

Cross-Topic

Nature of Dark Matter

Theory overview

Geraldine Servant, DESY-HH

ALPS

Axel Lindner, DESY-HH

Direct DM search with Ge

Klaus Eitel, KIT

Indirect DM search with
gamma rays and neutrinos

Johannes Knapp, DESY-Zeuthen

CASPER and GNOME

Arne Wickenbrock and
Dmitry Budker, HI Mainz

MU Tag 2016 HIM

Cross-Topic

Nature of Dark Matter

Studying generalized dark matter interactions with extended halo-independent methods F. Kahlhoefer, DESY-HH

Axion dark matter from topological defects K. Saikawa, DESY-HH

FUNK dark photon search: status and perspectives D. Verberic, KIT

Dark photon searches at MAMI and MESA A. Denig, JGU & HIMainz

Dark matter models with two mediators M. Duerr, DESY-HH

Phenomenology of flavoured dark matter S. Kast, KIT

Studying generalised dark matter interactions with extended halo-independent methods.

Felix Kahlhoefer

MUTAG2016

Helmholtz Institute Mainz

12-13 December 2016

Based on

arXiv:1607.04418 with Sebastian Wild

Motivation

- > Differential event rates in direct detection experiments are given by

$$\frac{dR}{dE_R} = \frac{\rho}{m_T m_\chi} \int_{v_{\min}}^{\infty} v f(\mathbf{v} + \mathbf{v}_E(t)) \frac{d\sigma}{dE_R} d^3v \quad \text{with} \quad v_{\min}(E_R) = \sqrt{\frac{m_T E_R}{2 \mu_{T\chi}^2}}$$

Dark matter (DM) mass and scattering cross section:
This is what we (particle physicists) would like to find out

Local DM density and velocity distribution:
These are uncertainties that enter experimental predictions in a complicated way

- > How reliably can we get to the particle physics properties of DM from future direct detection signals in the face of astrophysical uncertainties?

There are many models of DM

- There are many plausible models for the interactions of DM beyond standard spin-independent (SI) and spin-dependent (SD) scattering:

- DM with an Anapole moment (AM):

$$\mathcal{A} \bar{\chi} \gamma^\mu \gamma^5 \chi \partial^\nu F_{\mu\nu}$$

This is the only possible interaction of Majorana fermions with photons.

- DM with a magnetic dipole moment (MDM):

$$\mathcal{D}_{\text{magn}} \bar{\chi} \sigma^{\mu\nu} \chi F_{\mu\nu}$$

This leads to long-range interactions due to the massless photon exchange.

- DM with a dark magnetic dipole moment (DMDM):

$$\mathcal{D}_{\text{magn}} \bar{\chi} \sigma^{\mu\nu} \chi F'_{\mu\nu}$$

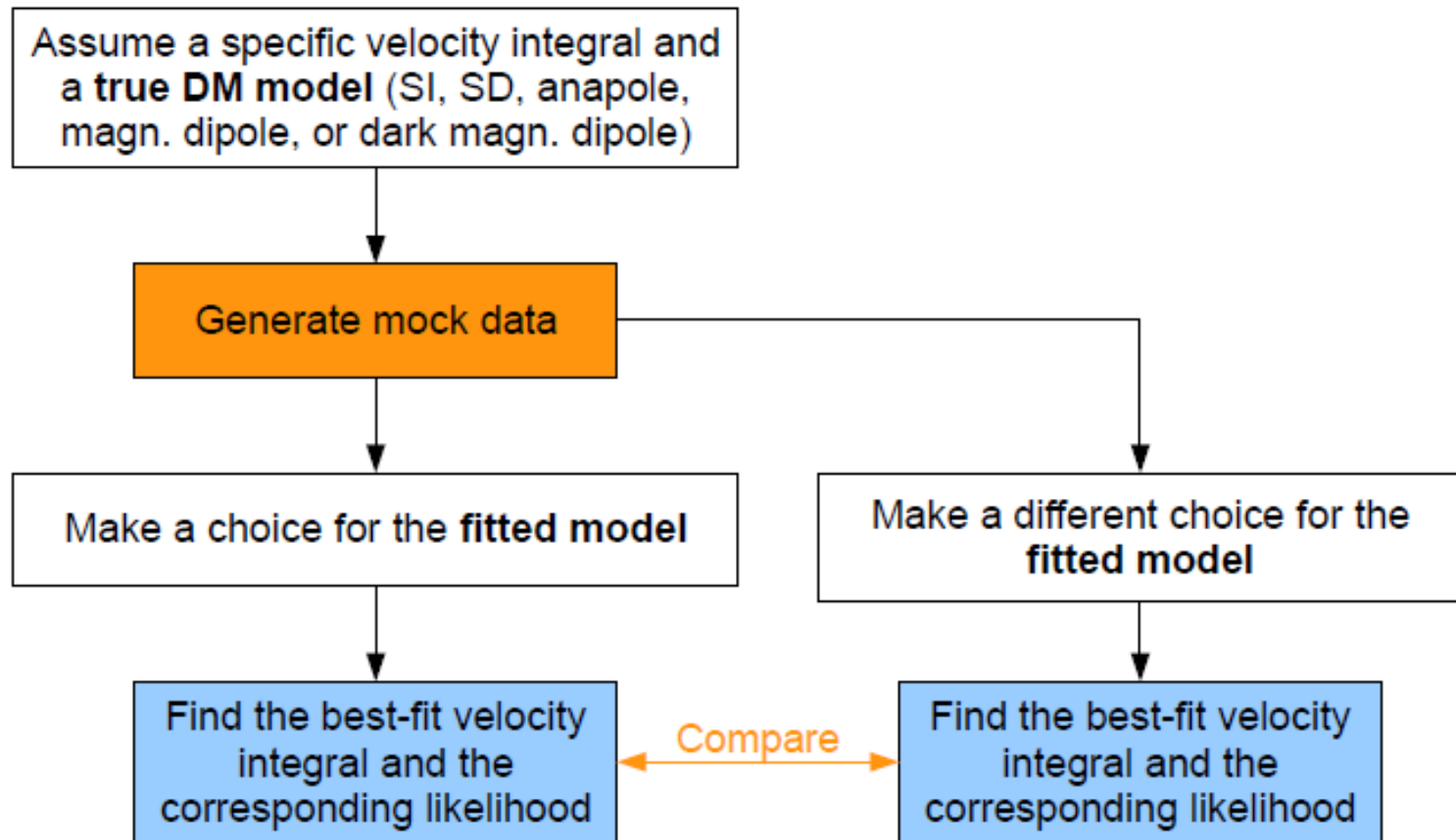
Same as the the MDM but for a massive hidden photon ($1/q^2 \rightarrow 1/M^2$).



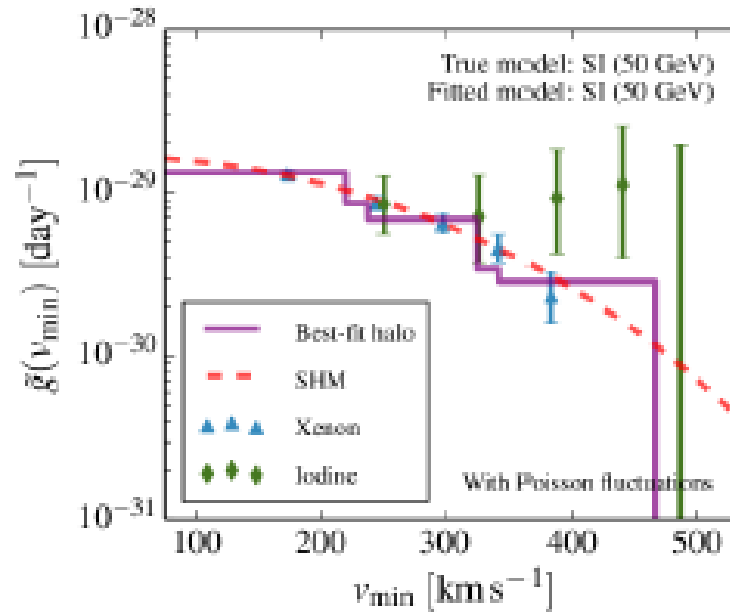
One can then construct a matrix D_{ij} such that the expected number of events in bin i (R_i) is given by a simple matrix multiplication:

