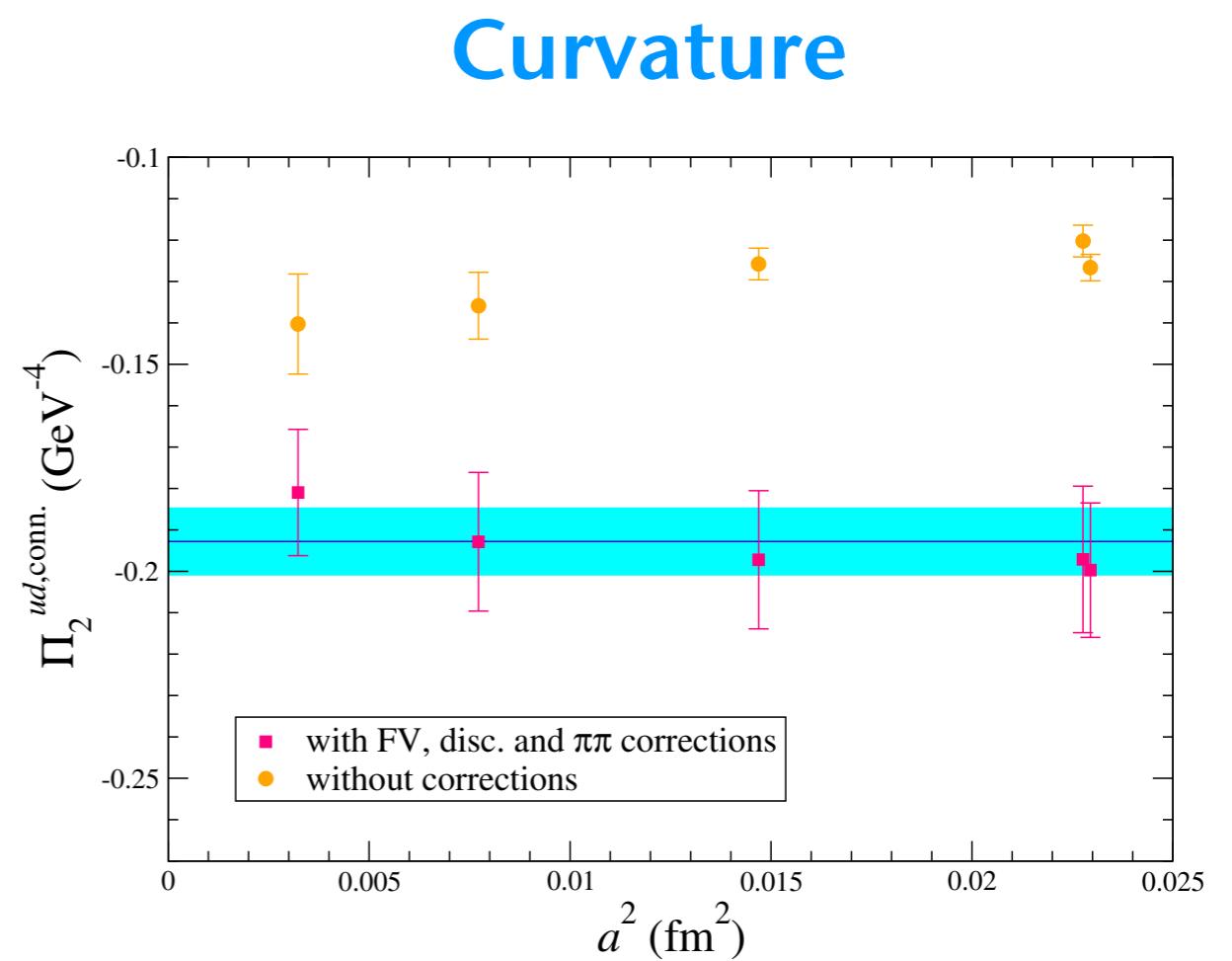
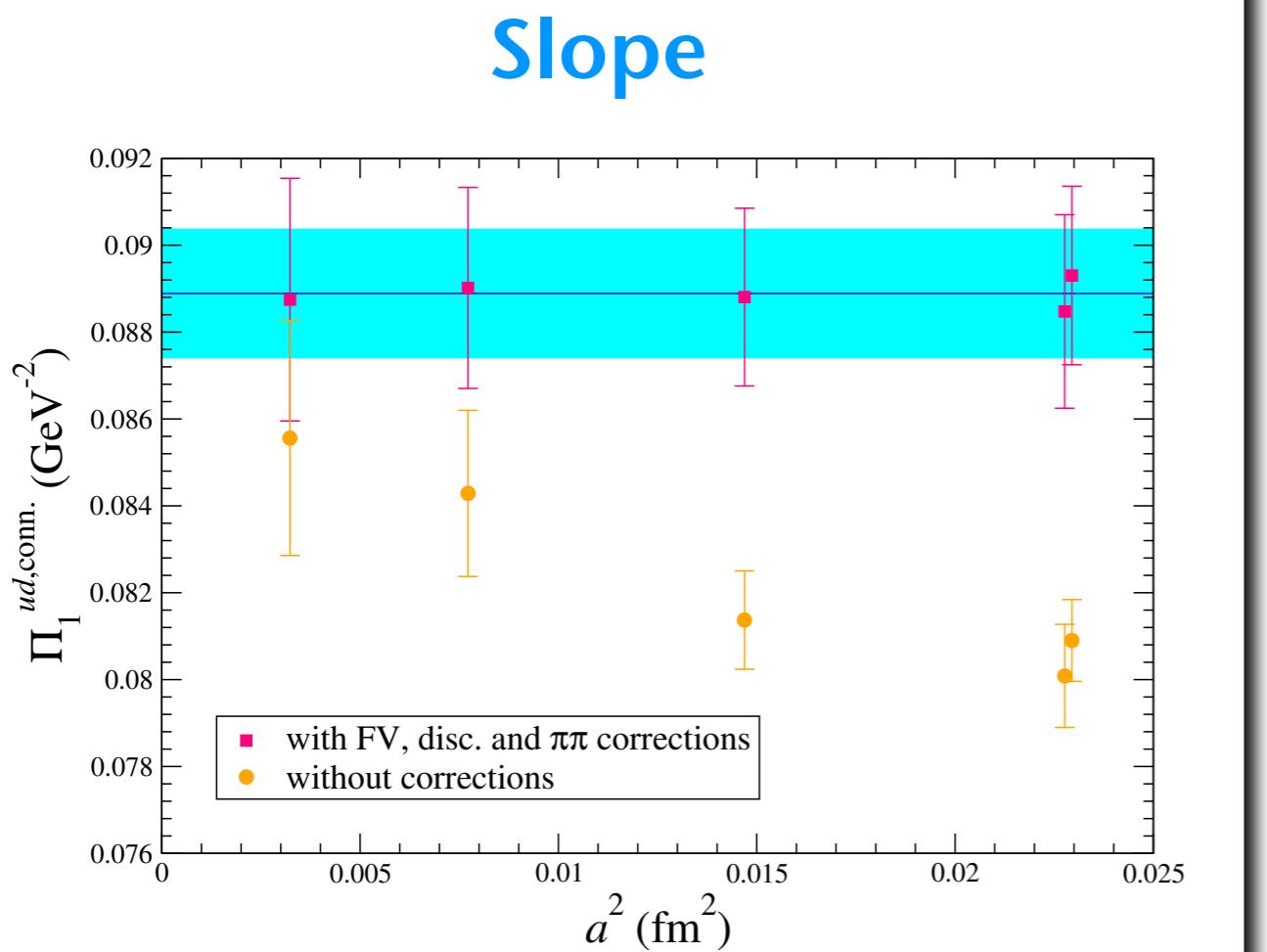
Total leading-order a_μ^{HVP}

