# The heavy tails of axion-photon conversion

## M.C. David Marsh **Stockholm University 11th August, 2022**

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### Collaborators





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[cajohare.github.io/axionlimits/]



[Wouters, Brun], [Conlon et al.], [Berg et al.], [Reynolds et al.], [Reynes et al.], [Chen, Conlon], [Day, Krippendorf], [Matthews et al.], [Schallmoser et al.]



### The photon disappearance channel





#### Final photon spectrum



### Galaxy clusters as axion-photon converters

- Largest gravitational bound objects (~100s kpc). Magnetised ( $\mu$ G). Long coherence lengths (~kpc). Luminous sources (AGNs, quasars).
- Unsuppressed *conversion ratios*:

$$P_{\gamma a} \sim \mathcal{O}\left(\frac{1}{2}\right)$$



 $g_{a\gamma}$  $) \times \left( \frac{3\omega}{10^{-11} \,\mathrm{GeV}} \right)$ 





#### [Reynolds, *DM*, et al.] [Sisk-Reynes et al.]

### **Precision spectra**



### Magnetic field models v1: "cell models"



### Status: standard practice



### Magnetic field models v2: Gaussian random fields



100  $x \; [kpc]$ 

Status: "state-of-the-art"



### How robust are these limits?



[Matthews et al.]



### **Perturbative formalism**

Small amplitude oscillations motivate perturbative solutions.

Simplest case:  $m_a > \omega_{\rm pl}$ 

 $P_{\gamma a}(\eta_a) =$ 

 $\eta_a = \frac{m_a^2}{2\omega}$ 

$$= \frac{g_{a\gamma}^2}{4} |\tilde{B}_i(\eta_a)|^2$$
$$\tilde{B}_i(\eta_a) = \int_{-L/2}^{L/2} dz \, B_i(z\hat{z}) e^{-iz\eta_a}$$

[Raffelt, Stodolsky]



### **Perturbative formalism: applicability**



#### [Matthews et al.]

### **Structure and phases**



[Maron, Goldreich]

Is ALP-photon conversion independent of MHD structure?

 $P_{\gamma a}(\eta_a) = \frac{g_{a\gamma}^2}{4} |\tilde{B}_i(\eta_a)|^2$ 



### **Dedicated MHD simulations: time-evolution**



 $L^3 = (200 \text{ kpc})^3$ #lattice points =  $512^3$ periodic bc, external forcing Dynamo-enhanced, turbulent magnetic field

[Carenza et al.]



## **MHD** magnetic field: cross-sections



100  $x \; [kpc]$ 

**Red:**  $|{\bf B}| > 3B_{\rm rms}$ 

### Want: statistical properties from ensemble of trajectories



## **GRF** magnetic field: cross-sections





100

### **Analytic GRF predictions**

$$\langle \widehat{B}_{a}(\mathbf{k})\widehat{B}_{b}^{*}(\mathbf{k}')\rangle = \delta^{3}(\mathbf{k} - \mathbf{k}') \left[\frac{P_{3D}(k)}{2} \left(\delta_{ab} - \frac{k_{a}k_{b}}{k^{2}}\right) - i\epsilon_{abc}\frac{k_{c}}{k}H(k)\right],$$

$$\int dk + k + \left(\sqrt{-1 - k^{2}}\right) \left(1 - \frac{k^{2}}{k}\right)$$

$$P_{1D}(\eta_a) = \int \frac{dk_{\perp}k_{\perp}}{2(2\pi)^3} P_{3D}\left(\sqrt{\eta_a^2 + k_{\perp}^2}\right) \left(1 - \frac{1}{2}\frac{k_{\perp}^2}{\eta_a^2 + k_{\perp}^2}\right)$$

#### Probability distribution of conversion ratios:

$$f_{P_{a\gamma}(\eta_a)}(p) = \frac{e^{-p/p_0}}{p_0}$$

$$p_{0} = \frac{g_{a\gamma}^{2}}{4} \frac{L}{2\pi} P_{1\mathrm{D}}(\eta_{a})$$

## **Heavy-tailed MHD distributions**



### Holds for arbitrary masses, polarisations



### **Non-Gaussianity**

Expectation value of conversion ratio depends only on power spectrum of B:

$$\langle P_{\gamma a}(\eta_a) \rangle = \frac{g_{a\gamma}^2}{4} \langle |\tilde{B}_i(\eta_a)|^2 \rangle = \frac{g_{a\gamma}^2}{4} \frac{L}{2\pi} P_{1\mathrm{D}}(\eta_a)$$

Heavy tails come from larger-than-Gaussian higher-order correlations functions, i.e.

$$\langle P_{\gamma a}(\eta_a)^2 \rangle, \quad \langle P_{\gamma a}(\eta_a)^3 \rangle, \quad \langle P_{\gamma a}(\eta_a)^4 \rangle \quad \text{etc.}$$

Structure does not directly affect conversion ratios, but non-Gaussianity does.

Same for MHD and GRF

Different for MHD and GRF



## **Non-Gaussianity**

### Two possible sources:







### MHD and GRF predictions agree for typical fluctuations.

MHD predicts heavy tails: larger probability of large conversion ratios.

Use full cluster MHD simulations with existing data.

Getting ready for the next-generation of X-ray satellites, e.g. Athena.

### Conclusions

- Suggests existing limits conservative.

- **Future directions**