$$R_i = \sum_j D_{ij} g_j$$

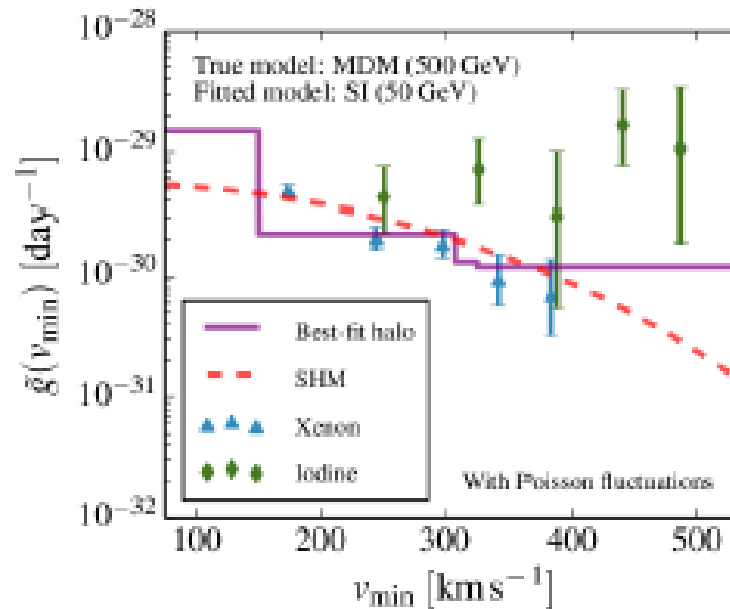
Best-fit velocity integrals



DM toy model Xenon & Iodine

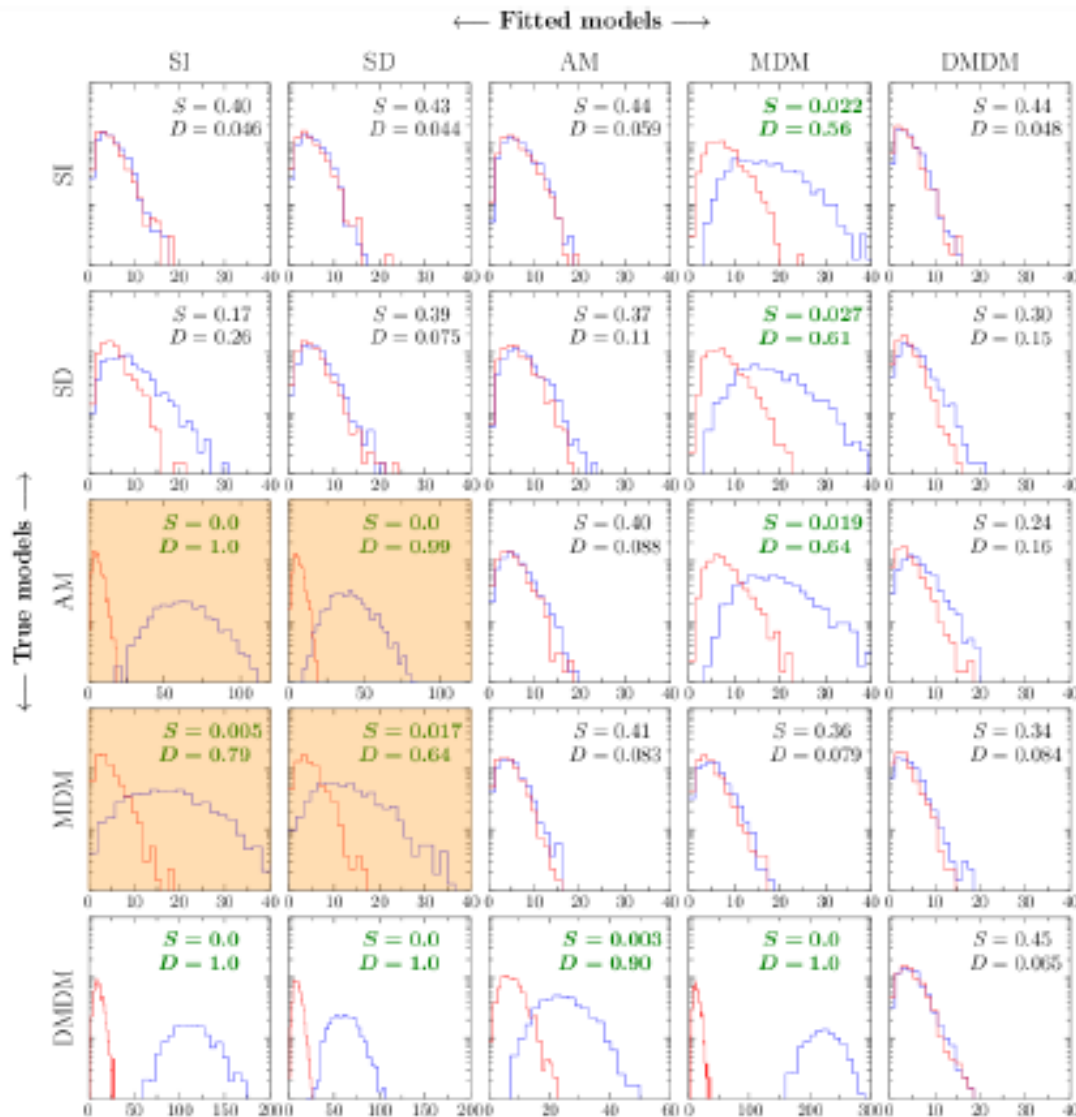


true model
=
fitted model



true model
 \neq
fitted model

Results for xenon+iodine experiments



S : p-value of a typical realisation

D : fraction of realisations that are excluded at 95% CL

- The presence of a second experiment with an iodine target makes it possible to rule out SI or SD interactions if DM has an anapole moment or a magnetic dipole moment.
- The reason is that iodine has a much higher sensitivity to anapole and dipole interactions.

arXiv:1607.04418



Axion dark matter from topological defects

Ken'ichi Saikawa (DESY)



Based on

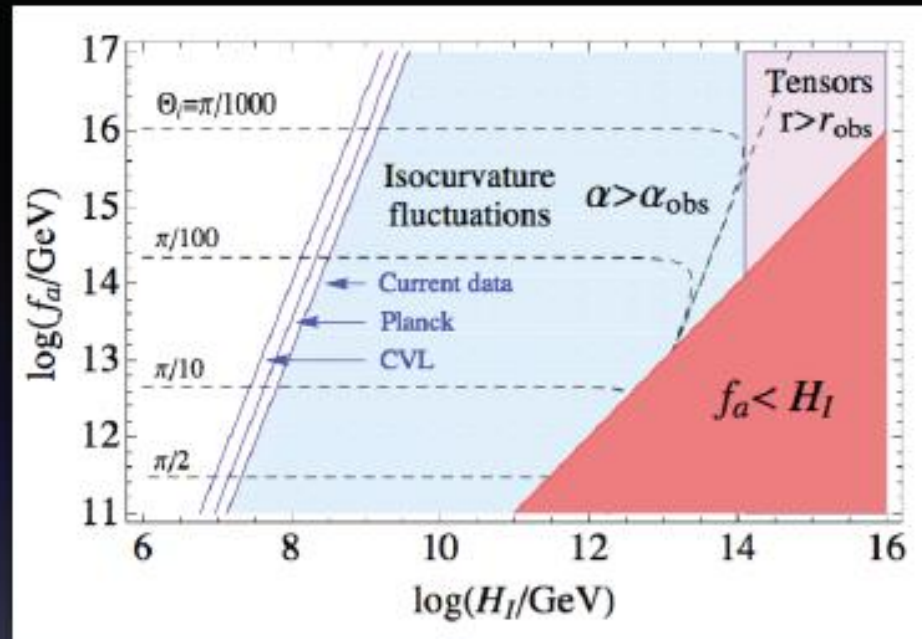
T. Hiramatsu, M. Kawasaki, KS, T. Sekiguchi, PRD85, 105020 (2012) [1202.5851]

T. Hiramatsu, M. Kawasaki, KS, T. Sekiguchi, JCAP01, 001 (2013) [1207.3166]

M. Kawasaki, KS, T. Sekiguchi, PRD91, 065014 (2015) [1412.0789]

A. Ringwald, KS, PRD93, 085031 (2016) [1512.06436]

Axions in the inflationary universe: two scenarios

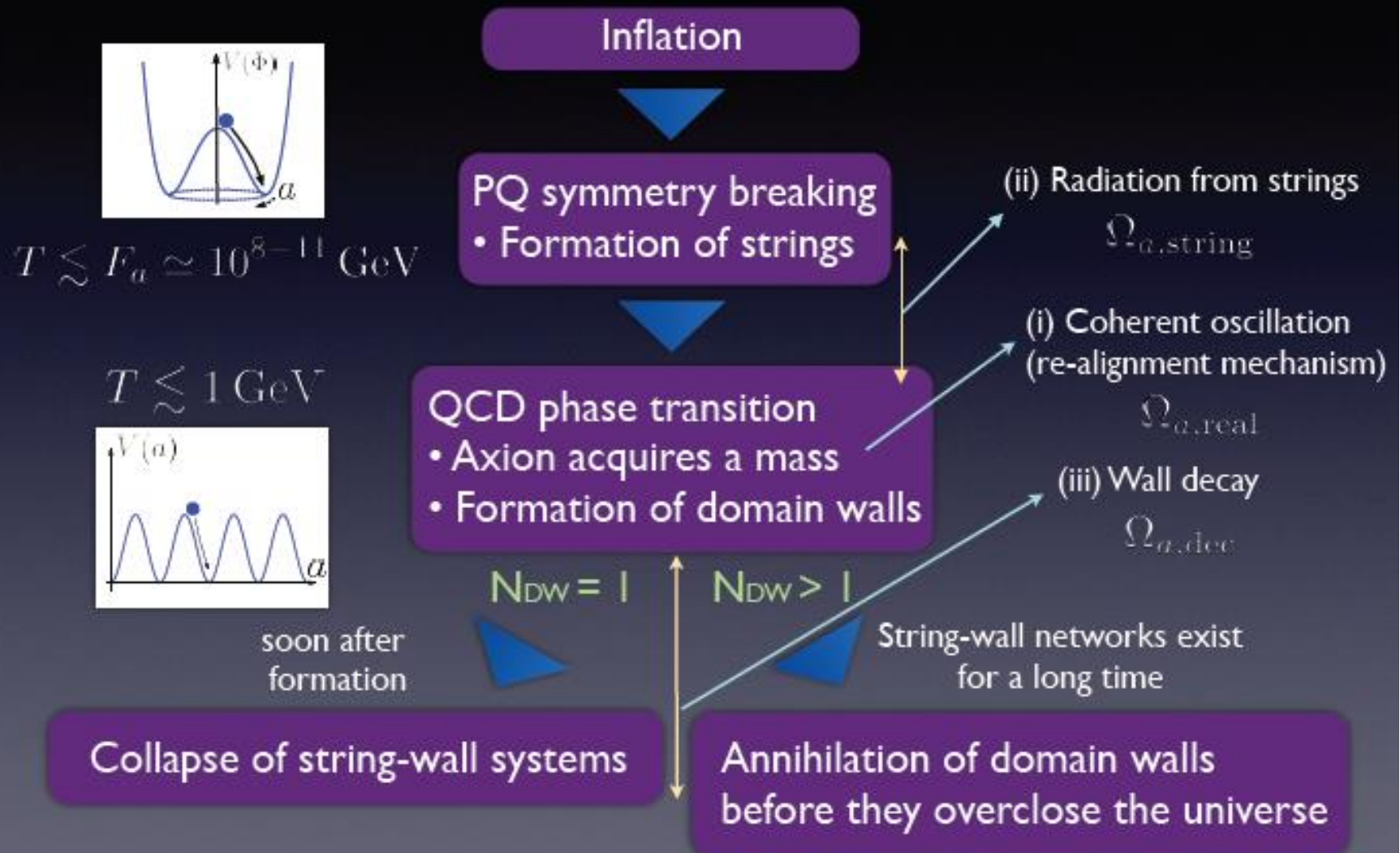


Hamann, Hannestad, Raffelt and Wong (2009)

- Pre-inflationary PQ symmetry breaking
 - Severe isocurvature constraints
 - Tuning of the initial field value (“anthropic window”)
- Post-inflationary PQ symmetry breaking ← this talk
 - Formation of topological defects