- ♦ As in [HPQCD+RV \[PRD96 \(2017\) no.3, 034516\]](#), add to light-quark connected contribution:
 - (1) [s, c-, and b-quark connected contributions from HPQCD](#) [PRD 91, no. 7, 074514 (2015)]
 - (2) [Quark-disconnected contribution from HPQCD+HadSpec](#) [PRD 93, no. 7, 074509 (2016)]
 - (3) [1% \$\oplus\$ 1% = 1.4% error to account for omission of QED and isospin breaking](#)
- ★ wait to include direct estimate of strong-isospin-breaking correction when we also have QED contribution, since cancellations are expected

Total leading-order a_μ^{HVP}



New! Taylor coefficients of $\Pi(q^2)$



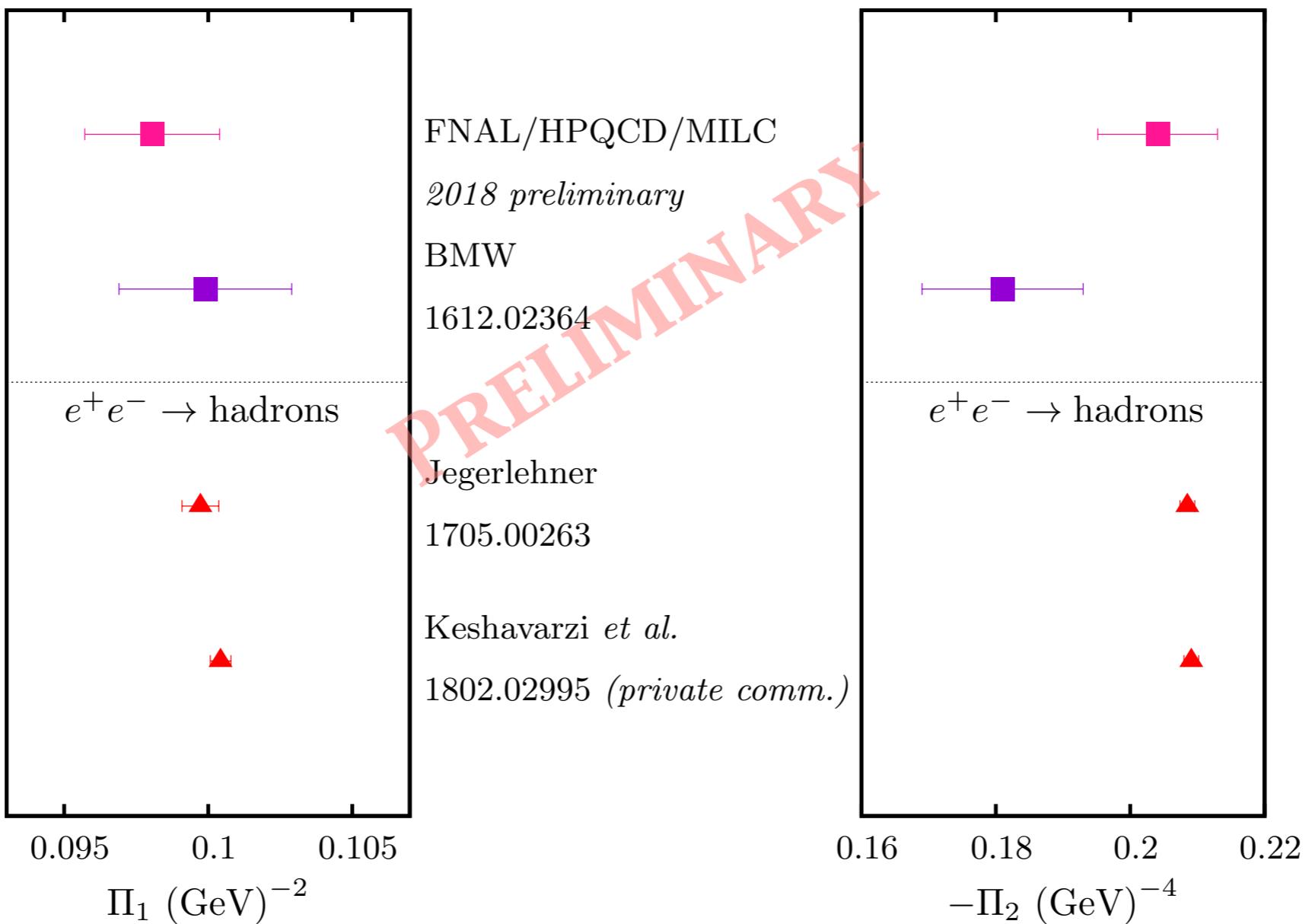
$$\chi^2/\text{dof}[\text{dof}] = 0.05[2]$$

$$p = 0.95$$

$$\chi^2/\text{dof}[\text{dof}] = 0.46[2]$$

$$p = 0.63$$

Comparison with other results



consistent with R-ratio determinations,
but tension with BMW for Π_2

Work in progress

- ♦ Current and proposed running focused on reducing leading sources of error from:

(1) Statistics

- ❖ Better statistics will allow increased t^* and reduce low-energy $\pi\pi$ uncertainty
- ❖ More precise $a \sim 0.09$ fm and 0.06 fm data will reduce (*dominant*) continuum-extrapolation error

(2) Electromagnetic effects

- ❖ Need direct calculation to replace (*and reduce error from*) phenomenological estimate

(3) Quark-disconnected contributions

- ❖ Contributes roughly half of total uncertainty on a_μ^{HVPLO} !

(4) Low-energy $\pi\pi$ contributions

- ❖ BlueWaters allocation to calculate $\langle \rho | \rho \rangle$, $\langle \rho | \pi\pi \rangle$, and $\langle \pi\pi | \pi\pi \rangle$ on $a \sim 0.15$ fm physical-mass ensemble and calculate levels using GEVP
- ❖ Will enable check of chiral-perturbation-theory estimate

QED+QCD ensemble

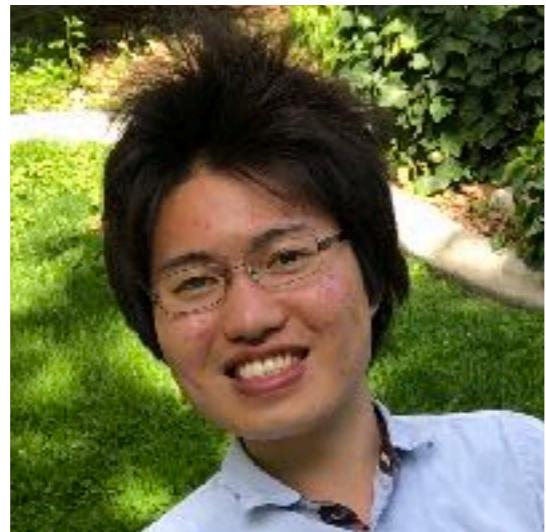
- ◆ Generation of dynamical HISQ QCD + QED ensemble in progress
 - ❖ Noncompact QED_L with formulation of Hayakawa & Uno [Prog.Theor.Phys.120:413-441,2008] for the zero-mode subtraction
 - ❖ **a~0.15 fm with physical quark masses & EM charges**
(same parameters as $N_f = 1+1+1+1$ QCD ensemble)
 - ❖ About 300 thermalized configurations available
- ◆ Presently checking spectrum code; will then compare with previous quenched-QED results
- ◆ **Theoretical work to understand and estimate finite-volume corrections still needed**