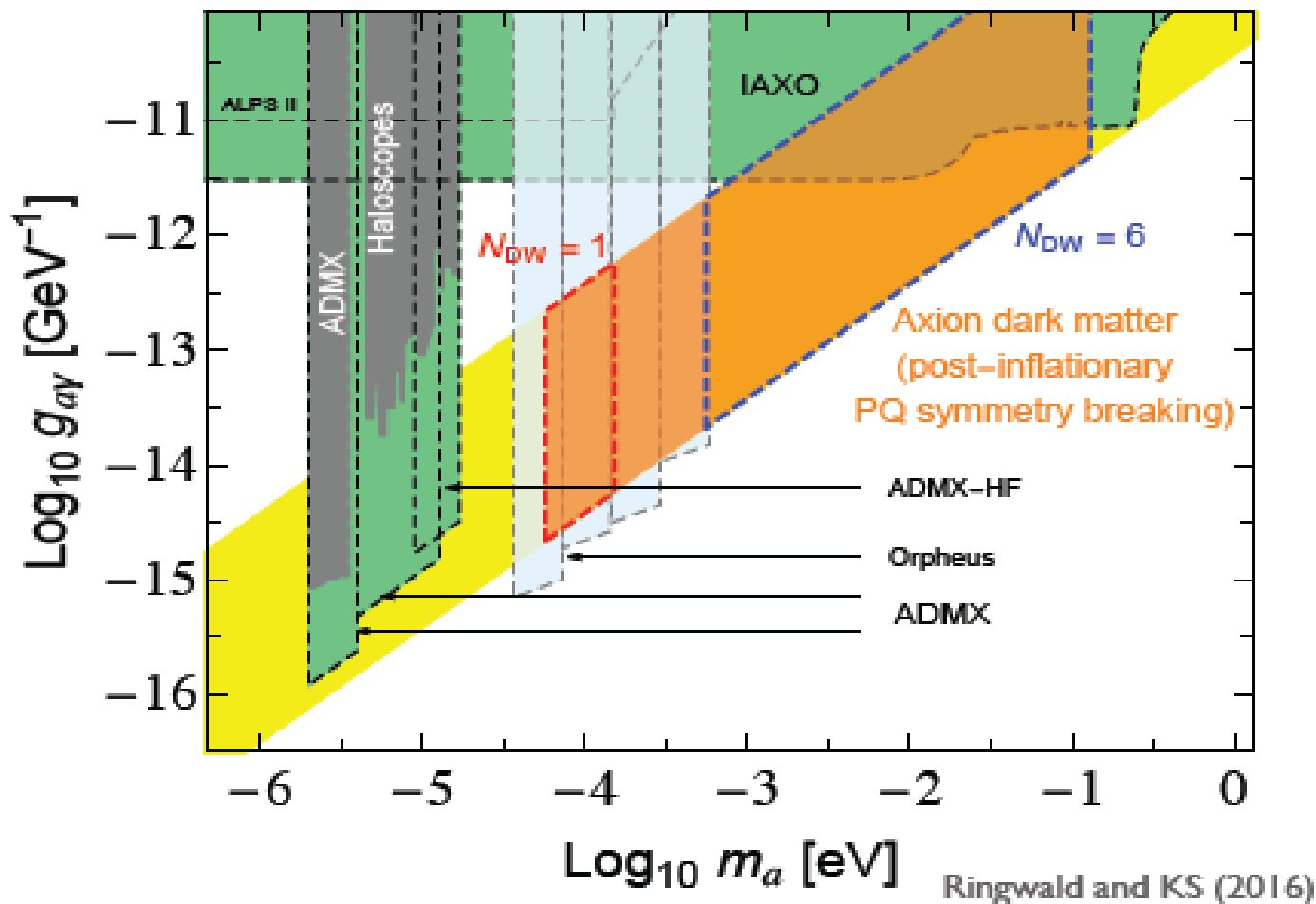
Production of axions in the early universe

(post-inflationary PQ symmetry breaking scenario)



Search for axion dark matter

Search space in photon coupling $g_{a\gamma} \sim \alpha/(2\pi F_a)$ vs. mass m_a





Search for Hidden-Photon Dark Matter

Status & Perspectives

K. Daumiller, B. Döbrich, R. Engel, L. Gastaldo, J. Jaeckel, M. Kowalski,
A. Lindner, H.-J. Mathes, A.-S. Müller, J. Redondo, M. Roth, C. Schäfer,
T. Schwetz-Mangold, R. Ulrich, D. Veberic



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Hidden Photons

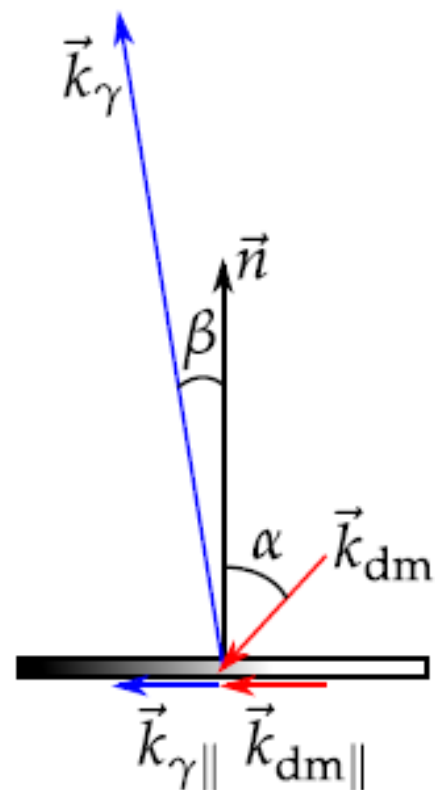
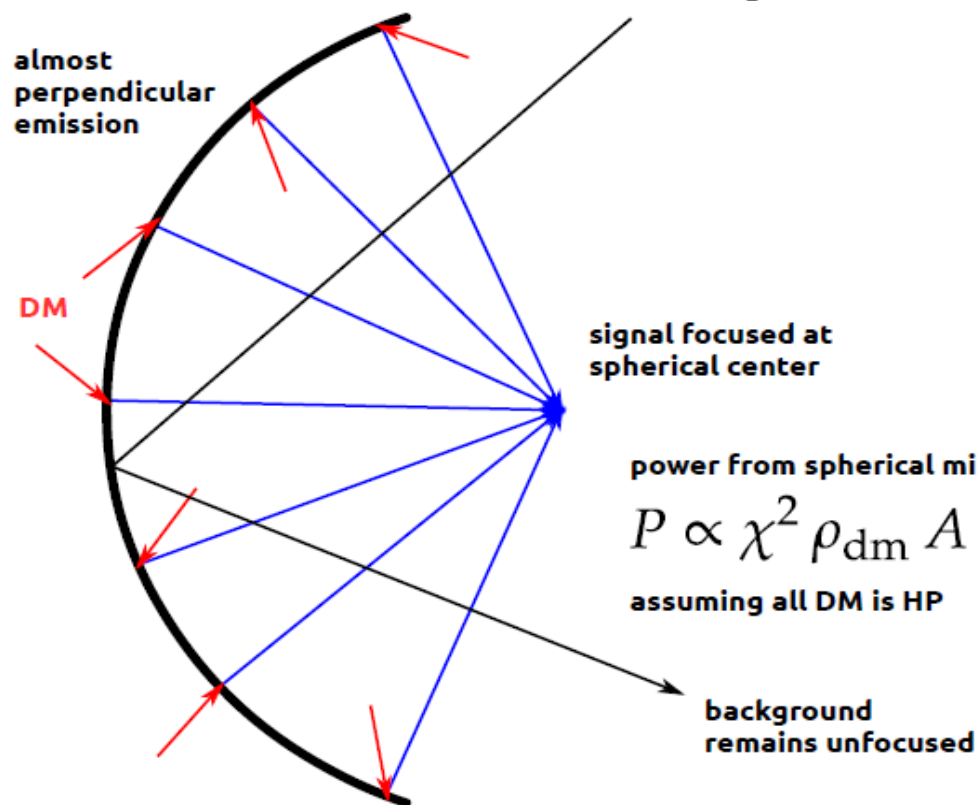
$$\mathcal{L} = -\frac{1}{4}(F_{\mu\nu}F^{\mu\nu} + \tilde{X}_{\mu\nu}\tilde{X}^{\mu\nu}) - \frac{\chi}{2}F_{\mu\nu}\tilde{X}^{\mu\nu} + \frac{m^2}{2}\tilde{X}_\mu\tilde{X}^\mu + J^\mu A_\mu$$

light extra U(1)
gauge boson:
hidden photon

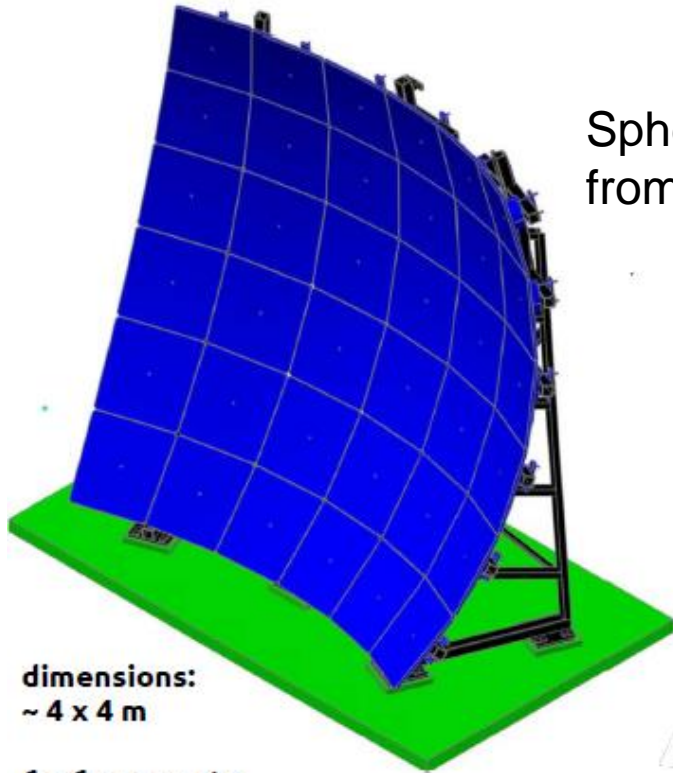
kinetic mixing
coupling

mass term

Horns et al., JCAP 04 (2013) 016



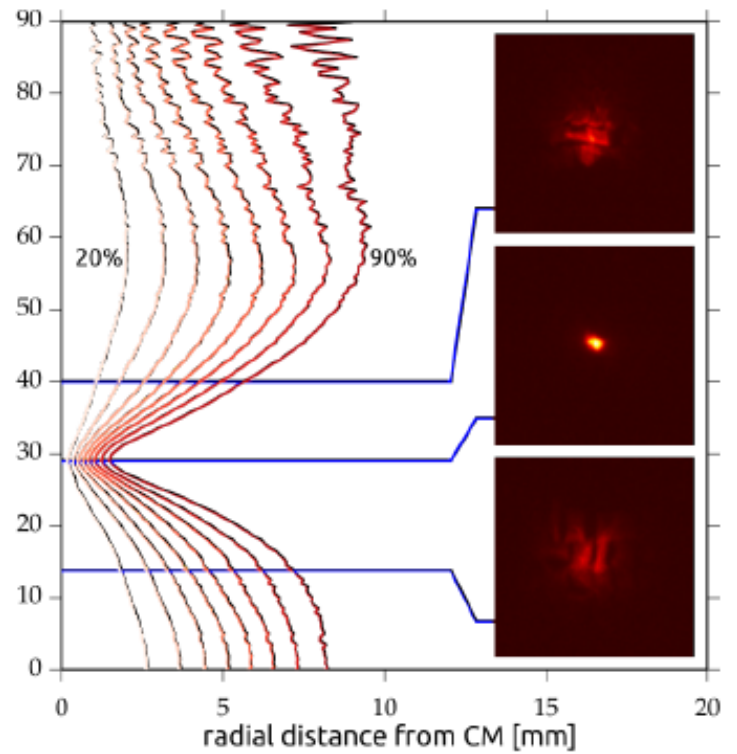
Spherical mirror from PAO



dimensions:
~ 4 x 4 m

6 x 6 segments

only 3 unique shapes



spot size ~2 mm

PCO Sensicam VGA
640x480 pixel CCD
internally Peltier cooled to -15°C

1. test: CCD camera



curent: PMT 160-630 nm

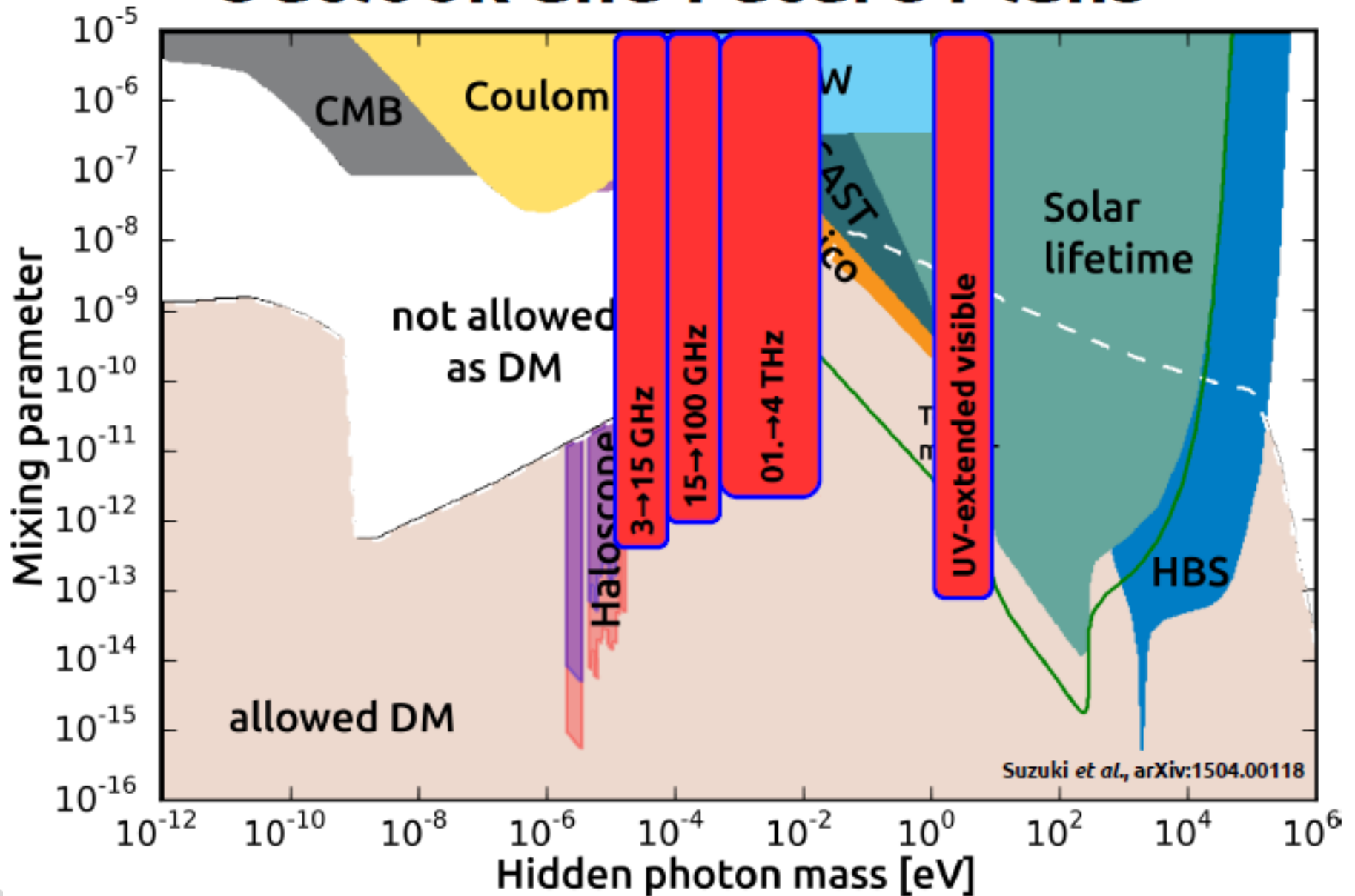
ET Enterprises
electron tubes



30 nm

1.13" window

Outlook and Future Plans



MU Programmtag 2016

12-13 December 2016
Helmholtz Institute Mainz

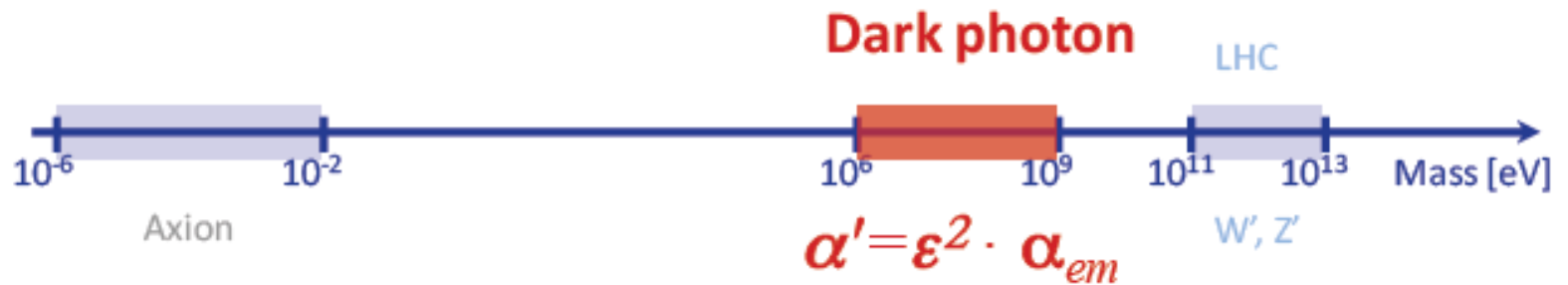


Dark Photon Searches at MAMI and MESA

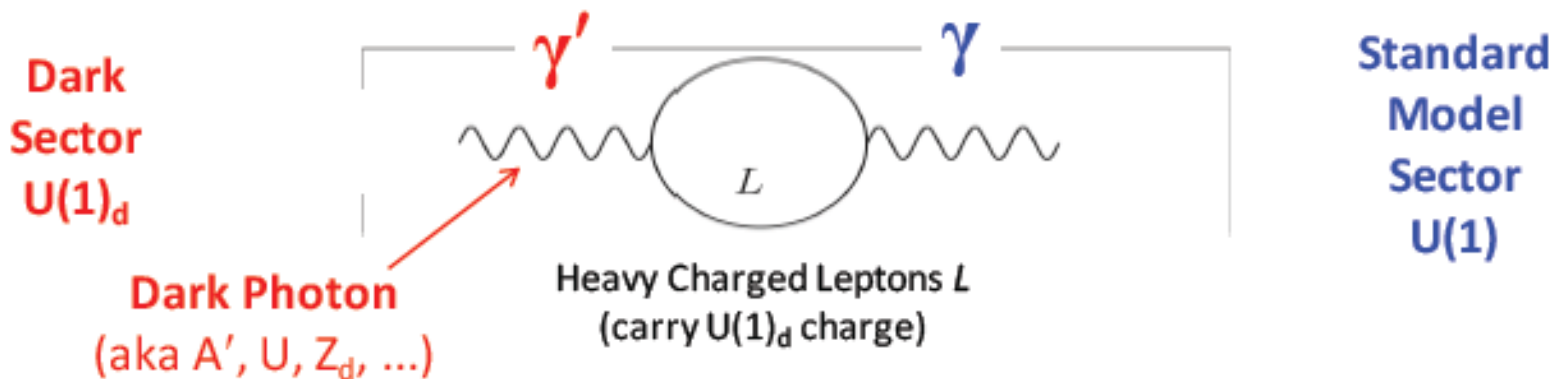
Achim Denig, Mainz



New massive force carrier of extra $U(1)_d$ gauge group;
 predicted in almost all string compactifications



A way to relate the dark sector to the SM (coupling $\sim \epsilon^2$)

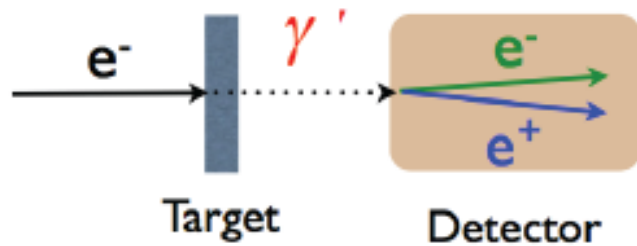


Results from A1 Pilot Run (2011)