Yuzhi Liu



Quark-disconnected contribution

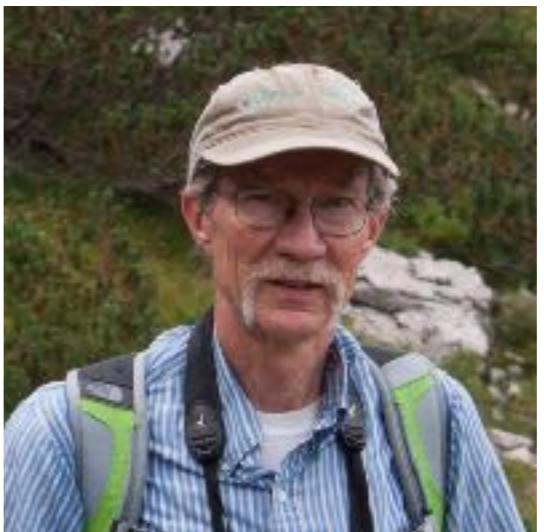
- ◆ Calculating quark-disconnected contributions on a~0.15 fm physical-mass ensemble
- ◆ One-link-split HISQ current density
- ◆ To reduce statistical errors, employ:
 - (1) Deflation
 - (2) Dilution
 - (3) Truncated Solver Method