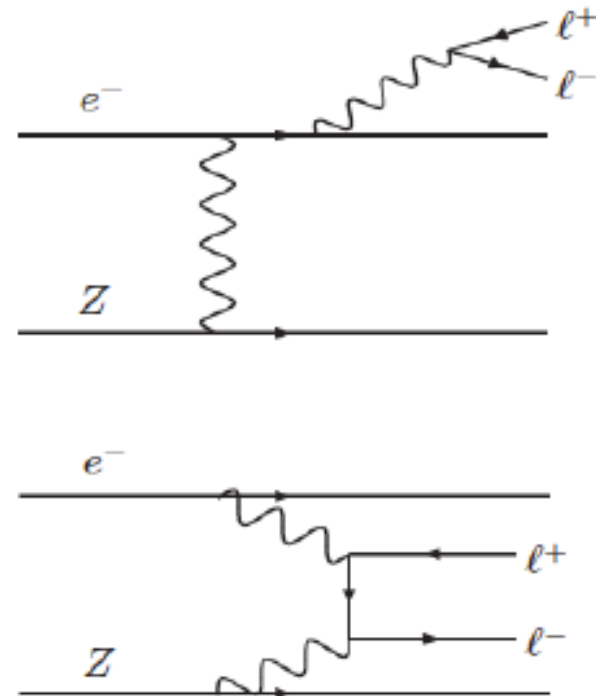
Low-Energy Electron Accelerator with high intensity suited for DP search

3 days @ MAMI

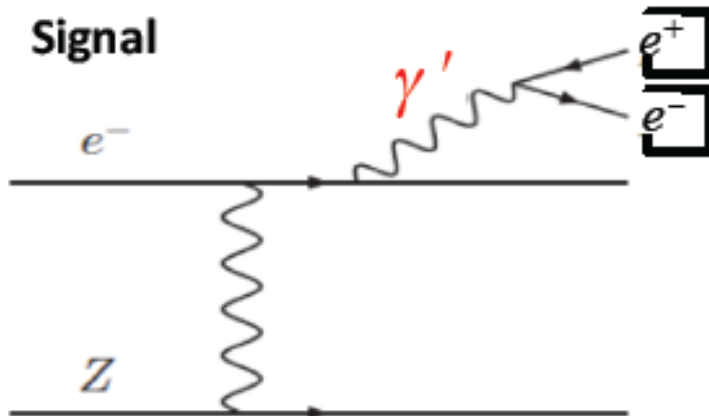
Bjorken, Essig, Schuster, Toro (2009)



Background



Signal

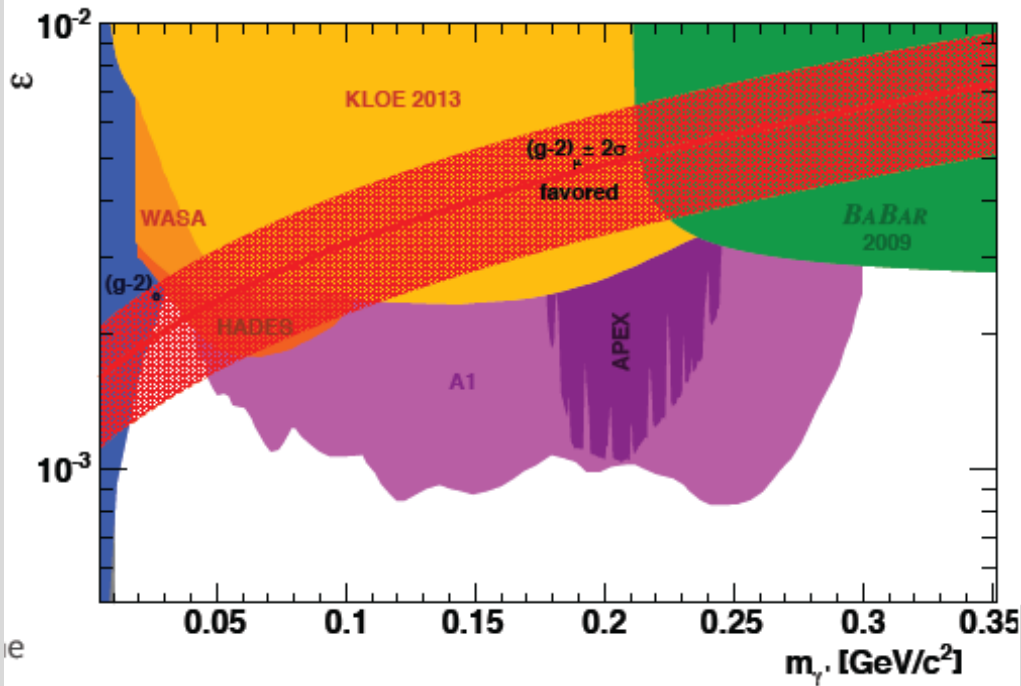
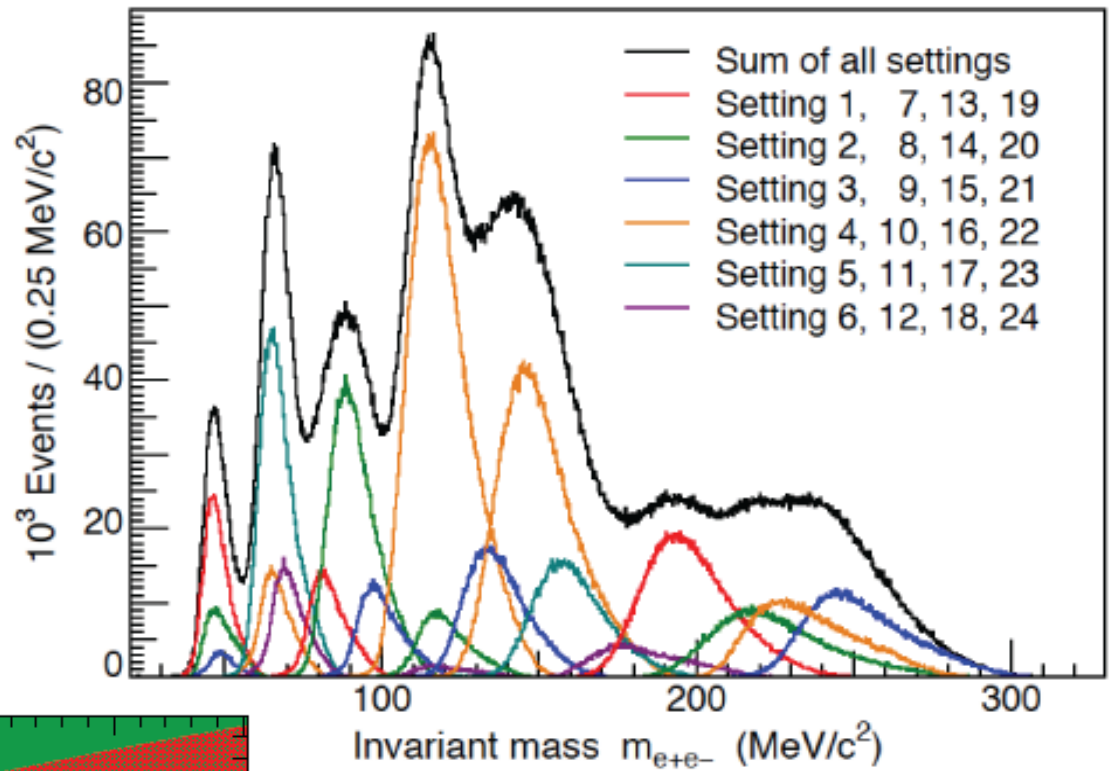


Bump search!

Merkel et al. [A1]
 PRL '11
 PRL '14

4 weeks @ MAMI

- E_{beam} 180 - 855 MeV
- 100 μA beam current
- Stack of Ta targets
- 22 kinematic settings
- O(1 month) of beam time



at time of publication most stringent limit ruling out major part of the parameter range motivated by $(g-2)_\mu$

New Tool for Low-Precision Physics: MESA

Mainz Energy-Recovering Superconducting Accelerator

$E_{\max} = 155 \text{ MeV}$

$I_{\max} > 1 \text{ mA (ERL)}$

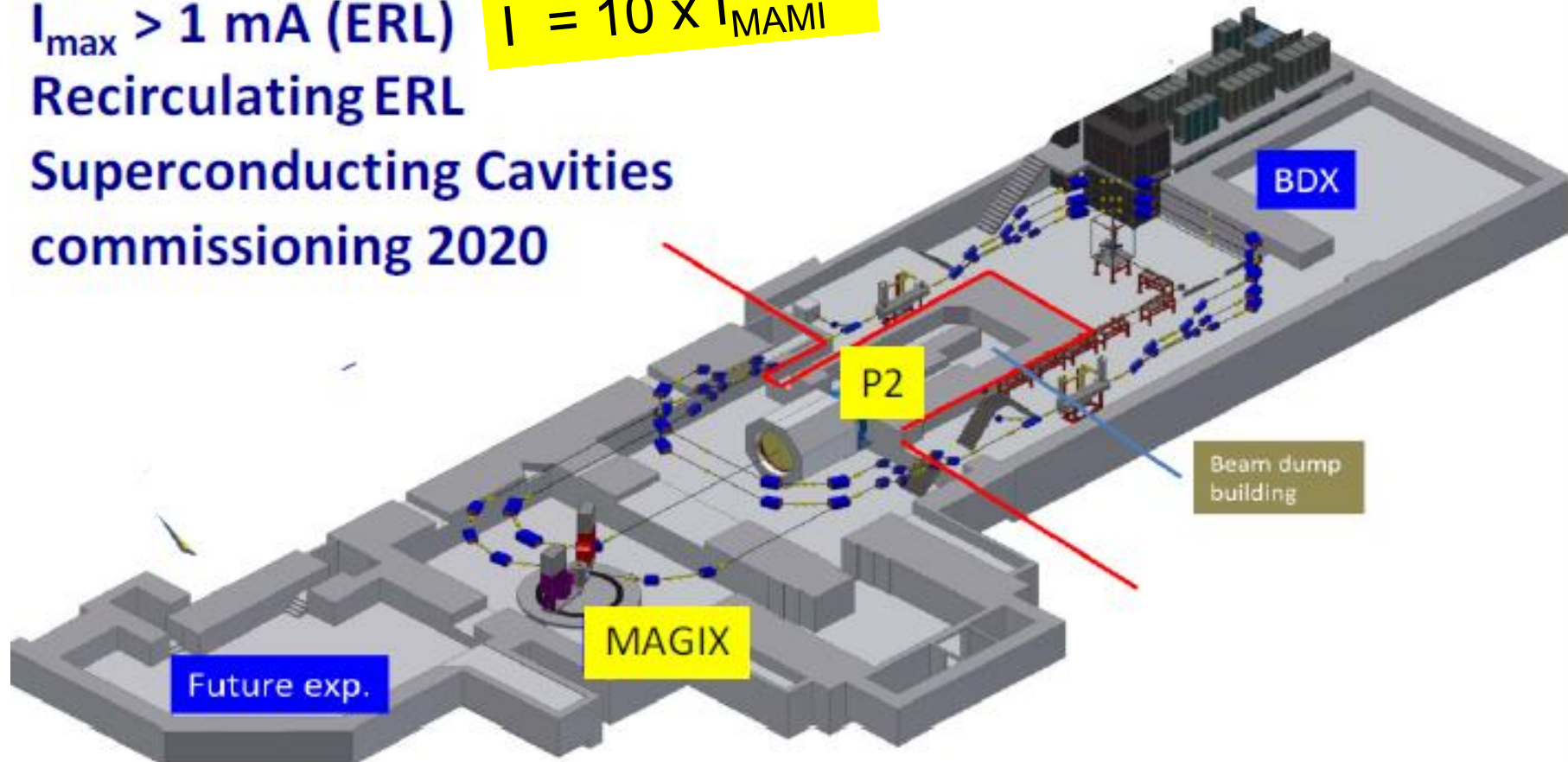
Recirculating ERL

Superconducting Cavities

commissioning 2020

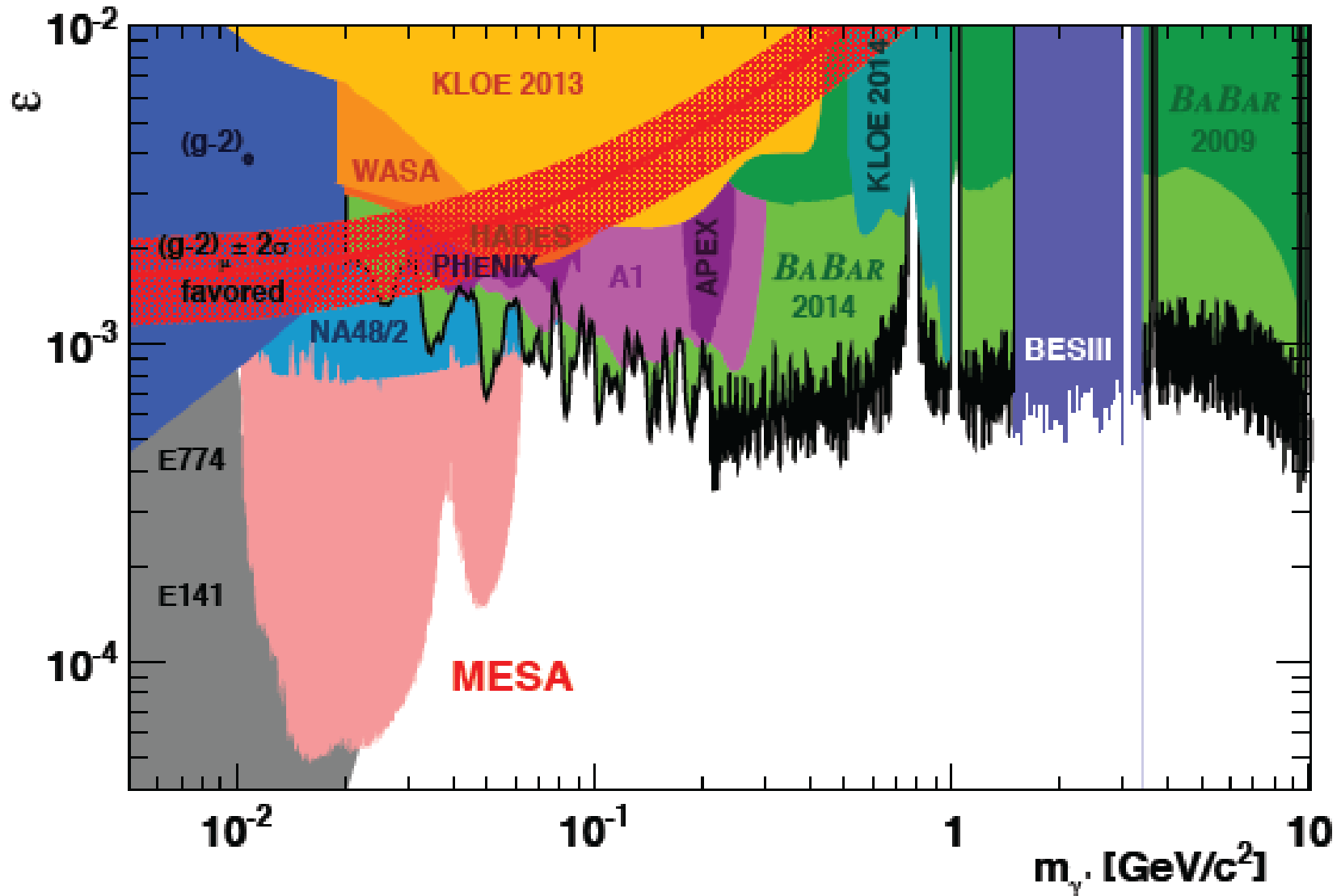
$$E = 1/10 E_{\text{MAMI}}$$

$$I = 10 \times I_{\text{MAMI}}$$



Potential of MESA w/ MAGIX after 6 months data taking

Karlsruhe Institute of Technology

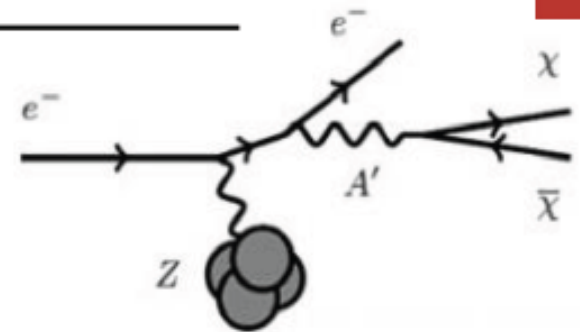


Invisible Decay of Dark Photon

So far Kinetic Mixing: $M(\gamma') < 2M(\chi)$

Consider now: $M(\chi) < 0.5M(\gamma')$

Dark Photon decays dominantly into
Light Dark Matter (LDM), which is not yet constraint experimentally!

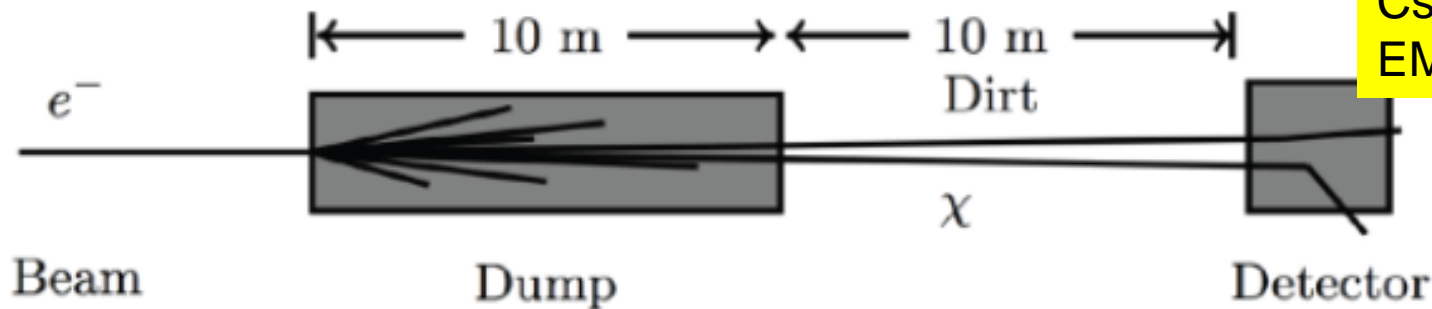


Beam Dump Experiment (BDX)

Electron Scattering on Beam Dump

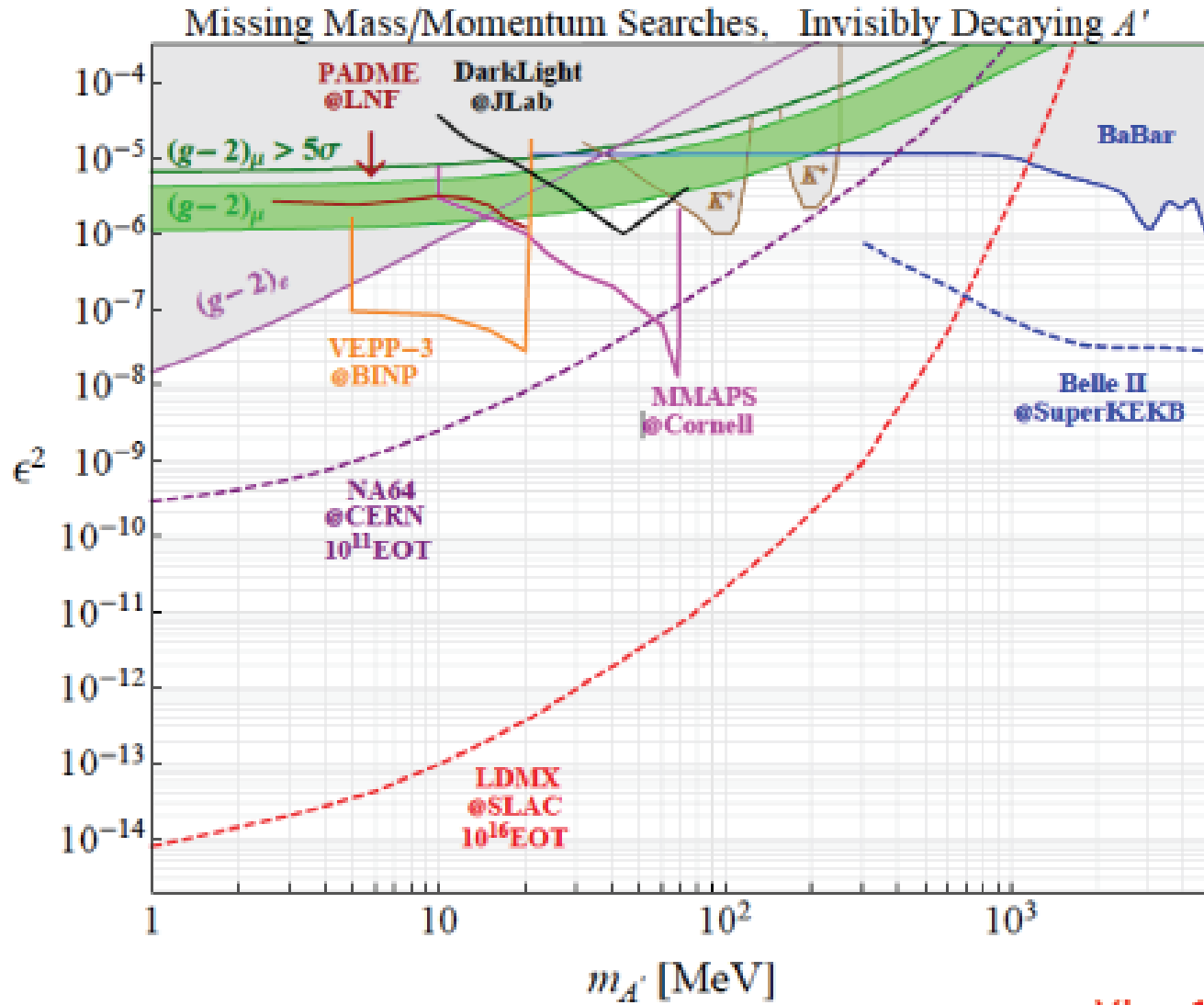
→ Collimated and boosed (!) pair of Dark Matter particles

1m³(1-5t)
CsI(Tl)
EM shower det.