Shuhei Yamamoto



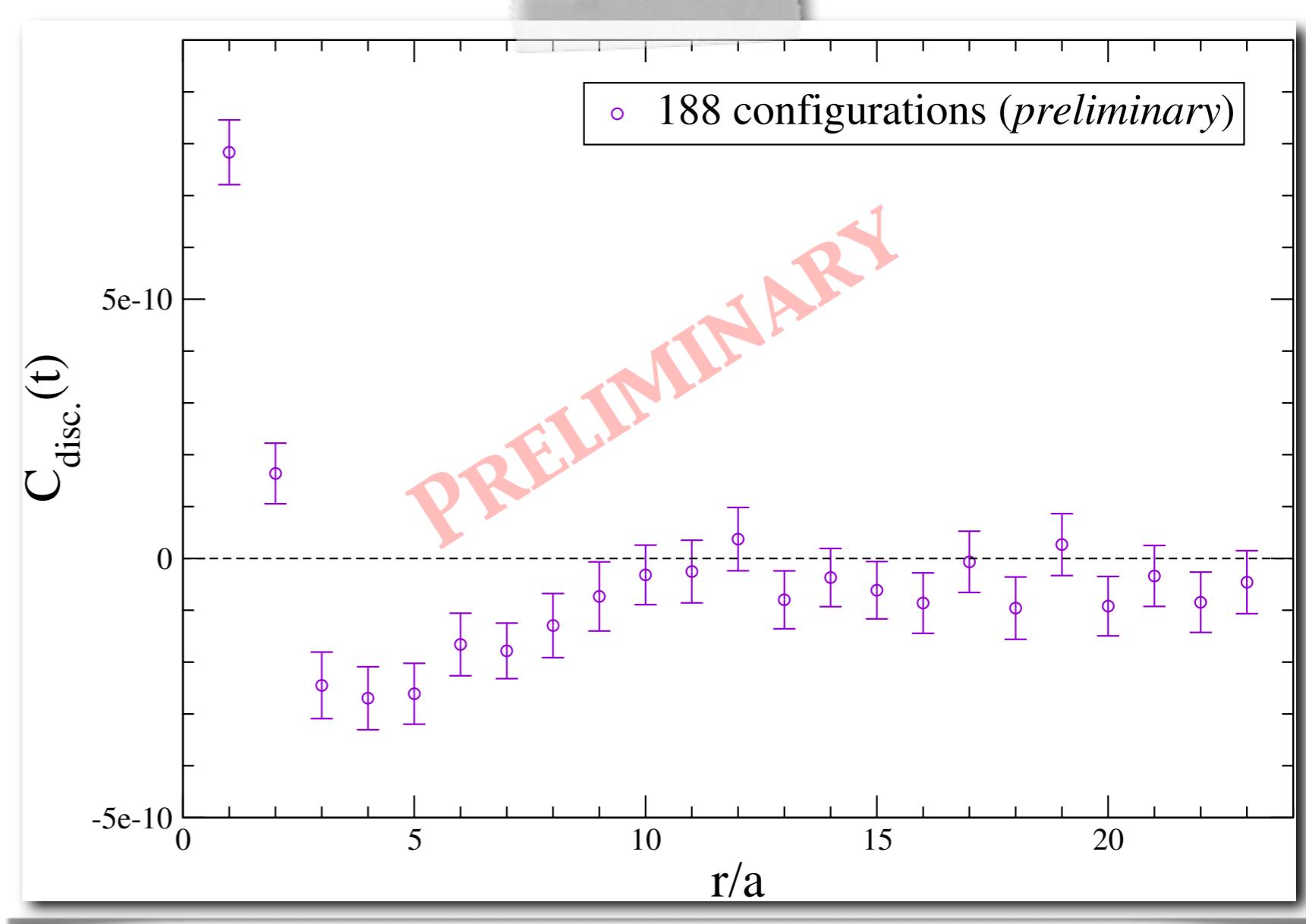
Craig



Carleton

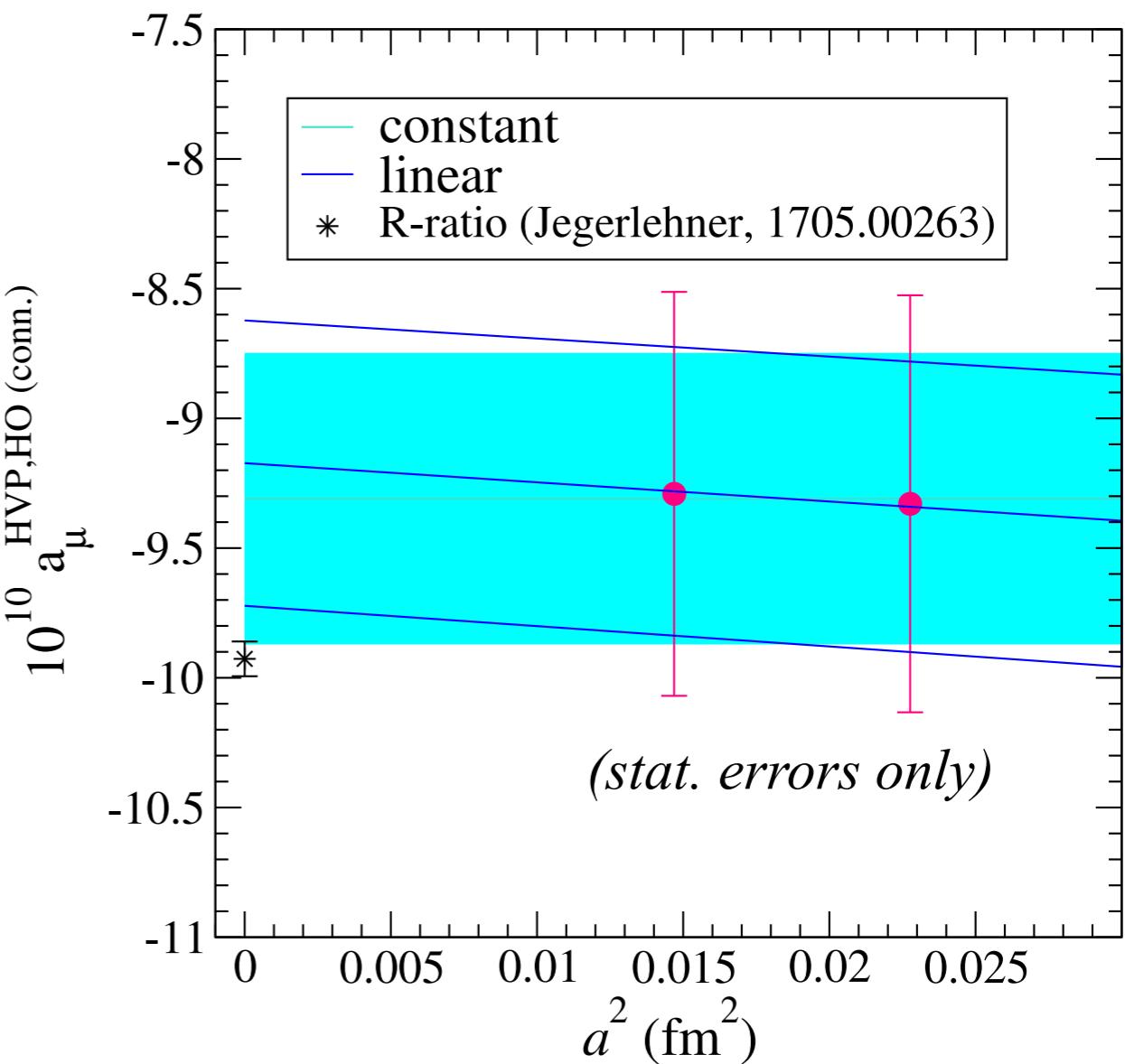
First results

- ♦ Include contributions of u, d, s, and c quarks
- ♦ **Obtain statistically-significant signal with ~200 configurations:**
hope to increase by factor of four within few weeks
- ♦ Next will employ same methodology on a ~ 0.12 physical-mass ensemble to take continuum limit



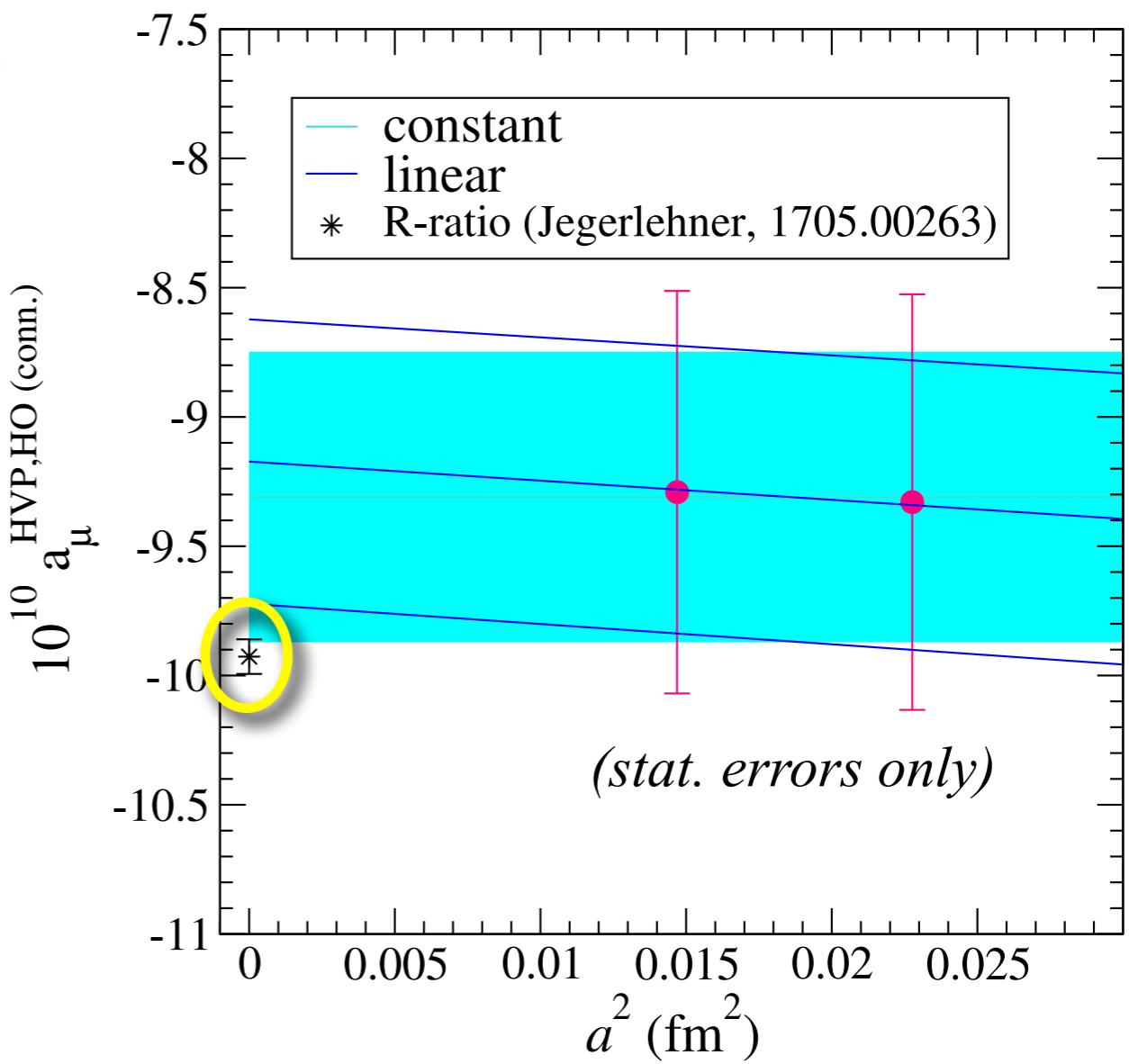
Advertisement! NLO HVP contribution

- ♦ Lattice QCD can also provide higher-order HVP contributions (inspired by first Muon $g-2$ Theory Initiative workshop a year ago)
 - ♦ Not critical to reach target experimental uncertainty, but good check of lattice methods
- ♦ Paper with first lattice-QCD result for NLO HVP contribution to appear soon
 - ♦ Employs published quark-connected Taylor coefficients at two lattice spacings from HPQCD
 - ♦ Provides new expressions suitable to compute $O(\alpha_{EM}^3)$ HVP contributions from renormalized $\Pi(q^2)$ or Euclidean vector-current correlator



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consistent — but not (yet)
competitive — with experiment

Summary

- ❖ Since last year, paper on first lattice-QCD calculation of strong isospin-breaking effects in LO HVP contribution at the physical pion mass [PRL120 (2018) no.15, 152001]
- ❖ Analysis of light-quark connected HVP contribution from 5 physical-mass ensembles with $a \sim 0.15 - 0.06$ fm is in progress
 - ❖ New high-statistics $a \sim 0.15$ fm ensemble enables explicit check of replacing data with fit and estimate of missing low-energy $\pi\pi$ contribution above $t^* = 1.5$ fm
 - ❖ Both a_μ^{HVP} & leading Taylor coefficients of vacuum polarization function agree with determinations from R-ratio data (**preliminary**)
- ❖ Current focus is improving statistics & reducing leading sources of error in 2016 result for $a_\mu^{\text{HVP,LO}}$ from **omission of electromagnetism & quark-disconnected contribution**
 - ❖ First signal for quark-disconnected contribution!