JLAB, Mainz, SLAC Cornell

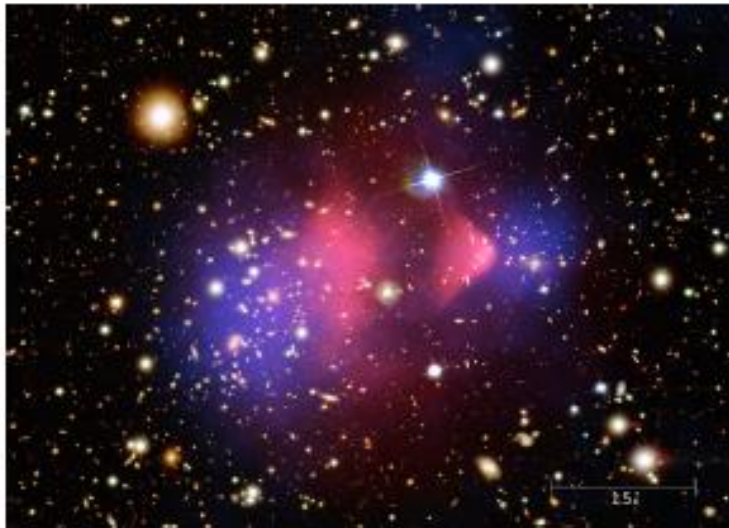
Future of Invisible Dark Photon Searches



arXiv: 1608.08632

DM models with two mediators.

How to save the WIMP



© NASA

Michael Duerr

MU Programmtag 2016

Mainz, 12 December 2016

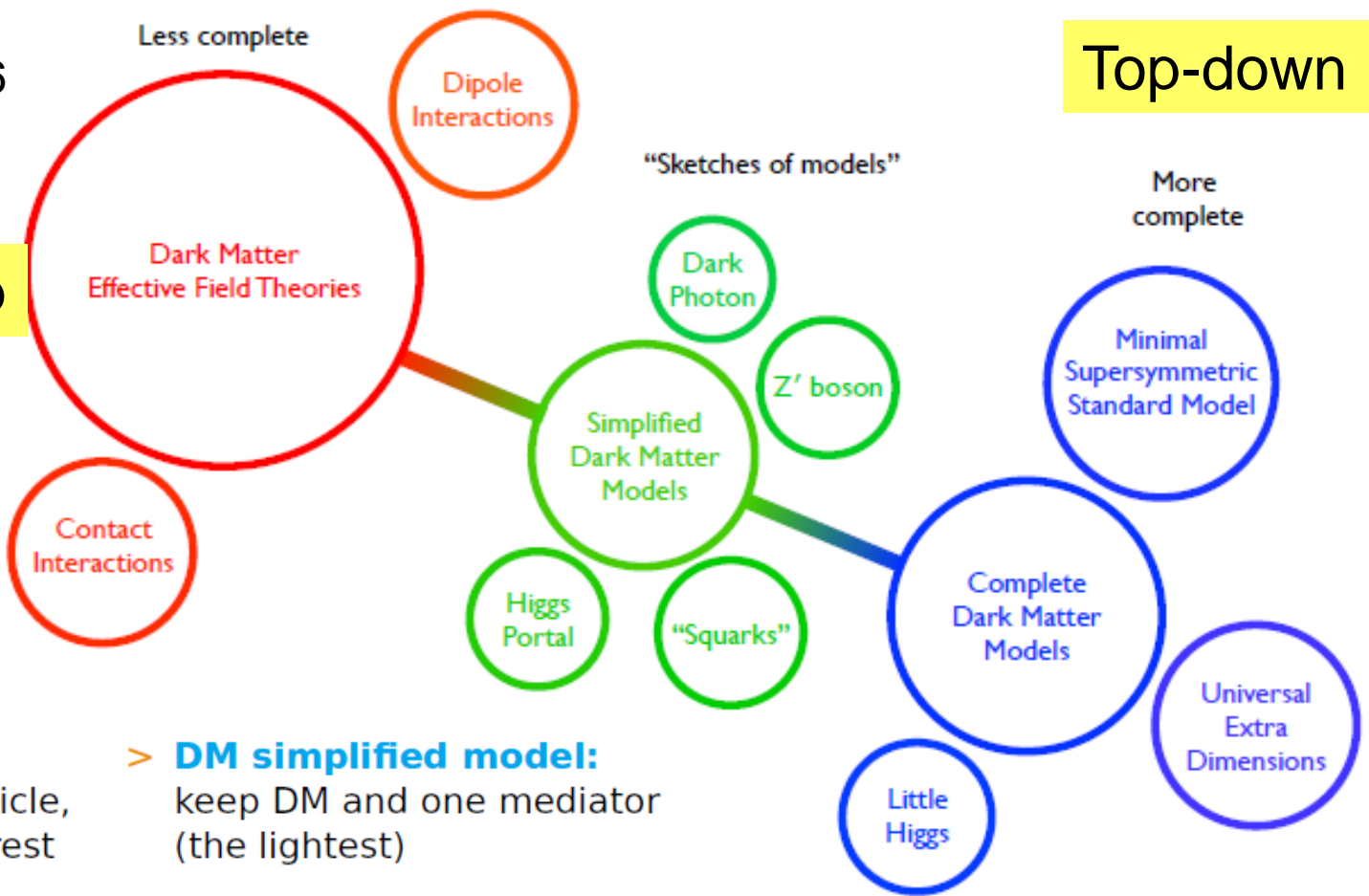
based on:

[arXiv:1304.0576](#), [arXiv:1309.3970](#),
[arXiv:1409.8165](#), [arXiv:1508.01425](#), and
[arXiv:1606.07609](#)

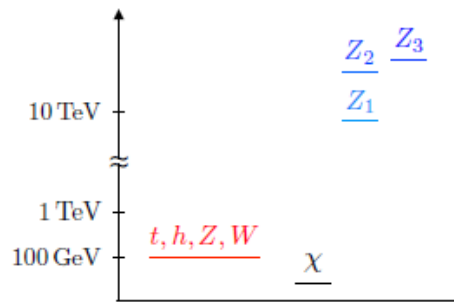
in collaboration with:

[P. Fileviez Pérez](#), [F. Kahlhoefer](#), [K. Schmidt-Hoberg](#),
[Th. Schwetz](#), [J. Smirnov](#), [S. Vogl](#), [M. B. Wise](#)

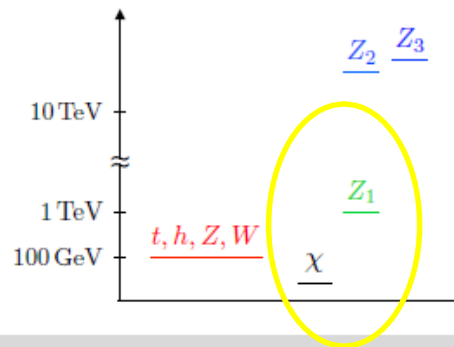
Bottom-up



> **DM EFT:**
only keep DM particle,
integrate out the rest



> **DM simplified model:**
keep DM and one mediator
(the lightest)



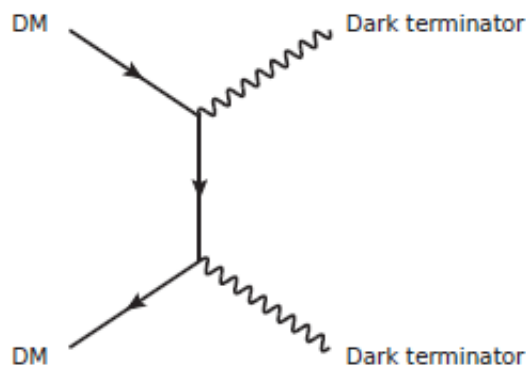
Dark matter model with two mediators.

> 6 independent parameters:

particle masses		coupling constants	
DM mass	m_χ	dark-sector coupling	g_χ or y_χ
Z' mass	$m_{Z'}$	quark-Z' coupling	g_q
dark Higgs mass	m_S	Higgs mixing angle	θ

> A combination of different simplified models:

	$g_q \gg \sin \theta$	$g_q \sim \sin \theta$	$\sin \theta \gg g_q$
$m_S \gg m_{Z'}$	Spin-1 mediator simplified model		Spin-0 mediator with spin-1 terminator
$m_{Z'} \sim m_S$		Two-mediator model	
$m_{Z'} \gg m_S$	Spin-1 mediator with spin-0 terminator		Spin-0 mediator simplified model



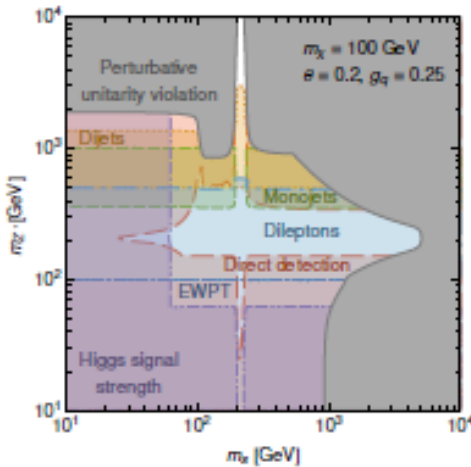
Dark terminator
new final state for
DM annihilation

Two mediators: results.