► With direct determinations of these corrections / contributions in hand, plus additional statistics, hope to obtain sub-per cent precision for total LO HVP contribution to muon g-2!

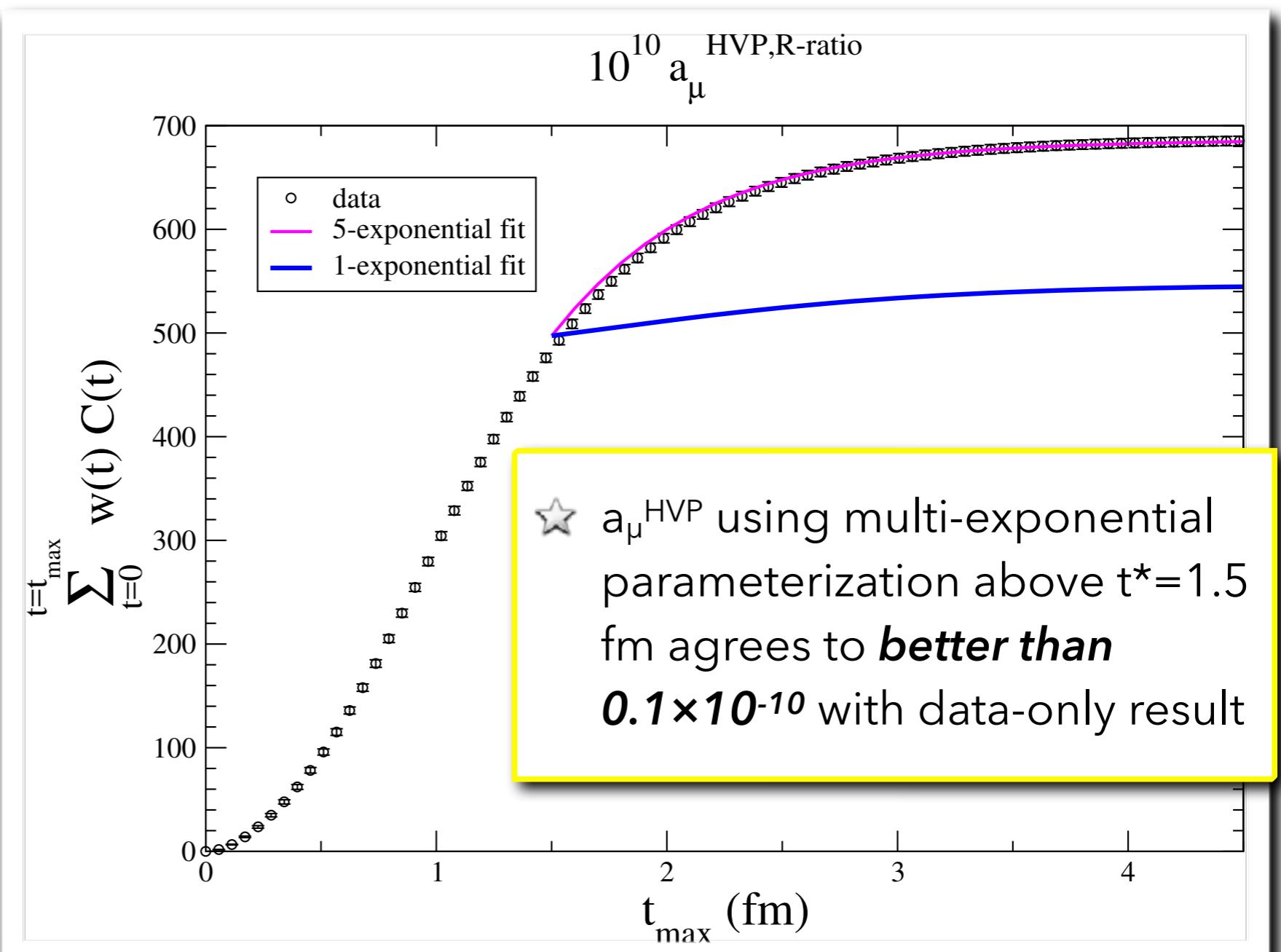
A photograph of a modern, multi-story building with a glass facade, reflected perfectly in the calm water in front of it. The sky is filled with long, wispy clouds colored in shades of blue, orange, and yellow, suggesting either sunrise or sunset. The overall scene is peaceful and architectural.

Tests with R-ratio data

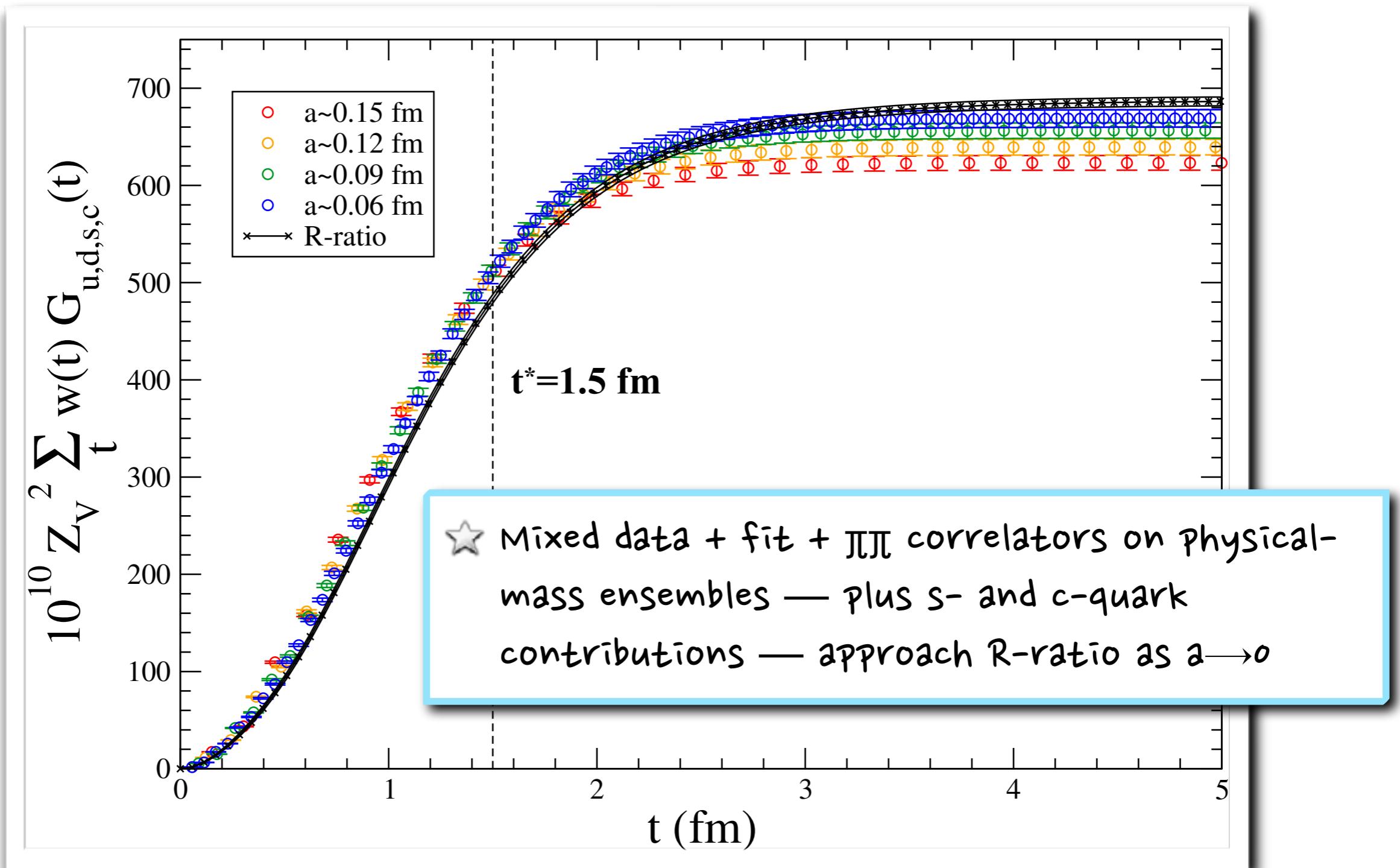
Multi-exponential 2-point fits

- ♦ Fit e^+e^- data to multi-exponential parameterization and compute a_μ using data below 1.5 fm and fit above 1.5 fm
- Method yields correct result provided use of multi-exponential parameterization that accurately describes data above transition time