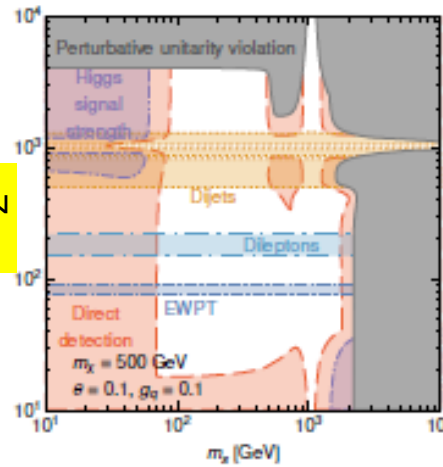
$M_\chi = 100 \text{ GeV}$

$M_\chi = 500 \text{ GeV}$

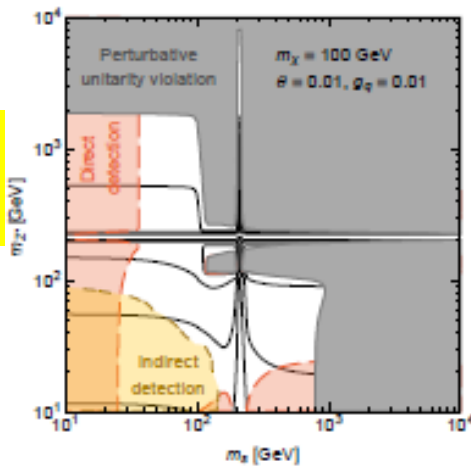
$\theta = 0.2,$
 $g_q = 0.25$



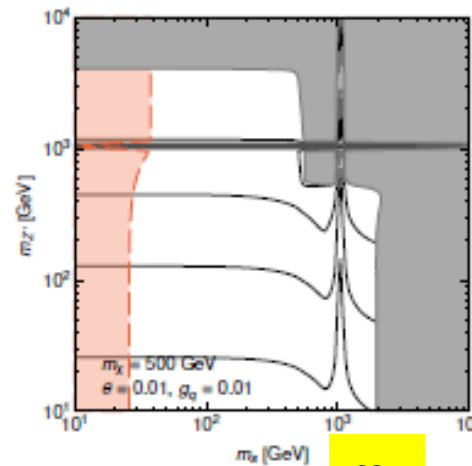
$m_{Z'}$



$\theta = 0.01,$
 $g_q = 0.01$



m_s



- > sizeable g_q and $\sin \theta$:
- > for $m_\chi = 100 \text{ GeV}$, only small regions close to the resonances remain viable
- > for $m_\chi = 500 \text{ GeV}$, larger regions are allowed because s or Z' can be terminators without being strongly constrained
- > secluded from the SM:
- > region with $m_{Z'}, m_s > m_\chi$ is tightly constrained because annihilations into SM final states cannot reproduce the relic abundance with perturbative couplings
- > for $m_{Z'}, m_s < m_\chi$, annihilation into dark terminators typically dominates
- > experimental constraints can be suppressed since g_q and θ can be small \rightarrow difficult to probe
- > for small masses, set-up can still be probed by indirect detection



Phenomenology of flavoured dark matter

MU Programtag 2016, Mainz

Monika Blanke, Simon Kast | Dez 12, 2016

KARLSRUHE INSTITUTE OF TECHNOLOGY



www.kit.edu

The Flavour Gate to Dark Matter

Assume an analogy to the SM fermions \rightarrow dark flavour triplet χ_i .

Flavoured Dark Matter coupling to SM right-handed up-quark triplet:

$$\mathcal{L}_{\text{NP,int}} = -\lambda_{ij} \bar{U}_{Ri} \chi_j \phi + h.c.$$

Dark Minimal Flavour Violation

[Agrawal, Blanke, Gemmler '14]

Flavour Symmetry

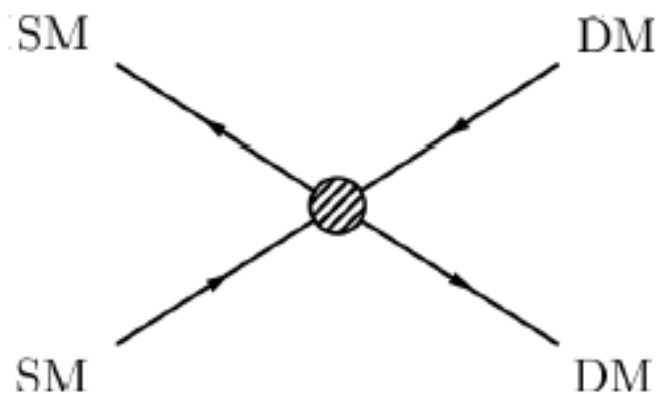
$$U(3)_u \times U(3)_d \times U(3)_q \times U(3)_\chi$$

is only broken by SM Yukawa couplings and the **DM-quark coupling** λ_{ij} (Dark Minimal Flavour Violation).

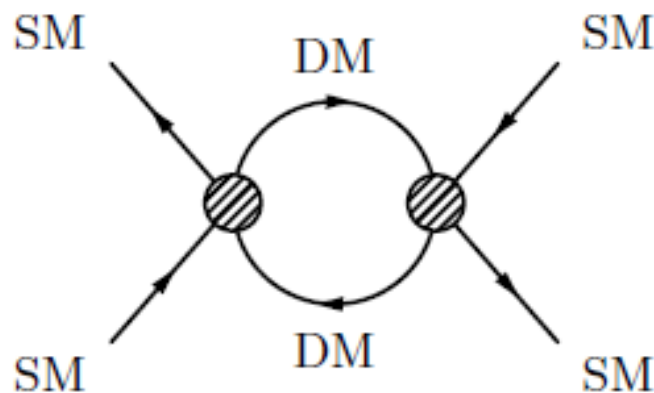
\Rightarrow only DM mass splitting comes from RG running:

$$m_{ij} = m_\chi (\mathbb{1} + \eta \lambda^\dagger \lambda + \dots)_{ij}.$$

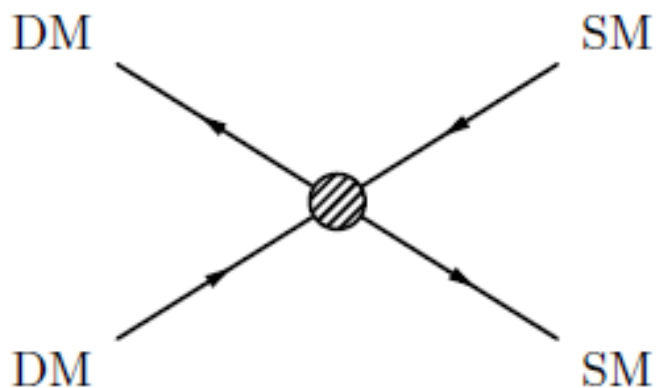
How to Detect Flavoured Dark Matter?



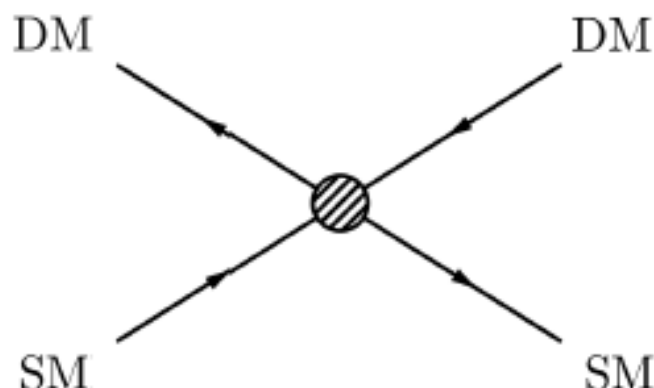
collider searches



precision flavour data



indirect detection / freeze-out



direct detection

DM Bounds from Direct Detection Experiments

[LUX collaboration '16]

- Many contributions to total WIMP-nucleon cross section, only Z-penguin with neutron is negative. \Rightarrow saves the day
- Tree level and neutron Z-penguin have to nearly cancel each other. \Rightarrow serious constraints on Θ_{13}
- For too large couplings the cancellation is no longer possible \rightarrow excluded.
- Top-flavoured DM is the natural choice.

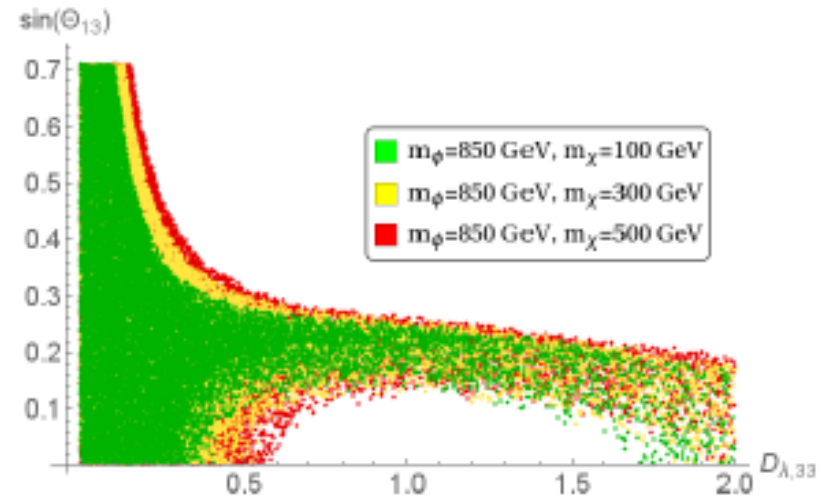


Figure : Valid mixing angle Θ_{13} vs $D_{\lambda,33}$.



Figure : Cancellation of tree-level and neutron Z-penguin contributions (symbolic).

Conclusion and Outlook

- All kinds of different constraints → multitude of effects and interesting interplay.
- Especially interesting effect on mixing angle θ_{13} due to DD and RA constraints.
 - ⇒ Future measurements of direct detection experiments will test a large part of the parameter space.
 - ⇒ Ongoing Xenon experiments or experiments with other noble gases well motivated.
- Simplified models are powerful tool to study diversity of constraints.
- Going beyond Minimal Flavour Violation is worth the effort.
 - Dark Minimal Flavour Violation as guidance.

Cross-Topic Nature of Dark Matter

Studying generalized dark matter interactions
with extended halo-independent methods

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Axion dark matter from topological defects

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FUNK dark photon search: status and perspectives

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A. Denig, JGU & HI Mainz

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Phenomenology of flavoured dark matter

S. Kast, KIT

- large variety of activities (theory & experiments)
- lively discussion, science-driven
- large attendance
- no PoF structural discussions, but
„MUTag very interesting exchange of ideas“