★ Thanks to F. Jegerlehner for compiling e^+e^- data in public [alphaQED fortran package!](#)

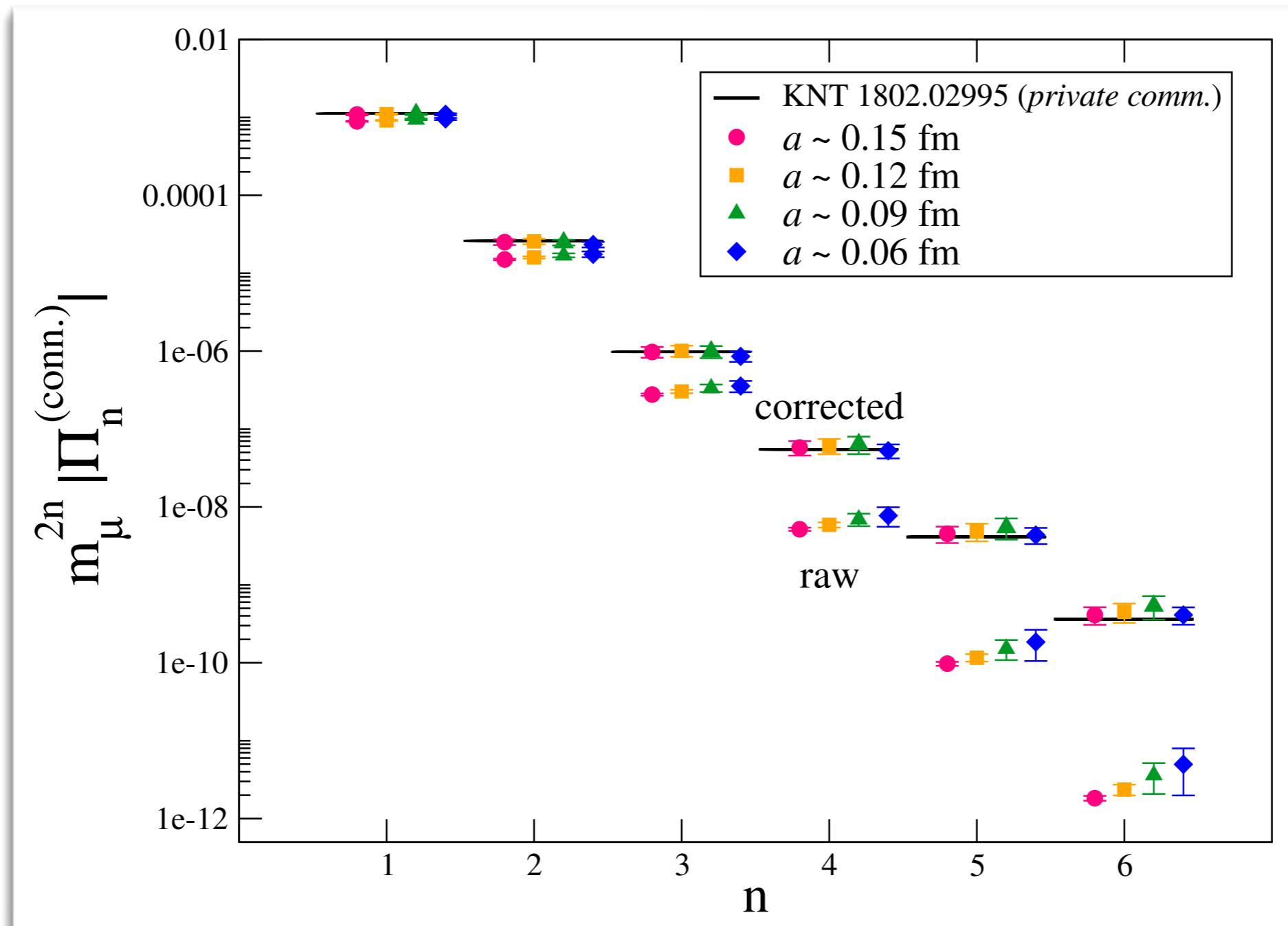


Mixed data + fit correlators



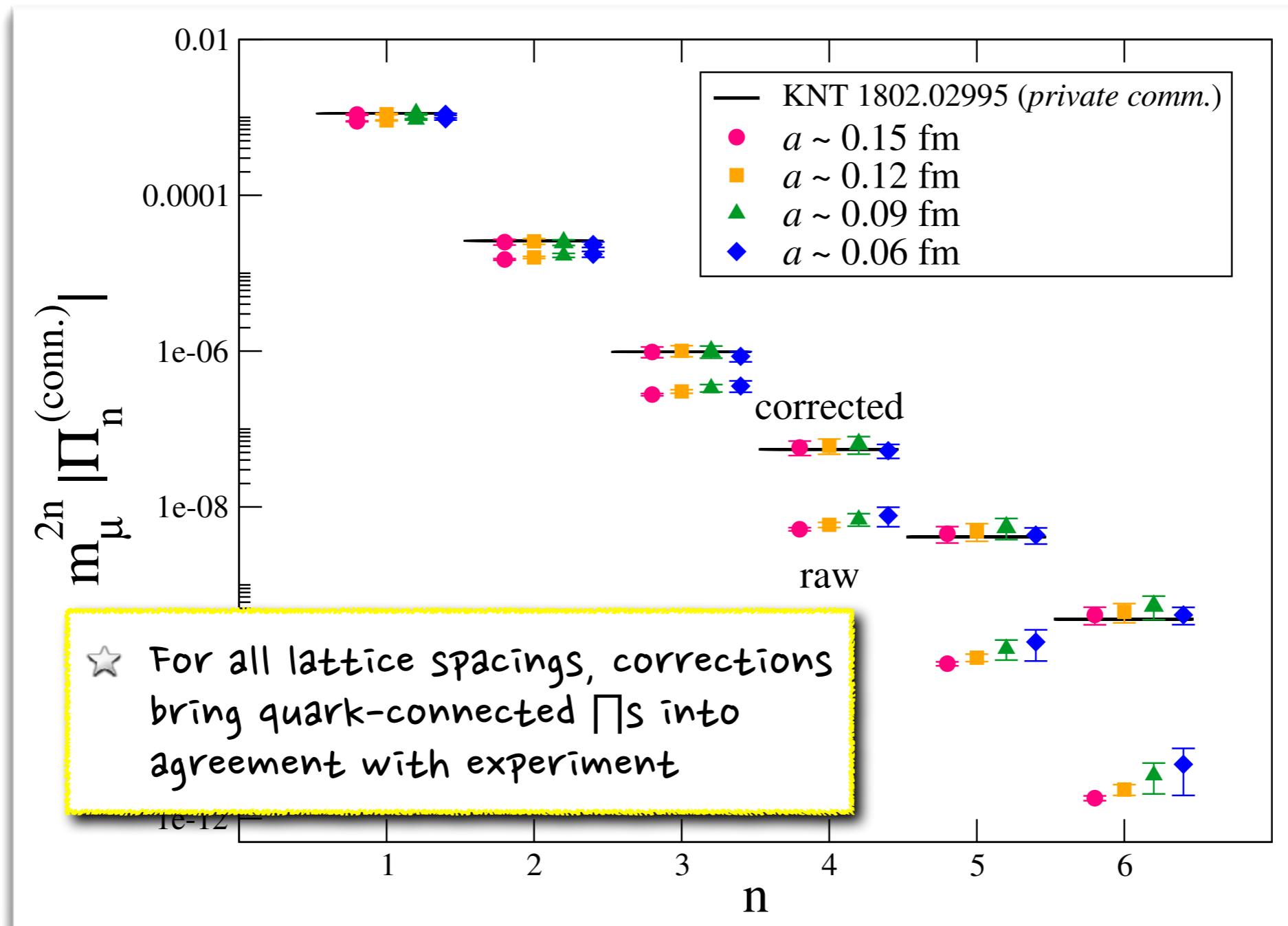
FV + discretization corrections

- ◆ Compare Taylor coefficients on physical-mass ensembles with R-ratio values



FV + discretization corrections

- ◆ Compare Taylor coefficients on physical-mass ensembles with R-ratio values



A photograph of a modern, multi-story building with a glass facade, reflected in a calm body of water. The sky is filled with dramatic, colorful clouds at sunset or sunrise. The text "Supplementary material" is overlaid in the center-left area.

Supplementary material

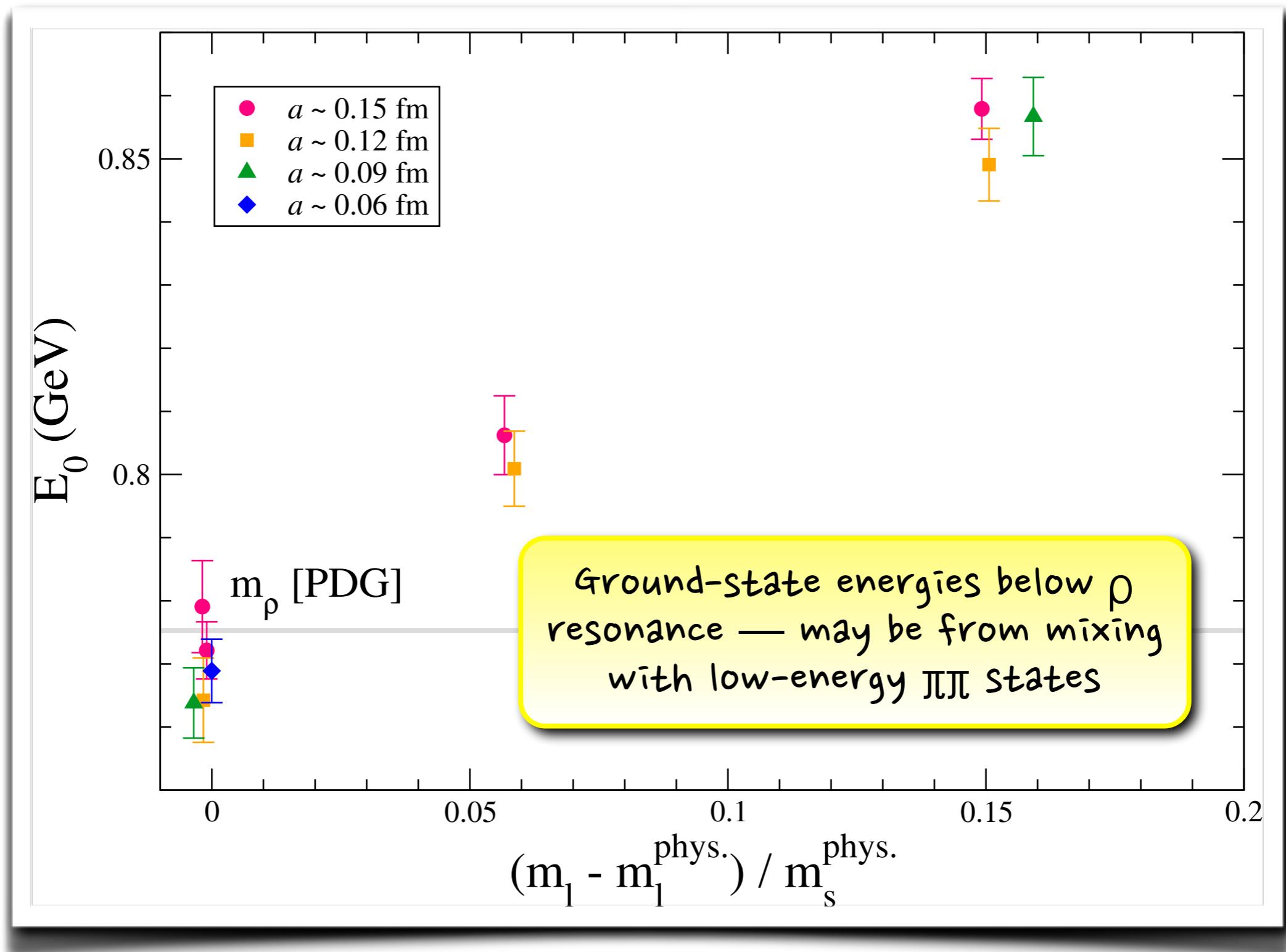
Correlator fits

Simultaneous fit of
four combinations
of (local, smeared)
correlators

$$G_{ij}(t) = a^3 \sum_{k=0}^{N-1} b_i^{(k)} b_j^{(k)} \left(e^{-E^{(k)} t} + e^{-E^{(k)} (T-t)} \right)$$
$$- (-1)^t a^3 \sum_{k=0}^{N-1} d_i^{(k)} d_j^{(k)} \left(e^{-\tilde{E}^{(k)} t} + e^{-\tilde{E}^{(k)} (T-t)} \right)$$

- ♦ Constrain energies & amplitudes with Gaussian priors
- ♦ Employ SVD cuts to reduce d.o.f. and improve reliability of correlation matrix
 - ❖ Conservative approach replaces eigenvalues of correlation matrix below SVD cut by the SVD cut times the maximum eigenvalue, thereby increasing fit error
- ♦ Choose number of states & fit range based on stability of E_0 , A_0 , E_1 , & goodness-of-fit
 - ❖ Fit between $(t_{\min}, T-t_{\min})$ to ensure that fit describes correlator over entire lattice time extent
 - ❖ Employ same $t_{\min} \sim 0.45$ fm for all lattice spacings
 - ❖ Obtain good correlated fits and stable results with $N_{\text{states}} \geq 3$

Ground-state energies



Chiral-continuum fit function

$$a_\mu = a_\mu^{\text{LO}} \times \left(1 + c_\ell \frac{\delta m_\ell}{\Lambda} + c_s \frac{\delta m_s}{\Lambda} + \tilde{c}_\ell \frac{\delta m_\ell}{m_\ell} + c_{a^2} \frac{(a\Lambda)^2}{\pi^2} \right)$$

$$\delta m_f \equiv m_f - m_f^{\text{phys.}}, \quad \Lambda = 500 \text{ MeV}$$

Constraints

$$a_\mu^{\text{LO}} = 600(200) \times 10^{-10}$$

$$c_\ell, c_{a^2} = 0(1), \quad c_s = 0.0(0.3), \quad \tilde{c}_\ell = 0.00(0.03)$$

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$$\delta m_f \equiv m_f - m_f^{\text{phys.}}, \quad \Lambda = 500 \text{ MeV}$$

- ◆ Correct for quark-mass mistuning
- ◆ c_s small because only enters through sea, and m_s well tuned

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- ◆ Correct for quark-mass mistuning
- ◆ c_s small because only enters through sea, and m_s well tuned

- ◆ Correct for small residual quark-mass dependence after $\pi^+\pi^-$ correction

Constraints

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- ◆ Account for residual generic and taste-breaking discretization errors

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$\delta m_f \equiv m_f - m_f^{\text{phys.}}, \quad \Lambda = 500 \text{ MeV}$

- ◆ Correct for quark-mass mistuning
- ◆ c_s small because one enters through sea, m_s well tuned

◆ Correct for small
  No m_l dependent terms with only physical-mass analysis correction

◆ Account for residual generic and taste-breaking discretization errors

Constraints

$$a_\mu^{\text{LO}} = 600(200) \times 10^{-10}$$

$$c_\ell, c_{a^2} = 0(1), \quad c_s = 0.0(0.3), \quad \tilde{c}_\ell = 0.00(0.03)$$

Isospin-breaking & QED effects

- ◆ Data-driven phenomenological analyses & model calculations suggest that contributions to connected u/d contribution to leading-order a_μ^{HVP} are at or below the ~1% level
 - ▶ Dominant EM effect from $\pi^0\gamma$ vacuum polarization bubbles estimated from $e^+e^- \rightarrow \pi^0\gamma$ data to be $\Delta(a_\mu^{\text{HVP}}) = 4.6(2) \times 10^{-10}$ [Hagiwara et al., PRD69, 093003 (2004)]
 - ▶ Dominant isospin-breaking contribution from $\rho\omega$ mixing estimated from $e^+e^- \rightarrow \pi^+\pi^-$ data to be $\Delta(a_\mu^{\text{HVP}}) \sim 2-5 \times 10^{-10}$ [Wolfe & Maltman, PRD83 (2011) 077301]

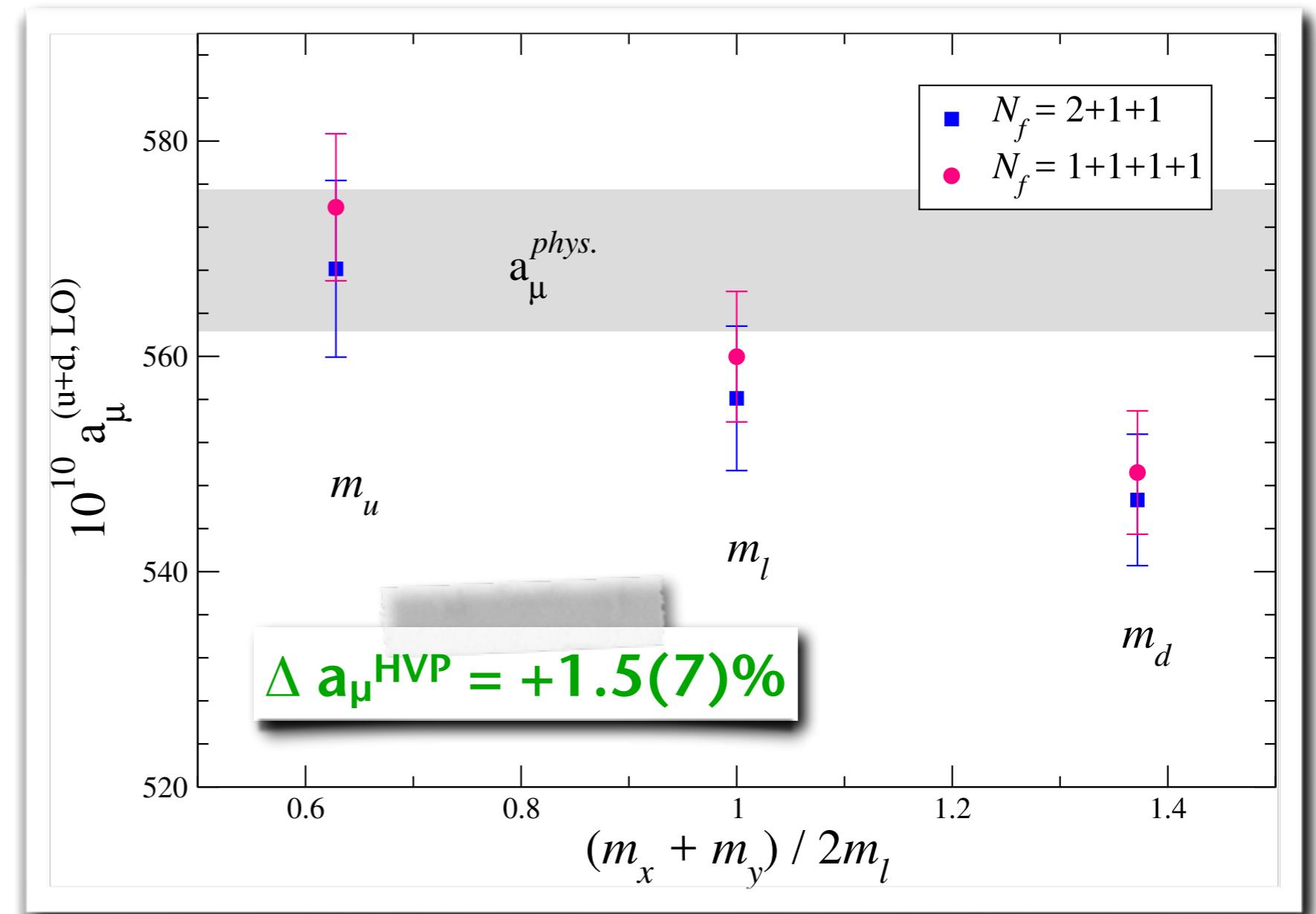
► Consistent with Jegerlehner compilation of missing effects in lattice QCD simulations performed in the isospin limit and without QED effects [*private communication, to appear in “Springer Tracts in Modern Physics”*]

Missing effect	$10^{10} \delta a_\mu^{\text{HVP,LO}}$
$\rho\omega$ mixing	+2.71
$\rho\gamma$ mixing	-2.74
Final-state radiation	+4.22
EM effects on $M_\pi, M_\rho, \Gamma_\rho$	+1.12
$\pi^0\gamma$ channel	+4.64
$\eta\gamma$ channel	+0.16
Total	+10.11

Strong isospin-breaking correction

- ◆ Directly calculate correction to a_μ^{HVP} from $m_u \neq m_d$ on $a \sim 0.15 \text{ fm}$ $N_f = 1+1+1+1$ (and $N_f = 2+1+1$) physical-mass ensemble [PRL120 (2018) no.15, 152001]

- ❖ Fix ratio $m_u/m_d = 0.4582$ to physical value from MILC quenched-QED analysis [1606.01228]
- ❖ Compute a_μ with valence-quark masses m_u , $m_l = (m_u + m_d)/2$, & m_d
- ❖ c.f. subsequent RBC/UKQCD preprint $a_\mu^{\text{SIB}}/a_\mu^{\text{ud,conn.}} \sim 1.6\%$ [1801.07224]



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